

SIGNATURES OF SUPERSYMMETRY IN
UNDERGROUND ν DETECTORS
FROM UHECR ?

→ ICECUBE?

DESY THEORY WORKSHOP '04

D. FARGION
WITH B. MELE
(IN PREPARATION)

Anindya Datta
INFN, Sezione di Roma
Università La Sapienza

- Observation of cosmic ray flux beyond GZK energies
AGASA ; Fly's Eye ...
One of the greatest puzzles in physics to be solved
- Many equally attractive models / proposals

- Top-down models
 - * primaries producing beyond GZK airshowers coming from
decays of some supermassive "X" particles - - -
- Important !!
- Decay life time & masses of these particles - - -

More on "X" particle decay.....

- Decay life time : comparable to the present age of universe .. .
- Masses of these heavy particles \sim M_{GUT} as they may be due to GUT symmetry breaking !!
- Density of "X" particles .. .
- Distance from our universe .. .
- as primaries of beyond GZK airshowers are mainly protons .. .
protons have energy loss length $\sim 50 \text{ Mpc}$ at these energies .. .

- Decays of supermassive "X" particles....
- need of a dynamics which governs the decays (cascades....)
- After long decay chains : end products $\rho, \gamma, \nu \dots$ stable particles..

Spectra : depends on "X" particle mass ...
Technique : well known; QCD fragmentation functions...)

For SM : Sarkar, Toldra Nuc. Phys B621

Supersymmetry driven "X" decays ?
 a well motivated extension of SM (do many important things....)
 also provides a candidate for cold dark matter : neutralino

A natural end product with ρ, ν, γ of "X" decay is thus neutralino ($\tilde{\chi}_1^0$)

Berezinsky, Kachelnikov PLB 422
 • Barbot & Drees PLB 533

Detection of χ^0 's in WHECR

$$\chi^0 + q \rightarrow \tilde{q}_{L,R} \rightarrow q \chi^0 \quad \text{Berezhinsky \& Kachelriess}$$

PLB 422 (1998)

from nucleons . . .

similar like $\nu N \rightarrow \nu X$, $\nu NC / DIS$ interaction . . .
final state hadrons have more energy . . .

- More recently in the context of EUSO . . . Anchordoqui, Goldberg, Nath
PRD 70 (2004)

production of rather heavy (~ 1 TeV) $\tilde{q}_{L,R}$ resonances . . .

→ to get rid of ν -background . . .

Then what we want to do ?

- Event rate at ICECUBE

$$\text{In general : } \tilde{\chi}^0 + e^- \rightarrow \tilde{e}_{L,R} \rightarrow l_f \rightarrow m_\tau = m_e \\ \tilde{\chi}^0 + q \rightarrow \tilde{q}_{L,R} \rightarrow l_f \rightarrow m_\tau = m_p/m_n$$

$$\text{But remember : } \sigma_{\text{peak}} = \frac{8\pi}{m_T^2} \left(\frac{m_T^2}{m_T^2 - m_{\tilde{\chi}^0}^2} \right)^2 \text{Br}(\tilde{e}, \tilde{q} \rightarrow e, q + \tilde{\chi}^0) \text{Br}(\tilde{e}, \tilde{q} \rightarrow f)$$

We stick to mSUGRA

$\text{Br}(\tilde{e}_L, \tilde{q}_L \rightarrow \tilde{\chi}^0 + e, q)$ is considerably small compared to \tilde{e}_R, \tilde{q}_R

$\text{Br}(\tilde{e}_R \rightarrow e \tilde{\chi}^0) \approx 1$ always mass ordering of \tilde{e}_R / nature of $\tilde{\chi}^0$

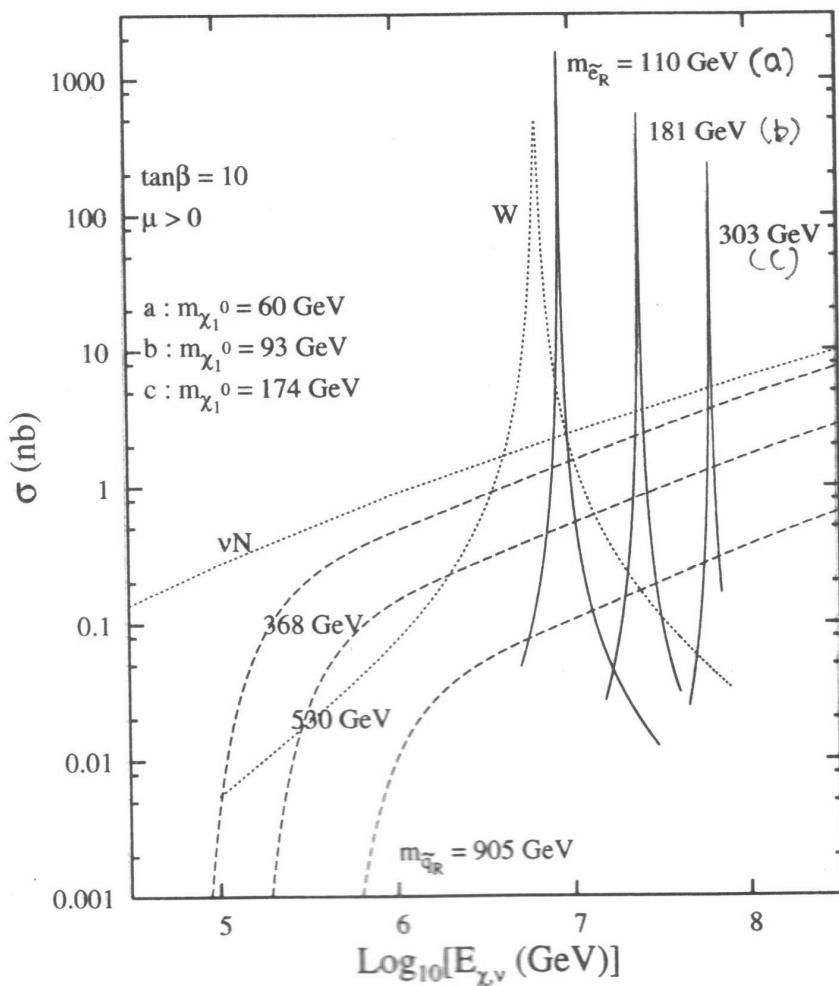
$\text{Br}(\tilde{q}_R \rightarrow q \chi^0) \approx 1$ $m_{\tilde{q}} < m_{\tilde{g}}$ (not a very unrealistic choice)

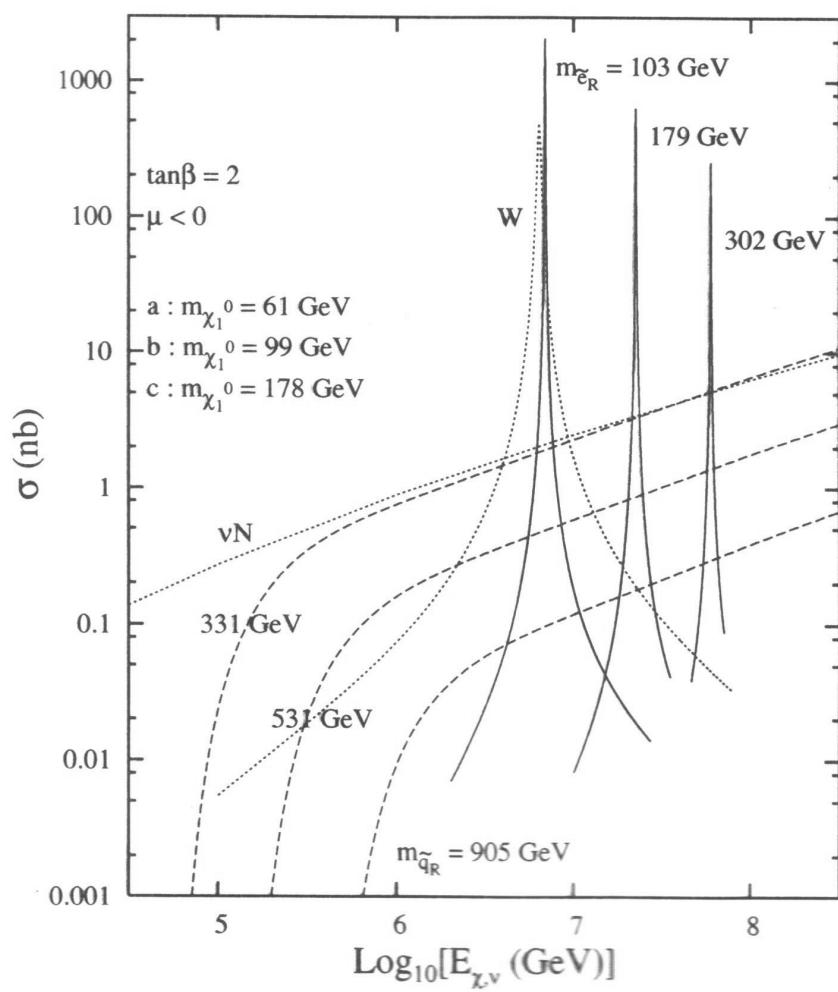
So our estimate of # of events are conservative side

$$E_x^{\text{peak}} = \frac{m_T^2 - m_{\tilde{\chi}^0}^2}{2m_T}$$

Let us look at some results...

- Height of the peaks are not only important qty.
- $\bar{\nu}_e + e^- \rightarrow W^- \rightarrow \text{visible}$ } backgrounds - - -
 $\nu + N \rightarrow \nu, l + X$ }
- $\sigma_{\nu} (\text{NC} + \text{CC}) = 8 \times 10^{-3} \text{ nb} \left(\frac{E_{\nu}}{1 \text{ GeV}} \right)^{0.36}$ Gandhi et al.





One caveat concerning $\tilde{\chi}^0$ -flux ...

No existing experimental bound on $\tilde{\chi}^0$ -flux ...

Only existing bound on UHE- ν flux ...

assuming : $\sigma(\nu N) \sim \sigma(x^0 N)$
Limits on $\phi_\nu \rightarrow \phi_{x^0}$

Apart from that :

ν -flux (e.g. Drees) : Highly model dependent ...

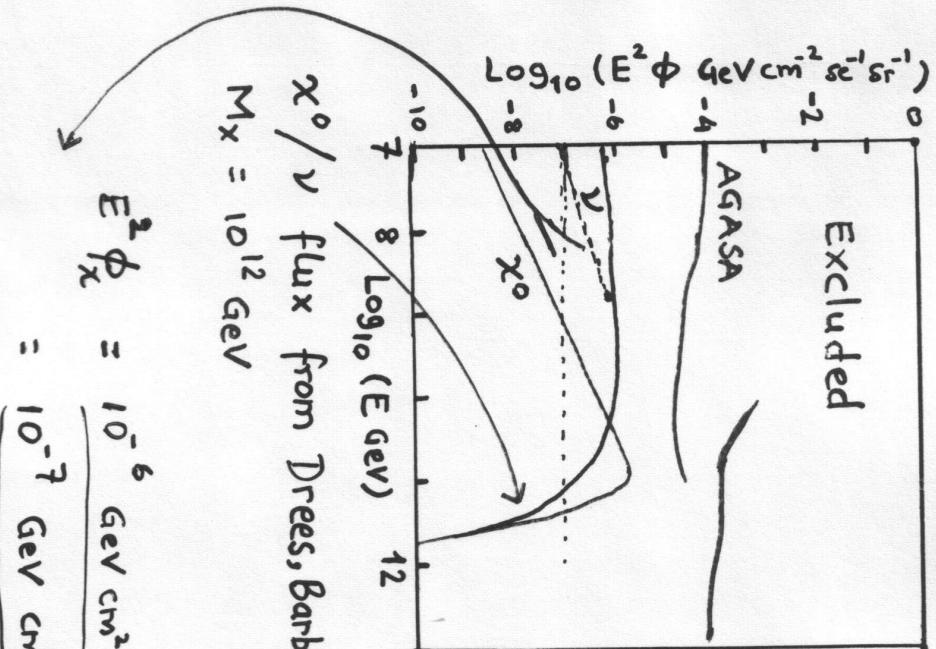
e.g. assumes no relevant physics between $m_X \nmid 1\text{TeV}$

χ^0/ν flux from Drees, Barbat

$$M_X = 10^{12} \text{ GeV}$$

$$\begin{aligned} E^2 \phi_x &= \frac{10^{-6} \text{ GeV cm}^2 \text{ sec}^{-1} \text{ sr}^{-1}}{\left(\frac{Ex}{10^4 \text{ GeV}} \right)^{0.5}} \\ &= 10^{-7} \text{ GeV cm}^2 \text{ sec}^{-1} \text{ sr}^{-1} \end{aligned}$$

We do not strictly follow
Drees flux ...



- Number of events

$(\text{km}^3 \text{ ice} / \text{year})$

$$\bar{N}_e = N_0 \sigma_{\text{peak}}^{e_R} \Gamma \left(\frac{m_{\tilde{\chi}}}{2m_e} \right) E_{\text{peak}}^{-\beta}$$

$$N_{\tilde{q}} = \tilde{N}_0 \int_{E_{\text{min}}} \sigma^{q_R}(E) E_x^{-\beta} dE_x$$

$m_{\tilde{\chi}}$	$m_{\tilde{q}}$	m_0	\bar{N}_e	$N_{\tilde{q}}$
103	331	61	23	4
			19	2.3
179	531	99	2.3	0.6
			3.4	0.5
302	905	178	0.3	0.07
			0.8	0.1
260	432	69	0.28	1.45
			0.7	1
363	611	100	0.08	0.37
			0.26	0.33

$$\lambda = 10^{-7} \text{ GeV cm}^2 \text{ sec}^{-1} \text{ sr}^{-1}$$

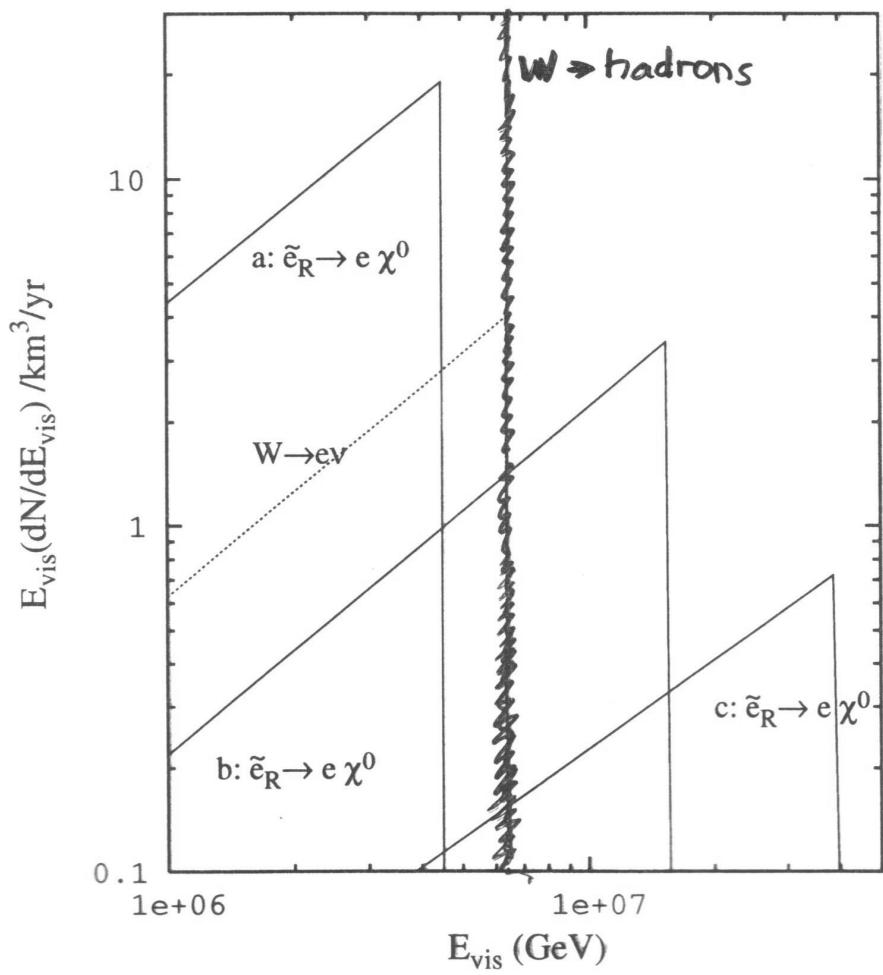
$$\phi_x = 10^{-6} \text{ GeV cm}^2 \text{ sec}^{-1} \text{ sr}^{-1} \left(\frac{E_x}{10^3 \text{ GeV}} \right)^{0.5}$$

$m_{\tilde{\chi}}$	$m_{\tilde{q}}$	m_0	\bar{N}_e	$N_{\tilde{q}}$
302	905	178	0.3	0.07
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$\tan \beta = 2$ (not very important!)

Backgrounds from W^- : 30 (24) hadronic channel 5 (4) in electronic channel...
from $V_N (cc + NC)$: 70 (50) events $E_x > 10^{5.5} \text{ GeV}$

How the visible spectra look like . . .
typical two-body decay . . .



Beyond GZK : A possible $\tilde{\nu}$ -burst?

- Similar to χ -burst :
but $\chi^0 \nu_r \rightarrow \tilde{\nu} \rightarrow$ visible particles

Possible hadronic decays of $\tilde{\nu}$

$$\begin{array}{c} \tilde{\nu} \rightarrow e \tilde{\chi}^\pm \\ \quad \quad \quad \downarrow \\ \tilde{\chi}^\pm \rightarrow q\bar{q}'\chi^0 \end{array}$$

$$\begin{array}{c} \tilde{\nu} \rightarrow \nu \tilde{\chi}_2^0 \\ \quad \quad \quad \downarrow \\ \tilde{\chi}_2^0 \rightarrow q\bar{q} \text{ hadrons} \end{array}$$

Possibly less energetic than hadrons from Z -decay....
(depends on masses of the particles in cascade)

Z -burst vis-a-vis $\tilde{\nu}$ -burst ...

- Crucial $m_{\tilde{\nu}} = 0.1 \text{ eV}$

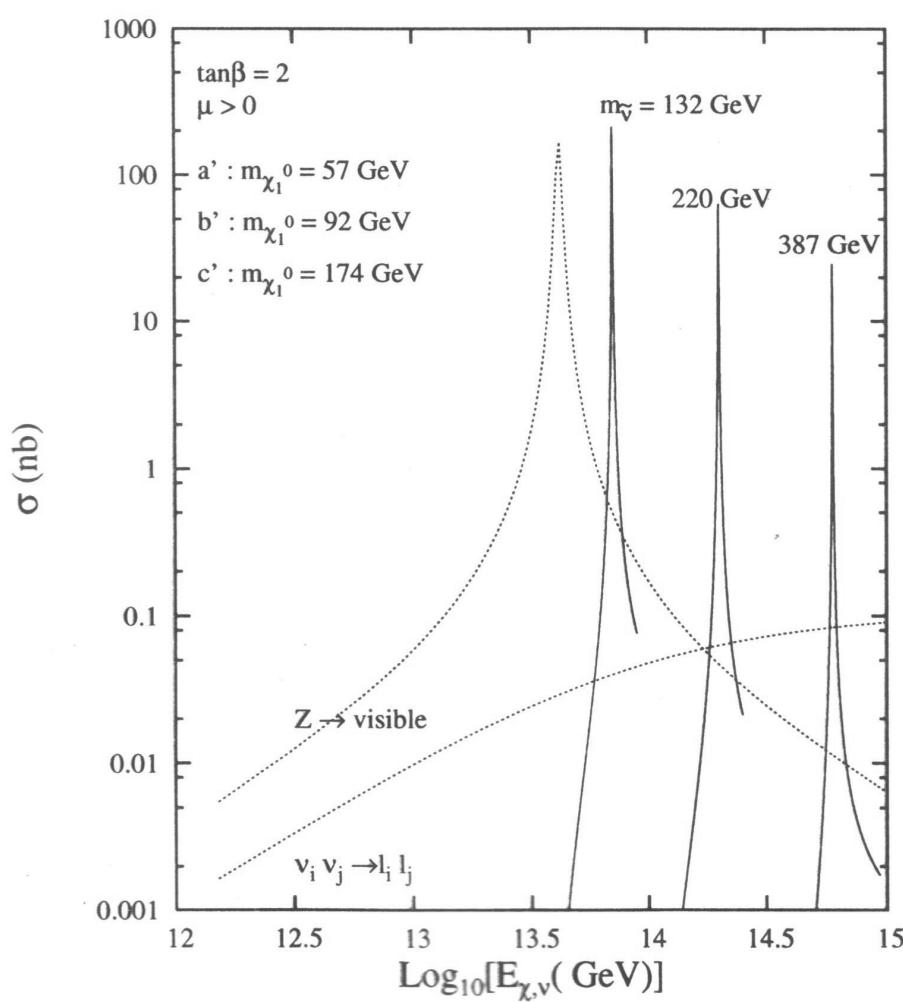
$m_{\tilde{\nu}}$	m_{χ^0}	σ^{peak}	$\Gamma_{\tilde{\nu}}$	B^{vis}
132	57	212 nb	0.37 GeV	0.46
220	92	64	0.60	0.47
387	174	25	0.59	0.42
269	170	22	2.53	0.72
383	93	11	3.29	0.70

For Z -decay

$$\sigma^{\text{peak}} = 200 \text{ nb}$$

$$\Gamma_Z = 2.5 \text{ GeV}$$

$$B^{\text{vis}} = 0.80$$



$$E_{\chi}^{\text{peak}} = \frac{m_{\tilde{\nu}}^2 - m_{\chi^0}^2}{2m_{\nu}}$$

$$E_{\nu}^Z = \frac{m_Z^2}{2m_{\nu}}$$

To summarise.....

- UHE neutralinos are natural end products of superheavy particle decay models of UHECR beyond GZK-energies....
- We have calculated events rate corresponding to \tilde{e}_R and \tilde{q}_R production at resonance from UHE- $\tilde{\chi}_0^0$ interaction in underground ν -detector like ICECUBE
- Events rates are small unless the \tilde{e}_R -masses are close to their present expt. limits....
(only shadows of SUSY ?)
- UHE- $\tilde{\chi}^0$ can produce $\tilde{\nu}$ -resonances interacting with relic- ν ... a similar process like Z-burst..
- Again $\tilde{\nu}$ -burst events are small compared to Z-burst, producing less energetic visible particles on the average....