

Absorption Dips - Detection of the Cosmic Neutrino Background

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see also: Relic Neutrino Absorption Spectroscopy, hep-ph/0401203
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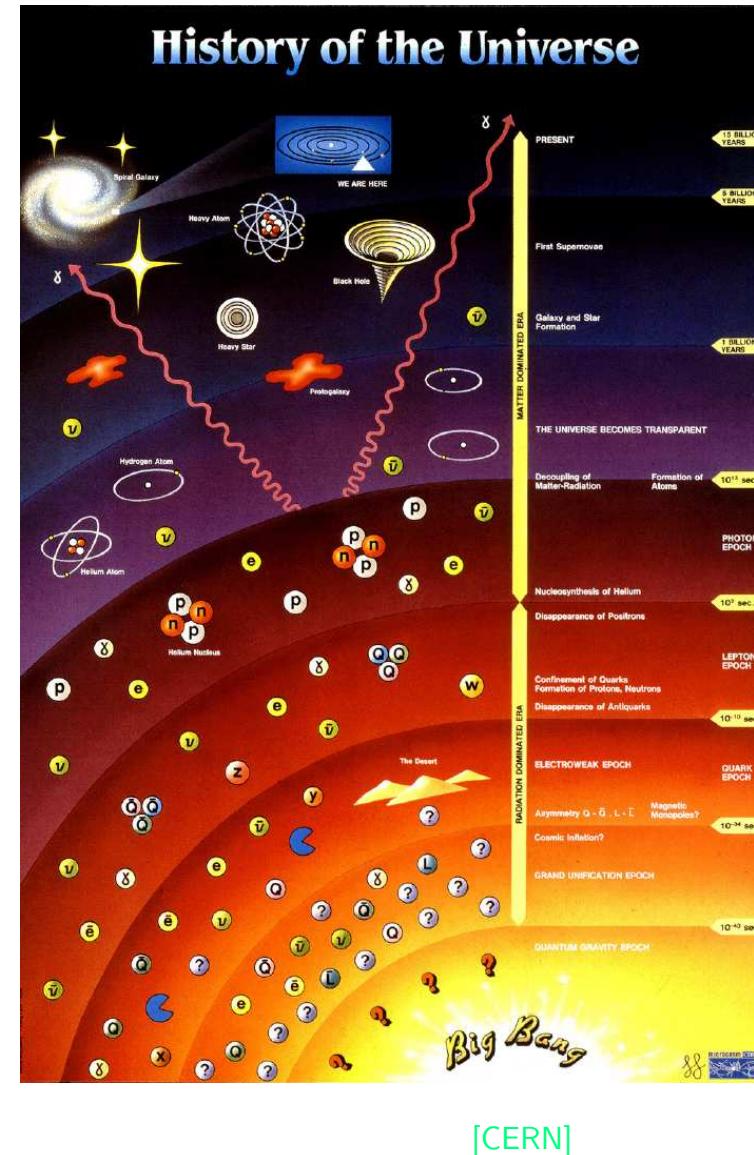
Big Bang - the relic ν story

- Big Bang: all particles in thermal equilibrium
- $\Gamma_{\text{INT.}} < \Gamma_{\text{EXP}}$ → decoupling; for **Cosmic neutrino Background (CνB)** at $t \sim 1$ s
→ CνB gives glimpse to **early universe!**
- since decoupling, wavelength of relic ν 's expanding in proportion to size of Universe
- subsequent e^+e^- annihilation reheats the photon gas, such that $(T_\nu/T_\gamma)^3 = 4/11$.

thus we can deduce for the **CνB**:

- **density** $\langle n_{\nu_i} \rangle_0 = \frac{3}{22} \langle n_\gamma \rangle_0 \simeq 56 \text{ cm}^{-3}$
- **momentum**

$$\langle |\vec{p}_{\nu_i}| \rangle_0 \simeq 3 \left(\frac{4}{11} \right)^{1/3} T_{\gamma 0} \simeq 5 \cdot 10^{-4} \text{ eV}$$



[CERN]

Well predicted - but no direct evidence!

C ν B - why so difficult to detect?

- ν 's interact only weakly - for $s \ll M_Z^2, M_W^2$: **small cross section**:

$$\sigma_{\nu_i}^{\text{NC,CC}} \sim G_F^2 s \simeq 10^{-55} \text{ cm}^2 (s/\text{eV}^2)$$

- ν 's have mass - $m_{\nu_3} > \langle |\vec{p}_\nu| \rangle \simeq 5 \cdot 10^{-4} \text{ eV}$, leading to a **small velocity**.
- ν mass is small - $m_{\nu_3} < 1 \text{ eV}$, **no substantial clustering** $n(r) \simeq \langle n_\nu \rangle_0$

Event rate for incoherent scattering

$$R_{\nu_i}^{\text{ic}} = V n_{\text{T/b}} \langle n_{\nu_i} \rangle_0 \langle |\vec{v}| \rangle \sigma_{\nu_i}^{\text{NC,CC}}$$

- ▷ detecting **C ν B flux** - incoherent scattering - kton targets

$$R_{\nu_i}^{\text{ic}} = N_T \langle |\vec{v}_{\nu_i}| \rangle_0 \langle n_{\nu_i} \rangle_0 \sigma_{\nu_i} \sim 3 \cdot 10^{-8} \text{ yr}^{-1} \left(\frac{m_{\nu_i}}{0.1 \text{ eV}} \right) \left(\frac{N_T}{10^{33}} \right)$$

- ▷ using **C ν B** as a **target** - collider experiment - TESLA, LHC, VLHC

$$R_{\nu_i \text{ beam}} \sim 10^{-11} \text{ yr}^{-1} \left(\frac{I}{\text{Ampere}} \right) \left(\frac{L}{10 \text{ km}} \right) \left(\frac{m_{\nu_i}}{0.1 \text{ eV}} \right) \left(\frac{E_{\text{beam}}}{1 \text{ TeV}} \right) \left(\frac{\langle n_{\nu_i} \rangle_0}{56 \text{ cm}^3} \right)$$

Mission impossible? ⇒ **ways out**:

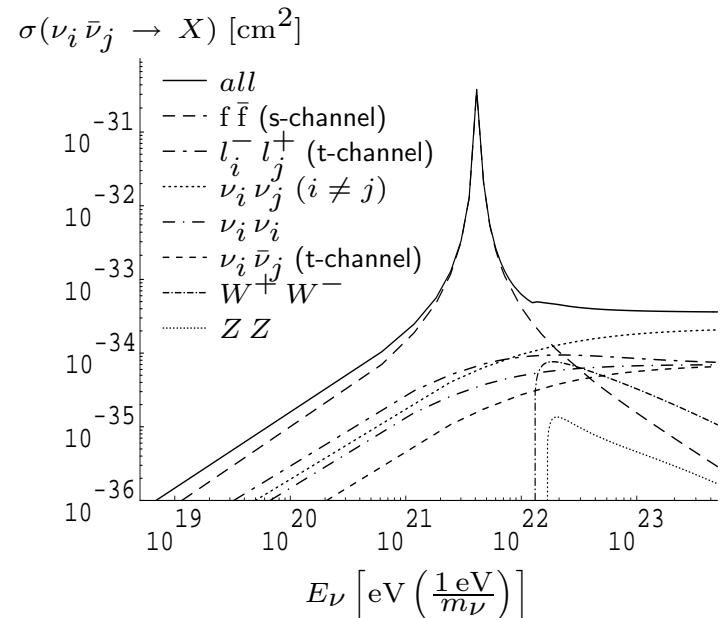
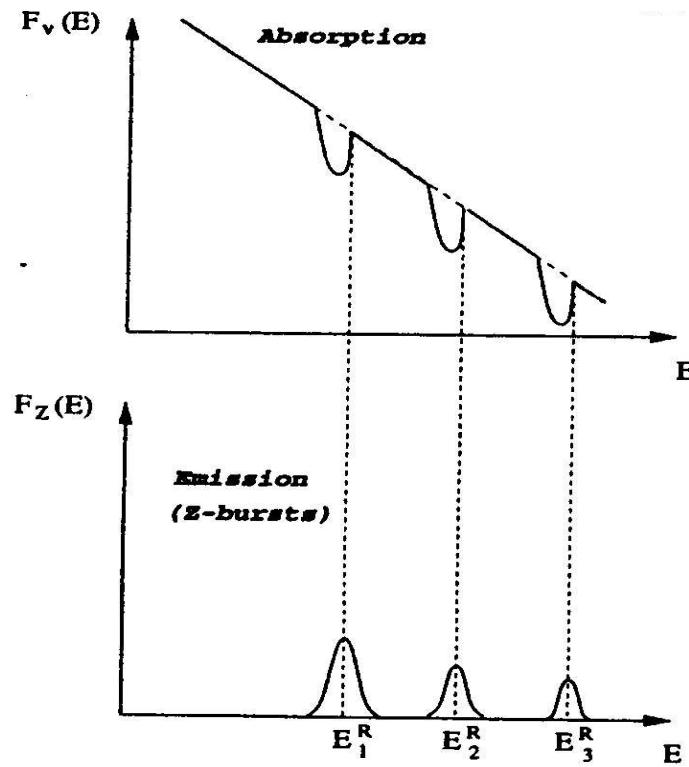
- **coherent scattering** - then: $R_\nu^{\text{coh}} \propto N_T^2$
- **C ν B as target** for **E**xtreme **E**nergy **C**osmic **R**ays (EECR), $E > 10^{19} \text{ eV}$

EEC ν scattering on the C ν B

The annihilation of a cosmic ray ν with a relic $\bar{\nu}$ (or vice versa) on the Z -resonance

⇒ **large cross-section at resonant energies**

$$\begin{aligned} E_{\nu_i}^{\text{res}} &= M_Z^2 / (2 m_{\nu_i}) \\ &= 4.2 \cdot 10^{21} \text{ eV} (1 \text{ eV}/m_{\nu_i}) \end{aligned}$$



- **Absorption features (dips)** in the extreme energy cosmic neutrino spectrum at the **resonant energies**

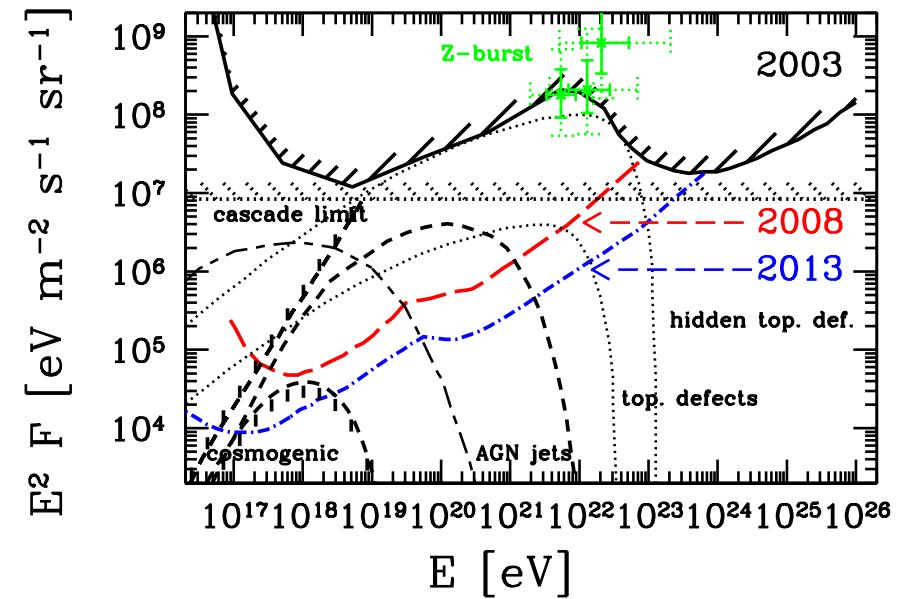
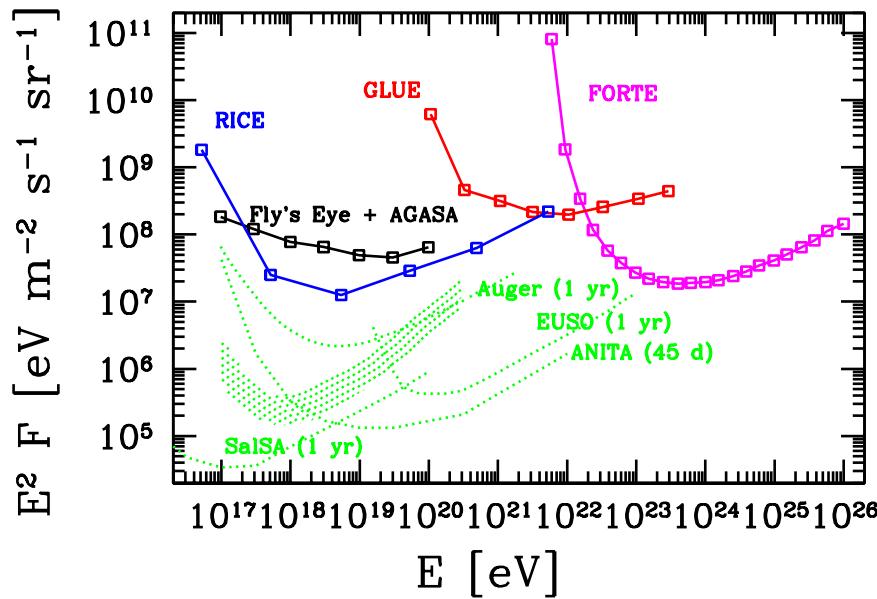
[Weiler '82; '84; Roulet '93; ...]

- **Emission features (Z-bursts)** as protons (or photons) above the Greisen-Zatsepin-Kuzmin-cutoff

[Fargion,Mele,Salis '99; Weiler '99; ... ; Fodor,Katz,Ringwald '01; '02]

At present time - best opportunities for absorption dips

- Precise best fits of **cosmological values** determining the expansion rate
- ν_3 -mass known to lie within an order of magnitude range
$$0.04 \text{ eV} \leq m_{\nu_3} \lesssim 0.4 \text{ eV} \quad \leftrightarrow \quad 1 \times 10^{22} \text{ eV} \lesssim E_{\nu_3}^{\text{res}} \leq 1 \times 10^{23} \text{ eV}$$
- **Experiments:** next decade - building phase for a lot of EEC ν detectors!



Outline of the talk

- Introduction
- Mathematical preparation:
formulas, simplifications and parameterizations
- Results
- Conclusions

Flux formulae I

$$F_{\nu_\alpha}(E) \equiv \frac{d^4 N_{\nu_\alpha}}{dE dA dt d\Omega} =$$

$$\frac{1}{4\pi} \int_0^\infty dE_i \int_0^\infty \underbrace{\frac{dz}{(1+z) H(z)}}_{dr} \sum_\beta \frac{-\partial P_{\nu_\alpha|\nu_\beta}(E; E_i, z)}{\partial E} \mathcal{L}_{\nu_\beta}(z, E_i)$$

EEC ν source emissivity distribution:

$$\mathcal{L}_{\nu_\beta}(z, E_i) \equiv \frac{d^3 N_{\nu_\beta}}{dV_{\text{comov.}} dt dE_i} \quad \text{injected at } z \text{ from Earth with energy } E_i$$

Propagation function:

$P_{\nu_\alpha|\nu_\beta}(E; E_i, z) \equiv$ expected N_{ν_α} above E if one ν_β started at z with E_i

Evolving Hubble parameter:

$$H(z) = H_0 \left[\Omega_M (1+z)^3 + \Omega_k (1+z)^2 + \Omega_\Lambda \right]^{1/2}$$

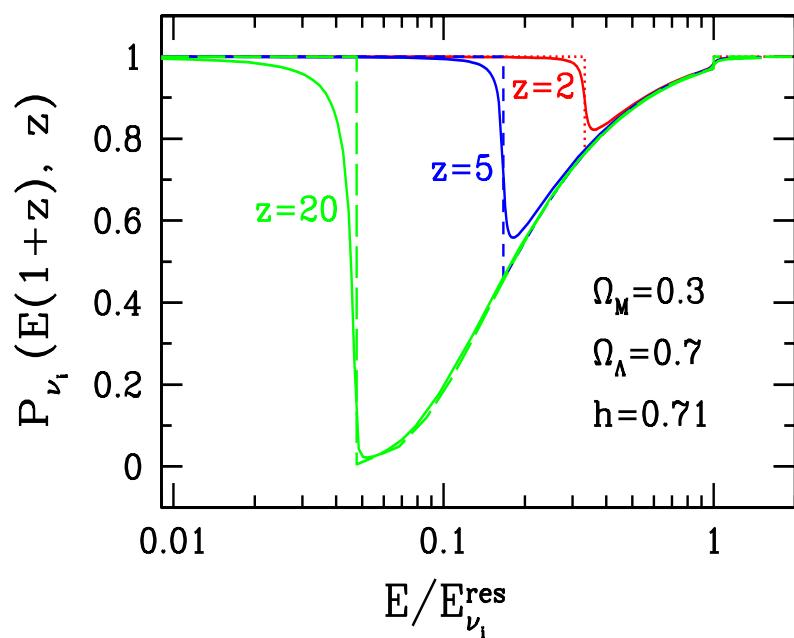
Assumptions and simplifications

- Cross section **Z-resonance dominated**
- **Interaction = absorption**
→ energy loss only due to expansion, $E_i = E(1 + z)$
- **Cosmological parameters:** $\Omega_M = 0.3$, $\Omega_\Lambda = 0.7$
- \sum of flux detected $\rightarrow \sum$ **leptonic mixing matrices** equal to **unity**

Flux formulae II - now simple

$$\sum_{\alpha} F_{\nu_{\alpha}}(E) \simeq \frac{1}{4\pi} \int_0^{\infty} \frac{dz}{H(z)} \frac{1}{3} \mathcal{L}_{\nu}^{\text{tot}}(z, E(1+z)) \times \sum_{j=1}^3 P_{\nu_j}(E(1+z), z)$$

with



$$P_{\nu_j}(E(1+z), z) \simeq \exp \left[- \left(\frac{\langle n_{\nu} \rangle_0 \langle \sigma_{\nu\bar{\nu}}^{\text{ann}} \rangle}{H_0} \right) \frac{\left(\frac{E_{\nu_j}^{\text{res}}}{E} \right)^3}{\left[\Omega_M \left(\frac{E_{\nu_j}^{\text{res}}}{E} \right)^3 + \Omega_{\Lambda} \right]^{1/2}} \right]$$

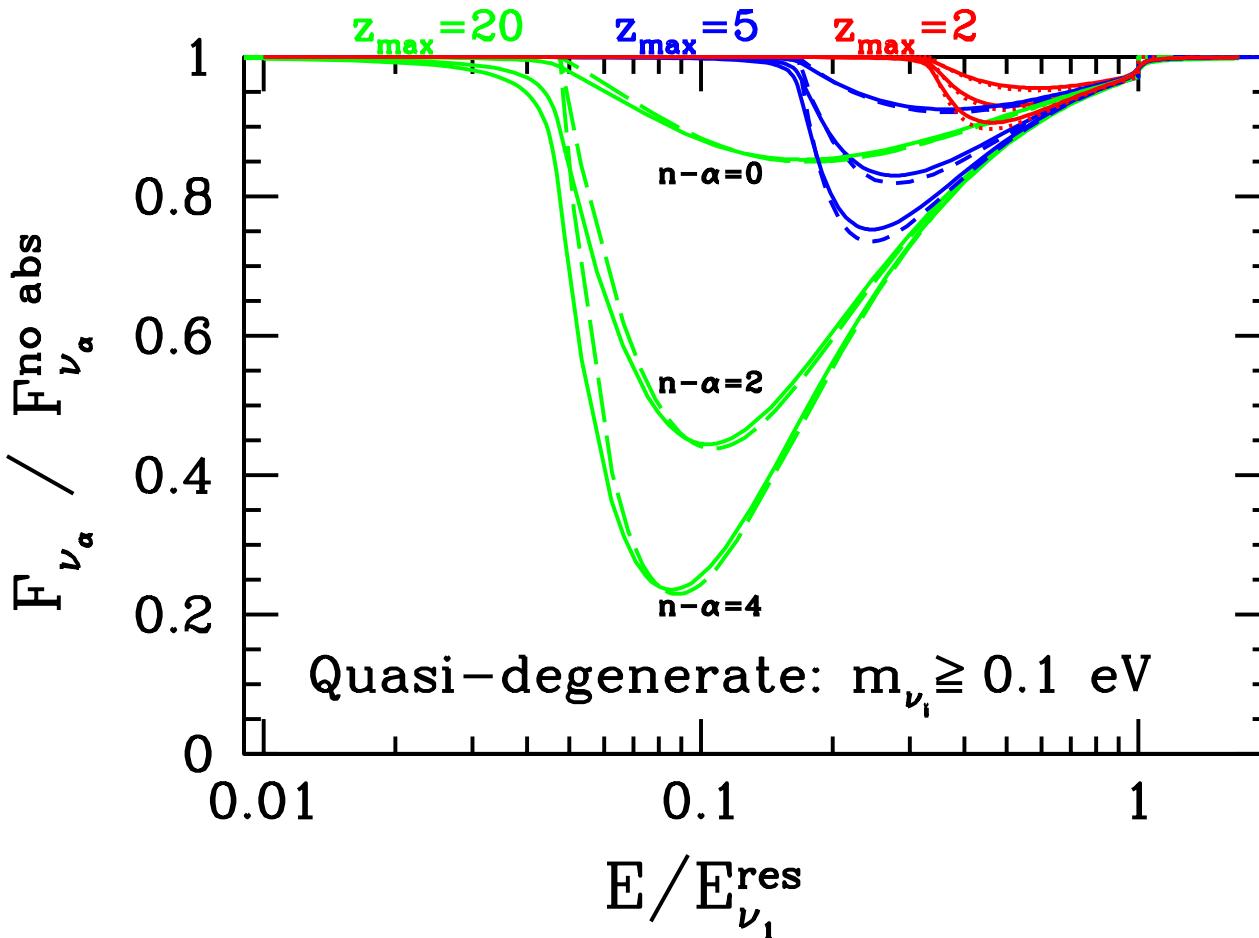
for $\frac{1}{1+z} < \frac{E}{E_{\nu_j}^{\text{res}}} < 1$,

1 otherwise

Open parameters

- **Source emissivity parameters**
- **ν -mass spectrum**

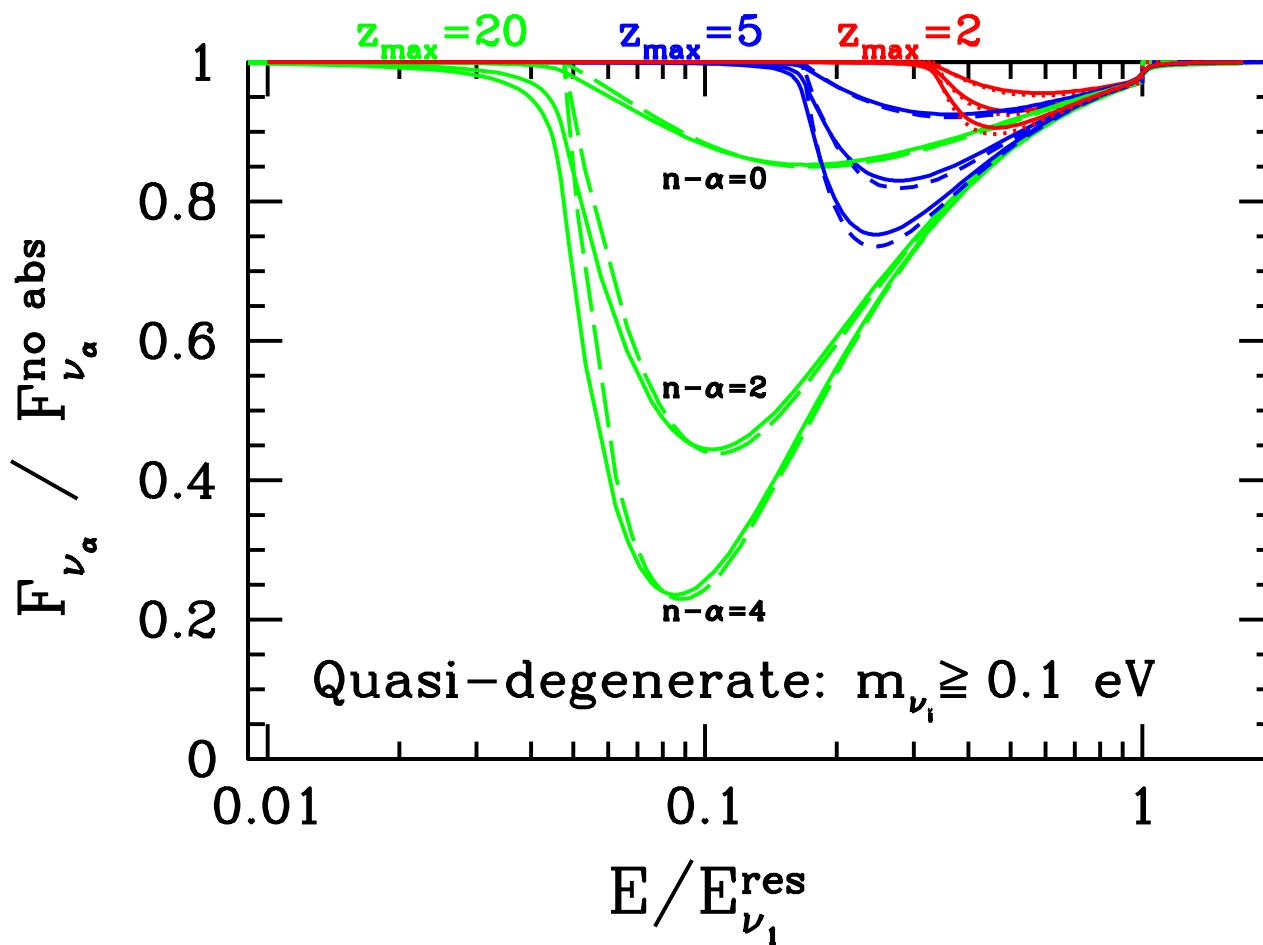
Case study - source emissivity



General features
of the dip:

- $\frac{\langle n_\nu \rangle_0 \langle \sigma_{\nu\bar{\nu}}^{\text{ann}} \rangle}{H_0} \sim 3.0 \%$
- **Depth:**
large source evolution
& flat energy spectrum
- **Magnification:**
starts at early times
& high cutoff energy

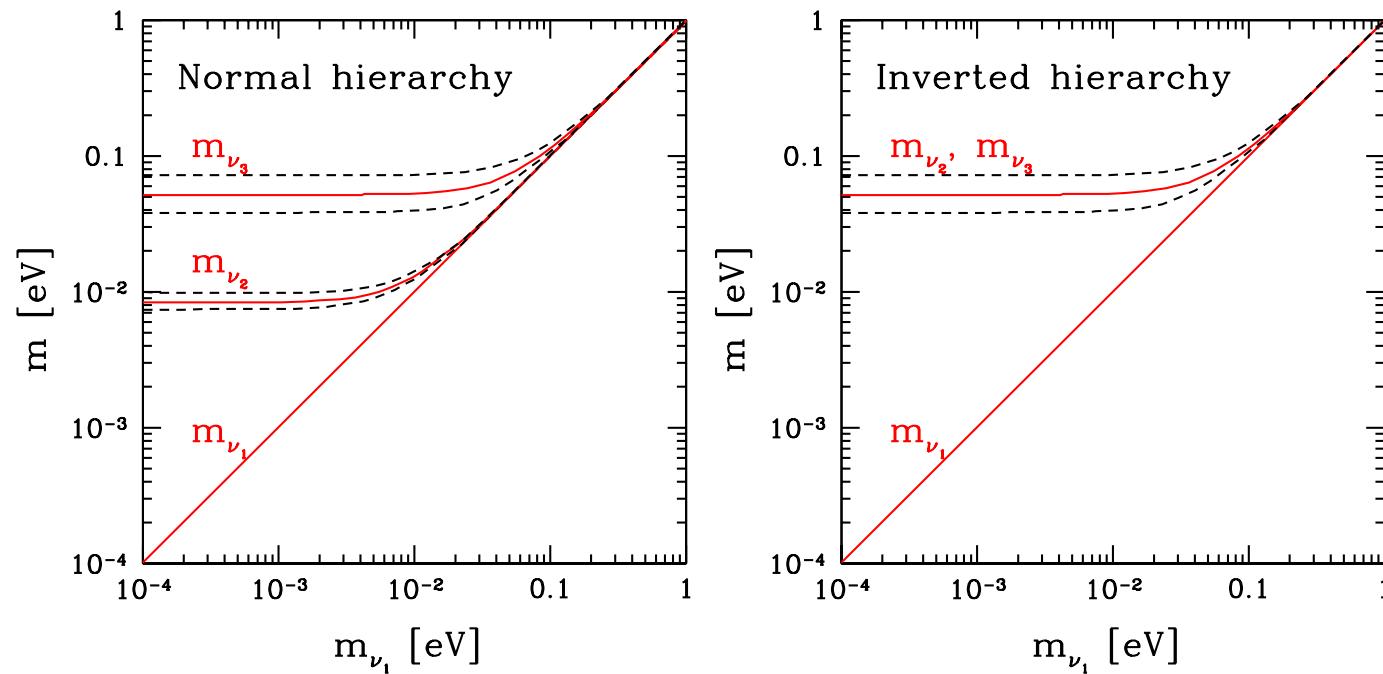
Open parameters - source emissivity: what is realized in Nature?



- **topological defects:**
 - low source evolution
 - starts at early times
 - cutoff energy large

- **astrophysical sources:**
 - large source evolution
 - starts at redshift time $z_{\max} = 2 \div 5$
 - cutoff energy low

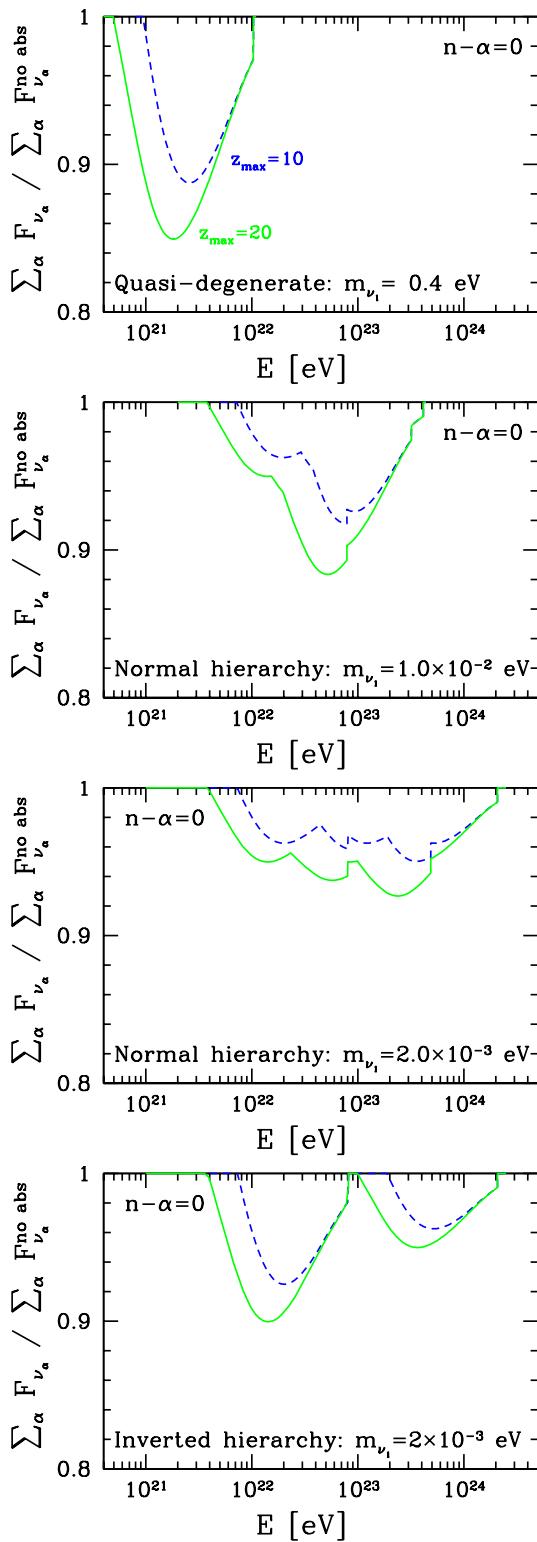
Open parameters - neutrino masses



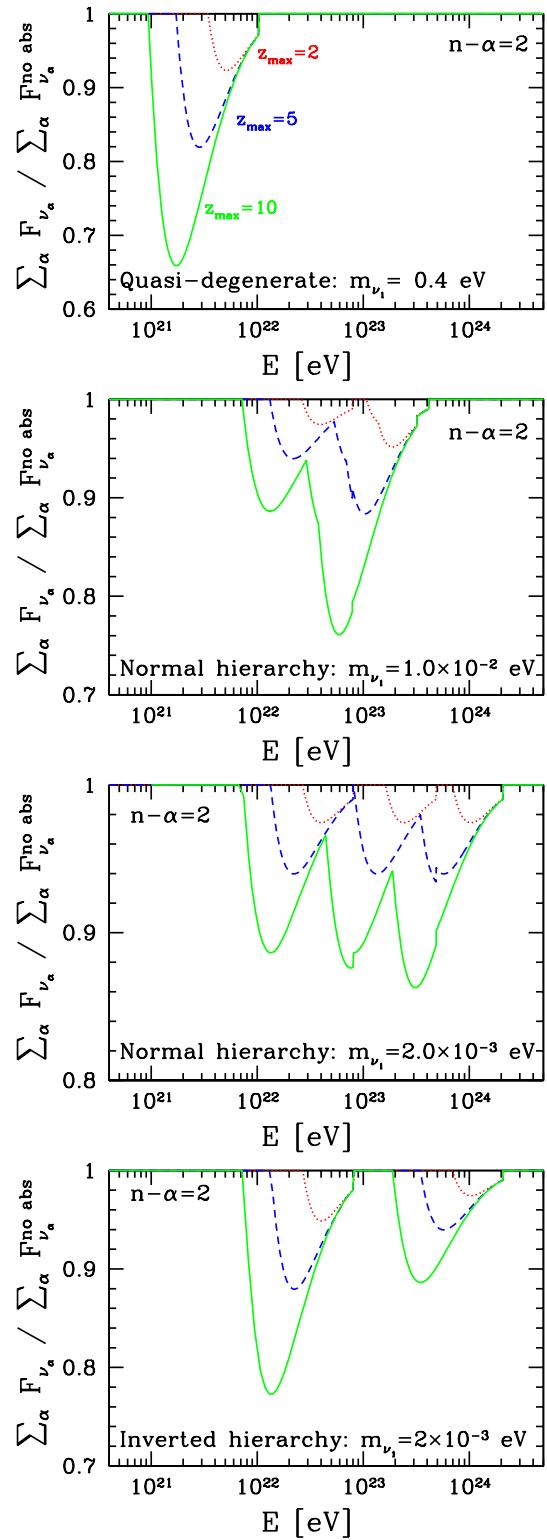
Allowed ranges for three neutrino masses in normal and inverted hierarchy

Case studies - different ν mass patterns

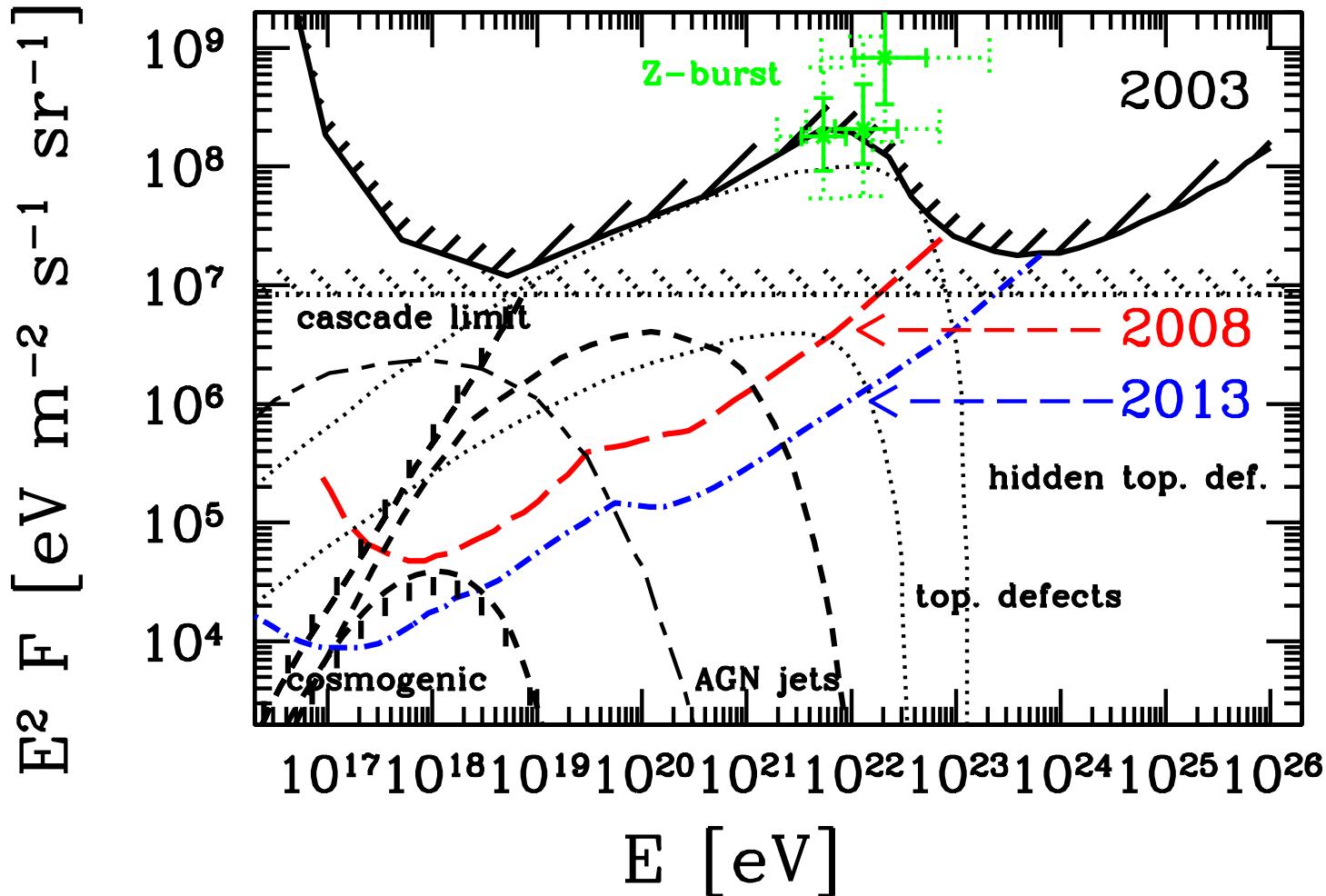
topol. defect source



acceleration source



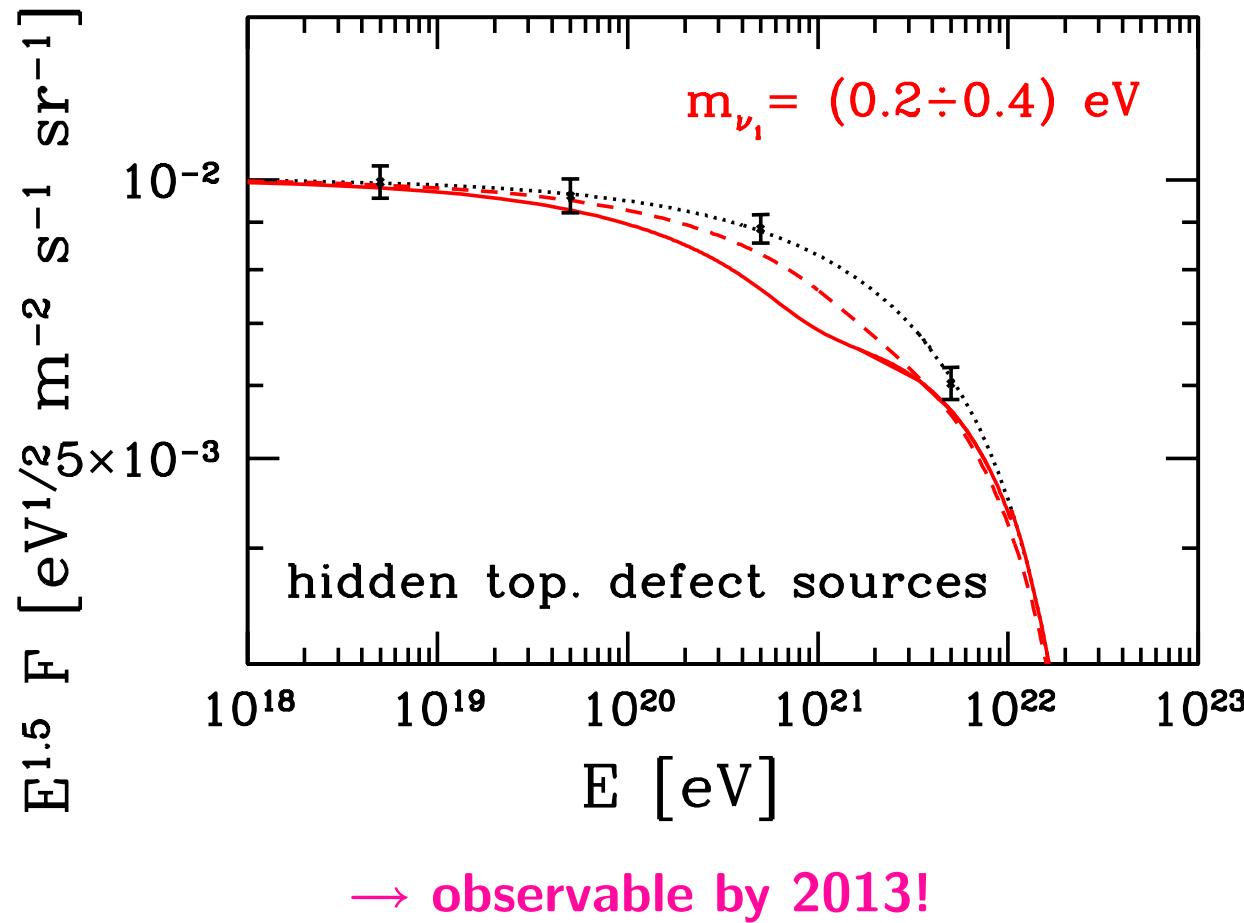
Open parameters - absolute scale of source emissivity



current observational upper bound, experimental sensitivities of 2008/2013,
predictions of various EEC ν fluxes

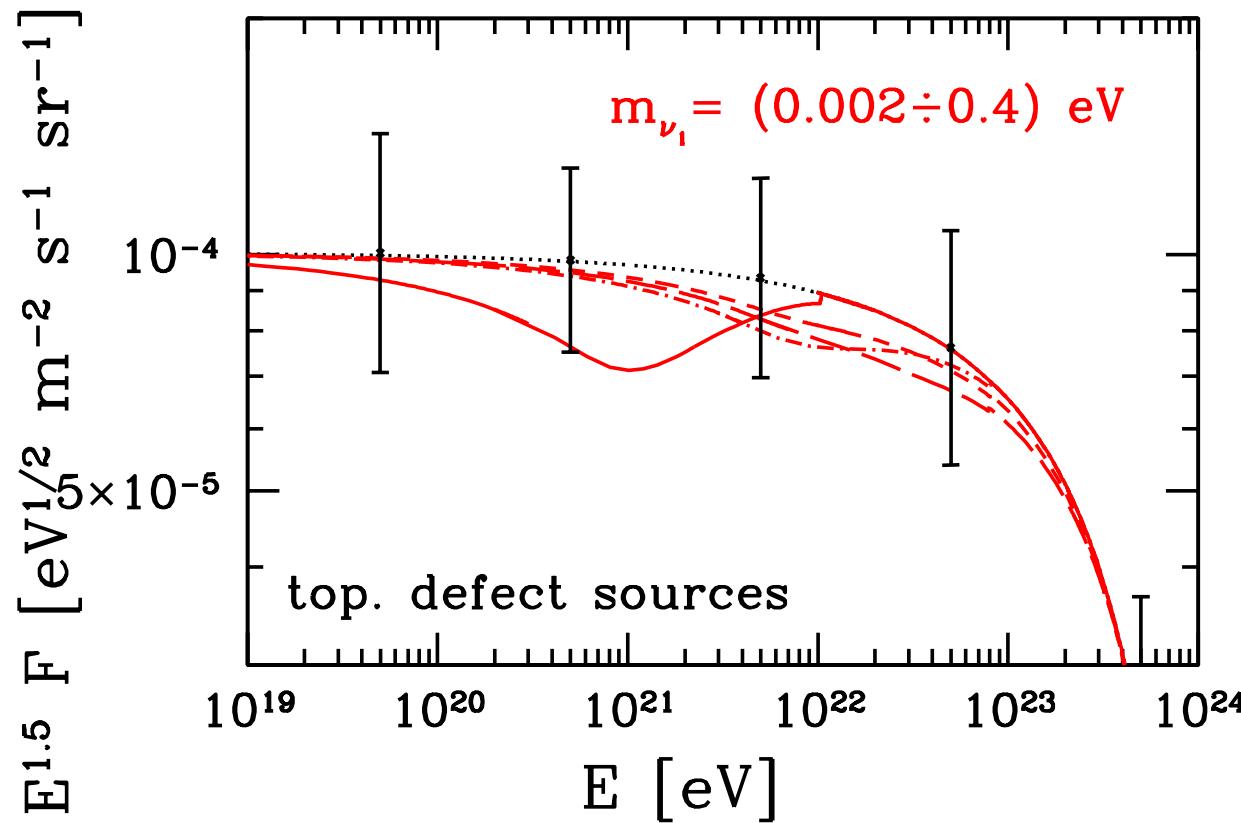
Benchmark fluxes - optimistic scenario

- "Hidden" source: ν flux touches the todays observational upper bound
- Quasi degenerated neutrino masses



Benchmark fluxes - less optimistic scenario

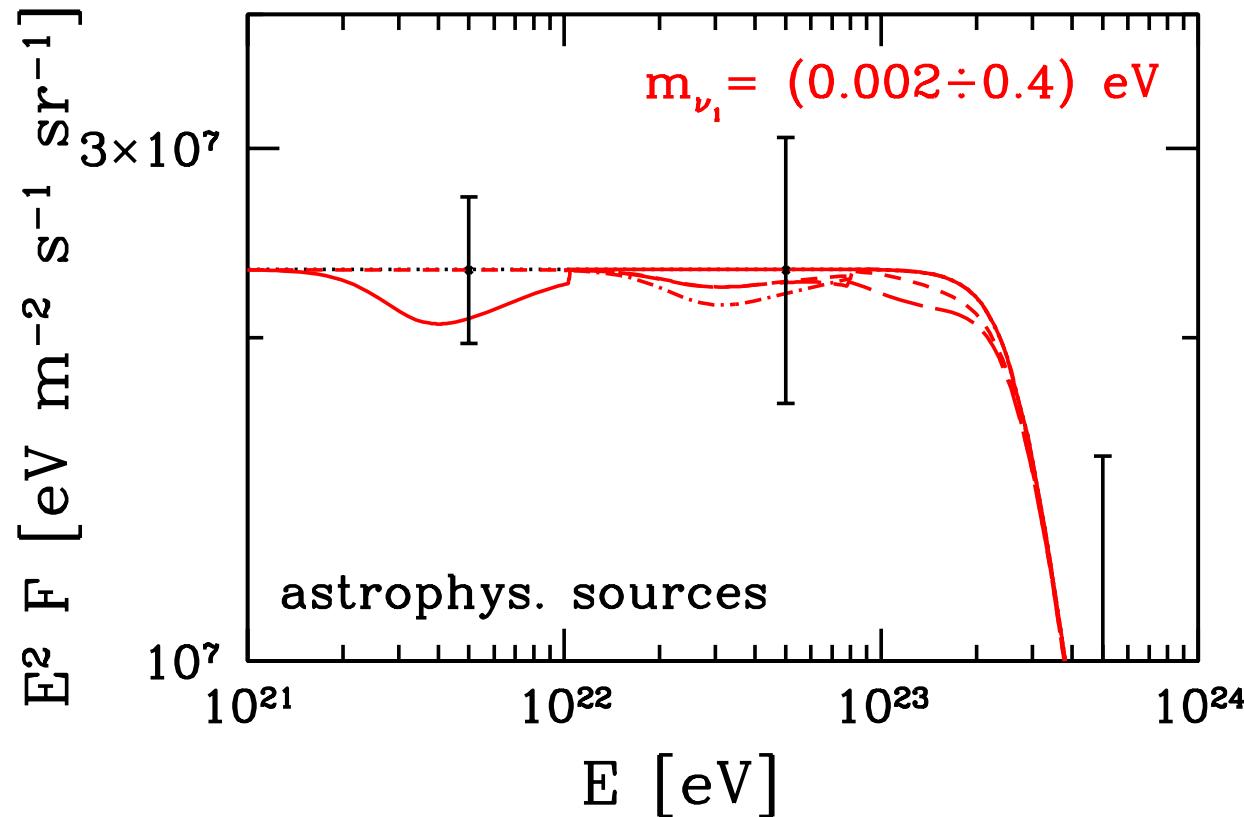
- **Transparent source:** ν flux touches the cascade limit
- Different mass scenarios



→ looks nice, but statistics worse

Benchmark fluxes - less optimistic scenario

- **Transparent source:** ν flux touches the cascade limit
- Different mass scenarios



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Conclusion

Detection of **C ν B** might be feasible, furthermore:
absorption dips rich in **particle and astrophysical information**

- High energy edge of dip: ν mass
- Shape, depth: Cosmol.param., activity, inj. spectrum and ν mass pattern
- Width: source evolutionary history and ν mass spectrum

However, detection strongly dependent on statistics:

On the magnitude of the ν flux:

Flux needed: should touch cascade limit - better higher flux
→ **Auger** should detect it within next few years

And on the depth of the dip:

ν mass spectrum - quasi deg. $\leftrightarrow m_{\nu_1} \gtrsim 0.1$ eV
→ measures **KATRIN** in upcoming decade

In any case - next decade **exciting** and decisive for C ν B detection!