Déflections of UHECRs Extra-Galactic Magnetic Fields

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Motivations

- EAS experiments provide some angular information about the arrival direction of UHECRs (with a resolution $\sim 1^{\circ}$)
- UHECRs ($E > 10^{19} \text{ eV}$) are, most likely, protons
- At such energies protons are expected to undergo tiny deflections in the Galactic Magnetic Fields
- If Extra Galactic MFs (EGMFs) are not too big the directional information is not lost and it may allow to identify UHECRs sources UHECR astronomy may then be possible

However

• We know very little about EGMFs from observations We need simulations What do we know about EGMFs from observations ?

MF in the Inter-Galactic Medium (IGMFs)

<u>Only upper limits are available</u> which are based on the Faraday Rotation Measurements (RMs) of Quasars' radio emission at cosmological distance $(z \sim 1)$



First limit by Kronberg'94 who, however assumed $n_e = \text{const.}$ and $\Omega_b = 1$!!

Then Blasi, Burles and Olinto, `99 accounted for n_e inhomogen. in Ly-alpha clouds

$$B_{H^{-1}} \leq 10^{-9} G$$

$$B_{50 Mpc} \leq 6 \times 10^{-9} G$$

$$B_{1 Mpc} \leq 10^{-8} G$$

proton deflections travelling over cosmological distances may be large if these limits are saturated !

Magnetic fields in galaxy clusters (ICMFs) **OBSERVATIONS**

Synchrotron Radio Halos

Minimum energy condition:

Faraday Rotation Measurements (RMs)

$$\langle B \rangle_{\rm V} = 0.1 \quad 1 \ \mu G$$

∜

$$\mathbf{RM} = \frac{\phi}{\lambda_{obs}^2} \propto \int_0^L \mathbf{n}_e^{\mathrm{T}}(1) \left(\frac{\lambda_l}{\lambda_{obs}}\right)^2 \mathbf{B}(1) \cdot \mathrm{d1} \qquad \mathrm{rad} \ \mathrm{m}^{-2}$$

Non thermal X-ray emission

if due to IC sc

 $\langle B \rangle_{V} = 0.1 \quad 1 \ \mu G$

if due to IC scattering
$$\Rightarrow$$

$$\frac{L_{Syn}}{L_{IC}} \propto \frac{B^2 / 8\pi}{\rho_{CMB}} \Rightarrow$$

$$B_{cell} = 1 \quad 10 \ \mu G$$

 $l_{c} = 10 \quad 100 \ kpc$

$B_{cell} = 1$	10 µG
$l_{\rm C} = 10$	100 kpc

Unknown origin, but

MSPH simulations in galaxy clusters

MSPH: Magnetic Smooth Particle Hydrodynamics [Dolag, Bartelmann & Lesch, astro-ph/0109541, 0202272]

N-body simulations of DM + gas hydrodynamics (SPH) + MHD

Lagrangian simulation (adaptive resolution)

A pre-existing seed field is invoked which is amplified by

Adiabatic compression:

$$\Phi_{\rm B} = \text{const} \implies B \propto R^{-2} \propto n_{\rm gas}^{2/3}$$

Cosmologically this implies: $B(z) \propto (1 + z)^2$

Non linear magnetic induction driven by the gas shear flows:

$$\frac{\partial \mathbf{B}}{\partial \mathbf{t}} = \vec{\nabla} \times \vec{\mathbf{v}} \times \vec{\mathbf{B}} + \vec{\mathbf{B}}_{\rm in}$$

Possible origin of the seed field

Primordial origin (z > 1000); turbulence at the reionization; starburst galaxies or early AGNs (z > 4) can account for the required seed



In the simulation B_{in} is switched-on at $z_{in} = 20$

ICMF evolution is highly non-linear for z < 3 \Rightarrow the final field does not change significantly if the same comoving field is switched-on even at $z \sim 3$

The geometrical structure of the seed field is almost irrelevant <u>The memory of the seed field structure is wiped-out in the</u> <u>clusters</u>

A uniform seed field can be assumed

Simulated RMs vs observations



	I	I
	$B_0(G)$	<b<sub>fin>_{core}</b<sub>
low	0.5 x 10 ⁻¹²	3 x 10 ⁻⁷
medium	2.5 x 10 ⁻¹²	8 x 10 ⁻⁷
high	1 x 10 ⁻¹¹	2 x 10 ⁻⁶

RMs intensity and profile are reasonably reproduced for

$$5 \times 10^{-13}$$
 < B_0 < 1×10^{-11} G

$$B_0 = B(z_{in}) (1 + z_{in})^{-2}$$

Best for $B_0 \sim 2 \ 10^{-12} \text{ G}$ especially for A119

The strategy:

- We assume that the MF in galaxy clusters have been originated from a cosmological seed field generated at high z. This should maximize UHECR deflections (largest filling)
- We combine MSPH with a constrained N-body simulation of the DM in the local universe (R ~ 100 Mpc)
- We choose a seed field strength which best reproduce (or maximize) RMs. We assume a uniform seed field.
- We compute proton deflection along random directions and construct deflections maps of UHE protons

Constrained simulations of the local universe

<u>Constrained Simulations</u>: initial conditions (density fluctuations) are chosen randomly from a Gaussian field with a power spectrum compatible with Λ CMD cosmology but constrained so that the smoothed density field coincide with that observed.



[Mathis, <u>Springel</u>, White et al., astro-ph/0111099]: Lagrangian code

IRAS 1.2 galaxy catalogue was used to constrain initial conditions Smoothing length = 7 Mpc

Our simulations: Constrained sim. for the DM + MSPH

[Dolag, D.G., Springel, & Tkachev, astro-ph/0310902, JETP Lett. '04 long paper submitted to JCAP '04]

• <u>Simulation volume</u>:

high resolution sphere radius : 115 Mpc embedded in a low resolution box of side \sim 350 Mpc

• <u>Resolution</u>

 $5 \ge 10^{7}$ particles with mass: $5 \ge 10^{9} M_{\odot}$ (DM); $7 \ge 10^{8} M_{\odot}$ (gas); max spatial resolution (<u>adaptative</u>) ~ 10 kpc

- Initial redshift: z = 60
- <u>Initial magnetic field</u> (comoving): $B_0 = 1 \ge 10^{-11}$ G (run 1: mhd_y) $B_0 = 2 \ge 10^{-12}$ G (run 2: mhd_z) homogeneous in both case but with \perp orientation



MFs in galaxy clusters



B correlates with the cluster temperature (i.e. with the mass) which agrees with the correlation observed between the cluster radio power and T_X

RMs vs X-ray luminosity



$$\mathbf{S}_{\mathrm{X}} \propto \int n_{e}^{2} \sqrt{\mathbf{T}} \, \mathrm{dx}$$

 $\sigma_{\mathrm{RM}} \propto \int n_{e} \, \mathbf{B}_{\parallel} \, \mathrm{dx}$

Observations: $\sigma_{\rm RM} \propto S_{\rm X}^{\alpha}$

$$\alpha \cong 1 \implies B \propto n_e$$

RMs radial profile



 $B_0 = 2 \times 10^{-12} G (run 2: mhd_z)$

MFs in the filaments and in the voids



50 Mpc



Filament





Across a filament

IGMFs in the Local Universe



Construction of deflection maps

- Deflection are summed along straight lines converging from randomly distributed points at maximal distance (110 Mpc) to the observer
- Trajectories with $\delta > 5^{\circ}$ trajectories are ignored (small pitch angle only)
- We consider protons with $E = 4 \times 10^{19} \text{ eV}$ and $E = 10^{20} \text{ eV}$ (arrival energy)
- Continuous energy losses are considered only for $E = 10^{20} eV$

$E_{obs} = 10^{20} eV$ with energy losses

Defelction on the sky $B_0 = 2e10^{-12} G$ 100 mhd_z mhd_y . mhd_z * 5.15 80 60 10 1.00e+00 [Degrees] 0.00e+00 40 Defelction on the sky 20 0 2 3 ϕ_{tresh} [Degree] 0 1 4 5 $B_0 = 10^{-11} G$ mhd_y **5**° **5.0** [Degrees] 0.0





Sky fraction covered by observable deflections extrapolation at large distances

 $E = 4 \times 10^{19} eV$

 $B_0 = 2e10^{-12}G \text{ (mhd_z)}$



This behaviour is consistent with a uniform density of deflectors (filaments)

There is a considerable fraction of the sky where correlation with sources is preserved !!

Sky fraction covered by observable deflections extrapolation at large distances

 $E = 4 \times 10^{19} eV$ $B_0 = 10^{-11} G$ ("maximal")



Conclusions

- We constructed the first simulated maps of UHE proton deflection produced by MF in the LSS of the local universe
- These maps have to be intended as maps of <u>maximal deflections</u> produced under the hypothesis that the seed field was originated before major cluster accretion mergers (deflections are smaller if ICMFs are produced locally)
- Pointing of UHECR sources should not be prevented over almost the entire sky at 10²⁰ eV independently on the structure of the seed field
- At smaller energies ~ $4 \times 10^{19} \text{ eV}$ deflections are mainly produced in filaments and sheets. Observable deflections may be produced over a significant fraction of the sky the exact amount depending on the structure of the seed field.
- The study of correlation with sources (e.g. BL Lacs) may allow to learn something on this structure, hence on the origin of EGMFs.