

# A Brane model with two asymptotic regions.

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Desy, Hambourg 2004.

## Plan of the talk

- *The Kip Thorne Interpretation Of The Melvin Solution.*
  - $\rho = \infty$ : asymptotic region near the symmetry axis.
- *Inclusion Of A Complex Higgs Field.*
  - divergence of the energy.
- Solution Interpolating Between The Cosmic String And The Melvin Solution.
  - potential.
  - causal structure.
- No Trapping of Gravity.

## References.

- Kip.S.Thorne, Phys.Rev. 138(1965) B251.
- Gary W. Gibbons, Daisuke Ida, Tetsuya Shiromizu, Phys.Rev.D66(2002)044010.
- J. Louko, D.L. Wiltshire, JHEP 0202:007,2002.
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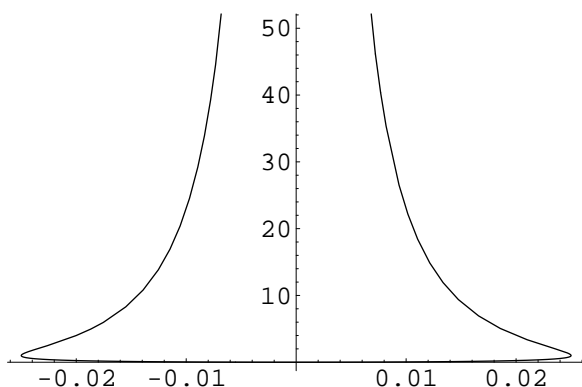
## The Kip Thorne Interpretation Of The Melvin Solution.

$$S = \int d^d x \sqrt{|g|} \left[ -\frac{1}{4} F_{ab} F^{ab} + \frac{R}{16\pi G} \right].$$

$$ds^2 = \left( 1 + \frac{\rho^2}{a^2} \right)^{2/(d-3)} \left( \eta_{\mu\nu} dx^\mu dx^\nu - d\rho^2 \right)$$

$$- \rho^2 \left( 1 + \frac{\rho^2}{a^2} \right)^{-2} d\phi^2, \quad ,$$

$$F = B_0 \left( 1 + \frac{\rho^2}{a^2} \right)^{-2} \rho d\rho \wedge d\phi, \quad ,$$



$\rho = \infty$ : asymptotic region near the symmetry axis.

## Inclusion Of A Complex Higgs Field.

Gibbons et al: real scalar field.

$$S = \int d^6x \sqrt{-g} \left[ \frac{1}{2} D_\mu \Phi D^\mu \Phi^* - \frac{\lambda}{4} (\Phi \Phi^* - v^2)^2 \right].$$

$$ds^2 = \beta^2(\rho) \eta_{\mu\nu} dx^\mu dx^\nu - \gamma^2(\rho) d\rho^2 - \alpha^2(\rho) d\phi^2$$

where  $\mu = 0, \dots, 3$  ,

$$\Phi = v f(\rho) e^{i\phi} \quad \text{and} \quad A_\phi = \frac{1}{e} (1 - p(\rho)) .$$

Energy density

$$\begin{aligned} \epsilon(\rho) = & \sqrt{|g|} \left[ \frac{1}{2} g^{\rho\rho} |D_\rho \Phi|^2 + \frac{1}{2} g^{\phi\phi} |D_\phi \Phi|^2 + \frac{1}{4} F_{\rho\phi} F^{\rho\phi} \right. \\ & \left. + \frac{\lambda}{4} (\Phi^* \Phi - v^2)^2 \right]. \end{aligned}$$

Higgs potential

$$\begin{aligned} \rho = \infty \quad \text{"symmetry axis"} & \Rightarrow f(\infty) = 0 \\ V(\Phi) & \sim \rho^{7/3} \lambda v^4 \end{aligned}$$

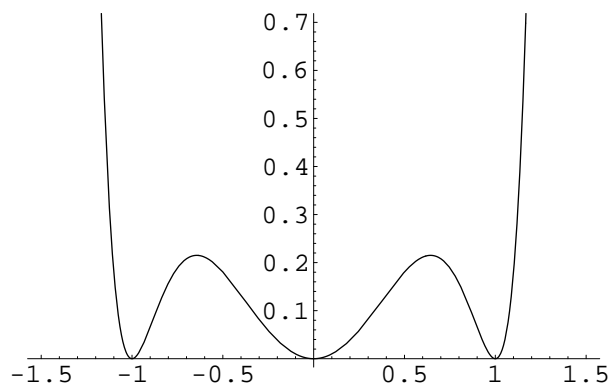
divergence.

## Solution Interpolating Between The Cosmic String And The Melvin Solution.

$$V(0) = 0 \quad .$$

Susy like potential

$$V(\Phi) = \lambda e^{w^2 \Phi \Phi^*} \Phi \Phi^* (\Phi \Phi^* - v^2)^2 \quad .$$



$$\mu = Gv^2 \quad , \quad \nu = \lambda^2 G^{-3} \quad , \quad \tau = e^2 v \quad \text{and} \quad \sigma = w^2 v^2 \quad .$$

$$\rho = L x \quad \text{where} \quad L = \frac{1}{2\sqrt{\pi}\mu^{5/4}\nu^{1/4}} \frac{1}{\sqrt{v}}$$

## Boundary conditions

$\rho$ : a length(cartesian) coordinate along the meridian.

$\rho \rightarrow -\infty$ : far region of the cosmic string solution:

$$\begin{aligned} A(x) &\sim x \quad , \quad B(x) \sim 1 \quad , \quad \bar{\gamma} \sim 1 \quad , \\ f(x) &\sim 1 \quad , \quad p(x) \sim 0 \end{aligned}$$

$\rho \rightarrow \infty$ : asymptotic region of the Melvin solution:

$$\begin{aligned} A(x) &\sim \frac{a^2}{x} \quad , \quad B(x) \sim \bar{\gamma}(x) \sim \left(\frac{x^2}{a^2}\right)^{1/3} \quad , \\ f(x) &\sim 0 \quad , \quad p(x) \sim 1 \quad . \end{aligned}$$

True radial coordinate.

$$r = \int_0^\rho \gamma(y) dy$$

## Field Equations

$$\begin{aligned}
& e^{\sigma F^2(x)} B(x) F^2(x) \bar{\gamma}^2(x) (1 - F^2(x))^2 + \frac{A'(x) B'(x)}{A(x)} \\
+ & 3 \frac{B'^2(x)}{B(x)} - \frac{B'(x) \bar{\gamma}'(x)}{\bar{\gamma}(x)} - 2 \frac{B(x) P'(x)^2}{A^2(x)} + B''(x) = 0 \quad , \\
& \frac{1}{4} e^{\sigma F^2(x)} B(x) F^2(x) \bar{\gamma}^2(x) (1 - F^2(x))^2 + 2\pi\mu B(x) F'^2(x) \\
- & \frac{B(x) A'(x) \bar{\gamma}'(x)}{4A(x) \bar{\gamma}(x)} - \frac{B'(x) \bar{\gamma}'(x)}{\bar{\gamma}(x)} + 3 \frac{B(x) P'^2(x)}{2A^2(x)} \\
+ & \frac{B(x) A''(x)}{4A(x)} + B''(x) = 0 \quad , \\
& e^{\sigma F^2(x)} B(x) F^2(x) \bar{\gamma}^2(x) (1 - F^2(x))^2 \\
+ & \frac{2\tau}{\pi\mu^{5/2} \sqrt{\nu}} \frac{F^2(x) \bar{\gamma}^2(x) P^2(x)}{A(x)} + 4 \frac{A'(x) B'(x)}{B(x)} - \frac{A'(x) \bar{\gamma}'(x)}{\bar{\gamma}(x)} \\
& + 6 \frac{P'^2(x)}{A(x)} + A''(x) = 0 \quad , \\
& \frac{\tau}{2\pi\mu^{5/2} \sqrt{\nu}} F^2(x) \bar{\gamma}^2(x) P(x) - \frac{A'(x) P'(x)}{A(x)} + 4 \frac{B'(x) P'(x)}{B(x)} \\
- & \frac{\bar{\gamma}'(x) P'(x)}{\bar{\gamma}(x)} + P''(x) = 0 \\
& \frac{1}{2\pi\mu} e^{\sigma F^2(x)} F(x) \bar{\gamma}^2(x) (1 + (-4 + \sigma) F^2(x)) \quad +
\end{aligned}$$

$$\begin{aligned}
& (3 - 2\sigma)F^4(x) + \sigma F^6(x) - \frac{\tau}{4\pi^2\mu^{7/2}\sqrt{\nu}} \frac{F(x)\bar{\gamma}^2(x)P(x)}{A^2(x)} \\
& \frac{A'(x)F'(x)}{A(x)} + 4\frac{B'(x)F'(x)}{B(x)} - \frac{F'(x)\bar{\gamma}'(x)}{\bar{\gamma}(x)} + F''(x) = 0 \quad .
\end{aligned}$$

## The causal structure.

$$\left. \begin{array}{l} \bar{u} \\ \bar{v} \end{array} \right\} = \arctan[t \mp \sigma(\rho)] \quad ,$$

$$\sigma(\rho) = \int_0^\rho d\xi \frac{\gamma(\xi)}{\beta(\xi)}$$

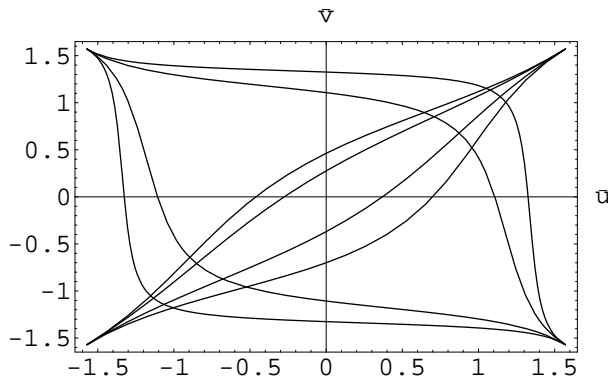


Figure 1: The causal structure of the solution is given. The time-like infinities are  $I^+(\bar{u} = \bar{v} = \pi/2)$  and  $I^-(\bar{u} = \bar{v} = -\pi/2)$ . As  $\rho$  can change sign, there are two space-like infinities :  $I_o^>(-\pi/2, \pi/2)$  and  $I_o^<(\pi/2, -\pi/2)$ . The curves which begin at  $I^-$  and end at  $I^+$  correspond to fixed values of  $\rho$  while the others correspond to fixed  $t$ .



## Trapping Of Gravity.

$$\int dx^4 dx^5 g^{00} \sqrt{|g|} < \infty$$

not satisfied.

### Conclusions.

- Melvin solution,  $\rho \rightarrow \infty$ , vanishing complex scalar field. Higgs potential excluded.
- Potential with one maximum and two minima.
- Solution interpolating between The Melvin and the cosmic string solutions.
- No trapping of gravity.