

Can dark energy
evolve to the Phantom!

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astro-ph/0407107

Motivation

A) Late acceleration of the Universe. (DE)

- 1) Observations of the transition from $w > -1$ to $w < -1$?
Alam, Sahni, Starobinsky 2003/4
- 2) Dynamical DE + $w \approx -1 \rightarrow$
Can DE evolve through $w = -1$?
- 3) if $w_{\text{now}} < -1$ observed,
coincidence problem
- 4) "big rip" singularity
for $w < -1$
- 5) Stability

3) Early Universe

- 1) "big crunch" singularity
in pre-big bang scenarios
- 2) exit from superexponential
inflation

C) General ?

Dynamical violation of DEC
• or WEC

Model

4

GR + 1 scalar field φ
(dominance of DE)

$$S = \int d^4x \sqrt{-g} \left[-\frac{R}{2} + p(\varphi, X) \right]$$

$$X \equiv \frac{1}{2} \nabla_\mu \varphi \nabla^\mu \varphi$$

$$T_{\mu\nu} = p_{,X} \nabla_\mu \varphi \nabla_\nu \varphi - p g_{\mu\nu} \rightarrow$$

„perfect fluid“

$$\text{pressure} = p(\varphi, X)$$

$$\text{energy density } \mathcal{E} = 2X p_{,X} - p$$

Armendáriz-Picón, Damour, Mukhanov

Armendáriz-Picón, Mukhanov, Steinhardt¹⁹⁹⁹
2000

EE for spatially flat FRW Univers

$$\frac{\ddot{a}}{a} = -\frac{1}{6} (\epsilon + 3p)$$

$$\underline{H^2 = \frac{\epsilon}{3}} \quad \rightarrow \quad \begin{array}{l} \epsilon \geq 0 \text{ + dominant} \\ \rightarrow \epsilon > 0 \end{array}$$

continuity equation

$$\dot{\epsilon} = -3H(\epsilon + p)$$

Whole dynamics in Eq. of Motion

$$\epsilon_{,x} \ddot{\varphi} + \dot{\varphi} \sqrt{3\epsilon} p_{,x} + \epsilon_{,\varphi} = C$$

$(\varphi, \dot{\varphi})$ - Phase space

↑
Phase diagram needed

$$W \equiv \frac{P}{\mathcal{E}} = -1 + \frac{2X}{\mathcal{E}} p_{,x}$$

$$X \geq 0 \quad \mathcal{E} > 0 \quad \rightarrow \quad p_{,x}$$

should change the sign
for the dynamical evolution
from $W \geq -1$ to $W < -1$

We do not live in an ideal FRW
Universe! \rightarrow Perturbations

sound speed

$$c_s^2 \equiv \frac{p_{,x}}{\mathcal{E}_{,x}} = \frac{(W+1)\mathcal{E}}{2X\mathcal{E}_{,x}}$$

Garriga, Mukhanov
1999

$$C_s^2 = \frac{(\omega+1)\epsilon}{2X\epsilon_{,X}} \rightarrow \begin{array}{l} \text{if } X \neq 0 \\ \epsilon_{,X} \neq 0 \end{array}$$

$(\omega+1)$ changes the sign \leftrightarrow

C_s^2 changes the sign!

For Stability $C_s^2 > 0!$ \rightarrow

Transition destroys \approx FRW
Universe!

Consider the transition

via $X=0$. C_s^2 ?

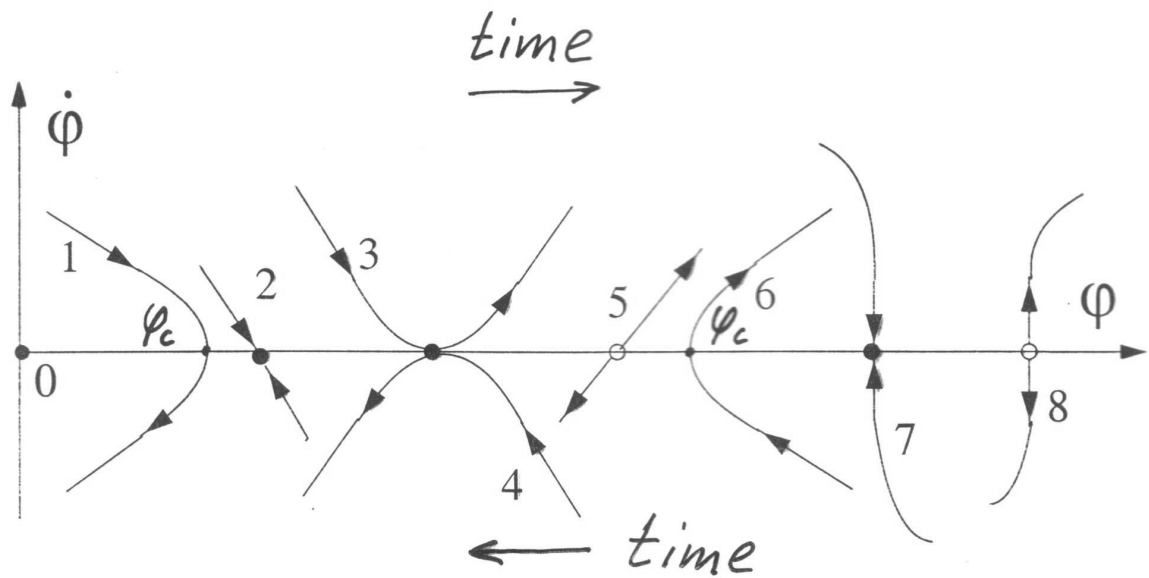


Figure 1: Possible phase curves in the neighborhood of the φ -axis. Only on the curves 1 and 6 the system pass the φ -axis. The curves 2,3,4,7 have an attractor as a shared point with the φ -axis, whereas the curves 5,8 have an repulsor. These attractors and repulsors can be a fixed point solutions or singularities as well.

$$\frac{d\dot{\varphi}}{d\varphi} = \frac{\ddot{\varphi}}{\dot{\varphi}}$$

- For phase curves 1 and 6

$$\varphi - \varphi_c \sim \dot{\varphi}^{2n} = X^n \quad n \geq 1$$

thus $x_{p,x}(X^n, X)$ cannot change the sign.

- For phase curves 3 and 4

$$\dot{\varphi} \sim (\varphi - \varphi_c)^{2n} \quad \text{and} \quad n \geq 1$$

$$t = \int_{\varphi_{in}}^{\varphi_c} \frac{d\varphi}{\dot{\varphi}(\varphi)} \sim \int_{\varphi_{in}}^{\varphi_c} \frac{d\varphi}{(\varphi - \varphi_c)^{2n}} = \infty$$

The last possibility for transition
the points $\Psi_c \equiv (\psi_c, \dot{\psi}_c)$ where

$$p_{,\psi}(\Psi_c) = 0, \quad \mathcal{E}_{,\psi}(\Psi_c) = 0 \\ \chi_c \neq 0$$

Eq. of Motion Leads to the
Eq. for phase curves

$$\frac{d\dot{\psi}}{d\psi} = - \frac{\dot{\psi} p_{,\psi} \sqrt{3\mathcal{E}} + \mathcal{E}_{,\psi}}{\dot{\psi} \mathcal{E}_{,\psi}}$$

$$\exists \dot{\psi}/\psi_c \rightarrow \mathcal{E}_{,\psi}(\Psi_c) = 0$$

if $\mathcal{E}_{,\psi}(\Psi_c) \neq 0$ then $\nexists \dot{\psi}$ and

$$\frac{d^2\psi}{d^2\dot{\psi}} = - \frac{\dot{\psi}^2}{\mathcal{E}_{,\psi}} \mathcal{E}_{,\psi\psi} \neq 0 \rightarrow$$

$\mathcal{E}_{,\psi}$ Transition impossible

Ψ_c is a singularity of Eq. Motion

Let us study the type of the singularity.

Use QTDE. Linearization.

- Why it is important to investigate this type?
- The form of the Eq. Motion excludes nodal point!
- The only mathematically allowed transition is absolutely unstable. The measure of the solutions which realize the transition is zero in $(\varphi, \dot{\varphi})$

Conclusions

- For dominating DE with $p(\varphi, \nabla_\mu \varphi)$
- 1) the transitions from $w \geq -1$ to $w < -1$ (or vice versa) are physically implausible.
- 2) no exit from superexp. infl.
- 3) smooth bounce cannot be realized
- DE with $p(\varphi, \nabla_\mu \varphi) = \frac{1}{2} (\nabla_\mu \varphi)^2 f(\varphi) - V$ cannot evolve from $w \geq -1$ to $w < -1$ (or vice versa).

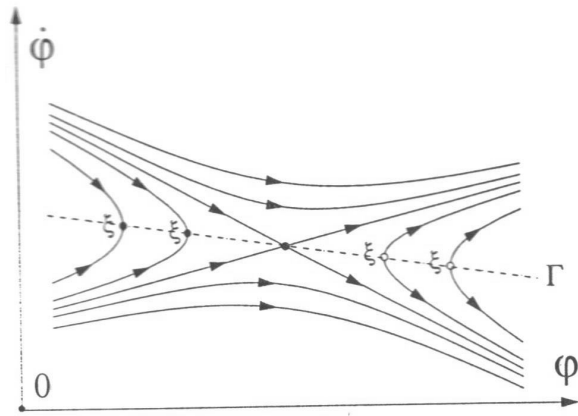


Figure 2: Phase curves in the neighborhood of the singular point Ψ_c^+ are plotted for the case of the real λ . At the points ξ the solutions $\varphi(t)$ do not exist. These points together with Ψ_c^+ build the curve Γ on which $\varepsilon_{,X}(\Gamma) = 0$.

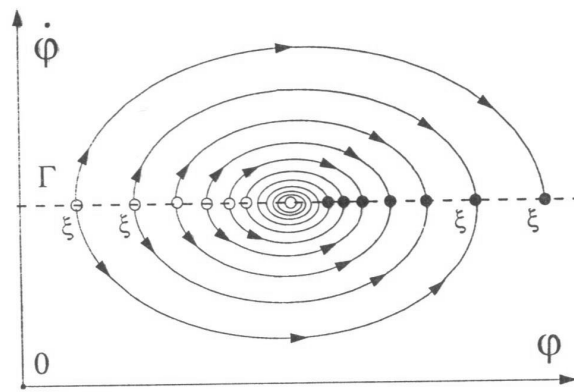


Figure 3: Phase curves in the neighborhood of the singular point Ψ_c^+ are plotted for the case of the pure imaginary λ . Here we assume that the singular point is a focus.

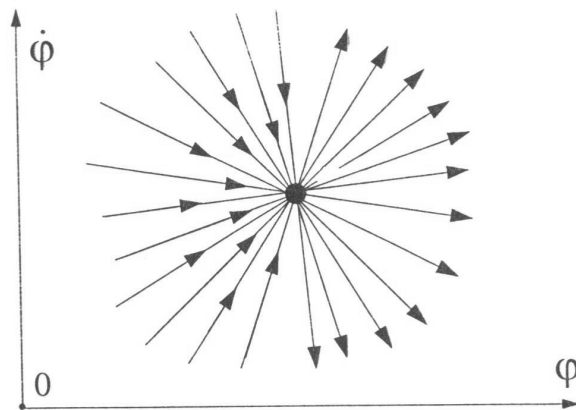


Figure 4: If \mathbf{A} had the eigenvalues $\lambda_1 = \lambda_2$ then the singular point Ψ_c^+ would be a nodal point and there would be a continuous set of trajectories passing through it. To illustrate this we plot here the phase curves in the particular case of an degenerate nodal point. The form of the equation of motion excludes such type of the singular points and therefore prevents the possibility of such transitions.

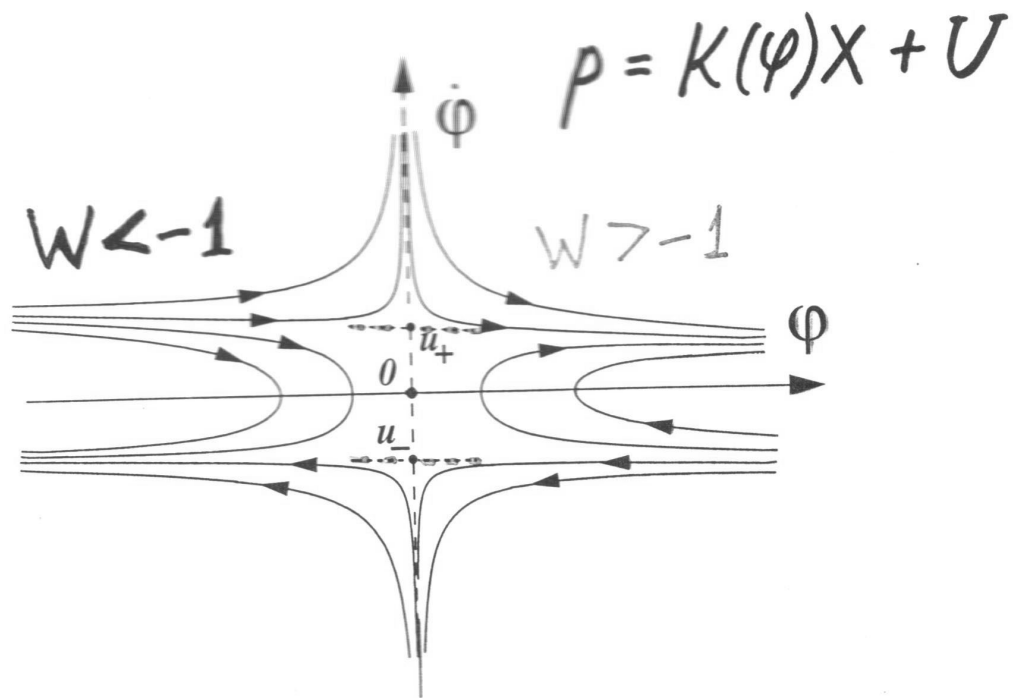


Figure 5: The typical behavior of the phase curves in the neighborhood of the "critical" line where $K(\varphi) = 0$ (here $\dot{\varphi}$ -axis), is plotted for the case when $K'_c > 0$, $V'_c < 0$. Horizontal dashed lines are the analytically obtained separatrices $\dot{\varphi}_{\pm}$, $(0, u_{\pm})$ are the points of transition.

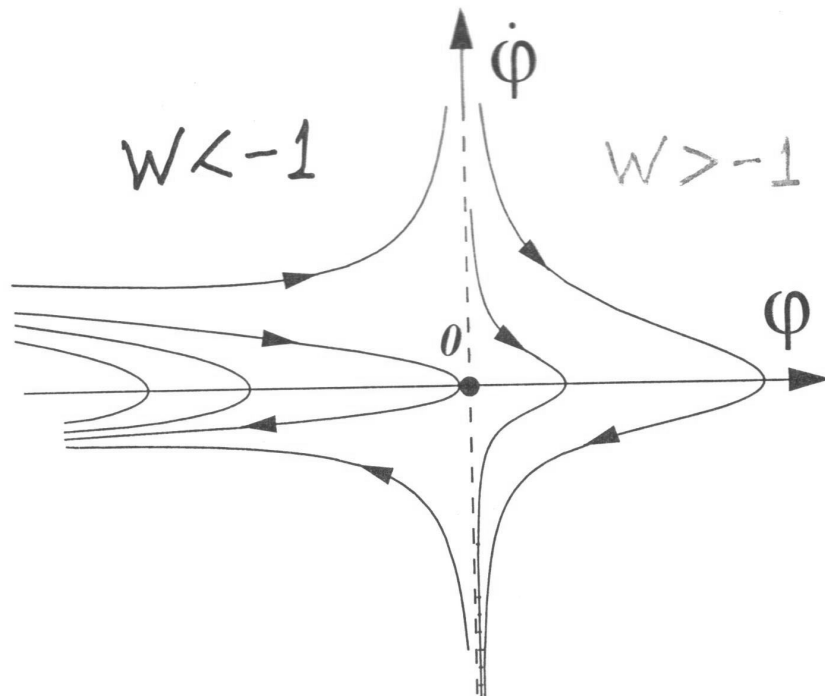


Figure 6: The typical behavior of the phase curves in the neighborhood of the "critical" line where $K(\varphi) = 0$ (here $\dot{\varphi}$ -axis), is plotted for the case when $K'_c > 0$, $V'_c = 0$, $V''_c > 0$.

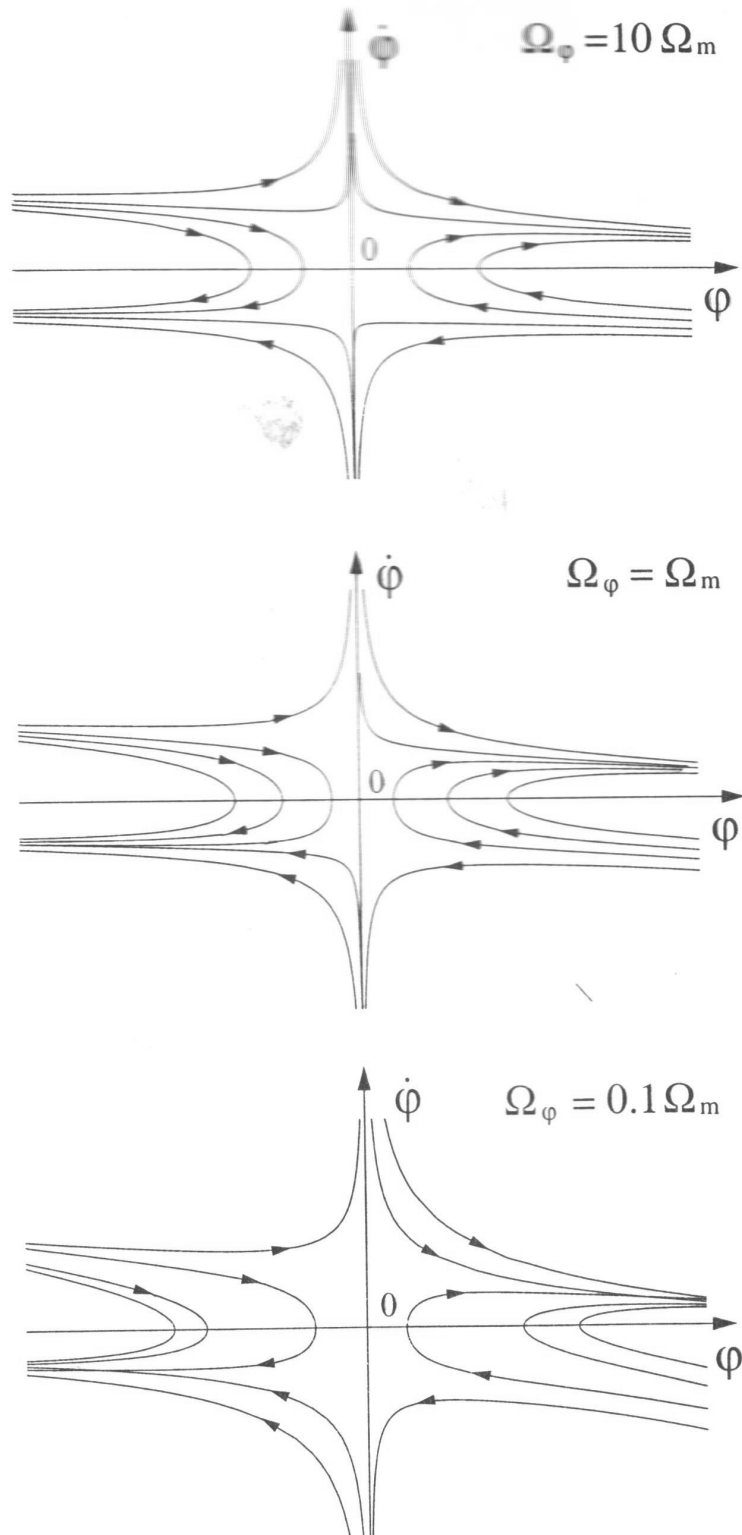


Figure 7: Numerically obtained trajectories of the dark energy described by a Lagrangian linear in X are plotted for the cases $\Omega_\phi = 10\Omega_m$, $\Omega_\phi = \Omega_m$ and $\Omega_\phi = 0.1\Omega_m$.