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## **Electroweak Form Factor in Sudakov and Threshold Regimes with Effective Field Theories**

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# Electroweak Form Factor in Sudakov and Threshold Regimes with Effective Field Theories

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**ABSTRACT:** We compute the massive gauge and scalar corrections to form factors in both the Sudakov and threshold regimes up to and including two-loop orders. The corrections are calculated for processes involving two external fermions and scalars in the spontaneously broken  $SU(N)$ -Higgs model, examining a range of composite operators. Our results are general, so we discuss how our form factors are mappable from our model to the Standard Model and beyond. The effective theory formalism deployed in our work extends previous studies based on infrared evolution equations, which either neglect scalar contributions or are restricted to the Sudakov regime.

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# 1 Introduction

The often addressed form factor is a crucial building block in the perturbative analysis of scattering processes occurring at the LHC and future colliders [19, 15]. It is also the simplest amplitude which can be used to study the infrared (IR) structure of the Standard Model and beyond. For reference, the QCD form factors of massless quarks have been evaluated through the three-loop approximation [31] and even to four loops in the leading-colour approximation [53]. On the other hand for massive quarks in QCD three-loop results are available thus far [10, 2, 1]. In our study, we consider massive gauge and Higgs corrections to the form factor for scalar, fermion and mixed external particles on a range of operators, checking and extending the results of [17] to two-loop orders. Furthermore, unlike previous work which focus solely on the Sudakov regime relevant for Hadron collider (LHC) studies, we also consider the threshold regime appropriate for future high precision colliders [13, 1, 51].

When considering the Sudakov regime, one tends to refer to partonic LHC processes since their centre of mass-energy is  $\sqrt{s} \sim 14$  TeV. This energy scale is more than an order of magnitude larger than the masses of the massive electroweak (EW) bosons. Radiative corrections to scattering processes depend on the ratio of mass scales, and such corrections at high energy depend on factors of the form  $\log(s/M_{W,Z}^2)$ . Such radiative corrections are further enhanced by two logarithmic powers in exclusive processes perturbatively at each order, and they are often referred to as Sudakov (double) logarithms. Electro-weak Sudakov corrections are significant at LHC energies, as each double logarithmic contribution is of the  $\mathcal{O}(10^{-1})$ . Thus, fixed-order perturbation theory breaks down, and one needs to employ resummation at all orders. Double logarithms even appear for inclusive processes, e.g. the total  $e^+e^-$  cross-section at large angles, since the colliding particles are not EW gauge singlets [20]. The literature on EW Sudakov effects in most cases focuses on employing infrared (IR) evolution equations to deal with computations [21, 27, 43, 28, 40, 25, 38, 39]. The Sudakov logarithm,  $\log(s/M_{W,Z}^2)$ , can be seen as an IR logarithm in EW theory, as it diverges as  $M_{W,Z} \rightarrow 0$ . In an effective field theory (EFT) formalism, IR logarithms in the original theory are convertible to ultraviolet (UV) logarithms in the effective theory, and then summable using standard renormalisation group (RG) techniques. The effective theories needed are soft-collinear effective theory (SCET) and heavy particle effective theory (HPET) for both fermions and scalars [7, 6], which have been previously used to study high energy EW Sudakov corrections [17, 18], and to perform resummation. This paper studies high energy EW Sudakov corrections using SCET and HPET, expanding on previous work employing the EFT formalism by other authors [17, 18, 16].

With regards to studies in the high energy (Sudakov) regime, an impressive level of accuracy has been achieved with many observables. For instance, the uncertainties on the predictions of the inclusive production cross-section of a top quark pair, now are at around 3 – 5% for a fixed top quark mass of  $m_t = 172.5$  GeV [47]. On the other hand, while these precise measurements provide a firm ground for testing the predictions within the SM, beyond the Standard Model (BSM) physics scenarios can just as well hide between said small uncertainties. To find a hint of BSM physics or to rule some hypotheses out, we need more precision than what can be provided by the LHC, and indeed, a future high precision collider which operates at threshold energies along with theoretical studies can achieve that [54, 61]. In the threshold regime, the processes we consider have a centre of mass energy,  $\sqrt{s}$ , near equal to the sum of the on-shell masses of the particles produced. Radiative corrections to scattering processes at threshold depend on the large on-shell external particle masses, as well as EW masses which are significant and, again, must be taken into account. We note further that in the threshold case, we take the gauge and Higgs masses to be IR as in the Sudakov case. Although there is extensive literature on QCD corrections at threshold, there is much more that needs to be achieved when considering EW and even BSM physics. The effective theory we employ

at threshold is HPET along with standard RG techniques to perform logarithmic resummation.

In essence, we generalise previous results in a gauge-invariant fashion to massive scalars and fermions, including radiative corrections due to the Higgs sector. Moreover, we study the threshold regime, as in previous works only the Sudakov regime was considered. Our results are computed without assuming that the Higgs and EW gauge bosons are degenerate in mass, as in previous calculations [40, 39], and we take the EFT analysis to two-loop orders to match the highest precision IR evolution results. We discuss the form factor computation in detail, checking and expounding on previous results. Although the form factor itself is not of direct relevance to collider processes, it still allows us to illustrate the EFT method for operators involving two external particles. More crucially, the form factor is known to be a building block for a vast array of processes. For example, our findings here can be used to compute corrections to processes relevant for the LHC and beyond, such as di-jet production,  $\bar{t}t$  production, squark pair production, or DM production in various models, which all involve operators with four external particles [16, 9, 57, 22]. Previous and future results on processes can be obtained from the computations given in this paper by summing over all pairs of external particles with the appropriate replacements of group-theoretic factors. The reason being is that the model we study, SU(N)-Higgs theory with spontaneous symmetry breaking (SSB), is selected for its generality. Moreover, the various set of composite operators, we look into allows future studies to be derived from our results. To illustrate such derivations, we apply our formalism to EW corrections in the SM for the case of light quarks, leptons and the top quark as external particles.

## 2 Full and Effective Theory Formalism

We begin by outlining the full theory employed in detail as well as including a primer for the effective theories which includes SCET, and HPET for both scalars and fermions.

### 2.1 SU(N)-Higgs Theory and the Standard Model

Our calculation is set in a spontaneously broken SU(2) gauge model, however we keep our results quite general, i.e. not substituting numerical colour factors and sticking with composite operators so that our results are more conveniently mapped to more specific models for phenomenological studies. In particular, with regards to the SM, the mapping of our model to the SM has been studied in detail previously [39, 17, 16], the sole difference being that in the SM, the isospin SU(2) group for left-handed fermions is mixed with the hypercharge U(1) group through the mass eigen-states of the Z-boson and the photon. In our model the electroweak fields,  $W^\pm$  and Z, are replaced with neutral SU(2) gauge bosons,  $W^a : a = \{1, 2, 3\}$ , with identical mass,  $M = M_W$ . The generators of SU(N) in the fundamental representation are labelled by  $T^a : a = \{1, \dots, N^2 - 1\}$ . The Lie algebra results in structure constants  $f^{abc}$  with Casimir operators for the fundamental and adjoint representations given by,  $C_F = (N^2 - 1)/2N$  and  $C_A = N$ , respectively. Moreover, we take the convention,  $\text{tr}(T^a T^b) = T_F \delta^{ab}$ , and even in the specific case of  $N = 2$  for SU(2), we remain with the general symbols rather than the specific values, which makes our results easily convertible, in particular for the case of the hypercharge U(1) gauge group. In the specific case of SU(2), the group generators are  $T^a = \sigma^a/2 : a = \{1, 2, 3\}$  where  $\sigma^a$  are the Pauli matrices, and  $f^{abc} = \epsilon^{abc}$ . With the above specifications we may now state the SU(2)-Higgs Lagrangian in the t'Hooft-Feynman gauge,

$$\mathcal{L} = \mathcal{L}_\psi + \mathcal{L}_\chi + \mathcal{L}_{YM} + \mathcal{L}_{GF} + \mathcal{L}_{gh} + \mathcal{L}_{Higgs} + \mathcal{L}_{Yuk}. \quad (2.1)$$

The Lagrangian is split into a few parts;  $\mathcal{L}_\psi$  and  $\mathcal{L}_\chi$  which describe the fermions and scalars (external particles), respectively;  $\mathcal{L}_{YM}$  and  $\mathcal{L}_{GF}$  corresponds to the massive Yang Mills (YM) and gauge-fixing (GF) terms, respectively;  $\mathcal{L}_{gh}$  describes the Faddeev-Popov (FP) ghost fields;  $\mathcal{L}_{Higgs}$

corresponds to the free Higgs Lagrangian which induces SSB and lastly,  $\mathcal{L}_{Yuk}$  entails the Yukawa interaction terms which provide mass to the external fermions and scalars.

**Fermions/Scalars** Let  $\psi_i(x)$  and  $\chi_i(x)$  correspond to Fermions and scalar fields with subscripts labelling fields as we consider different incoming outgoing external states for generality. The Dirac and scalar Lagrangians then have the following form,

$$\mathcal{L}_\psi = \bar{\psi}_i i \not{D} \psi_i, \quad \mathcal{L}_\chi = D_\mu \chi_i^\dagger D^\mu \chi_i, \quad (2.2)$$

where  $D_\mu = \partial_\mu - igW_\mu^a T^a$ ,  $W^a(x)$  is the gauge field as previously defined and  $g$  corresponds to the  $SU(N)_W$  gauge coupling.

**YM and Gauge-Fixing** The Yang-Mills and gauge-fixing Lagrangians have the usual form,

$$\mathcal{L}_{YM} = -\frac{1}{4} F_{\mu\nu}^a F^{\mu\nu,a}, \quad \mathcal{L}_{GF} = -\frac{1}{2\xi_W} F_W^2, \quad (2.3)$$

such that  $F_{\mu\nu}^a = \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + gf^{abc} W_\mu^b W_\nu^c$  and  $F_W = (\partial^\mu W_\mu^a - \xi_W M_W \phi^a) T^a$  where  $\phi^a$  is the Goldstone boson field and  $\xi_W$  the linear t'Hooft gauge fixing parameter.

**FP-ghosts** The gauge-fixing Lagrangian,  $\mathcal{L}_{GF}$ , involves the unphysical components of gauge fields. In order to compensate for their effects, one introduces the Lagrangian,

$$\mathcal{L}_{gh} = -i(\partial^\mu \bar{c}^a) D_\mu^{ab} c^b - \xi_W m_W^2 \bar{c}^a c^a, \quad (2.4)$$

with FP-ghosts,  $c^a(x)$ ,  $\bar{c}^a(x)$ , and  $D_\mu^{ab} = \partial_\mu \delta^{ab} + gf^{abc} W_\mu^c$ .

**Higgs and Yukawa** The minimal Higgs sector consists of a single complex scalar field,  $\Phi(x)$ , which is coupled to the gauge fields with a covariant derivative and has a self-coupling, resulting in the Lagrangian,

$$\mathcal{L}_H = (D_\mu \Phi)^\dagger D^\mu \Phi - V(|\Phi|^2), \quad (2.5)$$

with the Higgs potential,  $V(|\Phi|^2) = \frac{\lambda}{2} (|\Phi|^2 - v^2/2)^2$ . The potential is constructed in such a way that it gives rise to spontaneous symmetry breaking. Meaning the parameters,  $\lambda$  and  $v$ , are chosen in such a way that the potential minimum occurs for a non-vanishing Higgs field. More specifically, the theory is constructed such that the classical ground state of the scalar field satisfies,

$$|\langle \Phi \rangle|^2 = \frac{v^2}{2} \neq 0. \quad (2.6)$$

In perturbation theory one has to expand around the ground state and the Higgs field is written as

$$\Phi = \frac{1}{\sqrt{2}} ((H + v) + i\phi^a T^a), \quad (2.7)$$

where  $H$  and  $\phi^a$  have zero vacuum expectation value and are real. The field,  $H$ , is the physical Higgs field and  $\phi^a$  are the Goldstone bosons which illustrate the unphysical degrees of freedom. Inserting (2.7) back into the full Lagrangian,  $\mathcal{L}$ , provides mass to the Higgs field and W-boson,

$$M_H = \sqrt{\frac{\lambda}{2}} v \quad \text{and} \quad M_W = \frac{gv}{2}, \quad (2.8)$$

respectively. As for fermion and scalar masses, these arise from the Yukawa-like interactions in the Lagrangian,

$$\mathcal{L}_{Yuk} = -y_{f,i} \bar{\psi}_i \Phi \psi_i - y_{s,i} \chi_i^\dagger \Phi \chi_i + h.c., \quad (2.9)$$



where  $y_{f,i}$  and  $y_{s,i}$  are the Yukawa couplings for the fermions and scalars, respectively. After spontaneous symmetry breaking, i.e. inserting (2.7) back into (2.9), results in mass terms for said fermions and scalars,

$$\mathcal{L}_{Yuk} = -\sqrt{2}(y_{f,i}\bar{\psi}_i\psi_i + y_{s,i}\chi_i^\dagger\chi_i)(H + v), \quad (2.10)$$

therefore we can re-write,

$$m_\psi = \sqrt{2}vy_f \Rightarrow y_f = \frac{g}{2\sqrt{2}}\frac{m_\psi}{M_W} \equiv \frac{g}{2\sqrt{2}}Y_f, \quad (2.11a)$$

$$m_\chi^2 = \sqrt{2}vy_s \Rightarrow y_s = \frac{g}{2\sqrt{2}}\frac{m_\chi^2}{M_W} \equiv \frac{g}{2\sqrt{2}}Y_s, \quad (2.11b)$$

and in this notation the Lagrangian becomes,

$$\mathcal{L}_{Yuk} = -m_{\psi_i}\bar{\psi}_i\psi_i - m_{\chi_i}^2\chi_i^\dagger\chi_i - \frac{g}{2}Y_{f,i}H\bar{\psi}_i\psi_i - \frac{g}{2}Y_{s,i}H\chi_i^\dagger\chi_i, \quad (2.12)$$

where  $h_{\psi,\chi}$  is conventionally used in Feynman rules, as given in Appendix A, which we attain by expanding each term in the full Lagrangian.

## 2.2 Heavy Particle Effective Theory

In the case of fermions we deploy heavy quark effective theory (HQET), which we describe briefly in this section but refer to other works for more detail [33, 34, 52]. HQET is used in calculations involving a bound state of a heavy quark with mass  $m \gg \Lambda_{QCD}$ , and light quarks with mass smaller than the colour confinement scale,  $\Lambda_{QCD}$ . The energy scale of the interactions between the light and heavy quark is of order  $\Lambda_{QCD}$ , which is small compared to the mass of the heavy quark. The momentum,  $p$ , of the system can therefore be decomposed in the following way,

$$p^\mu = mv^\mu + k^\mu, \quad (2.13)$$

such that  $v$  is the velocity of the heavy quark, which is usually normalised such that  $v^2 = 1$ , and  $k$  is the small residual momentum corresponding to light quark interactions. More precisely, the first part of (2.13) represents the energy of the heavy quark and is approximately conserved in interactions. The second part corresponds to a parameterisation of the remaining momentum, which is due to the motion of the light quarks and interactions between the light and heavy quarks, such that,

$$|k| \sim \mathcal{O}(\Lambda_{QCD}) \quad \text{and} \quad m \gg \Lambda_{QCD} \quad (2.14)$$

A hierarchy of scales is thus present, whence one can organize an effective theory founded upon hierarchy. An interesting feature of HQET, as will be seeing below, is that its propagating degrees of freedom are massless and the propagating degrees of freedom carry the residual momentum,  $k$ . We now outline the derivation of the HQET Lagrangian for a quark coupled to our SU(N) gauge and Higgs fields. Our starting point is the Lagrangian,

$$\mathcal{L} = \bar{\psi}(i\not{D} - m)\psi - \frac{g}{2}Y_f H\bar{\psi}\psi, \quad (2.15)$$

such that  $D^\mu = \partial^\mu - igW^\mu$  and  $Y_f$  is the Yukawa coupling as previously defined. Next, following the pedagogical derivation in [60], by introducing the projection operators,

$$P_\pm = \frac{1 \pm \not{v}}{2}, \quad (2.16)$$

and two eigen-functions of these operators,

$$h_f = e^{imv \cdot x} P_+ \psi, \quad (2.17a)$$

$$H_f = e^{imv \cdot x} P_- \psi. \quad (2.17b)$$

This allows us to decompose the spinor field as follows,

$$\psi = \frac{1 + \not{v}}{2}\psi + \frac{1 - \not{v}}{2}\psi = e^{-imv \cdot x}(h_f + H_f), \quad (2.18)$$

where the field and anti-field are given by  $H_f$  and  $h_f$ , respectively and they satisfy the relations,  $\not{v}h_f = h_f$  and  $\not{v}H_f = -H_f$ . The details of the external states of the heavy fields are explained in previous work [52]. Now, substituting into (2.18) into (2.16), using some simple gamma matrix and projection operator identities, and integrating out the anti-field,  $H_f$ , using its equation of motion, we arrive at our HQET Lagrangian,

$$\mathcal{L}_{HQET} = \bar{h}_f i v \cdot D h_f - \frac{g}{2} Y_f H \bar{h}_f h_f + \mathcal{O}(1/m), \quad (2.19)$$

where we neglect terms of  $\mathcal{O}(1/m)$  in our derivation as they are heavily suppressed. The heavy quark propagator is thus,

$$S(k) = \frac{1}{k \cdot v + i\delta} \frac{1 + \not{v}}{2}, \quad (2.20)$$

where  $k$  is the residual momentum defined earlier and the vertex coupling is given in Appendix A.

As for the derivation of the heavy-field limit of a real scalar, or spin-0, field, it is very similar to the fermionic, or spin-1/2, derivation. One starts by considering the Lagrangian of a complex scalar field,  $\chi$ , with mass,  $m$ , coupled again to our SU(N) gauge and Higgs fields,

$$\mathcal{L} = D_\mu \chi^\dagger D^\mu \chi - m^2 \chi^\dagger \chi - \frac{g}{2} Y_s H \chi^\dagger \chi \quad (2.21)$$

Motivated by earlier studies [12, 55, 35], we then decompose the scalar field in the following way,

$$\chi = \frac{e^{-imv \cdot x}}{\sqrt{2}} (h_s + H_s), \quad (2.22)$$

where again,  $H_s$  is the anti-field containing the heavy modes, which needs to be integrated out. More specifically,

$$h_s = \frac{e^{imv \cdot x}}{\sqrt{2m}} (iv \cdot \partial + m) \chi \quad (2.23a)$$

$$H_s = \frac{e^{imv \cdot x}}{\sqrt{2m}} (-iv \cdot \partial + m) \chi, \quad (2.23b)$$

and plugging (2.23) into (2.22) gives,

$$iv \cdot \partial H_s = (2m + iv \cdot \partial) h_s. \quad (2.24)$$

Hence, substituting (2.22) and (2.24) into the Lagrangian, (2.21), and using the equation of motion, one obtains the heavy scalar effective theory (HSET) Lagrangian in our model,

$$\mathcal{L}_{HSET} = h_s^\dagger i v \cdot D h_s - \frac{g}{2} Y_s H h_s^\dagger h_s + \mathcal{O}(1/m). \quad (2.25)$$

We again neglect terms of  $\mathcal{O}(1/m)$  in our derivation as they are heavily suppressed. The heavy scalar propagator is thus,

$$S(k) = \frac{1}{k \cdot v + i\delta}, \quad (2.26)$$

where  $k$  is again the residual momentum and the vertex coupling is also given in Appendix A. Comparing the HSET Feynman rules with the the full theory ones, we see that they are related by simply decomposing the momenta as in (2.13) and dividing by  $2m$ .

### 2.3 Soft-collinear effective theory

SCET is an effective theory for high-energy particles, with some energy of  $\mathcal{O}(Q)$ , where  $Q$  is a large scale that characterises the scattering process under consideration. SCET preserves the modes of the full theory which have an invariant mass much smaller than  $Q^2$ . The SCET fields and Lagrangian depend on the null vectors,  $n$  and  $\bar{n}$ , where  $n = (1, \mathbf{n})$  and  $\bar{n} = (1, -\mathbf{n})$ . The three-vector,  $\mathbf{n}$ , is chosen to be a unit vector, thus,  $\bar{n} \cdot n = 2$ .

When calculating the Sudakov form factor, we work in the so-called Breit frame, with  $n$  chosen to be along the  $p_2$  direction and  $\bar{n}$  is then along the  $p_1$  direction. The momentum transfer,  $q = p_2 - p_1$ , then has time component,  $q^0 = 0$ . We work with light-cone components, which for a four-vector,  $p$ , are defined by  $p^+ \equiv n \cdot p$  and  $p^- \equiv \bar{n} \cdot p$ . In our problem,  $p_1^- = p_{1\perp} = p_2^+ = p_{2\perp} = 0$ , and  $Q^2 = p_1^+ p_2^-$ , which is reflected in our Feynman rules, see Appendix A. When a fermion moves in a direction close to  $n$ , it is describable by an  $n$ -collinear SCET field,  $\xi_{n,p}(x)$ , where  $p$  is a label momentum, and has components  $\bar{n} \cdot p$  and  $p_\perp$  [7, 6]. Kinematically, the field,  $\xi_{n,p}(x)$ , describes a particle (both on or off-shell) with  $2E = \bar{n} \cdot p$  and  $p^2 \ll Q^2$ . The SCET power counting is then as follows,

$$p^- \sim Q, \quad p^+ \sim Q^2\lambda, \quad p_\perp \sim Q\lambda, \quad (2.27)$$

where  $\lambda$  is the parameter used for power counting in the EFT expansion. The total momentum of the SCET field,  $\xi_{n,p}(x)$ , is  $p + k$ , where as in HPET,  $k$  is the residual momentum, except in this effective theory,  $k$  is of order  $Q\lambda^2$ , and is obtained from a Fourier transform of the position vector,  $x$ . Note that the label momentum,  $p$ , only contributes to the minus and perpendicular components of the total momentum.

On the other hand, the gauge field in the effective theory is represented in many ways: Labelled  $n$ -collinear fields,  $W_{n,p}(x)$ , and  $\bar{n}$ -collinear fields,  $W_{\bar{n},p}(x)$ , and unlabelled ultrasoft (US) fields,  $W(x)$ , which are analogous to the soft and US fields introduced in NRQCD [58, 46]. The  $n$ -collinear field contains gauge fields with momentum near the  $n$ -direction, and momentum scaling given by,

$$\bar{n} \cdot p \sim Q, \quad n \cdot p \sim Q^2\lambda, \quad p_\perp \sim Q\lambda, \quad (2.28)$$

and the  $\bar{n}$ -collinear fields contain gluons moving near the  $\bar{n}$ -direction, with momentum scaling,

$$n \cdot p \sim Q, \quad \bar{n} \cdot p \sim Q^2\lambda, \quad p_\perp \sim Q\lambda. \quad (2.29)$$

Lastly, the ultrasoft field represents gauge bosons with all momentum components scaling as  $Q\lambda^2$ . There are a particularly useful identity we quote that holds for the SCET fermion field,

$$\frac{n\bar{n}}{4}\xi_{n,p} = \xi_{n,p}, \quad (2.30)$$

where

$$P_n = \frac{n\bar{n}}{4}, \quad P_{\bar{n}} = \frac{\bar{n}n}{4}, \quad P_{\bar{n}} + P_n = 1 \quad (2.31)$$

are projection operators. The leading order fermion Lagrangian is [7],

$$\bar{\xi}_{n,p} \frac{\bar{n}}{4} \left( i n \cdot D + \frac{p_\perp^2}{2\bar{n} \cdot p} \right) \xi_{n,p} \frac{\bar{n}}{4} \quad (2.32)$$

where  $iD^\mu = i\partial^\mu + gW^\mu$  is the ultra-soft covariant derivative, and we neglect terms involving the collinear gauge field. The fermionic SCET propagator is then given by,

$$S(p) = \frac{\not{n} \bar{n} \cdot p}{2 p^2} \quad (2.33)$$

SCET knows about the large momentum scale,  $Q$ , through labels,  $\bar{n} \cdot p_2$  and  $n \cdot p_1$ , attached to the fields,  $\xi_{n,p_2}$  and  $\xi_{\bar{n},p_1}$ , for the outgoing and incoming particles, respectively. As a result, SCET

anomalous dimensions can depend on  $Q$ . However, there are no modes in SCET which couple  $\bar{n} \cdot p_2$  to  $n \cdot p_1$ , so that SCET does not contain modes with off-shellness of  $\mathcal{O}(Q^2)$ , which are of course present in the full theory. We also require SCET fields for scalar particles, such as Higgs-like fields for instance. Let  $\Phi_{n,p}$  be the scalar analogue of  $\xi_{n,p}$  for fermions, which describes the  $n$ -collinear field for a scalar particle moving in a direction near  $n$ . One normalises the SCET field,  $\Phi_{n,p}$ , in the same way as the full theory field,  $\phi$ , producing scalar particles with unit amplitude. The scalar field kinetic energy term in the Lagrangian then becomes,

$$D_\mu \phi^\dagger D^\mu \phi \rightarrow \Phi_{n,p}^\dagger \left( (\bar{n} \cdot p)(in \cdot D) + p_\perp^2 \right) \Phi_{n,p} \quad (2.34)$$

in SCET. It is also convenient to re-define the scalar field as follows,

$$\phi_{n,p} = \sqrt{\bar{n} \cdot p} \Phi_{n,p} \quad (2.35)$$

in terms of which the kinetic term becomes,

$$\mathcal{L} = \phi_{n,p}^\dagger \left( in \cdot D + \frac{p_\perp^2}{(\bar{n} \cdot p)} \right) \phi_{n,p} \quad (2.36)$$

and then has the same normalisation as the fermion Lagrangian in (2.32). The re-scaled scalar propagator is given by,

$$\frac{1}{p^2} \rightarrow \frac{\bar{n} \cdot p}{p^2}. \quad (2.37)$$

Hence,  $\phi_{n,p}$  as defined, produces scalar particles moving in the  $n$ -direction with amplitude,  $\sqrt{\bar{n} \cdot p}$ .

### 3 The Form Factor

The physical quantity we consider in this work is the form factor in the Euclidean region, defined as the amplitude,  $F_E(Q^2) = \langle p_2 | \mathcal{O} | p_1 \rangle$  for the scattering of on-shell particles  $p_i^2 = m_i^2$  by an operator  $\mathcal{O}$ , with  $Q^2 = -(p_2 - p_1)^2 > 0$ . The time-like form factor is given by an analytic continuation,  $F(s) = F_E(-s - i0^+)$ , implying  $\log Q^2/\mu^2 \rightarrow \log s/\mu^2 - i\pi$ . We will compute  $F_E(Q^2)$  for fermion scattering with,  $\mathcal{O} = \bar{\psi}\gamma^\mu\psi, \bar{\psi}\psi, \bar{\psi}\sigma^{\mu\nu}\psi$ , scalar scattering with,  $\mathcal{O} = \chi^\dagger\chi, i(D^\mu\chi^\dagger\chi - \chi^\dagger D^\mu\chi)$ , and mixed scattering with,  $\mathcal{O} = \bar{\psi}\chi, \chi^\dagger\psi$ . All operators are taken to be gauge singlets and thus the external particles have the same gauge quantum numbers, but differing mass. We then compute the form factor,  $F_E(Q^2)$ , by employing a sequence of effective theories inspired by previous studies [17, 18]. In both the threshold and Sudakov regimes we consider, there are various widely separated scales and we switch to the relevant theory as we shift between scales.

To illustrate the matching, let us consider the Sudakov regime. At scales higher than  $Q^2$ , the model is the original Higgs-gauge theory or the so-called full theory in EFT terminology. As one shifts to scales below  $\mathcal{O}(Q^2)$ , we transition to an effective field theory (SCET) where degrees of freedom with off-shellness of  $\mathcal{O}(Q^2)$  are integrated out. The full and effective theory share identical infrared (IR) physics but differ in their ultraviolet (UV) behaviour. To ensure that operators in the full and effective theories have the same on-shell matrix elements, one must introduce a so-called multiplicative matching coefficient. If the full theory is matched onto SCET at  $\mu \sim Q$  then the matching coefficient selected,

$$\langle p_2 | \mathcal{O}(\mu) | p_1 \rangle = \exp[C(\mu)] \langle p_2 | \tilde{\mathcal{O}}(\mu) | p_1 \rangle, \quad (3.1)$$

where  $\exp[C(\mu)]$  is the matching coefficient at  $\mu \sim Q$  which is in exponential form for convenience and  $\tilde{\mathcal{O}}(\mu)$  is the effective theory version of the full theory operator,  $\mathcal{O}(\mu)$ . The matching coefficient is independent of IR physics and is computable if perturbation theory is valid at  $\mu \sim Q$ . In general,

a single operator,  $\mathcal{O}$ , can match onto a set of operators,  $\tilde{\mathcal{O}}_i$  in the EFT with identical quantum numbers [32, 36]. The matching coefficient  $C(\mu)$  contains logarithms,  $\log \mu^2/Q^2$ , and logarithms are not large if  $\mu \sim Q$ . Although we choose  $\mu = Q$ , any value of  $\mathcal{O}(Q)$  may be chosen as well, and all physical observables do not depend on the renormalisation scale,  $\mu$ . The convention we follow is to pick the coefficient,  $c(\mu)$ , of  $\mathcal{O}$  in the full theory, to be unity at  $\mu = Q$ . Our choice then provides the normalisation for  $F_E(Q^2)$ , and  $c(Q) = \exp[C(Q)]$  is the coefficient of  $\mathcal{O}$  in SCET at  $\mu = Q$ . Moreover, to do RGE for  $c(\mu)$  between scales we use the usual equation,

$$\mu \frac{dc(\mu)}{d\mu} = \gamma(\mu)c(\mu), \quad (3.2)$$

such that  $\gamma(\mu)$  is the anomalous dimension of  $\tilde{\mathcal{O}}$  in the effective theory. We then repeat these steps of matching and RGE as we shift between well-separated energy scales, integrating out the appropriate degrees of freedom along the way. The EFT approach is superior to IR evolution as it divides a multi-scale calculation into multiple single-scale pieces which are simpler to work with. One can then trivially identify so-called universal quantities which are independent of scale. Lastly, we re-state that in an EFT calculation, the IR divergences in the theory above a matching scale match the UV divergences in the theory below the matching scale. Thus, with regards to most of our results presented here, having checked the above and below UV-IR matching, we need only provide the physically relevant finite parts.

For reference, our notation is as follows, we use  $a(\mu) \equiv \alpha(\mu)/(4\pi)$ , and for applications to the SM,  $a_i(\mu) \equiv \alpha_i(\mu)/(4\pi)$  where  $i = \{s, 2, 1\}$  for the QCD, SU(2) and U(1) couplings. Hypercharge is taken to be normalised such that  $Q = T_3 + Y$ . Our various Logarithms are denoted by  $\mathcal{L}_A \equiv \log A^2/\mu^2$ , for  $A = Q, M_{W,H}, m_{1,2}$ .  $C_A, C_F$  and  $T_f$  are the SU(N) Casimir operators and index for external particles.

## 4 Renormalisation

### 4.1 Field Renormalisation

The on-shell renormalization of the external fermion/scalar fields in our form factor expansions require the multiplication of the vertex corrections by a factor of  $Z$ , where  $Z^{1/2}$  is the fermion/scalar wave function renormalization (WFR) constant. The factor,  $Z$ , is determined by the fermion/scalar self-energy corrections  $\Sigma$  at on-shell momentum  $p^2 = m^2$ , in specific ways we will describe below. In a perturbative expansion with the external fields we study,  $\{\psi, \chi, h_f, h_s, \xi_{n,p}, \phi_{n,p}\}$ , letting  $\{I, J\}$  denote these fields such that  $V_{IJ}$  and  $Z_{IJ} = \sqrt{Z_I Z_J}$  correspond to the vertex and wavefunction contributions, we have,

$$V_{IJ} = 1 + aV_{IJ}^{(1)} + a^2V_{IJ}^{(2)} + \mathcal{O}(a^3), \quad (4.1)$$

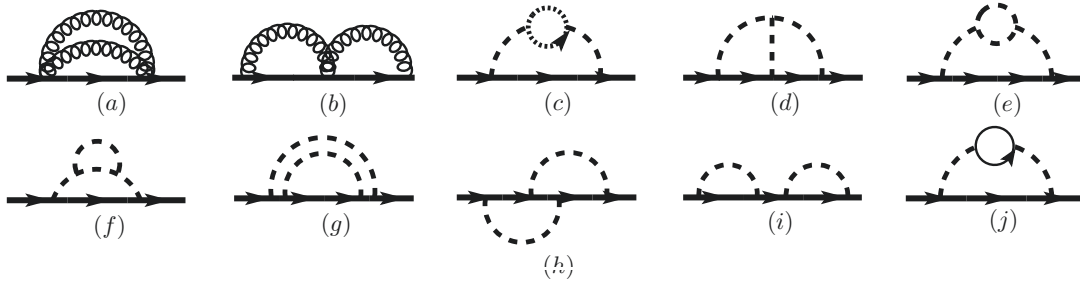
$$Z_I = 1 + a\delta Z_I^{(1)} + a^2\delta Z_I^{(2)} + \mathcal{O}(a^3). \quad (4.2)$$

Therefore, the WFR is given by,

$$Z_{IJ} = 1 + \frac{a}{2} \left( \delta Z_I^{(1)} + \delta Z_J^{(1)} \right) + \frac{a^2}{2} \left( \delta Z_I^{(2)} + \delta Z_J^{(2)} + \frac{1}{2} \delta Z_I^{(1)} \delta Z_J^{(1)} - \frac{1}{4} (\delta Z_I^{(1)})^2 - \frac{1}{4} (\delta Z_J^{(1)})^2 \right).$$

Whence, the total form factor,  $F_{IJ} = V_{IJ}Z_{IJ}$ , up to order  $\alpha^2$ , can be written as follows,

$$F_{IJ} = 1 + a \left\{ V_{IJ}^{(1)} + \frac{1}{2} \left( \delta Z_I^{(1)} + \delta Z_J^{(1)} \right) \right\} + a^2 \left\{ V_{IJ}^{(2)} + \frac{1}{2} \left( \delta Z_I^{(2)} + \delta Z_J^{(2)} \right) + \frac{1}{2} \left( \delta Z_I^{(1)} + \delta Z_J^{(1)} \right) V_{IJ}^{(1)} + \frac{1}{4} \delta Z_I^{(1)} \delta Z_J^{(1)} - \frac{1}{8} (\delta Z_I^{(1)})^2 - \frac{1}{8} (\delta Z_J^{(1)})^2 \right\}. \quad (4.3)$$



**Figure 1:** Two-loop wave-function correction graphs, arrowed lines represent all incoming-outgoing particles, dashed lines correspond to bosonic propagators. (a), (b) are seagull terms and only occur with scalar propagators, (h)-(j) and (c)-(g) represent Abelian and non-Abelian corrections, respectively.

With the above notation we may now discuss how to obtain the WFR constant,  $Z_I$ , for the spin- $\{0, 1/2\}$  fields we study. In all cases, the wavefunction corrections are garnered from self-energy amplitudes, denoted by  $\tilde{\Sigma}_I$ , which are quadratic matrices both in the spinor and in the isospin space [23, 60]. Moreover, we note here that the collinear correction to the particle propagator is the same as in the full theory [7, 17]. Therefore the wavefunction corrections are the same as in the full theory/non-collinear case. Whence, we only need to outline how to obtain the wavefunction contributions to the form factors for the full theory and HPET fields.

**Scalar field:** For massive scalars of momentum,  $p$ , and mass,  $m$ , the self-energy amplitudes, as shown in figure 1 are of the form,

$$\tilde{\Sigma}_\chi = -i\Sigma_\chi(p^2)\mathbf{1}. \quad (4.4)$$

From this we may extract the WFR contributions in the following way,

$$\delta Z_\chi = \frac{i}{4}\text{tr}(\partial_{p^2}\tilde{\Sigma}_\chi|_{p^2=m^2}). \quad (4.5)$$

The massless case is identical except one takes  $p^2 = 0$  instead.

**Fermion field:** In the case of fermions of momentum,  $p$ , and mass,  $m$ , the self-energy amplitudes are of the form,

$$\tilde{\Sigma}_\psi = -i(\Sigma_\psi^V(p^2)\not{p} + \Sigma_\psi^S(p^2)m)\mathbf{1}, \quad (4.6)$$

where the super-scripts,  $V$  and  $S$ , denote vector and scalar contributions, respectively. From this we may extract the WFR contributions,

$$\delta Z_\psi = \{\Sigma_\psi^V(m^2) + 2m^2\partial_{p^2}(\Sigma_\psi^V(p^2) + \Sigma_\psi^S(p^2))|_{p^2=m^2}\}. \quad (4.7)$$

The massless case simplifies as  $p^2 = 0$  instead and the terms proportional to  $m^2$  vanish.

**Heavy fields:** Lastly, for heavy scalars and fermions,  $h$ , of momentum,  $p$ , and velocity,  $v$ , the self-energy amplitudes are of the form,

$$\tilde{\Sigma}_h = -i\{\Sigma_h^F(v \cdot p) + \Sigma_h^R(M_B)\}\mathbf{1}, \quad (4.8)$$

for bosons of mass,  $M$ , coupling to the heavy fields, as in our case. We thus have an additional contribution to the heavy field residual mass term,  $\delta m$ , along with the usual wave function contribution,

$$\delta Z_h = i\partial_{v \cdot p}\tilde{\Sigma}_h|_{v \cdot p=0} \quad (4.9)$$

$$\delta m_I = -i\tilde{\Sigma}_h(v \cdot p = 0). \quad (4.10)$$

Field	$m$	$M$	$\delta Z^{(1)}$	$\delta Z^{(2)}$
$\psi$	0	0	$\frac{C_A}{2} \left( C_F + \frac{Y_f^2}{8} \right) \left\{ \frac{2}{\varepsilon_{\text{IR}}} - \frac{2}{\varepsilon_{\text{UV}}} \right\}$	$F_\psi^{(0,0)}$
$\psi$	0	$M$	$\frac{C_A}{2} \left\{ C_F \left( -\frac{2}{\varepsilon_{\text{UV}}} + 1 + 2\mathcal{L}_{M_W} \right) + \frac{Y_f^2}{8} \left( -\frac{2}{\varepsilon_{\text{UV}}} - 1 + 2\mathcal{L}_{M_H} \right) \right\}$	$F_\psi^{(0,M)}$
$\psi$	$m$	0	$\frac{C_A}{2} \left\{ C_F \left( -\frac{2}{\varepsilon_{\text{UV}}} - \frac{4}{\varepsilon_{\text{IR}}} - 4 + 3\mathcal{L}_m \right) + \frac{Y_f^2}{8} \left( -\frac{2}{\varepsilon_{\text{UV}}} + \frac{8}{\varepsilon_{\text{IR}}} + 14 - 6\mathcal{L}_m \right) \right\}$	$F_\psi^{(m,0)}$
$\psi$	$m$	$M$	$\frac{C_A}{2} \left\{ C_F \left( -\frac{2}{\varepsilon_{\text{UV}}} - 8 + 2\mathcal{L}_{M_W} - P(m/M_W) \right) + \frac{Y_f^2}{8} \left( -\frac{2}{\varepsilon_{\text{UV}}} + 14 + 2\mathcal{L}_{M_H} - P'(m/M_H) \right) \right\}$	$F_\psi^{(m,M)}$
$\chi$	0	0	$\frac{C_A C_F}{2} \left\{ \frac{4}{\varepsilon_{\text{UV}}} - \frac{4}{\varepsilon_{\text{IR}}} \right\}$	$F_\chi^{(0,0)}$
$\chi$	0	$M$	$\frac{C_A}{2} \left\{ -\frac{Y_s^2}{4M_H^2} + C_F \left( \frac{4}{\varepsilon_{\text{UV}}} + 3 - 4\mathcal{L}_{M_W} \right) \right\}$	$F_\chi^{(m,0)}$
$\chi$	$m$	0	$\frac{C_A}{2} \left\{ C_F \left( \frac{4}{\varepsilon_{\text{UV}}} - \frac{4}{\varepsilon_{\text{IR}}} \right) + \frac{Y_s^2}{2m^2} \left( \frac{1}{2\varepsilon_{\text{IR}}} + 1 - \mathcal{L}_m \right) \right\}$	$F_\chi^{(0,M)}$
$\chi$	$m$	$M$	$\frac{C_A}{2} \left\{ -\frac{Y_s^2}{2M_H^2} S'(m/M_H) + C_F \left( \frac{4}{\varepsilon_{\text{UV}}} - 4\mathcal{L}_{M_W} + S(m/M_W) \right) \right\}$	$F_\chi^{(m,M)}$
$h_{f,s}$		0	$\frac{C_A}{2} \left\{ C_F \left( \frac{4}{\varepsilon_{\text{UV}}} - \frac{4}{\varepsilon_{\text{IR}}} \right) + \frac{Y_s^2}{2} \left( -\frac{2}{\varepsilon_{\text{UV}}} + \frac{2}{\varepsilon_{\text{IR}}} \right) \right\}$	$F_h^{(0)}$
$h_{f,s}$		$M$	$\frac{C_A}{2} \left\{ C_F \left( \frac{4}{\varepsilon_{\text{UV}}} - 4\mathcal{L}_{M_W} \right) + \frac{Y_s^2}{2} \left( -\frac{2}{\varepsilon_{\text{UV}}} + 2\mathcal{L}_{M_H} \right) \right\}$	$F_h^{(M)}$

**Table 1:** Contributions to on-shell wavefunction renormalization. The exchanged boson masses are  $M = M_{W,H}$  where  $\mathcal{L}_M = \log M^2/\mu^2$ , and the external particle (fermion or scalar) mass is  $m$ . The two-loop wave-function corrections,  $F_I^{(i,j)}$ , and the parametric integral functions,  $P, P'$  and  $S, S'$ , are given in Appendix D.

The residual shift in the heavy particle mass,  $\delta m$ , is non-analytic in the boson mass squared, these arise from loop integrals which diverge as an odd power of loop momenta,  $l$ . Such non-analytic contributions are known to occur in mass corrections to particles with  $l \cdot v$  propagators [41, 37]. Such integrals are finite, but non-analytic, in dimensional regularization. There are two mass parameters for the heavy particle Lagrangian, the expansion parameter,  $m_0$ , and the residual mass term,  $\delta m$ . The two parameters are not independent; one can make the redefinition  $m_0 \rightarrow m_0 + \Delta m$ ,  $\delta m \rightarrow \delta m - \Delta m$ . A particularly convenient choice is to adjust  $m_0$  so that the residual mass term  $\delta m$  vanishes, when picking  $m_0$  such that  $\delta m = 0$  this choice is known as the pole mass [11], and we will follow this practice here.

## 4.2 Mass and Coupling Renormalisation

Our loop calculations up to two loop order have been performed using the unrenormalized Feynman rules. Introducing now the renormalized coupling constant and mass instead of the bare couplings does not change the two loop results at  $\mathcal{O}(\alpha^2)$ . However, the coupling and mass in the one-loop result have to be regarded as the bare parameters and must be replaced by the renormalised ones. In our work we employ the  $\overline{\text{MS}}$  scheme for the coupling renormalisation and the on-shell scheme for mass renormalisation. In the on-shell scheme, the square of the physical, renormalized mass is defined to be the real part of the pole of the propagator. In the case of coupling renormalisation the replacement can be applied naively as shown below. However, in the case of mass renormalisation, say given a mass  $M$ , with replacement (we denote the bare quantities with index, 0),

$$M_0^2 = M^2 + \delta M^2 + \mathcal{O}(\alpha^2), \quad (4.11)$$

in which  $\delta M^2$  corresponds to the mass contribution, the masses to be renormalised often appear in terms of the form  $(\mu^2/M^2)^\epsilon$  or in powers of logarithms. Thus the substitutions at one-loop are,

$$\left( \frac{M^2}{\mu^2} \right)^\epsilon \rightarrow \left( \frac{M^2}{\mu^2} \right)^\epsilon \left( 1 + \epsilon \frac{\delta M^2}{M^2} \right) + \mathcal{O}(\alpha^2), \quad (4.12)$$

$$\mathcal{L}_M^n = \mathcal{L}_M^n + n \mathcal{L}_M^{n-1} \frac{\delta M^2}{M^2} + \mathcal{O}(\alpha^2), \quad (4.13)$$

which, when applied provides corrections of  $\mathcal{O}(\alpha^2)$ . For the particles we are considering below, the renormalized quantities and renormalization constants are defined as follows,

$$\alpha_0 = (1 + \delta Z_\alpha)\alpha \quad (4.14a)$$

$$M_{W,0}^2 = M_W^2 + \delta M_W^2 \quad (4.14b)$$

$$M_{H,0}^2 = M_H^2 + \delta M_H^2 \quad (4.14c)$$

$$m_{\chi,0}^2 = m_\chi^2 + \delta m_\chi^2 \quad (4.14d)$$

$$m_{\psi,0} = m_\psi + \delta m_\psi, \quad (4.14e)$$

where the subscripts  $\psi$  and  $\chi$  indicate that the masses belong to fermion and scalar fields, respectively, that appear externally in the form factor.

#### 4.2.1 Coupling Renormalisation

According to the prescription of the  $\overline{\text{MS}}$  scheme, the unrenormalized coupling constant  $\alpha_0$  is replaced by the renormalized coupling  $\alpha$  via,

$$\alpha_0 = (1 + \delta Z_\alpha)\alpha = \alpha \left( 1 - \frac{\alpha}{4\pi} \frac{\beta_0}{\varepsilon_{\text{UV}}} \right) + \mathcal{O}(\alpha^3), \quad (4.15)$$

such that  $\beta_0$  is the leading (one-loop) coefficient of the renormalisation group  $\beta$ -function. We note that  $\beta_0$  has the following form,

$$\beta_0 = \frac{11}{3}C_A - \frac{4}{3}T_f n_f - \frac{1}{6}, \quad (4.16)$$

where the terms proportional to  $C_A$  and  $n_f$  correspond to the non-Abelian and fermionic contributions, respectively, while the last term corresponds to a Higgs contribution. Thus by applying the substitution (4.15) to our one-loop form factors, we get additional contributions of order  $\alpha^2$ .

#### 4.2.2 Gauge Mass Renormalisation

As this is the first case of mass renormalisation we consider we will discuss this in detail, at the amplitude level. The relation between the bare gauge boson mass,  $M_{W,0}$ , and the renormalized mass,  $M_W$ , is determined by the gauge boson self-energy corrections, which have the form,

$$\tilde{\Pi}^{\mu\nu,ab}(p) = i\delta^{ab}g^{\mu\nu}p^2\Pi(p^2)\mathbf{1} + \text{terms} \propto p^\mu p^\nu. \quad (4.17)$$

After extracting  $\Pi(p^2)$  from the amplitudes with the help of the projection operator,  $P_{\mu\nu} = g_{\mu\nu} - \frac{p_\mu p_\nu}{p^2}$ , the renormalised mass is given by setting  $\delta M_W^2 = -M_W^2\Pi(M_W^2)$ , and we may check various contributions at one-loop, up to  $\mathcal{O}(\varepsilon)$ , where  $\varepsilon$  are UV divergences. The results up to  $\mathcal{O}(\varepsilon)^2$ , needed for mass renormalisation contributing at two-loop orders is provided in Appendix B. We begin with the self-energy contributions from the fermion loop,

$$\Pi(M_W^2)_{n_f} = -\frac{4a}{9}T_f n_f \left\{ 5 + 3i\pi + \frac{3}{\varepsilon} - 3\mathcal{L}_{M_W} \right\} + \mathcal{O}(\varepsilon), \quad (4.18)$$

where as we stated before,  $a(\mu) = \alpha(\mu)/4\pi$ . Next, we have contributions from the non-Abelian gauge boson and ghost field loops,

$$\Pi(M_W^2)_{WW,cc} = \frac{a}{9}C_A \left\{ 82 - 12\sqrt{2}\pi + \frac{51}{\varepsilon} - 51\mathcal{L}_{M_W} \right\} + \mathcal{O}(\varepsilon), \quad (4.19)$$

from the loop with gauge and Higgs boson,

$$\Pi(M_W^2)_{WH} = a \left\{ -2 - \frac{1}{\varepsilon} + \mathcal{L}_{M_W} + \frac{sr}{M_W} \log w + r^2 \log r \right\} + \mathcal{O}(\varepsilon), \quad (4.20)$$



where we define  $r = M_H/M_W$ ,  $s = \sqrt{M_H^2 - 4M_W^2}$  and  $w = \frac{2M_W}{M_H + s}$ , and finally a contribution from the loops with Higgs and Goldstone bosons,

$$\Pi(M_W^2)_{\phi\phi} = \frac{a}{72} C_A \left\{ 34 - 3\sqrt{3} + \frac{15}{\epsilon} - 15\mathcal{L}_{M_W} \right\} + \mathcal{O}(\epsilon), \quad (4.21)$$

$$\begin{aligned} \Pi(M_W^2)_{H\phi} = a \left\{ \frac{1}{18} \left( 5 + \frac{3}{\epsilon} - 3\mathcal{L}_{M_W} \right) + \frac{r^2}{2} \left( \log r + \frac{3}{2} + \frac{1}{2\epsilon} - \frac{1}{3}\mathcal{L}_{M_H} \right) \right. \\ \left. - \frac{r^4}{12} - \frac{r^4}{2} \log r + \frac{r^5 s}{12M_W} \log w - \frac{r^3 s}{3M_W} \log w + \frac{r^6}{12} \log r \right\} + \mathcal{O}(\epsilon). \end{aligned} \quad (4.22)$$

Thus combining all terms provides one with the gauge boson mass correction in the replacement rules. Note that in the above contributions there are no terms from massive fermion and scalar loops, this makes sense as in the EFT formalism the scale where the bosons are no longer IR, i.e. where their masses are no longer zero, is the same scale where the fermions and scalars are taken to be UV or static. Moreover, we note that the self-energy diagrams with tadpoles have been omitted. They do not depend on the momentum of the gauge boson, so their contribution to the mass renormalization cancels exactly the corresponding vertex correction and field renormalization diagrams which are also dropped out.

### 4.2.3 Higgs Mass Renormalisation

As we were explicit in the previous section and broke down each contribution we will be brief now as the above still applies and we simply state the correction. The relation between the bare Higgs mass,  $M_{H_0}$ , and the renormalized mass,  $M_H$ , is determined by the Higgs self-energy corrections,  $\tilde{\Sigma}(p^2) = i\Sigma(p^2)\mathbf{1}$ . Extracting  $\Sigma(p^2)$  gives the renormalized mass by setting  $\delta M_H^2 = \Sigma(M_H^2)$ , which has the following form after combining all contributions,

$$\begin{aligned} \delta M_H^2 = a C_A C_F \frac{M_W}{64r} \left\{ -2M_W r(r^4 - 16r^2 + 36) + M_W r(r^4 - 16r^2 + 48) \left( \mathcal{L}_{M_W} - \frac{1}{\epsilon} \right) \right. \\ \left. - s(r^4 - 16r^2 + 56) \log \left( \frac{r(s - M_H)}{2M_W} + 1 \right) \right\} - ar^4 \frac{9M_W^2}{32} \left\{ 2 - \frac{\pi}{\sqrt{3}} \right. \\ \left. + \frac{1}{\epsilon} - \mathcal{L}_{M_W} - \log r \right\} + \mathcal{O}(\epsilon), \end{aligned} \quad (4.23)$$

up to  $\mathcal{O}(\epsilon)$ , where  $\epsilon$  are UV divergences. The results up to  $\mathcal{O}(\epsilon^2)$ , needed for mass renormalisation contributing at two-loop orders is provided in Appendix B. Whence the above provide us with the Higgs mass correction at two loop order. Note again the lack of contributions from massive fermion and scalar loops due to the corrections being applied at a scale where fermions and scalars are integrated out. Moreover, we note that the self-energy diagrams with tadpoles have been omitted for the same reason previously described.

### 4.2.4 Fermion and Scalar Mass Renormalisation

Lastly we discuss the mass renormalisation of the massive external fermion and scalar fields we consider. These masses appear and the corrections contribute at two loop order in the threshold regime at the scale where the Higgs and gauge masses are taken to be IR and vanish. We begin with the scalar contributions; the relation between the bare scalar mass,  $m_{\chi_0}$ , and the renormalised mass,  $m_\chi$ , is determined by the scalar self-energy corrections,  $\tilde{\Sigma}(p^2) = i\Sigma(p^2)\mathbf{1}$ . Extracting  $\Sigma(p^2)$  gives the renormalised mass by setting  $\delta m_\chi^2 = \Sigma(m_\chi^2)$ , which has the following form after combining all contributions,

$$\delta m_\chi^2 = ae^{\gamma_E \epsilon} \left( \frac{\mu^2}{m_\chi^2} \right)^\epsilon C_A \left\{ C_F m_\chi^2 \frac{(2\epsilon - 3)\Gamma(\epsilon - 1)}{2\epsilon - 1} + Y_s^2 \frac{\Gamma(\epsilon)}{4 - 8\epsilon} \right\}, \quad (4.24)$$

where  $\epsilon$  are UV divergences. Note that the dimensions of the second term match the first by definition of  $Y_s$  in (2.11). On the other hand, the relation between the bare fermion mass,  $m_{\psi_0}$ , and the renormalised mass,  $m_\psi$ , is determined by the fermion self-energy corrections,

$$\tilde{\Sigma}(p^2) = i(\Sigma^V(p^2)\not{p} + \Sigma^S(p^2)m_\psi)\mathbf{1}, \quad (4.25)$$

where the superscripts,  $S$  and  $V$ , label the scalar and vector contributions. Extracting  $\Sigma^{S,V}(p^2)$  gives the renormalized mass by setting  $\delta m_\psi = m_\psi(\Sigma^V(m_\chi^2) + \Sigma^S(m_\chi^2))$ , which has the following form after combining all contributions,

$$\delta m_\psi = ae^{\gamma_E \epsilon} \left(\frac{\mu^2}{m_\psi^2}\right)^\epsilon C_A \left\{ C_F \frac{(2\epsilon - 3)\Gamma(\epsilon)}{2\epsilon - 1} - \frac{Y_f^2}{8} \left( \Gamma(\epsilon - 1) + \frac{4\Gamma(\epsilon)}{1 - 2\epsilon} \right) \right\}, \quad (4.26)$$

where again  $\epsilon$  are UV divergences. In this case dimensions hold since  $Y_f$  is dimensionless as shown in (2.11). Note that the expansion up to  $\mathcal{O}(\epsilon)^3$  are needed for mass renormalisation. Now we have all the one-loop terms that arise in our problem which, when replacement rules are applied, contribute at the two-loop level.

### 4.3 Operator Renormalisation

Composite operators like ours require subsequent subtractions beyond wave-function renormalisation [52]. This holds for both full and effective theory operators, to illustrate, let us take, for instance, the bare heavy-light fermion operator from HPET,

$$\mathcal{O}^{(0)} = \bar{\psi}^{(0)}\Gamma h_f^{(0)} = \sqrt{Z_f Z_h}\bar{\psi}\Gamma h_f, \quad (4.27)$$

where  $\Gamma$  is an arbitrary Dirac matrix. The renormalised composite operator is then,

$$\begin{aligned} \mathcal{O} &= Z_{\mathcal{O}}^{-1}\mathcal{O}^{(0)} = \frac{\sqrt{Z_f Z_h}}{Z_{\mathcal{O}}}\bar{\psi}\Gamma h_f \\ &= \bar{\psi}\Gamma h_f + \text{counter term}, \end{aligned} \quad (4.28)$$

such that the additional operator,  $Z_{\mathcal{O}}$ , is determinable by computing a Green's function with an insertion of  $\mathcal{O}$ . Therefore,  $Z_{\mathcal{O}}$  can be found by taking the one particle irreducible Green's function of  $\bar{\psi}$ ,  $h_f$  and  $\mathcal{O}$ , where the counter term in (4.28) contributes,

$$\left( \frac{\sqrt{Z_f Z_h}}{Z_{\mathcal{O}}} - 1 \right) \Gamma, \quad (4.29)$$

to this time-ordered product. The vertex contribution also provides a UV divergent contribution to the time-ordered product. Consequently, the counter term, (4.28), must eliminate the divergences present in the vertex contribution and thus, (4.28) must be finite as  $\epsilon_{\text{UV}} \rightarrow 0$ . Plugging in the wave function contributions,  $\sqrt{Z_f Z_h}$  then gives  $Z_{\mathcal{O}}$  by the finiteness requirement. We then may obtain the anomalous dimension of the composite operator,

$$\gamma_{\mathcal{O}} = \frac{\mu}{Z_{\mathcal{O}}} \left( \frac{dZ_{\mathcal{O}}}{d\mu} \right), \quad (4.30)$$

from the renormalisation constant,

$$Z_{\mathcal{O}} = 1 + \delta Z_{\mathcal{O}} = 1 - \frac{1}{\epsilon_{\text{UV}}}\gamma_{\mathcal{O}} \quad (4.31)$$

Note in this case that the renormalisation of  $\mathcal{O}$  is independent of the gamma matrix of choice,  $\Gamma$ , in the composite operator. This is a consequence of heavy fermion spin symmetry and light fermion

chiral symmetry. In fact, this independence holds for all our effective operators in the threshold regime as they include operators with heavy/light fermions/scalars [52]. On the other hand, in the full theory as well as SCET the gamma matrix plays a role and  $Z_{\mathcal{O}}$  varies for different operators. In particular, in the full theory for both scalars and fermions, the scalar and tensor currents require renormalisation while the vector currents, at all orders, do not, meaning  $\delta Z_{\mathcal{O}}$  is null [60].

## 5 Radiative Corrections in Sudakov Limit

In this section, we calculate the form factor,  $\log F_E(Q^2)$ , in the large  $Q^2$ , or Sudakov, limit. We perform calculations up to two-loop order, extending previous studies and refraining from including computational details which have been presented in other works [17, 39].

### 5.1 Massless External Particles

Let us begin by considering the case of massless external particles in a fair amount of detail. The limit we consider is thus,  $Q^2 \gg M^2 \gg m^2$ , where  $M$  and  $m$  denote the bosonic and external masses, respectively. Schematically, in this case, the matching and running steps can be illustrated as follows,

$$\mathcal{O} \xleftarrow[m, M=0]{\mu \sim Q} e^C \tilde{\mathcal{O}}_1 \xrightarrow{\gamma_1} e^C \tilde{\mathcal{O}}_1 \xleftarrow[m=0]{\mu \sim M} e^{C+D} \tilde{\mathcal{O}}_2,$$

where  $C$  and  $D$  are multiplicative matching coefficients,  $\gamma_1$  the effective theory anomalous dimension and  $\tilde{\mathcal{O}}_{1,2}$  the effective theory operators at each scale. At scale,  $\mu > Q$ , we use the full theory, and at scale,  $\mu < Q$ , we match down to SCET with Wilson coefficient,  $c(\mu)$ . The RGE of  $c(\mu)$  is given by,

$$\mu \frac{dc(\mu)}{d\mu} = \gamma_F(a(\mu))c(\mu), \quad (5.1)$$

where  $\gamma_F(\mu)$  is the full theory anomalous dimension for a composite operator,  $\mathcal{O}$ . The full theory is matched onto SCET at a scale  $\mu \sim Q$ . The effective theory has modes with off-shellness of  $\mathcal{O}(Q)$  integrated out, so the matching coefficient depends on  $\mathcal{L}_Q$ , and these logarithms are not large if  $\mu \sim Q$ . The operator,  $\mathcal{O}$  in the full theory matches to the operator,  $\tilde{\mathcal{O}}$ , in SCET. More specifically,

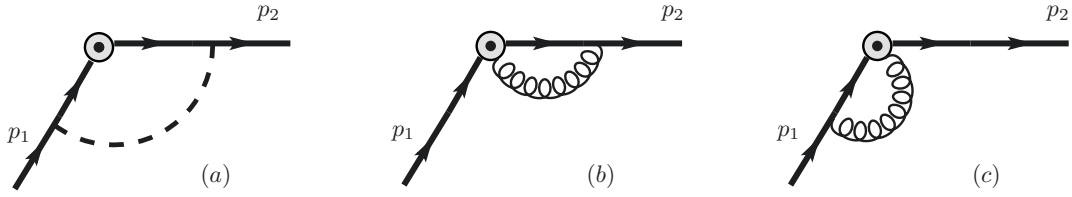
$$\bar{\psi}\Gamma\psi \rightarrow e^C(\bar{\xi}_{n,p_2} W_n)\Gamma(W_{\bar{n}}^\dagger \xi_{\bar{n},p_1}), \quad (5.2a)$$

$$\chi^\dagger\chi \rightarrow e^C(\Phi_{n,p_2}^\dagger W_n)(W_{\bar{n}}^\dagger \Phi_{\bar{n},p_1}), \quad (5.2b)$$

$$i\chi^\dagger \overleftrightarrow{D}_\mu \chi \rightarrow e^C(\Phi_{n,p_2}^\dagger W_n)[i\mathcal{D}_1 + i\mathcal{D}_2]_\mu(W_{\bar{n}}^\dagger \Phi_{\bar{n},p_1}), \quad (5.2c)$$

$$\bar{\psi}\chi \rightarrow e^C(\bar{\xi}_{n,p_2} W_n)(W_{\bar{n}}^\dagger \Phi_{\bar{n},p_1}), \quad (5.2d)$$

where  $i\mathcal{D}_1 = \mathcal{P} + g(n \cdot A_{\bar{n},q}) \frac{\bar{n}}{2}$ ,  $i\mathcal{D}_2 = \mathcal{P}^\dagger + g(\bar{n} \cdot A_{n,-q}) \frac{n}{2}$ ,  $\mathcal{P}$  are label operators in SCET and  $W_n$  is a Wilson line containing  $n$ -collinear gauge fields obtained by integrating over a path in the  $\bar{n}$ -direction [7]. Of course  $C(\mu)$  differs for each operator and we have written the multiplicative matching coefficient as  $\exp[C(\mu)]$  rather than  $C(\mu)$  for convenience. As is well-known, the matching coefficient can be computed as the finite part of the full theory matrix element, evaluated on-shell, with all IR scales, which in our case are the gauge and Higgs boson masses, are set to zero [51, 50, 49]. To illustrate the computation let us consider the one-loop result. The full and effective theory graphs to be evaluated are those in figure 2, except in SCET the external lines are both taken to be collinear and graphs (b) and (c) are no longer identical. After combining the vertex graphs with the wavefunction and tree-level graphs, one obtains the value of the full and effective theory matrix elements,  $\langle p_2 | \mathcal{O} | p_1 \rangle$  and  $\langle p_2 | \tilde{\mathcal{O}} | p_1 \rangle$ , respectively. The gauge and external particle masses are IR scales and can be set to zero in the matching, thus leading to scaleless integrals, for the EFT and wavefunction contributions. One then combines the vertex and wave-function contributions



**Figure 2:** One-loop vertex corrections, bulls-eye represents composite operator, arrowed lines represent all incoming-outgoing particles we consider, dashed lines correspond to bosonic propagators. (b), (c) only exists with the operator,  $\mathcal{O} = i\phi^\dagger \overleftrightarrow{D}_\mu \phi$ , and EFT equivalents,  $\tilde{\mathcal{O}}$ .

$\mathcal{O}$	$\gamma_F^{(1)}$	$C^{(1)}(\mu)$	$C^{(2)}(\mu)$	$\gamma_1^{(1)}$
$\bar{\psi}\psi$	$-3C_{ACF} + \frac{c_A Y_f^2}{8}$	$\frac{C_{ACF}}{6} \{-6\mathcal{L}_Q^2 + \pi^2 - 12\} + \frac{C_A Y_f^2}{4} \{\mathcal{L}_Q - 2\}$	$V_1^{(Q)}$	$C_{ACF}(4\mathcal{L}_Q - 6)$
$\bar{\psi}\gamma^\mu\psi$	$-\frac{c_A Y_f^2}{4}$	$\frac{C_{ACF}}{6} \{-6\mathcal{L}_Q(\mathcal{L}_Q - 3)\mathcal{L}_Q + \pi^2 - 48\} - \frac{C_A Y_f^2}{8} \{\mathcal{L}_Q - 1\}$	$V_2^{(Q)}$	$C_{ACF}(4\mathcal{L}_Q - 6)$
$\bar{\psi}\sigma^{\mu\nu}\psi$	$C_{ACF} - \frac{c_A Y_f^2}{8}$	$\frac{C_{ACF}}{6} \{-6\mathcal{L}_Q(\mathcal{L}_Q - 4)\mathcal{L}_Q + \pi^2 - 48\} + \frac{C_A Y_f^2}{4}$	$V_3^{(Q)}$	$C_{ACF}(4\mathcal{L}_Q - 6)$
$\chi^\dagger\chi$	$-3C_{ACF}$	$\frac{C_{ACF}}{6} \{-6\mathcal{L}_Q(\mathcal{L}_Q - 1)\mathcal{L}_Q + \pi^2 - 12\}$	$V_4^{(Q)}$	$C_{ACF}(4\mathcal{L}_Q - 8)$
$i\chi^\dagger \overleftrightarrow{D}_\mu \chi$	0	$\frac{C_{ACF}}{6} \{-6\mathcal{L}_Q(\mathcal{L}_Q - 4)\mathcal{L}_Q + \pi^2 - 48\}$	$V_5^{(Q)}$	$C_{ACF}(4\mathcal{L}_Q - 8)$
$\bar{\psi}\chi, \chi^\dagger\psi$	$-\frac{3}{2}C_{ACF} - \frac{c_A Y_f^2}{16}$	$\frac{C_{ACF}}{6} \{-6\mathcal{L}_Q(\mathcal{L}_Q - 3)\mathcal{L}_Q + \pi^2 - 36\}$	$V_6^{(Q)}$	$C_{ACF}(4\mathcal{L}_Q - 7)$

**Table 2:** Matching corrections,  $C(\mu)$ , to the Sudakov form-factor at  $\mu \sim Q$ .  $V_i^{(Q)}$  are two loop vertex corrections, given in Appendix C.  $\gamma_F$  and  $\gamma_1$  are the full theory and SCET anomalous dimension,  $a \equiv \alpha/(4\pi)$ , and  $\mathcal{L}_Q \equiv \log Q^2/\mu^2$ .

as prescribed in (4.3) to obtain the one and two-loop order results. Moreover, as the masses are zero there are no two-loop contributions from mass renormalisation, only coupling renormalisation contributes. Scaleless integrals are set to zero in dimensional regularization, so the EFT matrix element is equal to its tree-level value. The full theory and EFT operators,  $\mathcal{O}$  and  $\tilde{\mathcal{O}}$ , are normalised to have the same tree-level value [51], thus,

$$\exp[C(\mu)] = \frac{\langle p_2 | \mathcal{O} | p_1 \rangle}{\langle p_2 | \mathcal{O} | p_1 \rangle_{tree}}. \quad (5.3)$$

When computing the one loop graphs for  $\mathcal{O}$ ,  $\exp[C(\mu)]$  is given by the on-shell full theory matrix element, normalised by its tree-level value. The particle masses are all much smaller than  $Q^2$ , and only contribute  $M^2/Q^2$  (where  $M$  corresponds to the gauge and Higgs masses) power corrections at the large scale,  $Q$ , which are being neglected. The one-loop values of  $C(\mu)$  for the other cases are computed similarly, and are given in table 2, where in the loop expansion,  $C(\mu) = aC^{(1)}(\mu) + a^2C^{(2)}(\mu) + \mathcal{O}(a^3)$ . Large logarithms do not appear if the matching scale  $\mu \sim Q$ , in this work we choose  $\mu = Q$  and the RGE of  $c(\mu)$  in the EFT is given by the anomalous dimension,  $\gamma_1$ , of  $\tilde{\mathcal{O}}$  in SCET. The full theory anomalous dimension,  $\gamma_F$ , of  $\tilde{\mathcal{O}}$  is also given in 2, we avoid presenting the two loop result as this has been previously found for a number of operators [48]. On the other hand, the SCET anomalous dimension,  $\gamma_1$ , is used to evolve  $c(\mu)$  from  $\mu = Q \rightarrow M$ . As previously defined, the anomalous dimension is given by the UV counter terms for the SCET graphs, and can depend on  $Q$ , the largest scale. UV divergences are independent of IR properties and  $\gamma_1$  is linear in  $\log \mu^2/Q^2$  to all order [49, 4], so one can always write,

$$\gamma_1(\mu) = A(\alpha(\mu)) \log \frac{\mu^2}{Q^2} + B(\alpha(\mu)). \quad (5.4)$$

$\mathcal{O}$	$\Delta C^{(2)}(\mu)$
$\bar{\psi}\psi$	$\frac{C_A C_F}{36} \{-2\mathcal{L}_Q^3 + (\pi^2 - 12)\mathcal{L}_Q - 28\zeta_3 + 24\} (C_A(3C_A + 22) - 8n_f T_f - 4) + \frac{C_A Y_f^2}{288} \{6\mathcal{L}_Q(\mathcal{L}_Q(-2C_A C_F \log \mathcal{L}_Q + C_A(6C_F + 22) - 8n_f T_f - 1) + C_A((\pi^2 - 36)C_F - 88) + 32n_f T_f + 4) - 168\zeta_3 C_A C_F + 48C_A(9C_F + 22) - 2\pi^2 C_A(3C_F + 11) + (\pi^2 - 48)(8n_f T_f + 1)\}$
$\bar{\psi}\gamma^\mu\psi$	$\frac{C_A C_F}{72} \{2\mathcal{L}_Q((9 - 2\mathcal{L}_Q)\mathcal{L}_Q + \pi^2 - 48) - 56\zeta_3 - 3\pi^2 + 192\} (C_A(3C_A + 22) - 8n_f T_f - 4) + \frac{C_A Y_f^2}{576} \{6\mathcal{L}_Q^2(2C_A(6C_F - 11) + 8n_f T_f + 1) + 12\mathcal{L}_Q(C_A((\pi^2 - 42)C_F + 22) - 8n_f T_f - 1) - 24C_A C_F \mathcal{L}_Q^3 - 12(28\zeta_3 - 90 + \pi^2)C_A C_F + 22(\pi^2 - 12)C_A - (\pi^2 - 12)(8n_f T_f + 1)\}$
$\bar{\psi}\sigma^{\mu\nu}\psi$	$\frac{C_A C_F}{36} \{\mathcal{L}_Q(-2(\mathcal{L}_Q - 6)\mathcal{L}_Q + \pi^2 - 48) - 28\zeta_3 - 2\pi^2 + 84\} (C_A(3C_A + 22) - 8n_f T_f - 4) + \frac{C_A Y_f^2}{48} \{\mathcal{L}_Q(-2C_A C_F(\mathcal{L}_Q - 6)\mathcal{L}_Q + C_A((\pi^2 - 36)C_F + 44) - 2(8n_f T_f + 1)) - 28\zeta_3 C_A C_F - 2C_A((\pi^2 - 24)C_F + 66) + 48n_f T_f + 6\}$
$\chi^\dagger\chi$	$\frac{C_A C_F}{72} \{2\mathcal{L}_Q((3 - 2\mathcal{L}_Q)\mathcal{L}_Q + \pi^2 - 12) - 56\zeta_3 - \pi^2 + 48\} (-6C_A^2 + 22C_A - 8n_f T_f + 5)$
$i\chi^\dagger \overleftrightarrow{D}_\mu \chi$	$\frac{C_A C_F}{36} \{\mathcal{L}_Q(-2(\mathcal{L}_Q - 6)\mathcal{L}_Q + \pi^2 - 48) - 28\zeta_3 - 2\pi^2 + 96\} (-6C_A^2 + 22C_A - 8n_f T_f + 5)$
$\bar{\psi}\chi, \chi^\dagger\psi$	$\frac{C_A C_F}{576} \{-2\mathcal{L}_Q((9 - 2\mathcal{L}_Q)\mathcal{L}_Q + \pi^2 - 36) + 56\zeta_3 + 3\pi^2 - 144\} (C_A(-12C_A + 3Y_f^2 - 176) + 64n_f T_f + 20)$

**Table 3:** Mass and coupling corrections to matching which contribute at two-loop order,  $\Delta C^{(2)}(\mu)$  is the correction at  $\mu \sim Q$ .

The loop expansion of the anomalous dimension,  $\gamma_1 = a\gamma_1^{(1)} + a^2\gamma_2^{(2)} + \mathcal{O}(a^3)$ , is given for each operator in table 2. By inspection, the SCET anomalous dimension,  $\gamma_1$ , depends solely on the external fields for the operators, as in it is equal for the three fermion and two scalar operators, respectively, as well as being the average of the two field's result for the mixed operator. The reason being that the effective theory anomalous dimension depends on the IR divergence of the full theory graph, and the IR divergence is independent of the vertex factors.

The next step matching step occurs at the lower scale,  $\mu \sim M$ , where the massive bosons are integrated out. The matching is done from SCET with massive bosons ( $\mu > M$ ), to SCET without massive bosons ( $\mu < M$ ). In our model, this is a free theory, so there is no need for propagating bosonic modes below  $M$ . The matching coefficient at  $\mu \sim M$  is given by  $d(\mu) = \exp[D(\mu)]$  in table 4 and is found from the SCET vertex and wave-function corrections. More specifically, one matches in the following way,

$$e^C(\bar{\xi}_{n,p_2} W_n) \Gamma(W_n^\dagger \xi_{\bar{n},p_1}) \rightarrow e^{C+D} \bar{\xi}_{n,p_2} \Gamma \xi_{\bar{n},p_1}, \quad (5.5a)$$

$$e^C(\Phi_{n,p_2}^\dagger W_n)(W_n^\dagger \Phi_{\bar{n},p_1}) \rightarrow e^{C+D} \Phi_{n,p_2}^\dagger \Phi_{\bar{n},p_1}, \quad (5.5b)$$

$$e^C(\Phi_{n,p_2}^\dagger W_n)[i\mathcal{D}_1 + i\mathcal{D}_2]_\mu(W_n^\dagger \Phi_{\bar{n},p_1}) \rightarrow e^{C+D} \Phi_{n,p_2}^\dagger i(\mathcal{P}^\dagger + \mathcal{P})_\mu \Phi_{\bar{n},p_1}, \quad (5.5c)$$

$$e^C(\bar{\xi}_{n,p_2} W_n)(W_n^\dagger \Phi_{\bar{n},p_1}) \rightarrow e^{C+D} \bar{\xi}_{n,p_2} \Phi_{\bar{n},p_1}. \quad (5.5d)$$

As for the results, although we calculate up to two loops fully for  $c(\mu)$ , we do not calculate the two-loop vertex contribute to  $d(\mu)$  due to the complexity of massive SCET integrals. Hence, mass and coupling renormalisation is not necessary as they only affect the next order, nonetheless, we present these sub-divergent  $\mathcal{O}(a^2)$  contributions for both  $c(\mu)$  and  $d(\mu)$  in tables 3 and 4. Moreover, by inspection of table 1, we do not include the collinear correction to the particle propagator for each case as it is the same as in the full theory [7]. The ultrasoft correction vanishes, so the wavefunction corrections are the same as in the full theory and we have these up to two loops. For a more detailed description on the specific one-loop SCET integrals, we point to previous work [17, 8], and it would be interesting to calculate the SCET vertex contributions at two-loop to have a complete account at this order. The above matching steps are identical at each order, and the two-loop vertex and wavefunction graphs we calculated are shown in figures 3 and 1. Furthermore, we note that both in the massive and massless external particle cases of SCET, there is no Higgs contributions in the vertex corrections. This is because the fermion Yukawa vertex vanishes, as by construction, (2.30) implies that,

$$\bar{\xi}_{n,p} \xi_{n,p} = \bar{\xi}_{n,p} \frac{\not{p} \not{p}}{4} \frac{\not{p} \not{p}}{4} \xi_{n,p} = 0, \quad (5.6)$$

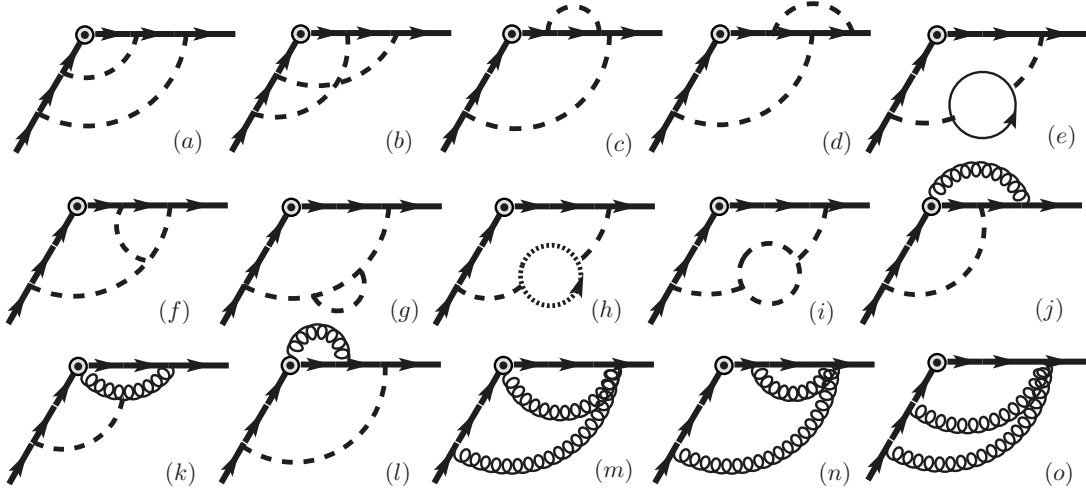
$\mathcal{O}$	$D^{(1)}(\mu)$	$\Delta D^{(2)}(\mu)$
$\bar{\psi}\Gamma\psi$	$\frac{C_A C_F}{6} \{-6\mathcal{L}_{M_W}^2 + 12\mathcal{L}_{M_W}\mathcal{L}_Q - 18\mathcal{L}_{M_W} + 27 - 5\pi^2\}$	$-\frac{C_A C_F}{72M_W^6} \{M_W^6 ((99\sqrt{3}\pi - 690)C_A + 32(5 + 3i\pi)n_f T_f + 124)$ $+ 3M_W^4 (\mathcal{L}_{M_W} (M_W^2 (141C_A - 32n_f T_f - 20) + 10M_H^2) - 8M_H^2 \mathcal{L}_{M_H})$ $- 54M_W^4 M_H^2 + 6M_H^2 M_H^2 - 6M_H s (12M_W^2 - 4M_W^2 M_H^2 + M_H^2) \log(w)$ $+ 6(8M_W^4 M_H^2 - 6M_W^2 M_H^2 + M_H^2) \log(M_W/M_H)\} (2\mathcal{L}_{M_W} - 2\mathcal{L}_Q + 3)$
$\chi^\dagger\chi, i\chi^\dagger\overleftrightarrow{D}_\mu\chi$	$\frac{C_A C_F}{6} \{-6\mathcal{L}_{M_W}^2 + 12\mathcal{L}_{M_W}\mathcal{L}_Q - 24\mathcal{L}_{M_W} + 21 - 5\pi^2\}$	$-\frac{C_A C_F}{36M_W^6} \{3M_W^4 (\mathcal{L}_{M_W} (M_W^2 (141C_A - 32n_f T_f - 20) + 10M_H^2) - 8M_H^2 \mathcal{L}_{M_H})$ $+ M_W^6 ((99\sqrt{3}\pi - 690)C_A + 32(5 + 3i\pi)n_f T_f + 124) + 6M_H^4 M_W^2$ $- 54M_H^2 M_W^4 - 6M_H s (M_H^2 - 4M_H^2 M_W^2 + 12M_W^2) \log(w)$ $+ 6(M_H^6 - 6M_H^4 M_W^2 + 8M_H^2 M_W^4) \log(M_W/M_H)\} (\mathcal{L}_{M_W} - \mathcal{L}_Q + 2)$
$\bar{\psi}\chi, \chi^\dagger\psi$	$\frac{C_A C_F}{6} \{-6\mathcal{L}_{M_W}^2 + 12\mathcal{L}_{M_W}\mathcal{L}_Q - 21\mathcal{L}_{M_W} + 24 - 5\pi^2\}$	$-\frac{C_A C_F}{144M_W^6} \{3M_W^4 (\mathcal{L}_{M_W} (M_W^2 (141C_A - 32n_f T_f - 20) + 10M_H^2) - 8M_H^2 \mathcal{L}_{M_H})$ $+ M_W^6 ((99\sqrt{3}\pi - 690)C_A + 32(5 + 3i\pi)n_f T_f + 124) + 6M_H^4 M_W^2$ $- 54M_H^2 M_W^4 - 6M_H s (M_H^2 - 4M_H^2 M_W^2 + 12M_W^2) \log(w)$ $+ 6(M_H^6 - 6M_H^4 M_W^2 + 8M_H^2 M_W^4) \log(M_W/M_H)\} (4\mathcal{L}_{M_W} - 4\mathcal{L}_Q + 7)$

**Table 4:** SCET contribution to the Sudakov form-factor at  $\mu \sim M$ , one-loop matching coefficient,  $D^{(1)}(\mu)$ , two-loop mass and coupling renormalisation correction,  $\Delta D^{(2)}(\mu)$ .

using the identity,  $\not{n}\not{n} = n^2 = 0$ . Moreover the tri-scalar couplings have dimension of mass and Higgs exchange corrections to the scalar operators are suppressed by powers of  $Y_s/Q$ , which is sub-leading in SCET power counting and we drop such terms. This is easily seen when using the re-scaled fields,  $\phi_{n,p}$ , which have a propagator of identical form to those of fermions. Then the Yukawa coupling becomes,

$$Y_s H \chi^\dagger \chi \rightarrow Y_s H \Phi_{n,p}^\dagger \Phi_{n,p} = \frac{Y_s}{\bar{n} \cdot p} H \phi_{n,p}^\dagger \phi_{n,p}, \quad (5.7)$$

which is  $\mathcal{O}(1/Q)$  as  $\bar{n} \cdot p$  is of order  $\mathcal{O}(Q)$  which suppresses any graph at each tri-scalar coupling. Thus, the only scalar graphs which appear are the matching at  $Q$ , which are full theory graphs as well as scalar contributions to the wave-function renormalisation in the effective theories.



**Figure 3:** Two-loop vertex correction graphs, (a)-(d) are Abelian corrections; (f)-(i) are non-Abelian, (j)-(m) only exists with the operator,  $\mathcal{O} = i\phi^\dagger \overleftrightarrow{D}_\mu \phi$  and EFT equivalents, (m)-(o) are seagull terms and occur only for scalar fields.

$\mathcal{O}$	$R^{(1)}(\mu)$	$\gamma_2^{(1)}(\mu)$	$T^{(1)}(\mu)$	$\gamma_3^{(1)}(\mu)$
$\bar{\psi}_2 \Gamma \psi_1$	$\frac{C_A C_F}{2} \left\{ \mathcal{L}_{m_2}^2 - \mathcal{L}_{m_2} + \frac{\pi^2}{6} + 4 \right\}$	$C_A C_F \{4\mathcal{L}_Q - 2\mathcal{L}_{m_2} - 5\}$	$\frac{C_A C_F}{2} \left\{ \mathcal{L}_{m_1}^2 - \mathcal{L}_{m_1} + \frac{\pi^2}{6} + 4 \right\}$	$4C_A C_F \{wh(w) - 1\}$
$\chi_2^\dagger \chi_1, i\chi_2^\dagger \overleftrightarrow{D}_\mu \chi_1$	$\frac{C_A C_F}{2} \left\{ \mathcal{L}_{m_2}^2 - 2\mathcal{L}_{m_2} + \frac{\pi^2}{6} + 4 \right\}$	$C_A C_F \{4\mathcal{L}_Q - 2\mathcal{L}_{m_2} - 6\}$	$\frac{C_A C_F}{2} \left\{ \mathcal{L}_{m_1}^2 - 2\mathcal{L}_{m_1} + \frac{\pi^2}{6} + 4 \right\}$	$4C_A C_F \{wh(w) - 1\}$
$\bar{\psi}_2 \chi_1$	$\frac{C_A C_F}{2} \left\{ \mathcal{L}_{m_2}^2 - \mathcal{L}_{m_2} + \frac{\pi^2}{6} + 4 \right\}$	$C_A C_F \{4\mathcal{L}_Q - 2\mathcal{L}_{m_2} - 6\}$	$\frac{C_A C_F}{2} \left\{ \mathcal{L}_{m_1}^2 - 2\mathcal{L}_{m_1} + \frac{\pi^2}{6} + 4 \right\}$	$4C_A C_F \{wh(w) - 1\}$
$\chi_2^\dagger \psi_1$	$\frac{C_A C_F}{2} \left\{ \mathcal{L}_{m_2}^2 - 2\mathcal{L}_{m_2} + \frac{\pi^2}{6} + 4 \right\}$	$C_A C_F \{4\mathcal{L}_Q - 2\mathcal{L}_{m_2} - 5\}$	$\frac{C_A C_F}{2} \left\{ \mathcal{L}_{m_1}^2 - \mathcal{L}_{m_1} + \frac{\pi^2}{6} + 4 \right\}$	$4C_A C_F \{wh(w) - 1\}$

**Table 5:** Matching and running results for  $Q \gg m_2 \gg m_1 \gg M$ .  $R$  and  $T$  are the matching coefficients at  $\mu \sim m_2$  and  $\mu \sim m_1$ ,  $\gamma_2$  is the anomalous dimension between  $m_2$  and  $m_1$ ,  $\gamma_3$  is the anomalous dimension between  $m_1$  and  $M$ .  $R$  and  $T$  only depend on whether the light particle is a fermion or scalar.

## 5.2 Massive External Particles

Thus far our results have been computed for external particles with masses,  $m_{1,2}$ , much smaller than the bosonic masses. In this section, we consider the Sudakov regime for massive external particles, extending previous results. Thus, we are primarily interested in the limits,  $Q \gg m_{1,2} \gg M$ , although we will discuss other cases that can be studied as well, in particular one that can be applied for LHC studies of the top quark.

There are two cases to consider,  $Q \gg m_2 \gg m_1 \gg M$  and  $Q \gg m_2 \sim m_1 \gg M$ , we begin with the former. Again, the Sudakov form-factor can be computed using a sequence of effective field theories [45]. One begins as in the massless external particle case by matching the full theory onto SCET with a single massive particle at the scale,  $\mu \sim Q$ . The same operators are matched to as in (5.2), except now the  $n$ -collinear SCET field,  $\xi_{n,p_2}$ , is taken to have mass,  $m_2$ . Again, this matching is independent of IR scales much smaller than  $Q$ , and thus is given by  $\exp[C(\mu)]$ , as presented in table 2. The next step is to run the operator from the scale  $Q$  to  $m_2$ , which can be done with the anomalous dimension,  $\gamma_1$ , given in 2, as the anomalous dimension is also independent of the lower mass scales. The matching steps that follow lie at scales  $\mu = m_2$ ,  $\mu = m_1$  and  $\mu = M$ . Schematically, the matching and running steps can be illustrated as follows,

$$e^C \tilde{\mathcal{O}}_1 \xleftarrow[\substack{\mu \sim m_2 \\ m_1, M=0}]{\mu \sim m_2} e^{C+R} \tilde{\mathcal{O}}_2 \xrightarrow{\gamma_2} e^{C+R} \tilde{\mathcal{O}}_2 \xleftarrow[\substack{\mu \sim m_1 \\ M=0}]{\mu \sim m_1} e^{C+R+T} \tilde{\mathcal{O}}_3 \xrightarrow{\gamma_3} e^{C+R+T} \tilde{\mathcal{O}}_3 \xleftarrow[\substack{\mu \sim M \\ M \neq 0}]{\mu \sim M} e^{C+R+T+U} \tilde{\mathcal{O}}_4,$$

where the exponents are multiplicative matching coefficients,  $\gamma_i$  the effective theory anomalous dimensions and  $\tilde{\mathcal{O}}_i$  the effective theory operators at each scale.

Firstly, for the matching step at  $\mu = m_2$ , one switches from SCET to a new EFT with the massive particle described by a heavy field [52],  $h_{f,s}$ , with a velocity,  $v_2$ , such that  $v_2^2 = 1$ . Whereas, the other particle remaining massless continues to be described by the  $\bar{n}$ -collinear SCET field,  $\xi_{\bar{n},p_1}$ . The fermionic operators, for instance, are then given by  $\bar{h}_{f,1} \Gamma W_{\bar{n}}^\dagger \xi_{\bar{n},p_1}$ , and similarly for other operators [30]. The matching correction at  $\mu = m_2$  is then given by the difference between the vertex graphs in figures 2 and 3, for the corresponding external particles in the effective theories above and below  $m_2$ . More specifically, in the fermion example, the difference between graphs where  $\xi_{n,p_2}$  and  $h_{f,2}$ , for the particle with mass,  $m_2$ . Note that in the theory below  $m_2$  there are no graphs with collinear Wilson lines associated with  $h_{f,s}$  and thus such corrections do not appear. The graphs in the theory above  $m_2$  are evaluated with bosonic masses set to zero, as  $m_2 \gg M$ , and on-shell at  $p_2^2 = m_2^2$ . Below  $m_2$  the graphs in the effective theory are evaluated at  $M = 0$  as well, at the on-shell point,  $k_2 \cdot v_2 = 0$  where  $k_2$  is the residual momentum of the heavy particle.

As for the wave-function graphs, the  $\xi_{\bar{n},p_1}$  and HQET graphs both vanish on-shell. Hence, the matching is given by the vertex correction and the on-shell wavefunction graph for  $\xi_{n,p_2}$ , the results of which are discussed in detail in previous work [17], and show in tables 5 and 1, respectively. We proceed then with next matching step with coefficient,  $\exp[T(\mu)]$ , at the scale of the lower particle mass,  $\mu \sim m_1$ . At this scale, the theory above  $m_1$  is SCET with heavy field for particle with mass,  $m_2$ , and the theory below  $m_1$ , the  $\bar{n}$ -collinear SCET field,  $\xi_{\bar{n},p_1}$  is replaced by the heavy field,  $h_{f,s}$ , with velocity,  $v_1$ , such that  $v_1^2 = 1$  and  $v_1 \cdot v_2 = w$ . The fermionic operators, for example, are then given by  $\bar{h}_{f,2} \Gamma h_{f,1}$  instead of  $\bar{h}_{f,2} \Gamma W_{\bar{n}}^\dagger \xi_{\bar{n},p_1}$ .



$\mathcal{O}$	$U^{(1)}(\mu)$	$U^{(2)}(\mu)$
$\bar{\psi}\Gamma\psi$	$\frac{C_A}{4}\{8C_F(wh(w)-1)\mathcal{L}_{M_W} - Y_f^2(h(w)-1)\mathcal{L}_{M_H}\}$	$(V_1^{(M)} + F_h^{(M)}) + \frac{2}{3}wC_A^2C_F^2(12\mathcal{L}_{M_W}-1)\mathcal{L}_{M_W} + \pi^2)h(w) + \frac{1}{24}C_A^2Y_f^4(12\mathcal{L}_{M_H}-1)\mathcal{L}_{M_H} + \pi^2)h(w) - \frac{1}{6}C_A^2C_FY_f^2h(w)(6\mathcal{L}_{M_W}((w+1)\mathcal{L}_{M_H}-2) + 3(w+1)\mathcal{L}_{M_W}^2 + 3\mathcal{L}_{M_H}((w+1)\mathcal{L}_{M_H}-4w) + \pi^2(w+1))$
$\chi^\dagger\chi, i\chi^\dagger\overleftrightarrow{D}_\mu\chi$	$\frac{C_A}{4}\{8C_F(wh(w)-1)\mathcal{L}_{M_W} - Y_s^2(h(w)-1)\mathcal{L}_{M_H}\}$	$(V_2^{(M)} + F_h^{(M)}) + \frac{2}{3}wC_A^2C_F^2(12\mathcal{L}_{M_W}-1)\mathcal{L}_{M_W} + \pi^2)h(w) + \frac{1}{24}C_A^2Y_s^4(12\mathcal{L}_{M_H}-1)\mathcal{L}_{M_H} + \pi^2)h(w) - \frac{1}{6}C_A^2C_FY_s^2h(w)(6\mathcal{L}_{M_W}((w+1)\mathcal{L}_{M_H}-2) + 3(w+1)\mathcal{L}_{M_W}^2 + 3\mathcal{L}_{M_H}((w+1)\mathcal{L}_{M_H}-4w) + \pi^2(w+1))$
$\chi^\dagger\psi, \bar{\psi}\chi$	$\frac{C_A}{4}\{8C_F(wh(w)-1)\mathcal{L}_{M_W} - (Y_fY_s h(w) - \frac{1}{2}(Y_f^2 + Y_s^2))\mathcal{L}_{M_H}\}$	$(V_3^{(M)} + F_h^{(M)}) + \frac{2}{3}wC_A^2C_F^2(12\mathcal{L}_{M_W}-1)\mathcal{L}_{M_W} + \pi^2)h(w) + \frac{1}{48}C_A^2(Y_f^2Y_s + Y_s^2Y_f)(12\mathcal{L}_{M_H}-1)\mathcal{L}_{M_H} + \pi^2) - \frac{1}{6}C_A^2C_FY_fY_s h(w)(3(2\mathcal{L}_{M_W}(\mathcal{L}_{M_H}-2) + \mathcal{L}_{M_W}^2 + \mathcal{L}_{M_H}^2) + \pi^2) - \frac{1}{12}wC_A^2C_F(Y_f^2 + Y_s^2)h(w)(3(\mathcal{L}_{M_W} + \mathcal{L}_{M_H})^2 - 12\mathcal{L}_{M_H} + \pi^2)h(w)$

**Table 6:** One and two-loop matching contribution,  $U^{(1,2)}$ , at  $\mu \sim M$ .  $V_i^{(M)}$  and  $F_i$  are two loop vertex and wave-function corrections, given in Appendices C and D.

In the theory below  $m_1$  there are no vertex corrections due to collinear Wilson lines, as there are no collinear wilson lines,  $W$ , associated with heavy fields. The matching contribution is again given by the difference of the vertex and wavefunction graphs in the theories above and below  $m_1$ , setting all scales less than  $m_1$  to zero. Note that in the theory below  $m_1$  we have scaleless integrals which vanish trivially and thus, the sole non-zero contributions come from vertex contributions above  $m_1$  and the  $\bar{n}$ -collinear wavefunction graph. Conveniently, these are the same graphs that contribute to the matching condition at  $m_2$ , so  $T$  is given by  $R$  with  $m_2 \rightarrow m_1$ , and is presented in table 5.

The final contributions needed are the anomalous dimension,  $\gamma_3$ , for the running between  $m_1$  and  $M$ , and the matching condition,  $\exp[U(\mu)]$ , at  $M$ . These can be computed from the HPET graphs, evaluated on-shell, but now with bosonic masses,  $M$ , included as they are no longer IR. The one-loop contributions are presented in the last column of table 5, which by inspection are independent of the composite operator. Whence, they are identical and depend solely on whether the external particles are fermions or scalars, and their dependence lies in the appropriate Yukawa factors. The function,

$$h(w) = \frac{\log(w + \sqrt{w^2 - 1})}{\sqrt{w^2 - 1}}, \quad (5.8)$$

is the well-known factor which occurs in the velocity-dependent anomalous dimension in HQET [52]. Note further that in the Sudakov regime, the Higgs contribution in  $\exp U$  is sub-leading as,  $Q^2 \sim m_1 m_2 w$ , and in this limit,

$$h(w) \sim \frac{\log w}{w}, \quad (5.9)$$

thus the gauge contribution dominates in the Sudakov regime. We will see later on that in the threshold regime, the Higgs and gauge contributions turn out to be on equal footing. We also present the two-loop contribution to the matching contribution,  $\exp[U(\mu)]$ , in table 6 which combines the vertex and wave-function contribution listed in table 1.

As for remaining two-loop contributions, we present the mass and coupling renormalisation, which contribute at two-loop order for each matching coefficient in table 7. The situation is similar in the case  $Q \gg m_2 \sim m_1 \gg M$ , which is why we left this for last. The evolution down to the scale  $m_1 \sim m_2$  is the same as for the case where  $m_i = 0$ . The  $n$  and  $\bar{n}$  collinear graphs at the scale  $m_1 \sim m_2$  are independent

$\mathcal{O}$	$\Delta R^{(2)}$	$\Delta T^{(2)}$	$\Delta U^{(2)}$
$\bar{\psi}_2\Gamma\psi_1$	$\frac{C_A^2 C_F}{2}\left\{C_F(1-2\mathcal{L}_{m_2})(4-3\mathcal{L}_{m_2}) - \frac{Y_f^2}{8}(1-2\mathcal{L}_{m_2})(7-3\mathcal{L}_{m_2})\right\}$	$\frac{C_A^2 C_F}{2}\left\{C_F(1-2\mathcal{L}_{m_1})(4-3\mathcal{L}_{m_1}) - \frac{Y_f^2}{8}(1-2\mathcal{L}_{m_1})(7-3\mathcal{L}_{m_1})\right\}$	$\Delta U_1^{(2)}$
$\chi_2^\dagger\chi_1, i\chi_2^\dagger\overleftrightarrow{D}_\mu\chi_1$	$C_A^2 C_F\left\{C_F(1-\mathcal{L}_{m_2})(7-3\mathcal{L}_{m_2}) - \frac{Y_s^2}{4m_2^2}(1-2\mathcal{L}_{m_2})(2-\mathcal{L}_{m_2})\right\}$	$C_A^2 C_F\left\{C_F(1-\mathcal{L}_{m_1})(7-3\mathcal{L}_{m_1}) - \frac{Y_s^2}{4m_1^2}(1-2\mathcal{L}_{m_1})(2-\mathcal{L}_{m_1})\right\}$	$\Delta U_2^{(2)}$
$\bar{\psi}_2\chi_1$	$\frac{C_A^2 C_F}{2}\left\{C_F(1-2\mathcal{L}_{m_2})(4-3\mathcal{L}_{m_2}) - \frac{Y_f^2}{8}(1-2\mathcal{L}_{m_2})(7-3\mathcal{L}_{m_2})\right\}$	$C_A^2 C_F\left\{C_F(1-\mathcal{L}_{m_1})(7-3\mathcal{L}_{m_1}) - \frac{Y_s^2}{4m_1^2}(1-2\mathcal{L}_{m_1})(2-\mathcal{L}_{m_1})\right\}$	$\Delta U_3^{(2)}$
$\chi_2^\dagger\psi_1$	$C_A^2 C_F\left\{C_F(1-\mathcal{L}_{m_2})(7-3\mathcal{L}_{m_2}) - \frac{Y_s^2}{4m_2^2}(1-2\mathcal{L}_{m_2})(2-\mathcal{L}_{m_2})\right\}$	$\frac{C_A^2 C_F}{2}\left\{C_F(1-2\mathcal{L}_{m_1})(4-3\mathcal{L}_{m_1}) - \frac{Y_f^2}{8}(1-2\mathcal{L}_{m_1})(7-3\mathcal{L}_{m_1})\right\}$	$\Delta U_3^{(2)}$

**Table 7:** Matching and contribution due to mass and coupling renormalisation for  $Q \gg m_2 \gg m_1 \gg M$ .  $\Delta R^{(2)}$ ,  $\Delta T^{(2)}$  and  $\Delta U^{(2)}$  are the two-loop order matching coefficients at  $\mu \sim m_2$ ,  $\mu \sim m_1$  and  $\mu \sim M$ . The contributions,  $\Delta U_i^{(2)}$ , are presented in Appendix B.



of each other, so the matching is imply given by the sum of  $R$  and  $T$  at  $m_2$  and  $m_1$  respectively. Below  $m_1 \sim m_2$  the matching and running is identical to the previous case with anomalous dimension,  $\gamma_3$ , and matching coefficient,  $\exp[U(\mu)]$ . Lastly, if  $m_2 = m_1$ , then the case is identical to  $m_2 \sim m_1$ , except one sets  $m_2 = m_1$  in all matching and running contributions.

**Further Cases:** We note, as considered in previous work [17], that there are other cases one can compare for complete generality, in particular, one case resonates with regard to top quark physics in the high energy regime. The Sudakov limit being,  $Q \gg m_1 \sim m_2 \sim M$ , which involves one running step with  $\gamma_1$ , as the running from  $Q$  is independent of the IR scales, and two matching steps. The matching at  $\mu \sim Q$  is represented by the usual,  $\exp[C(\mu)]$ . On the other hand, the matching at  $\mu \sim m_{1,2} \sim M$  is the same as for the massless case, except the matching condition,  $\exp[D(\mu)]$ , now involves massive collinear propagators, which modifies the matching in the following way,

$$D(m_1, m_2) = D(m_{1,2} = 0) + (f_2(z_2) - \tilde{f}_2(z_2)/2) + (f_1(z_1) - \tilde{f}_1(z_1)/2), \quad (5.10)$$

where  $z_i = m_i/M_W$ ,  $f_{1,2}$  corresponds to the massive collinear contributions,

$$f_i(z_i) \equiv I_n(m_i) - I_n(0), \quad (5.11)$$

where  $I_n$  is the collinear vertex contribution and,

$$\tilde{f}_i(z_i) \equiv \delta Z_i(m_i, M) - \delta Z_i(0, M), \quad (5.12)$$

is the difference between the wave-function contribution with all mass scales non-zero and the external mass scales set to zero from table 1. Both the vertex and wave-function contributions depend solely on whether the corresponding particle is a fermion or a scalar. More specifically,  $f_i(z_i)$  maps to  $f_F(z_i)$  and  $f_S(z_i)$  in the case of fermions and scalars, respectively, and are given by,

$$f_F(z) = 2 + \left(\frac{1}{z^2} - 2\right) \log z^2 + \frac{2\sqrt{1-4z^2}}{z^2} \tanh^{-1} \sqrt{1-4z^2} + \frac{1}{2} \log^2(z^2) - 2(\tanh^{-1} \sqrt{1-4z^2})^2 \quad (5.13a)$$

$$f_S(z) = 1 - \left(1 - \frac{1}{2z^2}\right) \log z^2 + \frac{\sqrt{1-4z^2}}{z^2} \tanh^{-1} \sqrt{1-4z^2} + \frac{1}{2} \log^2(z^2) - 2(\tanh^{-1} \sqrt{1-4z^2})^2 \quad (5.13b)$$

as was also found in [17]. Thus, now that we have considered cases of interest in the Sudakov limit, we can shift to studying counterparts in the threshold limit.

## 6 Radiative Corrections in Threshold Limit

In this section, we calculate the form factor,  $\log F_E(Q^2)$ , in the opposite limit, i.e. small  $Q^2$  and large  $m^2$ , or threshold regime. Evidently, at threshold, the masses of the external particles are then taken to be the largest scale, resulting in two cases to consider,  $m_1 \sim m_2 \gg M \gg Q$  and  $m_2 \gg m_1 \gg M \gg Q$ . These cases have not been studied previously, and we provide the form factor up to and including two-loop order, which is computed using a sequence of effective field theories.

We begin by noting that at scales higher than  $m^2$ , the theory is the original Higgs-gauge theory, or so-called full theory. Moving to scales below  $m^2$ , we transition to HPET where degrees of freedom of off-shellness on the order  $m^2$  are integrated out. More specifically, let us commence with the simpler case,  $m_1 \sim m_2 \gg M^2 \gg Q^2$ , where  $m_{1,2}$  and  $M$  denote the external particle and bosonic masses, respectively. Schematically, we then have the following matching and running steps, illustrated as follows,

$$\mathcal{O} \xleftarrow[Q, M=0]{\mu \sim m_{1,2}} e^B \tilde{\mathcal{O}}_1 \xrightarrow{\gamma_3} e^B \tilde{\mathcal{O}}_1 \xleftarrow[Q=0]{\mu \sim M} e^{B+U} \tilde{\mathcal{O}}_2,$$

where  $B$  and  $U$  are multiplicative matching coefficients,  $\gamma_3$ , is the effective theory anomalous dimension, and  $\tilde{\mathcal{O}}_{1,2}$  the effective theory operators at each scale. At the scale  $\mu > m_{1,2}$ , we employ the full theory graphs and below, at  $\mu < m_{1,2}$ , we match down to HPET with matching coefficient,  $b(\mu)$ , and RGE given by,

$$\mu \frac{db(\mu)}{d\mu} = \gamma_F(a(\mu))b(\mu), \quad (6.1)$$

where  $\gamma_F$  is the full theory anomalous dimension for operator,  $\mathcal{O}$ , and is independent of energetic regime as given in table 2. The full theory is then matched onto HPET at  $\mu \sim m_{1,2}$ . The matching coefficient then depends on logarithms,  $\mathcal{L}_{m_{1,2}}$ , which are not divergent if  $\mu \sim m_{1,2}$ . The matching is done between full and effective theory operators as follows,

$$\bar{\psi}_2 \Gamma \psi_1 \rightarrow e^{\tilde{B}} \bar{h}_{f,2} \Gamma h_{f,1}, \quad (6.2a)$$

$$\chi_2^\dagger \chi_1 \rightarrow e^{\tilde{B}} h_{s,2}^\dagger h_{s,1}, \quad (6.2b)$$

$$i \chi_2^\dagger \overleftrightarrow{D}_\mu \chi_1 \rightarrow e^{\tilde{B}} h_{s,2}^\dagger [v_1 + v_2]_\mu h_{s,1}, \quad (6.2c)$$

$$\bar{\psi}_2 \chi_1 \rightarrow e^{\tilde{B}} \bar{h}_{f,2} h_{s,1}, \quad \chi_2^\dagger \psi_1 \rightarrow e^{\tilde{B}} h_{s,2}^\dagger h_{s,1}. \quad (6.2d)$$

We can then calculate the matching coefficient,  $\exp[B(\mu)]$ , as the full theory vertex and wave-function corrections with IR scales,  $M$  and  $Q$ , set to zero. The results of which at one and two loop order are given in table 8. Note that for the two-loop results, since  $m_1 \sim m_2$  and we want to evaluate the master integrals analytically, this can only be achieved with master integrals at a single scale, whence, we expand the two-loop contributions about the difference of  $m_1$  and  $m_2$  to first order. This is an accurate representation as the scale we are considering is where  $m_1 \sim m_2$  and although we chose to expand to first order as is conventionally done one can expand to any order and perform the single-scale two-loop master integrals as they are independent of expansion order. As for the remaining two-loop contributions, we present the mass and coupling renormalisation contributions at two-loop order in table 8.

What remains is the anomalous dimension,  $\gamma_3$ , between  $m_{1,2}$  and  $M$ , and the matching coefficient,  $\exp[U(\mu)]$ , at  $\mu \sim M$ . These have been computed in the previously in tables 7 and 6. Again, these contributions are found by computing graphs in figures 3 and 1, evaluated on-shell, with bosonic masses,  $M$ , included and external lines taken to be heavy with incoming and outgoing velocities,  $v_1$  and  $v_2$ , respectively. The difference here being that in the threshold limit,

$$w \sim \frac{m_1^2 + m_2^2}{2m_1 m_2} \sim \mathcal{O}(1), \quad (6.3)$$

since we take  $m_1 \sim m_2$ , and thus,  $h(w) \sim \mathcal{O}(1)$ , by inspection of (5.8). Whence, the sub-leading Higgs contribution which was sub-leading in the Sudakov regime becomes of the same order as the gauge contribution in the threshold regime. The remaining contributions at two-loop order from mass and coupling renormalisation were presented previously in table 7. Finally, we consider the slightly more involved,  $m_2 \gg m_1 \gg M^2 \gg Q^2$  case, where  $m_{1,2}$  and  $M$  denote the external particle and bosonic masses, respectively. Schematically, we then have following matching and running steps, illustrated as follows,

$$\mathcal{O} \xleftarrow[Q, M, m_1=0]{\mu \sim m_2} e^{\tilde{B}} \tilde{\mathcal{O}}_1 \xrightarrow{\tilde{\gamma}_3} e^{\tilde{B}} \xleftarrow[Q, M=0]{\mu \sim m_1} e^{\tilde{B}+G} \tilde{\mathcal{O}}_2 \xrightarrow{\gamma_3} e^{\tilde{B}+G} \tilde{\mathcal{O}}_2 \xleftarrow[Q=0]{\mu \sim M} e^{\tilde{B}+G+U} \tilde{\mathcal{O}}_3,$$

where  $\tilde{B}$ ,  $G$  and  $U$  are multiplicative matching coefficients,  $\gamma_3$  and  $\tilde{\gamma}_3$ , are the effective theory anomalous dimensions, and  $\tilde{\mathcal{O}}_{1,2,3}$  the effective theory operators at each scale. At the scale  $\mu > m_2$ , we employ the full theory graphs and below, at  $\mu < m_2$ , we match down to an effective theory with a single heavy field of mass,  $m_2$ . Thus, the effective theory operator is given by the full theory operators with particle 2 represented by a heavy field,  $h_{f,s}$ , for instance in the fermionic case we have,  $\bar{h}_{f,2} \Gamma \psi_1$ , and similarly for the other operators. We can then calculate the matching coefficient,  $\exp[\tilde{B}(\mu)]$ , as the full theory vertex and wave-function corrections with IR scales,  $m_1$ ,  $M$  and  $Q$ , set to zero. The results of the vertex and wave-function contributions,  $\exp[\tilde{B}(\mu)]$ , as well as the anomalous dimension,  $\tilde{\gamma}_3$ , between  $m_2$  and  $m_1$  are given in table 9. Moreover, the coupling and mass renormalisation corrections that contribute at two-loop order are also given, in table 9. What remains then is to evaluate the matching at  $\mu \sim m_1$  as the final matching and running,  $\exp[U(\mu)]$  and  $\gamma_3$ , at  $M$  is identical to the previous case. The theory above,  $\mu > m_1$ , is the effective theory with particle 2 taken to be a heavy field and the theory below,  $\mu < m_1$ , is heavy particle effective theory where both particles 1 and 2 are taken to be heavy and the IR scale being the bosonic masses are set to zero. The theory below  $m_1$  is scaleless and thus does not contribute to the matching but the theory above  $m_1$  is one of two scales,  $m_1$  and  $w' = p_1 \cdot v_2$ . However,  $w'$  is integrated out at leading order in the threshold limit as,

$$w' = p_1 \cdot v_2 = \frac{m_1^2 + m_2^2 - Q^2}{2m_2} \sim \frac{m_2}{2}, \quad (6.4)$$

$\mathcal{O}$	$B^{(1)}(\mu)$	$B^{(2)}(\mu)$	$\Delta B^{(2)}(\mu)$
$\bar{\psi}_2 \psi_1$	$\frac{C_A C_F \mathcal{L}_{m_2}^2 (m_1^2 + m_2^2)}{2m_+ m_+} - \frac{C_A C_F \mathcal{L}_{m_2} (m_1^2 - 4m_1 m_2 - m_2^2)}{2m_+ m_+}$ $- \frac{C_A C_F \mathcal{L}_{m_1} (m_1^2 \mathcal{L}_{m_1} + m_2^2 \mathcal{L}_{m_1} + m_1^2 + 4m_1 m_2 - m_2^2)}{2m_+ m_+}$ $- \frac{Y_F^2 C_A (5m_1^2 \mathcal{L}_{m_1} - 9m_2^2 \mathcal{L}_{m_1} - 4m_1 m_2 \mathcal{L}_{m_1}^2 + 8m_1 m_2 \mathcal{L}_{m_1} - 6m_1^2 + 6m_2^2)}{16m_+ m_+}$ $- \frac{Y_F^2 C_A \mathcal{L}_{m_2} (9m_1^2 - 8m_1 m_2 - 5m_2^2)}{16m_+ m_+} - \frac{Y_F^2 m_1 m_2 C_A \mathcal{L}_{m_2}^2}{4m_+ m_+}$	$(V_1^{(m_1,2)} + \frac{1}{2} F_1^{(m_2)} + \frac{1}{2} F_1^{(m_1)}) +$ $C_A^2 C_F^2 (60m_1^2 \mathcal{L}_{m_1}^2 - 80m_1^2 \mathcal{L}_{m_1} - 12\mathcal{L}_{m_1}^2 + 28\mathcal{L}_{m_1} + 5\pi^2 m_1^2 + 104m_1^2 - \pi^2 - 36) +$ $C_A^2 C_F^2 (12m_1^2 \mathcal{L}_{m_1}^2 + 176m_1^2 \mathcal{L}_{m_1} + 60\mathcal{L}_{m_1}^2 - 188\mathcal{L}_{m_1} + \pi^2 m_1^2 - 64m_1^2 + 5\pi^2 + 236)(m_2 - m_1) -$ $C_A^2 C_F Y_F^2 (228m_1^2 \mathcal{L}_{m_1}^2 - 292m_1^2 \mathcal{L}_{m_1} - 60\mathcal{L}_{m_1}^2 + 128\mathcal{L}_{m_1} + 19\pi^2 m_1^2 + 386m_1^2 - 5\pi^2 - 160) -$ $C_A^2 C_F Y_F^2 (12m_1^2 \mathcal{L}_{m_1}^2 + 764m_1^2 \mathcal{L}_{m_1} + 252\mathcal{L}_{m_1}^2 - 708\mathcal{L}_{m_1} + \pi^2 m_1^2 - 366m_1^2 + 21\pi^2 + 912)(m_2 - m_1) +$ $C_A^2 Y_F^2 (72m_1^2 \mathcal{L}_{m_1}^2 - 144m_1^2 \mathcal{L}_{m_1} - 36\mathcal{L}_{m_1}^2 + 72\mathcal{L}_{m_1} + 6\pi^2 m_1^2 + 176m_1^2 - 3\pi^2 - 88) -$ $C_A^2 Y_F^2 (36m_1^2 \mathcal{L}_{m_1}^2 - 324m_1^2 \mathcal{L}_{m_1} - 108\mathcal{L}_{m_1}^2 + 360\mathcal{L}_{m_1} + 3\pi^2 m_1^2 + 322m_1^2 - 9\pi^2 - 408)(m_2 - m_1)$	$\Delta B_1^{(2)}$
$\bar{\psi}_2 \gamma^\mu \psi_1$	$\frac{C_A C_F \mathcal{L}_{m_2} (m_1^2 - 4m_1 m_2 + 5m_2^2)}{2m_+ m_+}$ $- \frac{C_A C_F}{2m_+ m_+} \{ M_1^2 \mathcal{L}_{m_1}^2 - 5m_1^2 \mathcal{L}_{m_1} + m_2^2 \mathcal{L}_{m_1}^2 - m_2^2 \mathcal{L}_{m_1} - 2m_1 m_2 \mathcal{L}_{m_1}^2$ $+ 4m_1 m_2 \mathcal{L}_{m_1} + 6m_1^2 - 6m_2^2 \}$ $+ \frac{m_+ C_A C_F \mathcal{L}_{m_2}^2}{16m_+ m_+} - \frac{Y_F^2 C_A \mathcal{L}_{m_2} (9m_1^2 + 9m_1^2 m_2 - 3m_1 m_2^2 - 11m_2^2)}{16m_+ m_+}$ $- \frac{Y_F^2 C_A}{16m_+ m_+} \{ 11m_1^2 \mathcal{L}_{m_1} + 11m_1^2 m_2 \mathcal{L}_{m_1} - 9m_2^2 \mathcal{L}_{m_1} - 17m_1 m_2^2 \mathcal{L}_{m_1} - 18m_1^2$ $- 26m_1^2 m_2 + 18m_1 m_2^2 + 26m_2^2 \}$	$(V_1^{(m_1,2)} + \frac{1}{2} F_1^{(m_2)} + \frac{1}{2} F_1^{(m_1)}) - \frac{C_A^2 C_F^2 (12m_1^2 \mathcal{L}_{m_1}^2 - 28m_1^2 \mathcal{L}_{m_1} + 12\mathcal{L}_{m_1}^2 - 16\mathcal{L}_{m_1} + \pi^2 m_1^2 + 30m_1^2 + \pi^2 + 16)}{4m_1(m_2 - m_1)}$ $- \frac{C_A^2 C_F^2 (36m_1^2 \mathcal{L}_{m_1}^2 + 36m_1^2 \mathcal{L}_{m_1} + 36\mathcal{L}_{m_1}^2 + 3\pi^2 m_1^2 - 74m_1^2 + 3\pi^2 + 40)}{8m_1^2} +$ $C_A^2 C_F Y_F^2 (12m_1^2 \mathcal{L}_{m_1}^2 - 16m_1^2 \mathcal{L}_{m_1} + 12\mathcal{L}_{m_1}^2 - 4\mathcal{L}_{m_1} + \pi^2 m_1^2 + 10m_1^2 + \pi^2 + 2) +$ $C_A^2 C_F Y_F^2 (36m_1^2 \mathcal{L}_{m_1}^2 + 72m_1^2 \mathcal{L}_{m_1} + 36\mathcal{L}_{m_1}^2 + 36\mathcal{L}_{m_1} + 3\pi^2 m_1^2 - 74m_1^2 + 3\pi^2 + 22) -$ $C_A^2 C_F Y_F^2 (324m_1^2 \mathcal{L}_{m_1}^2 - 1512m_1^2 \mathcal{L}_{m_1} + 900\mathcal{L}_{m_1}^2 - 1284\mathcal{L}_{m_1} + 27\pi^2 m_1^2 + 1278m_1^2 + 1202)(m_2 - m_1) +$ $C_A^2 Y_F^2 (36m_1^2 \mathcal{L}_{m_1}^2 - 108m_1^2 \mathcal{L}_{m_1} + 36\mathcal{L}_{m_1}^2 - 72\mathcal{L}_{m_1} + 3\pi^2 m_1^2 + 124m_1^2 + 3\pi^2 + 70) +$ $3C_A^2 (36m_1^2 \mathcal{L}_{m_1}^2 + 12m_1^2 \mathcal{L}_{m_1} + 36\mathcal{L}_{m_1}^2 - 24\mathcal{L}_{m_1} + 3\pi^2 m_1^2 - 80m_1^2 + 3\pi^2 + 46) +$ $C_A^2 (324m_1^2 \mathcal{L}_{m_1}^2 + 324m_1^2 \mathcal{L}_{m_1} - 252\mathcal{L}_{m_1}^2 + 840\mathcal{L}_{m_1} + 27\pi^2 m_1^2 - 36m_1^2 - 21\pi^2 - 778)(m_2 - m_1)$	$\Delta B_2^{(2)}$
$\bar{\psi}_2 \sigma^{\mu\nu} \psi_1$	$\frac{C_A C_F \mathcal{L}_{m_2}^2 (m_1^2 + m_2^2)}{2m_+ m_+} - \frac{C_A C_F \mathcal{L}_{m_2} (m_1^2 - 4m_1 m_2 + 7m_2^2)}{2m_+ m_+}$ $- \frac{C_A C_F (m_1^2 \mathcal{L}_{m_1}^2 - 7m_1^2 \mathcal{L}_{m_1} + m_2^2 \mathcal{L}_{m_1}^2 - m_2^2 \mathcal{L}_{m_1} + 4m_1 m_2 \mathcal{L}_{m_1} + 4m_1^2 - 4m_2^2)}{2m_+ m_+}$ $- \frac{Y_F^2 C_A (9m_1^2 \mathcal{L}_{m_1} - 9m_2^2 \mathcal{L}_{m_1} - 4m_1 m_2 \mathcal{L}_{m_1}^2 + 8m_1 m_2 \mathcal{L}_{m_1} - 14m_1^2 + 14m_2^2)}{16m_+ m_+}$ $- \frac{Y_F^2 C_A \mathcal{L}_{m_2} (9m_1^2 - 8m_1 m_2 - 9m_2^2)}{16m_+ m_+} - \frac{Y_F^2 m_1 m_2 C_A \mathcal{L}_{m_2}^2}{4m_+ m_+}$	$(V_1^{(m_1,2)} + \frac{1}{2} F_1^{(m_2)} + \frac{1}{2} F_1^{(m_1)}) + \frac{C_A^2 C_F^2 (12m_1^2 \mathcal{L}_{m_1}^2 - 58m_1^2 \mathcal{L}_{m_1} - 12\mathcal{L}_{m_1}^2 + 34\mathcal{L}_{m_1} + \pi^2 m_1^2 + 80m_1^2 - \pi^2 - 55)}{2m_1^2}$ $- \frac{C_A^2 C_F^2 (6m_1^2 \mathcal{L}_{m_1}^2 + 6\mathcal{L}_{m_1} - 10m_1^2 - 7)}{16m_1(m_2 - m_1)} + \frac{C_A^2 C_F Y_F^2 (12m_1^2 \mathcal{L}_{m_1}^2 - 4m_1^2 \mathcal{L}_{m_1} + 12\mathcal{L}_{m_1}^2 + 8\mathcal{L}_{m_1} + \pi^2 m_1^2 - 4m_1^2 + \pi^2 - 6)}{16m_1(m_2 - m_1)}$ $+ \frac{C_A^2 C_F^2 (12m_1^2 \mathcal{L}_{m_1}^2 - 22m_1^2 \mathcal{L}_{m_1} + 60\mathcal{L}_{m_1}^2 - 194\mathcal{L}_{m_1} + \pi^2 m_1^2 - 10m_1^2 + 5\pi^2 + 263)(m_2 - m_1) -$ $C_A^2 C_F Y_F^2 (84m_1^2 \mathcal{L}_{m_1}^2 - 388m_1^2 \mathcal{L}_{m_1} - 108\mathcal{L}_{m_1}^2 + 216\mathcal{L}_{m_1} + 7\pi^2 m_1^2 + 548m_1^2 - 9\pi^2 - 342) -$ $C_A^2 C_F Y_F^2 (36m_1^2 \mathcal{L}_{m_1}^2 + 132m_1^2 \mathcal{L}_{m_1} + 1476\mathcal{L}_{m_1}^2 - 4392\mathcal{L}_{m_1} + 3\pi^2 m_1^2 - 936m_1^2 + 123\pi^2 + 5378)(m_2 - m_1) -$ $C_A^2 Y_F^2 (36m_1^2 \mathcal{L}_{m_1}^2 - 108m_1^2 \mathcal{L}_{m_1} + 36\mathcal{L}_{m_1}^2 - 72\mathcal{L}_{m_1} + 3\pi^2 m_1^2 + 124m_1^2 + 3\pi^2 + 70) +$ $C_A^2 Y_F^2 (36m_1^2 \mathcal{L}_{m_1}^2 - 180m_1^2 \mathcal{L}_{m_1} - 108\mathcal{L}_{m_1}^2 + 72\mathcal{L}_{m_1} + 3\pi^2 m_1^2 + 376m_1^2 - 9\pi^2 - 138) +$ $C_A^2 (180m_1^2 \mathcal{L}_{m_1}^2 - 396m_1^2 \mathcal{L}_{m_1} - 396\mathcal{L}_{m_1}^2 + 1128\mathcal{L}_{m_1} + 15\pi^2 m_1^2 + 548m_1^2 - 33\pi^2 - 1130)(m_2 - m_1)$	$\Delta B_3^{(2)}$
$\chi_2^1 \chi_1$	$\frac{C_A C_F \mathcal{L}_{m_2}^2 (m_1^2 + m_2^2)}{2m_+ m_+} - \frac{m_2^2 C_A C_F \mathcal{L}_{m_2} - Y_F^2 C_A \mathcal{L}_{m_2} - Y_F^2 C_A \mathcal{L}_{m_2}^2}{8m_+^2} - \frac{Y_F^2 C_A \mathcal{L}_{m_2}^2}{16m_+ m_+}$ $- \frac{C_A C_F (m_1^2 \mathcal{L}_{m_1}^2 - 2m_1^2 \mathcal{L}_{m_1} + m_2^2 \mathcal{L}_{m_1}^2 - 6m_1^2 + 6m_2^2)}{2m_+ m_+}$ $+ \frac{Y_F^2 C_A (m_1^2 m_2^2 \mathcal{L}_{m_1}^2 - 2m_1^2 m_2^2 \mathcal{L}_{m_1} + 2m_2^2 \mathcal{L}_{m_1} + 2m_1^2 - 2m_2^2)}{16m_+^2 m_2^2 m_+}$	$(V_1^{(m_1,2)} + \frac{1}{2} F_1^{(m_2)} + \frac{1}{2} F_1^{(m_1)}) + \frac{C_A^2 C_F^2 (24m_1^2 \mathcal{L}_{m_1}^2 - 24m_1^2 \mathcal{L}_{m_1} - 12\mathcal{L}_{m_1}^2 + 12\mathcal{L}_{m_1} + 2\pi^2 m_1^2 + 24m_1^2 - \pi^2 - 12)}{3m_1^2}$ $- \frac{2C_A^2 C_F^2 (18m_1^2 \mathcal{L}_{m_1}^2 + 12\mathcal{L}_{m_1}^2 - 24\mathcal{L}_{m_1} - 6m_1^2 + \pi^2 + 18)(m_2 - m_1) -$ $C_A^2 C_F Y_F^2 (12m_1^2 \mathcal{L}_{m_1}^2 - 36m_1^2 \mathcal{L}_{m_1} - 12\mathcal{L}_{m_1}^2 + 32\mathcal{L}_{m_1} + \pi^2 m_1^2 + 30m_1^2 - \pi^2 - 38)(m_2 - m_1) -$ $C_A^2 C_F Y_F^2 (36m_1^2 \mathcal{L}_{m_1}^2 - 48m_1^2 \mathcal{L}_{m_1} - 12\mathcal{L}_{m_1}^2 + 24\mathcal{L}_{m_1} + 3\pi^2 m_1^2 + 60m_1^2 - \pi^2 - 30) -$ $C_A^2 Y_F^2 (12\mathcal{L}_{m_1}^2 - 12\mathcal{L}_{m_1} + \pi^2 + 12) - C_A^2 Y_F^2 (12\mathcal{L}_{m_1}^2 - 24\mathcal{L}_{m_1} + \pi^2 + 18)(m_2 - m_1)$	$\Delta B_4^{(2)}$
$i\chi_2^{\dagger 2} D_\mu \chi_1$	$\frac{2m_2^2 C_A C_F \mathcal{L}_{m_1} (m_1^2 + 8m_2^2)}{m_+ m_+ (m_1^2 + m_2^2)} + \frac{C_A C_F \mathcal{L}_{m_2}^2 (m_1^2 + m_2^2)}{2m_+ m_+}$ $- \frac{C_A C_F (m_1^2 \mathcal{L}_{m_1} - 32m_1^2 \mathcal{L}_{m_1} + 2m_1^2 m_2^2 \mathcal{L}_{m_1} - 4m_1^2 m_2^2 \mathcal{L}_{m_1} + m_2^2 \mathcal{L}_{m_1} + 18m_1^2 - 18m_2^2)}{2m_+^2 m_+ m_+}$ $- \frac{Y_F^2 C_A \mathcal{L}_{m_2} (m_1^2 - 2m_2^2)}{16m_+ m_+} - \frac{Y_F^2 C_A \mathcal{L}_{m_2}^2}{16m_+ m_+}$ $+ \frac{C_A}{16m_+^2 m_2^2 m_+} \{ M_1^2 m_2^2 \mathcal{L}_{m_1}^2 - 16m_1^2 m_2^2 \mathcal{L}_{m_1} + m_1^2 m_2^2 \mathcal{L}_{m_1}^2 - 8m_1^2 m_2^2 \mathcal{L}_{m_1}$ $+ 8m_1^2 \mathcal{L}_{m_1} + 6m_1^2 - 8m_1^2 m_2^2 + 8m_1^2 m_2^2 - 8m_2^2 \}$	$(V_1^{(m_1,2)} + \frac{1}{2} F_1^{(m_2)} + \frac{1}{2} F_1^{(m_1)}) +$ $C_A^2 C_F^2 (36m_1^2 \mathcal{L}_{m_1}^2 - 120m_1^2 \mathcal{L}_{m_1} + 60\mathcal{L}_{m_1}^2 - 72\mathcal{L}_{m_1} + 3\pi^2 m_1^2 + 108m_1^2 + 5\pi^2 + 48)(m_2 - m_1)$ $- \frac{C_A^2 C_F Y_F^2 (12m_1^2 \mathcal{L}_{m_1}^2 - 54m_1^2 \mathcal{L}_{m_1} + 72\mathcal{L}_{m_1}^2 - 132\mathcal{L}_{m_1} + \pi^2 m_1^2 + 78m_1^2 + 6\pi^2 + 126)(m_2 - m_1) -$ $C_A^2 C_F Y_F^2 (12m_1^2 \mathcal{L}_{m_1}^2 - 24m_1^2 \mathcal{L}_{m_1} - 36\mathcal{L}_{m_1}^2 + 36\mathcal{L}_{m_1} + \pi^2 m_1^2 + 24m_1^2 - 3\pi^2 - 36) -$ $C_A^2 (12m_1^2 \mathcal{L}_{m_1}^2 - 36m_1^2 \mathcal{L}_{m_1} - 12\mathcal{L}_{m_1}^2 + 24\mathcal{L}_{m_1} + \pi^2 m_1^2 + 48m_1^2 - \pi^2 - 30) -$ $C_A^2 (12m_1^2 \mathcal{L}_{m_1}^2 - 42m_1^2 \mathcal{L}_{m_1} - 24\mathcal{L}_{m_1}^2 + 60\mathcal{L}_{m_1} + \pi^2 m_1^2 + 57m_1^2 - 2\pi^2 - 72)(m_2 - m_1) -$ $C_A^2 C_F^2 (12m_1^2 \mathcal{L}_{m_1}^2 - 12m_1^2 \mathcal{L}_{m_1} + 12\mathcal{L}_{m_1}^2 + \pi^2 m_1^2 + 6m_1^2 + \pi^2)$	$\Delta B_5^{(2)}$
$\bar{\chi}_2^1 \psi_1$ , $\bar{\psi}_2 \chi_1 (1 \leftrightarrow 2)$	$\frac{C_A C_F \mathcal{L}_{m_2}^2 (m_1^2 + m_2^2)}{2m_+ m_+} - \frac{2m_2^2 C_A C_F \mathcal{L}_{m_2} - Y_F^2 m_1 m_2 C_A \mathcal{L}_{m_2}^2}{8m_+ m_+}$ $- \frac{C_A C_F}{2m_+ m_+} \{ M_1^2 \mathcal{L}_{m_1}^2 - 3m_1^2 \mathcal{L}_{m_1} + m_2^2 m_2 \mathcal{L}_{m_1} - 3m_1^2 m_2 \mathcal{L}_{m_1}$ $+ m_2^2 \mathcal{L}_{m_1} - m_2^2 \mathcal{L}_{m_1} + m_1 m_2^2 \mathcal{L}_{m_1}^2 + 3m_1 m_2^2 \mathcal{L}_{m_1} - 4m_1^2 + 4m_1 m_2^2 \}$ $+ \frac{C_A}{16m_+^2 m_+ m_+} \{ -9\gamma^2 m_1^2 m_2^2 \mathcal{L}_{m_1} - 9\gamma^2 m_1 m_2^2 \mathcal{L}_{m_1} + 9\gamma^2 m_2^2 \mathcal{L}_{m_1} + 9\gamma^2 m_1 m_2^2 \mathcal{L}_{m_1}$ $+ 9\gamma^2 m_1^2 m_2^2 - 9\gamma^2 m_1 m_2 - 9\gamma^2 m_2^2 + 2\gamma^2 Y_1 m_1 m_2^2 \mathcal{L}_{m_1}^2 + 2\gamma^2 Y_1 m_1 m_2^2 \mathcal{L}_{m_1}$ $- 4\gamma^2 Y_1 m_1 m_2^2 \mathcal{L}_{m_1} - 4\gamma^2 Y_1 m_1^2 m_2^2 + 4\gamma^2 Y_1 m_1^2 + 2\gamma^2 m_1^2 m_2 - 2\gamma^2 m_1 m_2^2 - 2\gamma^2 m_2^2 \}$ $+ \frac{Y_F C_A \mathcal{L}_{m_2} (2Y_1 m_1 m_2^2 - Y_1 m_1^2 - Y_1 m_2^2 + Y_1 m_1 m_2^2 + Y_1 m_2^2)}{8m_+^2 m_+ m_+}$	$(V_0^{(m_1,2)} + \frac{1}{2} F_1^{(m_2)} + \frac{1}{2} F_1^{(m_1)}) + \frac{C_A^2 C_F^2 (360m_1^2 \mathcal{L}_{m_1}^2 - 600m_1^2 \mathcal{L}_{m_1} - 60\mathcal{L}_{m_1}^2 + 108\mathcal{L}_{m_1} + 25\pi^2 m_1^2 + 744m_1^2 - 5\pi^2 - 132)}{24m_1^2}$ $- \frac{C_A^2 C_F^2 (60m_1^2 \mathcal{L}_{m_1}^2 + 312m_1^2 \mathcal{L}_{m_1} + 180\mathcal{L}_{m_1}^2 - 432\mathcal{L}_{m_1} + 5\pi^2 m_1^2 - 144m_1^2 + 15\pi^2 + 474)(m_2 - m_1) -$ $C_A^2 C_F Y_1 Y_2 (120m_1^2 \mathcal{L}_{m_1}^2 - 288m_1^2 \mathcal{L}_{m_1} - 60\mathcal{L}_{m_1}^2 + 162\mathcal{L}_{m_1} + 10\pi^2 m_1^2 + 228m_1^2 - 5\pi^2 - 171)(m_2 - m_1) -$ $C_A^2 C_F Y_1 Y_2 (180m_1^2 \mathcal{L}_{m_1}^2 - 264m_1^2 \mathcal{L}_{m_1} - 60\mathcal{L}_{m_1}^2 + 108\mathcal{L}_{m_1} + 15\pi^2 m_1^2 + 312m_1^2 - 5\pi^2 - 132) +$ $C_A^2 C_F Y_2^2 (108m_1^2 \mathcal{L}_{m_1}^2 - 396m_1^2 \mathcal{L}_{m_1} - 60\mathcal{L}_{m_1}^2 + 156\mathcal{L}_{m_1} + 9\pi^2 m_1^2 + 468m_1^2 - 5\pi^2 - 186)(m_2 - m_1) -$ $C_A^2 C_F Y_2^2 (60m_1^2 \mathcal{L}_{m_1}^2 - 132m_1^2 \mathcal{L}_{m_1} - 12\mathcal{L}_{m_1}^2 + 24\mathcal{L}_{m_1} + 5\pi^2 m_1^2 + 168m_1^2 - \pi^2 - 30) -$ $C_A^2 C_F Y_2^2 (180m_1^2 \mathcal{L}_{m_1}^2 - 396m_1^2 \mathcal{L}_{m_1} - 36\mathcal{L}_{m_1}^2 + 72\mathcal{L}_{m_1} + 15\pi^2 m_1^2 + 494m_1^2 - 3\pi^2 - 88) -$ $C_A^2 C_F Y_2^2 (36m_1^2 \mathcal{L}_{m_1}^2 + 36m_1^2 \mathcal{L}_{m_1} + 108\mathcal{L}_{m_1}^2 - 252\mathcal{L}_{m_1} + 3\pi^2 m_1^2 - 2m_1^2 + 9\pi^2 + 300)(m_2 - m_1) +$ $C_A^2 Y_1 Y_2 (108m_1^2 \mathcal{L}_{m_1}^2 - 180m_1^2 \mathcal{L}_{m_1} - 36\mathcal{L}_{m_1}^2 + 72\mathcal{L}_{m_1} + 9\pi^2 m_1^2 + 210m_1^2 - 3\pi^2 - 88) -$ $C_A^2 Y_1 Y_2 (72m_1^2 \mathcal{L}_{m_1}^2 - 144m_1^2 \mathcal{L}_{m_1} - 36\mathcal{L}_{m_1}^2 + 90\mathcal{L}_{m_1} + 6\pi^2 m_1^2 + 140m_1^2 - 3\pi^2 - 106)(m_2 - m_1) +$ $C_A^2 Y_2^2 Y_2 (36m_1^2 \mathcal{L}_{m_1}^2 - 60m_1^2 \mathcal{L}_{m_1} - 12\mathcal{L}_{m_1}^2 + 24\mathcal{L}_{m_1} + 3\pi^2 m_1^2 + 72m_1^2 - \pi^2 - 30) -$ $C_A^2 Y_2^2 Y_2 (60m_1^2 \mathcal{L}_{m_1}^2 - 144m_1^2 \mathcal{L}_{m_1} - 24\mathcal{L}_{m_1}^2 + 66\mathcal{L}_{m_1} + 5\pi^2 m_1^2 + 150m_1^2 - 2\pi^2 - 78)(m_2 - m_1)$	$\Delta B_6^{(2)}$

**Table 8:** Matching corrections,  $B(\mu)$ , to the threshold form-factor at  $\mu \sim m_{1,2}$ ,  $a \equiv \alpha/(4\pi)$ ,  $\mathcal{L}_{m_{1,2}} \equiv \log m_{1,2}^2/\mu^2$ , and  $m_\pm = m_1 \pm m_2$ .  $V_i^{(m_{1,2})}$  and  $F_I$  are two loop vertex and wave-function corrections given in Appendices C and D.  $\Delta B^{(2)}$  is the two-loop order contribution from mass and coupling renormalisation given in Appendix B.

and thus, we obtained the matching and wave-function contributions,  $\exp[G(\mu)]$ , with logarithms of a single

$\mathcal{O}$	$\tilde{B}^{(1)}(\mu)$	$\tilde{B}^{(2)}(\mu)$	$\tilde{\gamma}_3^{(1)}(\mu)$	$\Delta\tilde{B}^{(2)}(\mu)$
$\psi_2\psi_1$	$-\frac{1}{48}C_A(3\mathcal{L}_{m_2}(8C_F\mathcal{L}_{m_2}+8C_F+9Y_f^2)+4(24+\pi^2)C_F+9Y_f^2)$	$(V_1^{(m_2)}+\frac{1}{2}F_\psi^{(m_2,0)}+\frac{1}{2}F_\psi^{(0,0)})$ $+\frac{1}{6}C_A^2C_F^2(3\mathcal{L}_{m_2}(4(\mathcal{L}_{m_2}-2)\mathcal{L}_{m_2}+\pi^2+16)+6\zeta_3-2\pi^2-48)$ $-\frac{1}{6}C_A^2C_FY_f^2(2\mathcal{L}_{m_2}(6\mathcal{L}_{m_2}(2\mathcal{L}_{m_2}+1)+3\pi^2-46)+12\zeta_3+\pi^2+158)$ $+\frac{1}{64}C_A^2Y_f^4(36(\mathcal{L}_{m_2}-3)\mathcal{L}_{m_2}+3\pi^2+142)$	$-\frac{C_A}{8}(40C_F-Y_f^2)$	$\frac{1}{24}C_A C_F(-2\mathcal{L}_{m_2}(3\mathcal{L}_{m_2}(\beta_0-48C_A C_F)+120C_A C_F+(24+\pi^2)\beta_0+2\beta_0\mathcal{L}_{m_2}^2))+24(24+\pi^2)C_A C_F-(\pi^2-96)\beta_0-8\beta_0\zeta_3$ $-\frac{1}{192}Y_f^2 C_A(6\mathcal{L}_{m_2}(\mathcal{L}_{m_2}(48C_A C_F+5\beta_0)-208C_A C_F+6\beta_0)$ $+24(46+\pi^2)C_A C_F+(5\pi^2-132)\beta_0)+\frac{1}{32}Y_f^4(13-15\mathcal{L}_{m_2})C_A^2$
$\psi_2\gamma^\mu\psi_1$	$\frac{1}{48}C_A(-3\mathcal{L}_{m_2}(8C_F\mathcal{L}_{m_2}-40C_F+11Y_f^2)-4(60+\pi^2)C_F+27Y_f^2)$	$(V_1^{(m_2)}+\frac{1}{2}F_\psi^{(m_2,0)}+\frac{1}{2}F_\psi^{(0,0)})$ $+\frac{1}{12}C_A^2C_F^2(6\mathcal{L}_{m_2}(4\mathcal{L}_{m_2}^2-2\mathcal{L}_{m_2}+\pi^2+26)+12\zeta_3-\pi^2-246)$ $+\frac{1}{32}C_A^2C_FY_f^2(-6\mathcal{L}_{m_2}(4\mathcal{L}_{m_2}^2-2\mathcal{L}_{m_2}+\pi^2+30)+\pi^2+262-12\zeta_3)$ $-\frac{1}{256}C_A^2Y_f^4(36(\mathcal{L}_{m_2}-3)\mathcal{L}_{m_2}+3\pi^2+124)$	$-\frac{C_A}{4}(8C_F+Y_f^2)$	$\frac{1}{24}C_A C_F(-2\mathcal{L}_{m_2}(-3\mathcal{L}_{m_2}(48C_A C_F+5\beta_0)+552C_A C_F+(60+\pi^2)\beta_0+2\beta_0\mathcal{L}_{m_2}^2))+24(66+\pi^2)C_A C_F+(264+5\pi^2)\beta_0-8\beta_0\zeta_3$ $-\frac{1}{192}Y_f^2 C_A(6\mathcal{L}_{m_2}(48C_A C_F\mathcal{L}_{m_2}-496C_A C_F-18\beta_0+11\beta_0\mathcal{L}_{m_2})$ $+24(166+\pi^2)C_A C_F+(180+11\pi^2)\beta_0)+\frac{1}{32}Y_f^4(52-33\mathcal{L}_{m_2})C_A^2$
$\psi_2\sigma^{\mu\nu}\psi_1$	$\frac{1}{48}C_A(-3\mathcal{L}_{m_2}(8C_F\mathcal{L}_{m_2}-56C_F+9Y_f^2)-4(48+\pi^2)C_F+15Y_f^2)$	$(V_1^{(m_2)}+\frac{1}{2}F_\psi^{(m_2,0)}+\frac{1}{2}F_\psi^{(0,0)})$ $+\frac{1}{6}C_A^2C_F^2(3\mathcal{L}_{m_2}(4(\mathcal{L}_{m_2}-2)\mathcal{L}_{m_2}+\pi^2+28)+6\zeta_3-2\pi^2-108)$ $-\frac{1}{32}C_A^2C_FY_f^2(2\mathcal{L}_{m_2}(6\mathcal{L}_{m_2}(2\mathcal{L}_{m_2}+3)+\pi^2-50)$ $+3(4\zeta_3+50+\pi^2))+\frac{3}{192}C_A^2Y_f^4(36(\mathcal{L}_{m_2}-3)\mathcal{L}_{m_2}+3\pi^2+136)$	$-\frac{C_A}{8}(8C_F+Y_f^2)$	$\frac{1}{24}C_A C_F(-2\mathcal{L}_{m_2}(-3\mathcal{L}_{m_2}(48C_A C_F+7\beta_0)+696C_A C_F+(48+\pi^2)\beta_0+2\beta_0\mathcal{L}_{m_2}^2))+24(68+\pi^2)C_A C_F+(96+7\pi^2)\beta_0-8\beta_0\zeta_3$ $-\frac{1}{64}Y_f^2 C_A(2\mathcal{L}_{m_2}(48C_A C_F\mathcal{L}_{m_2}-496C_A C_F-10\beta_0+9\beta_0\mathcal{L}_{m_2})$ $+8(154+\pi^2)C_A C_F+(4+3\pi^2)\beta_0)+\frac{3}{32}Y_f^4(13-9\mathcal{L}_{m_2})C_A^2$
$\chi_2^T\chi_1$	$-\frac{1}{12}C_A C_F(6(\mathcal{L}_{m_2}-2)\mathcal{L}_{m_2}+\pi^2-12)$ $+\frac{Y_f^2 C_A(6(\mathcal{L}_{m_2}+6\log^2(m_2^2))+\pi^2+12)}{96m_2^2}$	$(V_1^{(m_2)}+\frac{1}{2}F_\chi^{(m_2,0)}+\frac{1}{2}F_\chi^{(0,0)})$ $+\frac{1}{6}C_A^2C_F^2(2\mathcal{L}_{m_2}(4\mathcal{L}_{m_2}^2+6\mathcal{L}_{m_2}+\pi^2-12)+4\zeta_3+\pi^2+24)$ $+C_A^2C_FY_f^2(-\mathcal{L}_{m_2}(4\mathcal{L}_{m_2}^2+\pi^2-3)-6+\psi^{(2)}(1))$ $+C_A^2Y_f^4(2\mathcal{L}_{m_2}(4\mathcal{L}_{m_2}^2-6\mathcal{L}_{m_2}+\pi^2+12)+4\zeta_3-\pi^2-24)$ $\frac{12m_2^2}{384m_2^2}$	$C_A C_F$	$\frac{1}{12}C_A C_F(6\mathcal{L}_{m_2}^2(12C_A C_F+\beta_0)-\mathcal{L}_{m_2}(240C_A C_F+(\pi^2-12)\beta_0)+6(38+\pi^2)C_A C_F-2\beta_0\mathcal{L}_{m_2}^2+\beta_0(\pi^2-4\zeta_3+9))$ $+\frac{Y_f^2 C_A}{96m_2^2}(6\mathcal{L}_{m_2}(-6\mathcal{L}_{m_2}(20C_A C_F+\beta_0)+384C_A C_F+(12+\pi^2)\beta_0+2\beta_0\mathcal{L}_{m_2}^2)-10(42+\pi^2)C_A C_F-(24+\pi^2)\beta_0+4\beta_0\zeta_3)$ $+Y_f^4(12\mathcal{L}_{m_2}-3)\mathcal{L}_{m_2}+\pi^2+42)C_A^2+\frac{1}{32}Y_f^4(13-15\mathcal{L}_{m_2})C_A^2$ $\frac{12m_2^2}{192m_2^2}$
$i\chi_2^T\overleftrightarrow{D}_\mu\chi_1$	$-\frac{1}{12}C_A C_F(6(\mathcal{L}_{m_2}-8)\mathcal{L}_{m_2}+\pi^2+54)$ $+\frac{Y_f^2 C_A(6(\mathcal{L}_{m_2}-4)\mathcal{L}_{m_2}+\pi^2+12)}{96m_2^2}$	$(V_1^{(m_2)}+\frac{1}{2}F_\chi^{(m_2,0)}+\frac{1}{2}F_\chi^{(0,0)})$ $+\frac{1}{12}C_A^2C_F^2(4\mathcal{L}_{m_2}(\mathcal{L}_{m_2}(4\mathcal{L}_{m_2}-9)+\pi^2+15)+8\zeta_3-3\pi^2-54)$ $+C_A^2C_FY_f^2(-4\mathcal{L}_{m_2}(8\mathcal{L}_{m_2}^2-9\mathcal{L}_{m_2}+2\pi^2+24)+3\pi^2+8(15+\psi^{(2)}(1)))$ $+C_A^2Y_f^4(4\mathcal{L}_{m_2}^3+(\pi^2-6)\mathcal{L}_{m_2}+2(\zeta_3+6))$ $\frac{12m_2^2}{192m_2^2}$	$4C_A C_F$	$\frac{1}{24}C_A C_F(-2\mathcal{L}_{m_2}(-24\mathcal{L}_{m_2}(3C_A C_F+\beta_0)+456C_A C_F+(51+\pi^2)\beta_0+2\beta_0\mathcal{L}_{m_2}^2))+6(32+\pi^2)C_A C_F-2\beta_0\mathcal{L}_{m_2}^2+\beta_0(-4\zeta_3-72+\pi^2)$ $+\frac{Y_f^2 C_A}{96m_2^2}(-2((387+5\pi^2)C_A C_F+\pi^2\beta_0)+\mathcal{L}_{m_2}(2\mathcal{L}_{m_2}(\beta_0\mathcal{L}_{m_2}-6(10C_A C_F+\beta_0))+600C_A C_F+(12+\pi^2)\beta_0+4\beta_0\zeta_3)$ $+Y_f^4(12(\mathcal{L}_{m_2}-4)\mathcal{L}_{m_2}+\pi^2+54)C_A^2$ $\frac{12m_2^2}{192m_2^2}$
$\chi_2^T\psi_1$	$-\frac{1}{12}C_A C_F(6(\mathcal{L}_{m_2}-2)\mathcal{L}_{m_2}+\pi^2-24)$ $-\frac{Y_f^2 C_A(\mathcal{L}_{m_2}-1)}{8m_2^2}$ $+\frac{Y_f Y_\psi C_A(6\mathcal{L}_{m_2}^2+\pi^2-12)}{48m_2}$	$(V_1^{(m_2)}+\frac{1}{2}F_\psi^{(m_2,0)}+\frac{1}{2}F_\psi^{(0,0)})$ $+\frac{1}{6}C_A^2C_F^2(\mathcal{L}_{m_2}(4\mathcal{L}_{m_2}(\mathcal{L}_{m_2}+3)+\pi^2-24)+2\zeta_3+\pi^2+24)$ $-C_A^2C_FY_fY_\psi(2\mathcal{L}_{m_2}(4\mathcal{L}_{m_2}^2+6\mathcal{L}_{m_2}+\pi^2-12)+4\zeta_3+\pi^2+24)$ $-\frac{C_A^2C_FY_f^2(2\mathcal{L}_{m_2}(4\mathcal{L}_{m_2}^2+6\mathcal{L}_{m_2}+\pi^2-24)+4\zeta_3+\pi^2+72)}{48m_2^2}$ $+\frac{C_A^2Y_f^2Y_f^2(4\mathcal{L}_{m_2}^3+(\pi^2-6)\mathcal{L}_{m_2}+2(\zeta_3+6))}{96m_2^2}$	$C_A C_F$	$\frac{1}{12}C_A C_F(6\mathcal{L}_{m_2}^2(12C_A C_F+\beta_0)-\mathcal{L}_{m_2}(240C_A C_F+(\pi^2-24)\beta_0)+6(32+\pi^2)C_A C_F-2\beta_0\mathcal{L}_{m_2}^2+\beta_0(-4\zeta_3-72+\pi^2))$ $-\frac{Y_f^2 C_A}{96m_2^2}(6\mathcal{L}_{m_2}(\mathcal{L}_{m_2}(3C_A C_F+\beta_0)-2(18C_A C_F+\beta_0))+4(54+\pi^2)C_A C_F+(24+\pi^2)\beta_0)+\frac{Y_f Y_\psi C_A}{48m_2}(\mathcal{L}_{m_2}(18C_A C_F+(\pi^2-12)\beta_0)-72C_A C_F\mathcal{L}_{m_2}-6(24+\pi^2)C_A C_F+2\beta_0\mathcal{L}_{m_2}^2+4\beta_0(\zeta_3+9))$ $Y_f Y_\psi^2(12(\mathcal{L}_{m_2}-2)\mathcal{L}_{m_2}+\pi^2+18)(2C_A C_F+1)-\frac{Y_f^2(2\mathcal{L}_{m_2}-3)C_A^2}{32m_2^2}$
$\psi_2\chi_1$	$-\frac{1}{12}C_A C_F(6(\mathcal{L}_{m_2}-5)\mathcal{L}_{m_2}+\pi^2+48)$ $-\frac{Y_f Y_\psi C_A}{4m_2}-\frac{9}{16}Y_f^2 C_A(\mathcal{L}_{m_2}-1)$	$(V_1^{(m_2)}+\frac{1}{2}F_\psi^{(m_2,0)}+\frac{1}{2}F_\psi^{(0,0)})$ $+\frac{1}{12}C_A^2C_F^2(6\mathcal{L}_{m_2}(2\mathcal{L}_{m_2}(2\mathcal{L}_{m_2}-7)+\pi^2+36)+12\zeta_3-7\pi^2-264)$ $-C_A^2C_FY_fY_\psi(4\mathcal{L}_{m_2}(3\mathcal{L}_{m_2}-10)+\pi^2+56)$ $+\frac{1}{16}C_A^2C_FY_f^2(-\mathcal{L}_{m_2}(12(\mathcal{L}_{m_2}-3)\mathcal{L}_{m_2}+3\pi^2+88)-6\zeta_3+3\pi^2+104)$ $+\frac{C_A^2Y_f^2Y_f^2(36(\mathcal{L}_{m_2}-3)\mathcal{L}_{m_2}+3\pi^2+142)}{128m_2^2}$	$C_A C_F$	$\frac{1}{12}C_A C_F(-2\mathcal{L}_{m_2}(-3\mathcal{L}_{m_2}(48C_A C_F+5\beta_0)+552C_A C_F+(48+\pi^2)\beta_0+2\beta_0\mathcal{L}_{m_2}^2))+24(60+\pi^2)C_A C_F+(192+5\pi^2)\beta_0-8\beta_0\zeta_3$ $-\frac{1}{64}Y_f^2 C_A(2\mathcal{L}_{m_2}(48C_A C_F\mathcal{L}_{m_2}-448C_A C_F-18\beta_0+9\beta_0\mathcal{L}_{m_2})$ $+8(152+\pi^2)C_A C_F+(68+3\pi^2)\beta_0)+\frac{9}{32}Y_f^4(5-3\mathcal{L}_{m_2})C_A^2$ $+\frac{Y_f Y_\psi C_A(-\mathcal{L}_{m_2}(6C_A C_F+\beta_0)+19C_A C_F+3\beta_0)}{4m_2}$ $+Y_f^2 Y_\psi(3\mathcal{L}_{m_2}-11)C_A^2$ $\frac{12m_2^2}{16m_2^2}$

**Table 9:** Matching and running,  $\tilde{B}(\mu)$  and  $\tilde{\gamma}_3(\mu)$ , to the threshold form-factor for  $m_2 \gg m_1 \gg M \gg Q$  at  $\mu \sim m_2$ ,  $a \equiv \alpha/(4\pi)$ , and  $\mathcal{L}_{m_2} \equiv \log m_2^2/\mu^2$ .  $\Delta\tilde{B}^{(2)}$  are the mass and coupling renormalisation contributions contributing at two-loop order.  $V_i^{(m_2)}$  and  $F_I$  are two loop vertex and wave-function corrections, given in Appendices C and D.

scale,  $m_1$ , and these are presented up to two-loops in table 10. As for the coupling and mass renormalisation corrections that also contribute at two-loop order, we present these results in table 10. With the above results, due to their generality one can map our results to operators in models that are similar to the SU(N)-Higgs model we discuss here, including those with spontaneous symmetry breaking at a certain scale.

## 7 Application to the Standard Model

The results we have obtained for our SU(N)-Higgs model can now be used to compute results for the SM. We dedicate this treatment to illustrating how the mapping from our to other models of a similar type, which may exhibit SSB, can be achieved. When considering the SM, one must select the correct coupling constants with care, since, it is a chiral gauge theory, and our model is vector-like. One can then obtain results for more than two external particles by combining the two-particle results computed in this paper with the appropriate gauge theory factors included. We focus here on how our results can be used to calculate the radiative corrections to quark and charged lepton production by gauge-invariant currents,  $\bar{Q}_i\gamma_\mu P_L Q_i$  and  $\bar{L}\gamma_\mu P_L L$ , respectively, where  $Q_i$  is the quark doublet for generation,  $i = u, c, t$ , with only the top quark mass,  $m_t$ , taken to be a non-zero fermion mass.

$\mathcal{O}$	$G^{(1)}(\mu)$	$G^{(2)}(\mu)$	$\Delta G^{(2)}(\mu)$
$\psi_2 \Gamma \psi_1$	$\frac{1}{16} C_A (\mathcal{L}_{m_1} (16C_F - 7Y_f^2) - 16C_F + 5Y_f^2)$	$(V_4^{(m_1)} + \frac{1}{2} F_\psi^{(m_1,0)} + \frac{1}{2} F_h^{(0)})$ $-\frac{1}{8} C_A^2 C_F^2 (4\mathcal{L}_{m_1} (3\mathcal{L}_{m_1} - 10) + \pi^2 + 56) +$ $\frac{C_A^2 C_F Y_f^2}{96m_1^2} \{-2\mathcal{L}_{m_1} (2\mathcal{L}_{m_1} (2\mathcal{L}_{m_1} - 9) + \pi^2 + 48)$ $-4\zeta_3 + 3\pi^2 + 120\} - \frac{1}{96} C_A^2 Y_f^4 (36(\mathcal{L}_{m_1} - 3)\mathcal{L}_{m_1} + 3\pi^2 + 142)$	$\frac{1}{12} C_A C_F \{6\mathcal{L}_{m_1} (\beta_0 \mathcal{L}_{m_1} - 2(12C_A C_F + \beta_0))$ $+168C_A C_F + (24 + \pi^2) \beta_0\} + \frac{1}{192} Y_f^2 C_A \{6\mathcal{L}_{m_1} (216C_A C_F$ $+10\beta_0 - 7\beta_0 \mathcal{L}_{m_1}) - 1512C_A C_F - (108 + 7\pi^2) \beta_0\}$ $+ \frac{1}{32} Y_f^4 (32 - 21\mathcal{L}_{m_1}) C_A^2\}$
$\chi_2^{\dagger} \chi_1$	$\frac{1}{24} C_A C_F (6\mathcal{L}_{m_1}^2 + \pi^2 + 48) - \frac{Y_f^2 C_A (\mathcal{L}_{m_1} - 1)}{8m_1^2}$	$(V_2^{(m_1)} + \frac{1}{2} F_\psi^{(m_1,0)} + \frac{1}{2} F_h^{(0)})$ $-\frac{1}{6} C_A^2 C_F^2 \{\mathcal{L}_{m_1} (4(\mathcal{L}_{m_1} - 3)\mathcal{L}_{m_1} + \pi^2 + 24)$ $+2\zeta_3 - \pi^2 - 24\} + \frac{C_A^2 C_F Y_f^2}{96m_1^2} \{-2\mathcal{L}_{m_1} (2\mathcal{L}_{m_1} (2\mathcal{L}_{m_1} - 9) + \pi^2$ $+48) - 4\zeta_3 + 3\pi^2 + 120\}$	$\frac{1}{24} C_A C_F \{-144C_A C_F \mathcal{L}_{m_1} + 2\beta_0 \mathcal{L}_{m_1}^2 + (48 + \pi^2) \beta_0 \mathcal{L}_{m_1}$ $+168C_A C_F + 4\beta_0 (\zeta_3 - 24)\} + \frac{Y_f^2 C_A}{96m_1^2} \{6\mathcal{L}_{m_1} (20C_A C_F + 2\beta_0$ $-\beta_0 \mathcal{L}_{m_1}) - 168C_A C_F - (24 + \pi^2) \beta_0\} - \frac{Y_f^4 (2\mathcal{L}_{m_1} - 3) C_A^2}{32m_1^2}$
$i\chi_2^{\dagger} \overleftrightarrow{D}_\mu \chi_1$	$\frac{1}{12} C_A C_F \{6(\mathcal{L}_{m_1} - 2)\mathcal{L}_{m_1} + \pi^2 + 48\}$ $-\frac{Y_f^2 C_A (\mathcal{L}_{m_1} - 1)}{8m_1^2}$	$(V_3^{(m_1)} + \frac{1}{2} F_\psi^{(m_1,0)} + \frac{1}{2} F_h^{(0)})$ $+\frac{1}{3} C_A^2 C_F^2 \{\mathcal{L}_{m_1} (4(\mathcal{L}_{m_1} - 3)\mathcal{L}_{m_1} + \pi^2 + 24)$ $+2\zeta_3 - \pi^2 - 24\} + \frac{C_A^2 C_F Y_f^2}{48m_1^2} \{-2\mathcal{L}_{m_1} (2\mathcal{L}_{m_1} (2\mathcal{L}_{m_1} - 9) + \pi^2$ $+48) - 4\zeta_3 + 3\pi^2 + 120\}$	$\frac{1}{12} C_A C_F \{\mathcal{L}_{m_1} ((48 + \pi^2) \beta_0 - 72C_A C_F) + 84C_A C_F +$ $2\beta_0 \mathcal{L}_{m_1}^2 - 6\beta_0 \mathcal{L}_{m_1} - \beta_0 (-4\zeta_3 + 96 + \pi^2)\}$ $+ \frac{Y_f^2 C_A}{96m_1^2} \{6\mathcal{L}_{m_1} (20C_A C_F + 2\beta_0 - \beta_0 \mathcal{L}_{m_1}) - 168C_A C_F - (24$ $+ \pi^2) \beta_0\} - \frac{Y_f^4 (2\mathcal{L}_{m_1} - 3) C_A^2}{32m_1^2}$
$\chi_2^{\dagger} \psi_1$	$\frac{1}{16} C_A \{\mathcal{L}_{m_1} (16C_F + Y_f (2Y_s - 9Y_f)) - 16C_F$ $+ Y_f (9Y_f - 4Y_s)\}$	$(V_4^{(m_1)} + \frac{1}{2} F_\psi^{(m_1,0)} + \frac{1}{2} F_h^{(0)})$ $-\frac{1}{12} C_A^2 C_F^2 \{12(\mathcal{L}_{m_1} - 2)\mathcal{L}_{m_1} + \pi^2 + 24\} + \frac{1}{48} C_A^2 C_F Y_f Y_s (\pi^2$ $+24 + 12(\mathcal{L}_{m_1} - 2)\mathcal{L}_{m_1}) + \frac{C_A^2 C_F Y_f^2}{96m_1^2} \{12(\mathcal{L}_{m_1} - 3)\mathcal{L}_{m_1} + \pi^2$ $+48\} - \frac{C_A^2 Y_f^3 Y_f (12(\mathcal{L}_{m_1} - 3)\mathcal{L}_{m_1} + \pi^2 + 48)}{384m_1^2}$	$\frac{1}{12} C_A C_F \{6\mathcal{L}_{m_1} (\beta_0 \mathcal{L}_{m_1} - 2(12C_A C_F + \beta_0)) + (24 + \pi^2) \beta_0$ $+168C_A C_F\} + \frac{1}{64} Y_f^2 C_A \{6\mathcal{L}_{m_1} (88C_A C_F + 6\beta_0 - 3\beta_0 \mathcal{L}_{m_1}) - (68 + 3\pi^2) \beta_0$ $-664C_A C_F\} + \frac{1}{96} Y_f Y_s C_A \{6\mathcal{L}_{m_1} (\beta_0 \mathcal{L}_{m_1} - 4(6C_A C_F + \beta_0)) + 240C_A C_F$ $+ (48 + \pi^2) \beta_0\} + \frac{1}{32} Y_f^3 Y_s \{6\mathcal{L}_{m_1} - 13\} C_A^2 + \frac{9}{32} Y_f^4 (5 - 3\mathcal{L}_{m_1}) C_A^2\}$
$\psi_2 \chi_1$	$\frac{C_A}{24m_1^2} \{6m_1^2 C_F \mathcal{L}_{m_1}^2 + (48 + \pi^2) m_1^2 C_F$ $- 3Y_s^2 (\mathcal{L}_{m_1} - 1)\}$	$(V_5^{(m_1)} + \frac{1}{2} F_\psi^{(m_1,0)} + \frac{1}{2} F_h^{(0)})$ $+\frac{1}{12} C_A^2 C_F^2 \{3\mathcal{L}_{m_1} (4(\mathcal{L}_{m_1} - 5)\mathcal{L}_{m_1} + \pi^2 + 56)$ $+6\zeta_3 - 5\pi^2 - 216\} + \frac{1}{64} C_A^2 C_F Y_f^2 \{-2\mathcal{L}_{m_1} (6\mathcal{L}_{m_1} (2\mathcal{L}_{m_1} - 9)$ $+3\pi^2 + 142) + 9\pi^2 + 350 + 6\psi^{(2)}(1)\}$	$\frac{1}{24} C_A C_F \{-144C_A C_F \mathcal{L}_{m_1} + 168C_A C_F + 2\beta_0 \mathcal{L}_{m_1}^2 + (48 + \pi^2) \beta_0 \mathcal{L}_{m_1}$ $+4\beta_0 (\zeta_3 - 24)\} + \frac{Y_f^2 C_A}{96m_1^2} \{6\mathcal{L}_{m_1} (20C_A C_F + 2\beta_0 - \beta_0 \mathcal{L}_{m_1})$ $-168C_A C_F - (24 + \pi^2) \beta_0\} - \frac{Y_f^4 (2\mathcal{L}_{m_1} - 3) C_A^2}{32m_1^2}$

**Table 10:** Matching corrections,  $G(\mu)$ , to the threshold form-factor for  $m_1 \gg m_2 \gg M \gg Q$  at  $\mu \sim m_1$ ,  $a \equiv \alpha/(4\pi)$ , and  $\mathcal{L}_{m_1} \equiv \log m_1^2/\mu^2$ .  $\Delta \tilde{G}^{(2)}$  are the mass and coupling renormalisation contributions contributing at two-loop order.  $V_i^{(m_1)}$  and  $F_I$  are two loop vertex and wave-function corrections, given in Appendices C and D.

## 7.1 Light Quarks

Let us begin by considering the representation of light quarks in the SM [60]. The first generation of the quark doublet in the mass eigenstate basis is given by,

$$Q_u = \begin{bmatrix} u \\ d' \end{bmatrix} = \begin{bmatrix} t \\ V_{ud}d + V_{us}s + V_{ub}b \end{bmatrix}. \quad (7.1)$$

At the scale,  $Q \gg m_q$ , in the full electro-weak theory, the operator coefficient is assumed to be unity. For the first generation, all quark masses and Yukawa couplings can be neglected, and so the matching is given by combining the gauge boson contributions computed earlier. The operator in SCET at the scale  $Q$  is,

$$\bar{Q}_u \gamma_\mu P_L Q_u \rightarrow c(Q) \left[ \bar{\xi}_{n,p_2}^{(Q_u)} W_n \right] \gamma_\mu P_L \left[ W_{\bar{n}}^\dagger \xi_{\bar{n},p_1}^{(Q_u)} \right], \quad (7.2)$$

where  $\xi^{(Q_u)}$  represents the left-handed EW quark doublet of (7.1) in SCET. Thus, the matching condition,  $c(\mu)$  at the scale  $\mu = Q$  with  $\mathcal{L}_Q = 0$  is,

$$\log c(Q) = a_{EW}(Q) \log c^{(1)}(Q) + a_{EW}(Q)^2 \log c^{(2)}(Q) + \mathcal{O}(a_{EW}^3), \quad (7.3)$$

where,

$$a_{EW}(\mu) = \left( \frac{\alpha_s(\mu)}{4\pi} \frac{4}{3} + \frac{\alpha_2(\mu)}{4\pi} \frac{3}{4} + \frac{\alpha_1(\mu)}{4\pi} \frac{1}{36} \right). \quad (7.4)$$

The gauge couplings have been multiplied by the corresponding gauge factors,  $C_F$ , which are 4/3 for an SU(3) triplet, 3/4 for an SU(2) doublet, and 1/36 for  $Y = 1/6$ . Moreover, the electroweak couplings, renormalized at  $\mu = M_Z$ , are given by [23],

$$\alpha_1(M_Z) = \frac{\alpha_{em}(M_Z)}{\cos^2 \theta_W},$$

$$\alpha_2(M_Z) = \frac{\alpha_{em}(M_Z)}{\sin^2 \theta_W}, \quad (7.5)$$

and their values at  $\mu \sim Q$  are obtained by the usual  $\beta$ -functions of the SM. The theory below  $Q$  is SCET with an  $SU(3) \times SU(2) \times U(1)$  gauge symmetry. In this regime, the SCET current in (7.2) is multiplicatively renormalised with the anomalous dimension  $\gamma(\mu)$  given by,

$$\log \gamma(\mu) = a_{EW}(\mu) \gamma_1^{(1)}(\mu) + a_{EW}(\mu)^2 \gamma_1^{(2)}(\mu) + \mathcal{O}(a_{EW}^3). \quad (7.6)$$

The anomalous dimension,  $\gamma(\mu)$ , is used to run  $c(\mu)$  down to a scale of order the gauge boson mass. One can integrate out the weak gauge bosons sequentially, by first integrating out the  $Z$ -boson at  $\mu = M_Z$ , followed by the  $W$ -boson at  $\mu = M_W$ . This is not a good choice to use for the SM, as  $M_W/M_Z$  is not negligible, and summing powers of  $M_W/M_Z$  is more important than summing  $\alpha \log^2 M_W/M_Z$  terms. Instead, we integrate out the  $W$  and  $Z$  at a common scale, chosen to be  $\mu = M_Z$ . In this way, we match directly from an  $SU(3) \times SU(2) \times U(1)$  gauge theory onto a  $SU(3) \times U(1)_{em}$  gauge theory of gluons and photons, which lacks the complications of an intermediate stage of broken EW symmetry where  $Z$  is integrated out, but not the  $W$ . Moreover, the Higgs corrections for light particles are sub-leading as the Yukawa coupling is proportional to the light mass and thus, are suppressed. At the scale  $\mu = M_Z$ , integrating out the  $W$  and  $Z$  bosons give a matching correction to the SCET operator,

$$\begin{aligned} \left[ \bar{\xi}_{n,p_2}^{(Q_u)} W_n \right] \gamma_\mu P_L \left[ W_n^\dagger \xi_{\bar{n},p_1}^{(Q_u)} \right] \rightarrow d^{(u)} \left[ \bar{\xi}_{n,p_2}^{(u)} W_n \right] \gamma_\mu P_L \left[ W_n^\dagger \xi_{\bar{n},p_1}^{(u)} \right] + \\ d^{(d')} \left[ \bar{\xi}_{n,p_2}^{(d')} W_n \right] \gamma_\mu P_L \left[ W_n^\dagger \xi_{\bar{n},p_1}^{(d')} \right]. \end{aligned} \quad (7.7)$$

Since the EW symmetry is broken, the  $u$  and  $d'$  parts of the operator get different matching corrections. The matching corrections are as follows,

$$\begin{aligned} \log d^{(u)}(M_Z) = a_1 \log d^{(1)}(M_W) + a_1^2 \log d^{(2)}(M_W) + \mathcal{O}(a_1^3) \\ + a_2 \log d^{(1)}(M_Z) + a_2^2 \log d^{(2)}(M_Z) + \mathcal{O}(a_2^3) \end{aligned} \quad (7.8)$$

where the terms proportional to  $a_1$  and  $a_2$  correspond the  $Z$  and  $W$  contributions, respectively, and,

$$a_1 = \frac{\alpha_{em}}{4\pi \sin^2 \theta_W \cos^2 \theta_W} \left( \frac{1}{2} - \frac{2}{3} \sin^2 \theta_W \right)^2 \quad (7.9)$$

$$a_2 = \frac{\alpha_{em}}{4\pi \sin^2 \theta_W \cos^2 \theta_W} \left( \frac{1}{2} \right)^2. \quad (7.10)$$

Below  $M_Z$ , the operators in (7.1) are multiplicatively renormalised, with anomalous dimensions,

$$\gamma^{(u)}(\mu) = \tilde{a}_1(\mu) \gamma_1^{(1)}(\mu) + \tilde{a}_1(\mu)^2 \gamma_1^{(2)}(\mu) + \mathcal{O}(\tilde{a}_1^3), \quad (7.11)$$

$$\gamma^{(d')}(\mu) = \tilde{a}_2(\mu) \gamma_1^{(1)}(\mu) + \tilde{a}_2(\mu)^2 \gamma_1^{(2)}(\mu) + \mathcal{O}(\tilde{a}_2^3) \quad (7.12)$$

such that

$$\tilde{a}_1(\mu) = \left\{ \frac{\alpha_s(\mu)}{4\pi} \frac{4}{3} + \frac{\alpha_{em}(\mu)}{4\pi} \frac{4}{9} \right\} \quad \text{and} \quad \tilde{a}_2(\mu) = \left\{ \frac{\alpha_s(\mu)}{4\pi} \frac{4}{3} + \frac{\alpha_{em}(\mu)}{4\pi} \frac{1}{9} \right\}, \quad (7.13)$$

for the  $u$  and  $d'$  terms. The final result for the operator at a low scale is thus,

$$\begin{aligned} \left[ \bar{\xi}_{n,p_2}^{(Q_u)} W_n \right] \gamma_\mu P_L \left[ W_n^\dagger \xi_{\bar{n},p_1}^{(Q_u)} \right] \rightarrow c^{(u)} \left[ \bar{\xi}_{n,p_2}^{(u)} W_n \right] \gamma_\mu P_L \left[ W_n^\dagger \xi_{\bar{n},p_1}^{(u)} \right] + \\ c^{(d')} \left[ \bar{\xi}_{n,p_2}^{(d')} W_n \right] \gamma_\mu P_L \left[ W_n^\dagger \xi_{\bar{n},p_1}^{(d')} \right], \end{aligned} \quad (7.14)$$

with

$$\log c^u(\mu) = \log c(Q) + \int_Q^{M_Z} \frac{d\mu}{\mu} \gamma(\mu) + \log d^{(u)}(M_Z) + \int_{M_Z}^\mu \frac{d\mu}{\mu} \gamma^{(u)}(\mu) \quad (7.15)$$

$$\log c^d(\mu) = \log c(Q) + \int_Q^{M_Z} \frac{d\mu}{\mu} \gamma(\mu) + \log d^{(d')}(M_Z) + \int_{M_Z}^\mu \frac{d\mu}{\mu} \gamma^{(d')}(\mu). \quad (7.16)$$

The EFT operator in (7.14) can then be used to compute processes such as dijet production using SCET [5]. For jet production, the renormalisation scale,  $\mu$ , would be chosen to be of order the jet invariant mass, or 30 GeV for jets at the LHC.

## 7.2 Leptons

The computation for the radiative corrections to the lepton current,  $\bar{L}\gamma_\mu P_L L$ , where  $L$  is the lepton doublet,

$$L = \begin{pmatrix} \nu \\ l \end{pmatrix}, \quad (7.17)$$

is identical to that for the quark doublet, aside from a few replacements. The full theory operator at the low scale,  $\mu$ , is,

$$\begin{aligned} \bar{L}\gamma_\mu P_L L \rightarrow & c^{(\nu)} \left[ \bar{\xi}_{n,p_2}^{(\nu)} W_n \right] \gamma_\mu P_L \left[ W_n^\dagger \xi_{n,p_1}^{(\nu)} \right] + \\ & c^{(l)} \left[ \bar{\xi}_{n,p_2}^{(l)} W_n \right] \gamma_\mu P_L \left[ W_n^\dagger \xi_{n,p_1}^{(l)} \right], \end{aligned} \quad (7.18)$$

with the coefficients given by (7.16) and replacements  $u \rightarrow \nu$ ,  $d' \rightarrow l$ , along with different gauge theory factors which implies the following coupling replacements,

$$a_{EW}(\mu) \rightarrow a'_{EW}(\mu) = \left( \frac{\alpha_2(\mu)}{4\pi} \frac{3}{4} + \frac{\alpha_1(\mu)}{4\pi} \frac{1}{4} \right), \quad (7.19)$$

$$a_1 \rightarrow a'_1 = \frac{\alpha_{em}}{4\pi \sin^2 \theta_W \cos^2 \theta_W} \left( \frac{1}{2} \right)^2, \quad (7.20)$$

$$a_2 \rightarrow a'_2 = \frac{\alpha_{em}}{4\pi \sin^2 \theta_W \cos^2 \theta_W} \left( -\frac{1}{2} + \sin^2 \theta_W \right)^2, \quad (7.21)$$

$$\tilde{a}_1(\mu) \rightarrow \tilde{a}'_1(\mu) = 0, \quad (7.22)$$

$$\tilde{a}_2(\mu) \rightarrow \tilde{a}'_2(\mu) = \frac{\alpha_{em}(\mu)}{4\pi}, \quad (7.23)$$

which provides us with the leptonic equivalent of the previous result.

## 7.3 Top Quarks

In this subsection, we show how our results can be used to compute the radiative corrections to  $t\bar{t}$ -production by a gauge-invariant vector current  $\bar{Q}_t \gamma_\mu P_L Q_t$ , where  $Q_t$  is the left-handed quark doublet in the SM. We may write the quark doublet in the mass eigenstate basis,

$$Q_t = \begin{pmatrix} t \\ b' \end{pmatrix} = \begin{bmatrix} t \\ V_{td}d + V_{ts}s + V_{tb}b \end{bmatrix}. \quad (7.24)$$

We will neglect all quark masses other than  $m_t$ . This example is useful as it illustrates how to use the fermion mass and Higgs exchange contributions computed in our model. We will examine both the Sudakov and threshold regimes in this case as they are both available to us in this example.

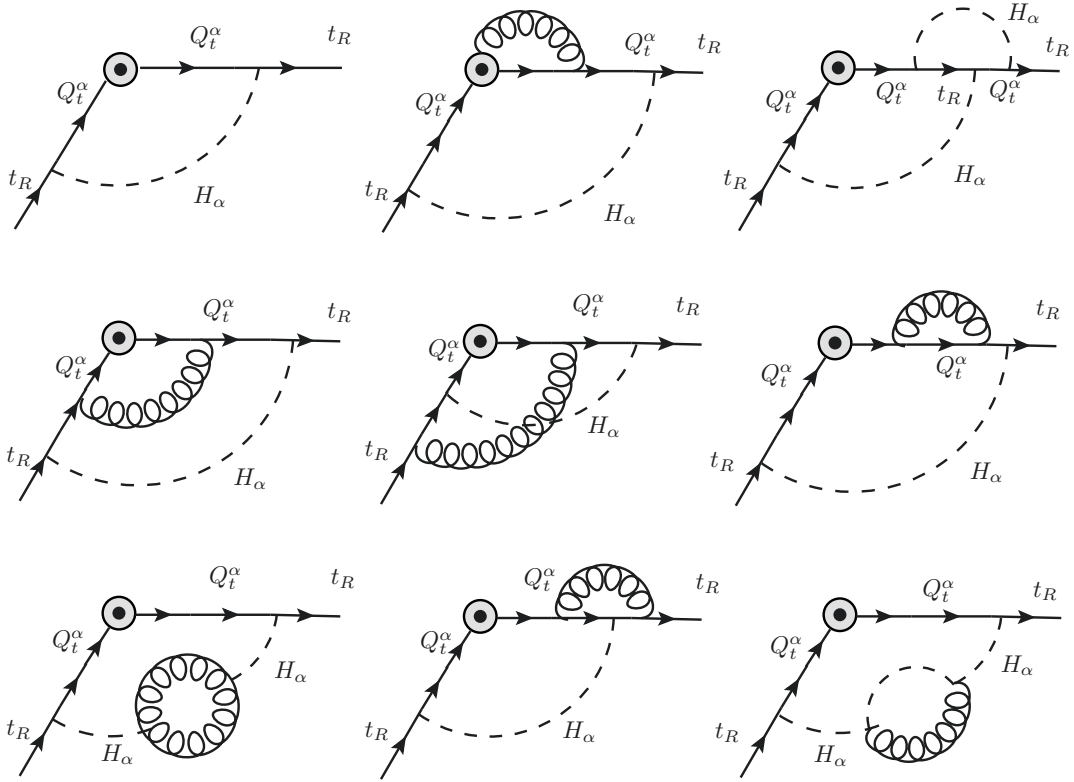
### 7.3.1 Sudakov Regime

In the Sudakov regime, the operator in SCET at the scale,  $\mu \sim Q$ , is as follows,

$$\begin{aligned} \bar{Q}_t \gamma_\mu P_L Q_t \rightarrow & c_L(Q) \left[ \bar{\xi}_{n,p_2}^{(Q_t)} W_n \right] \gamma_\mu P_L \left[ w_n^\dagger \xi_{n,p_1}^{(Q_t)} \right] + \\ & c_R(Q) \left[ \bar{\xi}_{n,p_2}^{(t)} W_n \right] \gamma_\mu P_R \left[ w_n^\dagger \xi_{n,p_1}^{(t)} \right], \end{aligned} \quad (7.25)$$

where  $\xi^{(Q_t)}$  and  $\xi^{(t)}$  represent the left-handed and right-handed  $t$ -quark doublet, (7.24), and singlet,  $t_R$ , respectively, in SCET with gauge indices suppressed. The reason  $t_R$  appears in this case is due to Higgs exchange graphs which are chiral in the SM, and have been computed in our model value which is a vector-like theory, thus when mapping to the SM we must plaster on the fact that the Yukawa vertex flips the fermion chirality. Practically, the Higgs exchange vertex correction mixes the  $Q_L$  operator with the  $t_R$  operator. The matching condition at the scale  $\mu = Q$  is then given by  $c_{L/R}(Q)$ , where one splits the left and right handed contributions of  $c(Q)$  which now has non-zero Yukawa couplings. Hence,  $c_R(Q)$  includes all terms which arise from Higgs exchange graphs of type illustrated in figure 4; and the remaining graphs contribute to  $c_L(Q)$ . Note further that one must include appropriate factors of two for terms in  $c_{L/R}$  arising





**Figure 4:** Vertex contributions to matching,  $c_R(\mu)$ , at one and two loop order. Higgs exchanges cause  $\bar{Q}_t \gamma_\mu P_L Q_t$  to mix with  $\bar{t}_R \gamma_\mu P_R t$  and the index,  $\alpha$ , is the fundamental SU(2) and index and is summed over.

from summing over each closed SU(2) index loop, i.e. because both the Higgs and  $Q_t$  are SU(2) doublets. As for the wavefunction correction, the  $t_L$  and  $t_R$  field renormalisation contributions which include Higgs exchange must also include appropriate factors of two from loops with SU(2) index summation.

The theory below  $Q$  is SCET with an  $SU(3) \times SU(2) \times U(1)$  gauge symmetry. In this regime the two operators in (7.25) are multiplicatively renormalised with anomalous dimensions (again splitting chiral contributions in the same way as for matching),

$$\frac{dc_L(\mu)}{d\mu} = \gamma_L(\mu)c_L \quad \text{and} \quad \frac{dc_R(\mu)}{d\mu} = \gamma_R(\mu)c_R. \quad (7.26)$$

At this scale, the Higgs vertex graph, which causes  $c_L/c_R$  mixing, is  $1/Q^2$  suppressed. The anomalous dimension  $\gamma$  is as usual used to run  $c_L$  and  $c_R$  down to a scale of order  $m_t$ . At  $\mu \sim m_t$  there are several different methods one can use. Since  $\{m_t, M_{W,Z}, M_H\}$  are not widely separated, one can integrate them all out together. In this way, one goes directly from an  $SU(3) \times SU(2) \times U(1)$  invariant theory to a  $SU(2) \times U(1)_{em}$  gauge theory, with broken  $SU(2) \times U(1)$  symmetry and no EW gauge bosons. This procedure keeps the entire mass dependence on the four mass scales. At the scale  $\mu = m_t$  the  $t$ -quark SCET field is replaced by the heavy quark field  $t_v$ , whereas the  $b'$  quark in the SCET field  $\xi^{(Q_t)}$  remains a SCET field,  $\xi^{(b')}$ . The operator matching is,

$$\left[ \bar{\xi}_{n,p_2}^{(Q_t)} W_n \right] \gamma_\mu P_L \left[ W_n^\dagger \xi_{n,p_1}^{(Q_t)} \right] \rightarrow a_1 \left[ \bar{\xi}_{n,p_2}^{(b')} W_n \right] \gamma_\mu P_L \left[ W_n^\dagger \xi_{n,p_1}^{(b')} \right] + a_2 \bar{t}_{v_2} t_{v_1} \quad (7.27)$$

$$\left[ \bar{\xi}_{n,p_2}^{(t)} W_n \right] \gamma_\mu P_R \left[ W_n^\dagger \xi_{n,p_1}^{(t)} \right] \rightarrow a_3 \bar{t}_{v_2} t_{v_1}, \quad (7.28)$$



where the matching coefficients,  $a_i$ , and are obtained using the results of the results from section 5.2. All running couplings are renormalized at  $\mu = m_t$ . The expressions are given by adding the contributions due to the  $W/Z$ ,  $g$ ,  $\gamma$  and  $H, h^{0,+}$ , where  $h^{0,+}$  are the unphysical Higgs scalars present in  $R_\xi = 1$  gauge. Lastly, below  $\mu = m_t$ , the  $\bar{t}_{v_2} t_{v_1}$  operator has the anomalous dimension,

$$\gamma_3 = a\gamma_3^{(1)} + a^2\gamma_3^{(2)} + \mathcal{O}(a^3), \quad a = \left( \frac{4}{3} \frac{\alpha_s}{4\pi} + \frac{4}{9} \frac{\alpha_{em}}{4\pi} \right) \quad (7.29)$$

from the fourth column of table 5, with the given group theory factor replacements. The radiative corrections to the  $\bar{t}t$  operator can then be combined with known methods to obtain  $t$ -quark decay distributions. The QCD corrections (the  $\alpha_s$  terms) have already been included in the analysis of previous work [30]. The new results in this paper are the additional two-loop EW radiative corrections, including Higgs effects.

### 7.3.2 Threshold Regime

In the Threshold regime at  $\mu \sim m_t$ , and  $m_t$  is the largest scale in the problem. Although the scales,  $\{m_t, M_{W,Z}, M_H\}$ , are not widely separated, one can no longer integrate them all out together as in the Sudakov regime. One can integrate out the scales  $m_t$ ,  $M_{W,Z}$  and  $M_H$  in various ways, e.g. one can integrate out each particle at a scale  $\mu$  equal to its mass, or integrate out one or more particles simultaneously at some common value of  $\mu$ , as was done in Sudakov regime. The most experimentally relevant way to integrate out the relevant scales is to first integrate out the top quark as  $m_t \sim 172$  GeV and  $m_t > m_H > M_{W,Z} \gg m_b$ , which leads to an effective theory that breaks  $SU(2) \times U(1)$  invariance as the  $b'$  quark remains along with dynamical  $W/Z$  bosons. From which one integrates out the Higgs first [26], as  $M_H \sim 125$  GeV, and then the  $W/Z$  bosons at a common scale,  $M_Z \sim 81$  GeV. Otherwise if one wants to avoid breaking  $SU(2) \times U(1)$  invariance, one is free to integrate out both  $t$  and  $b'$  at a common scale,  $\mu \sim m_t$ . Then one can either integrate all massive bosons at a common scale,  $M_Z$  or separate into integrating out the Higgs first then the  $W/Z$  at common scale. We consider the former here to illustrate and leave further analyses and numerical comparisons to upcoming work in the SM and beyond, as there is the further case of heavy-light currents to consider as well.

Integrating out both  $t$  and  $b'$  at the common scale  $\mu = m_t$ , below  $m_t$  the fields are replaced by their heavy quark counterparts  $t_v$  and  $b_v$ , respectively. The operator matching is,

$$\bar{Q}_t \gamma_\mu P_L Q_t \rightarrow b_1 \bar{t}_{v_2} t_{v_1} + b_2 \bar{b}_{v_2} b_{v_1} \quad \text{and} \quad \bar{t}_R \gamma_\mu P_R t_R \rightarrow b_3 \bar{t}_{v_2} t_{v_1}, \quad (7.30)$$

where the matching coefficients,  $b_i$ , are obtained using the matching coefficient,  $b(m_t)$ , from section 6 with the appropriate graphs and group theory factors for each part, adding the contributions due to both  $b'$  and  $t$  individually. All running couplings are then to be renormalized at  $\mu = m_t$ . As in the Sudakov case, below  $\mu = m_t$ , the  $\bar{t}_{v_2} t_{v_1}$  operator has the anomalous dimension for running from  $m_t \rightarrow M_Z$  and is given by,

$$\gamma_3 = a\gamma_3^{(1)} + a^2\gamma_3^{(2)} + \mathcal{O}(a^3), \quad a = \left( \frac{4}{3} \frac{\alpha_s}{4\pi} + \frac{4}{9} \frac{\alpha_{em}}{4\pi} \right) \quad (7.31)$$

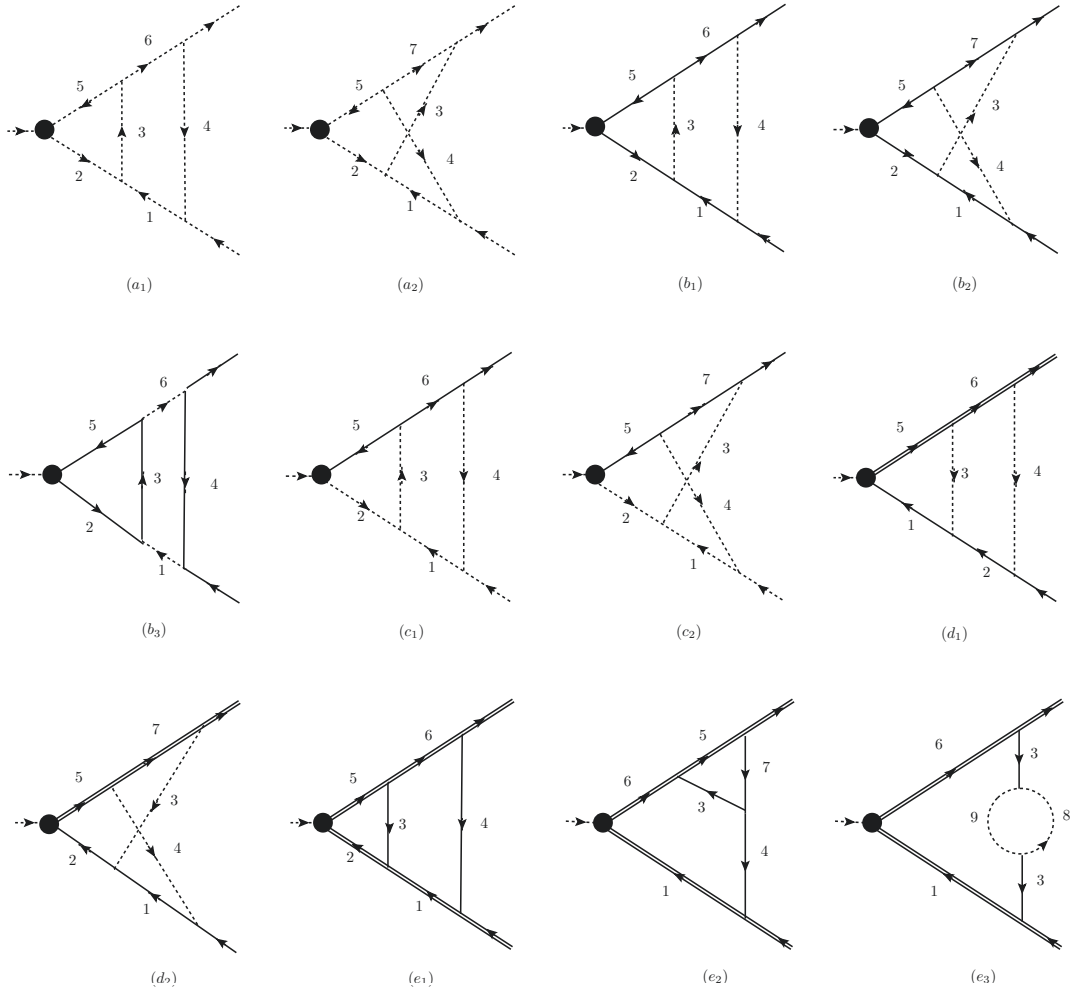
from the fourth column of table 5, with all running coupling renormalised at  $\mu = m_t$ . The remaining quantities needed are the matching contributions at  $\mu = M_Z$ , where the operators  $\bar{t}_{v_2} t_{v_1}$  and  $\bar{b}_{v_2} b_{v_1}$  above  $M_Z$  are matched to their counterparts below  $M_Z$  with massive bosons integrated out,

$$\bar{t}_{v_2} t_{v_1} \rightarrow u_1 \bar{t}_{v_2} t_{v_1} \quad \text{and} \quad \bar{b}_{v_2} b_{v_1} \rightarrow u_2 \bar{b}_{v_2} b_{v_1} \quad (7.32)$$

where the matching coefficients,  $u_i$ , are obtained using the matching coefficient,  $u(M_Z)$ , from combining contributions from tables in sections 5.2 and 6 with all running coupling renormalised at  $\mu = M_Z$ . The expressions are given by adding individual contributions due to each massive boson in the SM.

## 8 Technical Calculation

The Feynman diagrams we needed were generated using **QGRAF** [56], the output of which was then processed using **FORM** [63] to express the diagrams in terms of a linear combination of a set of scalar integrals. We then reduced these integrals to a much smaller set of so-called master integrals (MIs) using integration-by-parts identities (IBPs) [14], with the help of **LiteRed** [44] and home-grown tools. Our master integrals in



**Figure 5:** Topologies required for the calculation of two-loop form factors. Solid lines represent massive particles, double lines represent heavy particles, dashed lines correspond to massless propagators. Arrows represent direction of momentum.  $(a_i)$  represent topologies of MIs at  $\mu \sim Q$  (Sudakov),  $(b_i)$  represent topologies of MIs at  $\mu \sim m_{1,2}$  (threshold),  $(c_i)$  represent topologies of MIs at  $\mu \sim m_2$  (threshold),  $(d_i)$  represent topologies of MIs at  $\mu \sim m_1$  (threshold) and  $(e_i)$  represent topologies of MIs at  $\mu \sim M$  (Sudakov/threshold).

some cases are dependent on two mass scales taken to be not widely separated, either the external particle masses or the bosonic masses, respectively. One can perform these integrals numerically but to obtain analytic results we expand such amplitudes in the mass difference to NLO in said difference, leading to single scale integrals. Once the integrals are maximally reduced, all that remains is to evaluate the master integrals. As these procedures are well-known, we refrain from delving into too much detail.

We focus here on the calculation of MIs of the two-loop vertex and wave-function contributions. The full theory integrals have been computed analytically [1]. The result was found as a Laurent expansion in the dimensional parameter,  $\epsilon$ , using the differential equation method [42, 59]. We evaluate our MIs in the effective theory, HPET, which are not known analytically, in the same fashion. One requires the two-loop master integrals up to a sufficiently high order in  $\epsilon$  to obtain  $\mathcal{O}(\epsilon^2)$  accuracy in the form factors. As we are still evaluating the heavy-heavy and heavy-light master integrals, we will keep the details of this calculation for upcoming work. Instead, we briefly outline our classification scheme for reference and present the results

we do have explicitly. The master integrals are classified according to their underlying topology.

We begin by distinguishing between the vertex topologies for external full theory fields displayed in figures 5 (a-c). The master integrals for all topologies can be expressed in terms of a single integral family with seven propagators given by,

$$J_{\{\nu_1^{(m)} \dots \nu_7^{(m)}\}}^{(s)} = [(4\pi)^{2-\epsilon} e^{\gamma_E \epsilon}]^2 \int \not{d}^D l_1 \not{d}^D l_2 \frac{1}{D_1^{\nu_1}(m) \dots D_7^{\nu_7}(m)}, \quad (8.1)$$

where  $l_i : i = 1, 2$  are the loop momenta,  $s$  is the scale in the EFT formalism at which the MIs play a role, and,

$$\begin{aligned} D_1(m) &= l_1^2 - m^2, & D_2(m) &= l_2^2 - m^2, & D_3(m) &= (l_1 + l_2)^2 - m^2, \\ D_4(m) &= (l_1 - p_1)^2 - m^2, & D_5(m) &= (l_2 - q)^2 - m^2, \\ D_6(m) &= (l_1 + q)^2 - m^2, & D_7(m) &= (l_1 + l_2 - q - p_1)^2 - m^2. \end{aligned} \quad (8.2)$$

$$D_6(m) = (l_1 + q)^2 - m^2, \quad D_7(m) = (l_1 + l_2 - q - p_1)^2 - m^2. \quad (8.3)$$

Here the  $p_i : i = 1, 2$  are the external momenta, which are taken to be on-shell ( $p_i^2 = m_i^2$ ) and  $q = p_2 - p_1$  is the usual transfer momentum. We therefore label the MIs by the exponents,  $\nu_1 \dots \nu_7$  of the denominators,  $D_1 \dots D_7$ . Note the single mass scale in our denominators, this arises from the fact that for integrals involving two mass scales or more, we expand our results in the difference of mass scales up to NLO. For instance, for graphs that include propagators of both  $W$  and Higgs bosons, we expand about  $\Delta M = M_W - M_H$ , assuming them to be not widely separated. This is done purely so we can present our results analytically as any number of scales can be handled numerically, moreover our choice to expand to NLO is for presentability as there is no issue in expanding the amplitudes to higher orders in  $\Delta M$  computationally.

In the Sudakov regime the vertex matching contributions at the scale  $\mu \sim Q$ ,  $V_i^{(Q)}$ , has all mass scales set to zero as they are taken to be IR, and thus  $m = 0$  in cases below, in which case we have MIs with topology given by figures 5 (a). Post-reduction one is left with the following master integrals,

$$J_{1010100}^{(Q)} = \frac{Q^2 \left( 13\epsilon - 4\epsilon \log\left(\frac{Q^2}{\mu^2}\right) + 2 \right)}{8\pi^4 \epsilon}, \quad (8.4a)$$

$$J_{1100110}^{(Q)} = \frac{\epsilon \left( 2(\pi^2 - 72)\epsilon - 3 \right) + 6\epsilon \mathcal{L}_Q (\epsilon - 4\epsilon \mathcal{L}_Q + 4) - 12}{12\pi^4 \epsilon^2}, \quad (8.4b)$$

$$J_{1100101}^{(Q)} = -\frac{\epsilon \left( (114 + \pi^2)\epsilon + 30 \right) + 12\epsilon \mathcal{L}_Q (-5\epsilon + \epsilon \mathcal{L}_Q - 1) + 6}{12\pi^4 \epsilon^2}, \quad (8.4c)$$

$$\begin{aligned} J_{1111101}^{(Q)} &= \frac{59\pi^4 \epsilon^4 + 120\pi^2 \epsilon^2 - 80\epsilon \mathcal{L}_Q (3\pi^2 \epsilon^2 + \epsilon \mathcal{L}_Q (-3\pi^2 \epsilon^2 + \epsilon \mathcal{L}_Q (\epsilon \mathcal{L}_Q - 2) + 3) - 3)}{120\pi^4 \epsilon^4 Q^2} \\ &+ \frac{83\epsilon^3 \zeta_3 (1 - 2\epsilon \mathcal{L}_Q) - 3}{3\pi^4 \epsilon^4 Q^2}, \end{aligned} \quad (8.4d)$$

where  $\mathcal{L}_Q \equiv \log Q^2/\mu^2$  and  $\zeta_n$  denotes the Riemann  $\zeta$ -function,

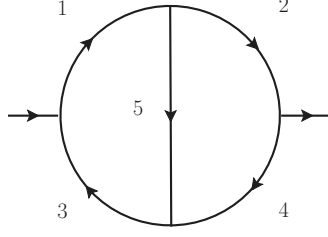
$$\zeta_m = \sum_{k=1}^{\infty} \frac{1}{k^n}, \quad n \geq 2 : n \in \mathbb{N}, \quad (8.5)$$

and these integrals have been verified from previous work [62]. On the other hand, in the threshold regime, the full theory MIs have topologies represented by figures 5 (b,c). Due to the threshold limit,  $Q \rightarrow 0$ , the MIs are further reduced down to two-loop self-energy topologies, as shown in figure 6, and analytic expressions can be obtained from the standard **ON-SHELL2** package [29]. The master integrals can therefore be expressed in terms of a single integral family with five propagators given by,

$$J_{\{\nu_1^{(m)} \dots \nu_5^{(m)}\}}^{(s)} = [(4\pi)^{2-\epsilon} e^{\gamma_E \epsilon}]^2 \int \not{d}^D l_1 \not{d}^D l_2 \frac{1}{D_1^{\nu_1}(m) \dots D_5^{\nu_5}(m)}, \quad (8.6)$$

where  $l_i : i = 1, 2$  are the loop momenta,  $s$  is the scale in the EFT formalism at which the MIs play a role, and,

$$\begin{aligned} D_1(m) &= l_1^2 - m^2, & D_2(m) &= l_2^2 - m^2, & D_3(m) &= (l_1 - p)^2 - m^2, \\ D_4(m) &= (l_2 - p)^2 - m^2, & D_5(m) &= (l_1 - l_2)^2 - m^2. \end{aligned} \quad (8.7)$$



**Figure 6:** Full theory self-energy topology. Arrows represent momentum direction. The MIs associated to other topologies are subsets of the MIs required for this topology.

To begin with, the vertex matching contributions at the scale  $\mu \sim m_2 \gg m_1, M$ , for matching from the full theory to heavy-light operators is represented by  $V_i^{(m_2)}$ . Post-reduction one is left with the following master integrals,

$$J_{1^{(m)}1^{(m)}000}^{(m_2)} = -\frac{m_2^4(2\epsilon\mathcal{L}_{m_2}(\epsilon\mathcal{L}_{m_2} - 2\epsilon - 1) + \epsilon(\epsilon(\zeta_2 + 3) + 2) + 1)}{\epsilon^2}, \quad (8.8a)$$

$$J_{11001^{(m)}}^{(m_2)} = -\frac{m_2^2(2\epsilon\mathcal{L}_{m_2}(\epsilon\mathcal{L}_{m_2} - 3\epsilon - 1) + \epsilon(\epsilon(3\zeta_2 + 7) + 3) + 1)}{2\epsilon^2}, \quad (8.8b)$$

$$J_{11^{(m)}100}^{(m_2)} = -\frac{m_2^2(2\epsilon\mathcal{L}_{m_2}(\epsilon\mathcal{L}_{m_2} - 3\epsilon - 1) + \epsilon(7\epsilon + 3) + 1)}{\epsilon^2}, \quad (8.8c)$$

$$J_{01^{(m)}101}^{(m_2)} = -\frac{m_2^2(4\epsilon\mathcal{L}_{m_2}(2\epsilon\mathcal{L}_{m_2} - 5\epsilon - 2) + \epsilon(\epsilon(20\zeta_2 + 11) + 10) + 4)}{8\epsilon^2}, \quad (8.8d)$$

$$J_{011^{(m)}02^{(m)}}^{(m_2)} = -\frac{-6\epsilon\mathcal{L}_{m_2}(\epsilon(-\mathcal{L}_{m_2}) + \epsilon + 1) + \epsilon(\epsilon(6S_1 + \zeta_2 - 3) + 3) + 3}{6\epsilon^2}, \quad (8.8e)$$

$$J_{01^{(m)}1^{(m)}01^{(m)}}^{(m_2)} = -\frac{m_2^2(4\epsilon\mathcal{L}_{m_2}(6\epsilon\mathcal{L}_{m_2} - 17\epsilon - 6) + \epsilon(\epsilon(12\zeta_2 + 59) + 34) + 12)}{8\epsilon^2}, \quad (8.8f)$$

$$J_{01101}^{(m_2)} = \frac{1}{8}m_2^2 \left( \frac{2}{\epsilon} - 4\mathcal{L}_{m_2} + 13 \right), \quad (8.8g)$$

$$J_{1^{(m)}001^{(m)}1}^{(m_2)} = -\frac{m_2^2(4\epsilon\mathcal{L}_{m_2}(4\epsilon\mathcal{L}_{m_2} - 11\epsilon - 4) + \epsilon(\epsilon(12S_1 + 8\zeta_2 + 35) + 22) + 8)}{8\epsilon^2}, \quad (8.8h)$$

$$J_{11110}^{(m_2)} = \frac{2\epsilon\mathcal{L}_{m_2}(\epsilon(-\mathcal{L}_{m_2}) + 4\epsilon + 1) + \epsilon(\epsilon(\zeta_2 - 12) - 4) - 1}{\epsilon^2}, \quad (8.8i)$$

$$J_{11^{(m)}111^{(m)}}^{(m_2)} = \frac{\frac{27}{2}S_1S_2 + i\pi\zeta_2 - 3\zeta_3}{m_2^2}, \quad (8.8j)$$

where  $S_1 = \frac{\pi}{\sqrt{3}}$  and  $S_2 = \frac{4}{9\sqrt{3}}Cl_2(\frac{\pi}{3})$  such that  $Cl_2$  denotes the second order Clusen function,

$$Cl_2(\theta) = -\int_0^\theta \log |2 \sin(\theta/2)| d\theta : \quad 0 < \theta < 2\pi. \quad (8.9)$$

Next, we consider the vertex contributions,  $V_i^{(m)}$ , at the scale  $\mu \sim m_{1,2} \gg M$ , for matching from the full theory to heavy-heavy operators. After reduction we have the following master integrals,

$$J_{001^{(m)}1^{(m)}0}^{(m)} = -\frac{m^4(2\epsilon\mathcal{L}_m(\epsilon\mathcal{L}_m - 2\epsilon - 1) + \epsilon(\epsilon(\zeta_2 + 3) + 2) + 1)}{\epsilon^2}, \quad (8.10a)$$

$$J_{1001^{(m)}1}^{(m)} = -\frac{m^2(4\epsilon\mathcal{L}_m(2\epsilon\mathcal{L}_m - 5\epsilon - 2) + \epsilon(\epsilon(20\zeta_2 + 11) + 10) + 4)}{8\epsilon^2}, \quad (8.10b)$$

$$J_{01^{(m)}1^{(m)}01^{(m)}}^{(m)} = \frac{m^2(-2\epsilon\mathcal{L}_m(\epsilon\mathcal{L}_m + \epsilon(S_1 - 3) - 1)}{\epsilon^2} \quad (8.10c)$$

$$+ \frac{\epsilon(-\epsilon(S_1(\log(3) - 3) - 9S_2 + \zeta_2 + 7) + S_1 - 3) - 1}{\epsilon^2}. \quad (8.10d)$$

As for the full theory two-loop wave-function contributions, present in Appendix D, it is well-known that they map to MIs illustrated by figure 6, and thus we refrain from going into detail.

With regards to the effective theory MIs, the HPET vertex and wave-function contributions have MIs with topologies represented by figures 5 (d,e) and 7, respectively, the analytic expressions of which we will present in upcoming work [3]. Instead we present our results in an attached ancillary file, in general dimension,  $d$ , with unevaluated master integrals, defined by the notation below. We begin with considering the heavy-light currents in figures 5 (d), the master integrals of which can be expressed in terms of a single integral family with seven propagators given by,

$$R_{\nu_1 \dots \nu_7}^{(s)} = [(4\pi)^{2-\epsilon} e^{\gamma_E \epsilon}]^2 \int \not{d}^D l_1 \not{d}^D l_2 \frac{1}{D_1^{\nu_1} \dots D_7^{\nu_7}}, \quad (8.11)$$

where  $v_2$  is the heavy particle velocity,  $p_1$  and  $m_1$  are the full theory field momentum and mass,  $p_1 \cdot v_2 \equiv w'$ , and thus,

$$\begin{aligned} D_1(m) &= (l_1 + p_1)^2 - m^2, & D_2(m) &= (l_2 + p_1)^2 - m^2, & D_3 &= (l_1 - l_2)^2, & D_4 &= l_2^2, \\ D_5 &= l_1 \cdot v_2, & D_6 &= l_2 \cdot v_2, & D_7 &= (l_1 - l_2) \cdot v_2 + w'. \end{aligned} \quad (8.12)$$

Similarly, for the heavy-heavy vertex contributions, all sub-topologies can be mapped to the largest unique two that are shown in figure 5 (e). We can again express all master integrals in terms of nine propagators given by,

$$K_{\nu_1 \dots \nu_9}^{(s)} = [(4\pi)^{2-\epsilon} e^{\gamma_E \epsilon}]^2 \int \not{d}^D l_1 \not{d}^D l_2 \frac{1}{D_1^{\nu_1} \dots D_9^{\nu_9}}, \quad (8.13)$$

where  $v_{1,2}$  are the heavy particle velocities,  $v_1 \cdot v_2 \equiv w$ ,  $M$  is the mass of exchanged bosons, and

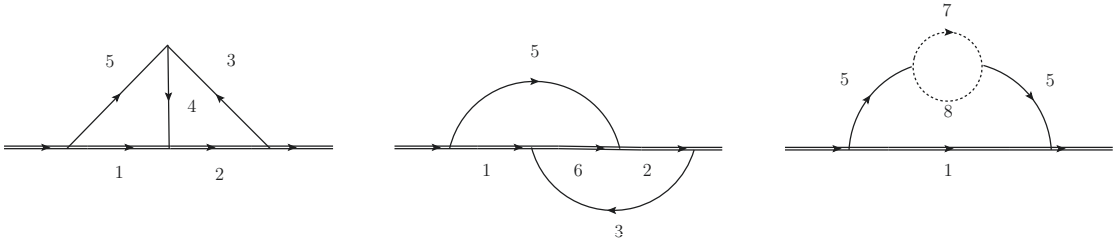
$$\begin{aligned} D_1 &= l_2 \cdot v_1, & D_2 &= l_1 \cdot v_1, & D_3 &= (l_1 - l_2)^2 - M^2, & D_4 &= l_1^2 - M^2, \\ D_5(M) &= l_1 \cdot v_2, & D_6(M) &= l_2 \cdot v_2, & D_7(M) &= l_2^2 - M^2, & D_8 &= l_2^2, & D_9 &= l_1^2 \end{aligned} \quad (8.14)$$

Finally, we examine the wave-function contributions of which all topologies are mapped to those shown in figure 7. In this case we can express all MIs in terms of six propagators given by,

$$L_{\nu_1 \dots \nu_8}^{(s)} = [(4\pi)^{2-\epsilon} e^{\gamma_E \epsilon}]^2 \int \not{d}^D l_1 \not{d}^D l_2 \frac{1}{D_1^{\nu_1} \dots D_8^{\nu_8}}, \quad (8.15)$$

where  $p$  and  $v$  are the heavy particle residual momentum and velocity,  $M$  is the mass of exchanged bosons, and

$$\begin{aligned} D_1 &= (p - l_1) \cdot v, & D_2 &= (p + l_2) \cdot v, & D_3(M) &= l_2^2 - M^2, & D_4(M) &= (l_1 + l_2)^2 - m^2, \\ D_5(M) &= l_1^2 - M^2, & D_6 &= (p + l_2 - l_1) \cdot v, & D_7 &= l_2^2, & D_8 &= (l_1 - l_2)^2. \end{aligned} \quad (8.16)$$



**Figure 7:** Heavy field self-energy topologies. The MIs associated to other topologies are subsets of the MIs required for topologies illustrated.

## 9 Discussion

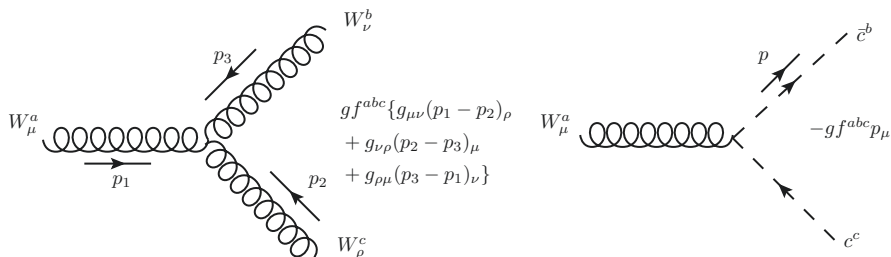
Both the massive and massless form factors are indispensable building blocks to a broad set of observables in both high and low energy regimes. Precisely studying these factors are crucial for shedding light on mysteries that remain in the standard model and beyond, such as the physical structure of the top quark, aspects of mass generation and the nature of dark matter. Our broad set of composite operators, choice of model for applicability, as well as consideration of two critical energetic regimes is emblematic of the breadth of the problem at hand. Our two-loop results are not complete, as we have not considered all possible regimes, for instance, the low-energy regime, nor have we calculated the vertex corrections at two-loop orders for SCET graphs. Continuing to map this space at two-loops and beyond is essential for our predicting power to be able to match the high precision potential of a future electron-positron collider and the LHC in its upcoming high luminosity operating phase.

At this point, the effective theory formalism is a central front of attack when it comes to tackling such complex problems by breaking them downscale by scale. By application to the SM, we have begun extending the work on EW corrections to high energy processes beyond NLO, as stated in the latest review [24]. Moreover, we are mapping other parts of the energetic landscape, aside from just the Sudakov regime, which itself opens the door for further investigation. Beyond the SM, the generality of the model and operators studied means that our results can be applied to BSM models by replacement of the proper coupling and group theory factors, which would be very interesting to examine further. For instance, one can apply our results to various models of dark matter [22, 57], where weak corrections are significant for indirect detection.

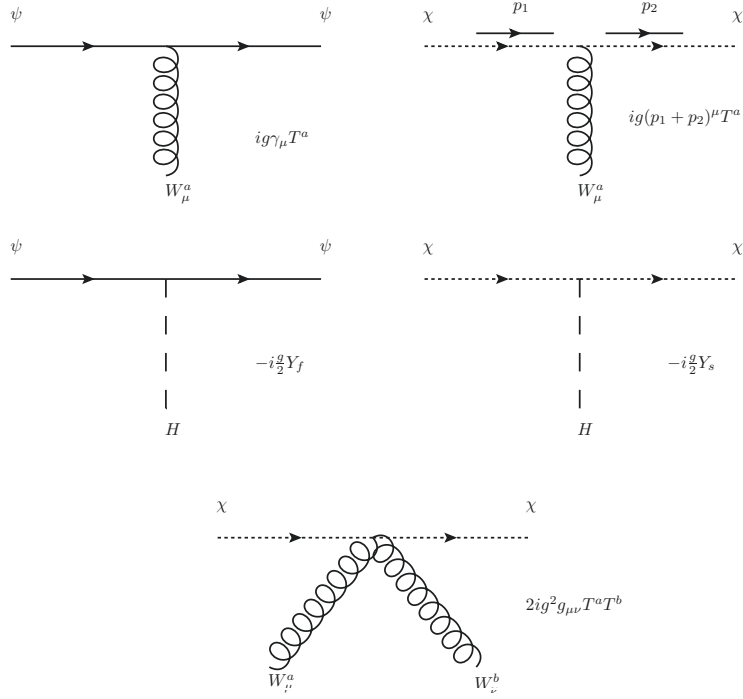
## A Feynman Rules

This appendix lists the Feynman rules of the vertices which are needed for the calculation of the form factor, as they follow from the Lagrangian of the spontaneously broken SU(2) gauge model described in section 2.1. The gauge boson fields of mass  $M = M_W$  are  $W_\mu^a : a = 1, 2, 3$  (with Lorentz vector index  $\mu$ ). To each  $W^a$  corresponds a ghost field  $c$  (and antighost  $\bar{c}$ ) and a Goldstone boson  $\phi^a$ , one of the unphysical components of the Higgs field. In the Feynman-t'Hooft gauge used by us, one sets  $M_\phi = M_W$ . The fields,  $\psi$  and  $\chi$  denote fermions and complex scalars, respectively. Lastly,  $g$  is the weak SU(2) coupling, and the labelling,  $\{1, 2\}$  differentiates between the particles on the grounds of mass, if two of the same kind exist in a vertex. Vertices which apply beyond two-loops are omitted here but should be included if one wants to venture beyond.

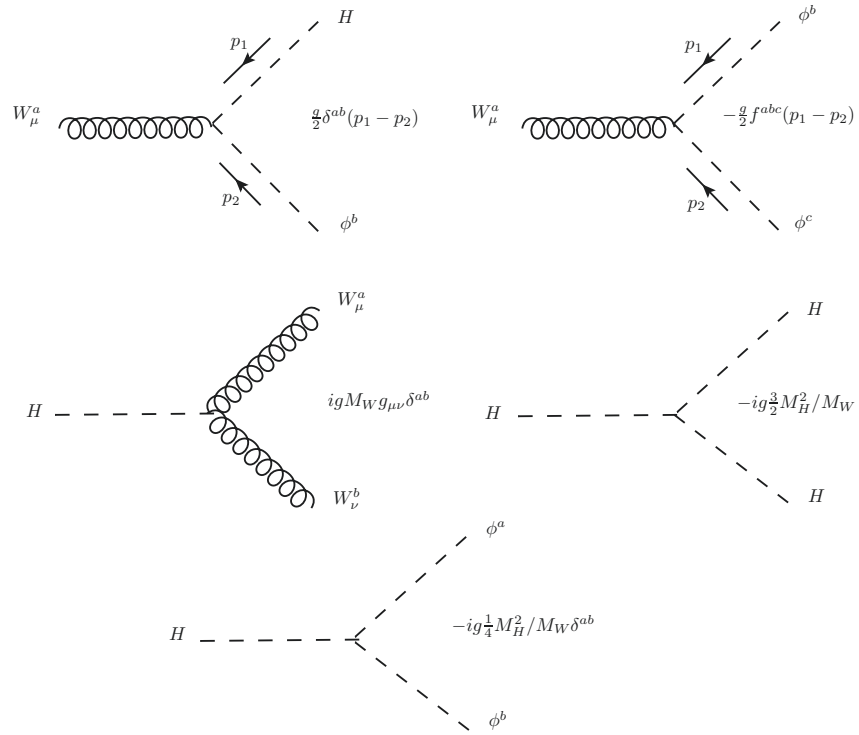
### A.1 Gauge boson self and Ghost couplings



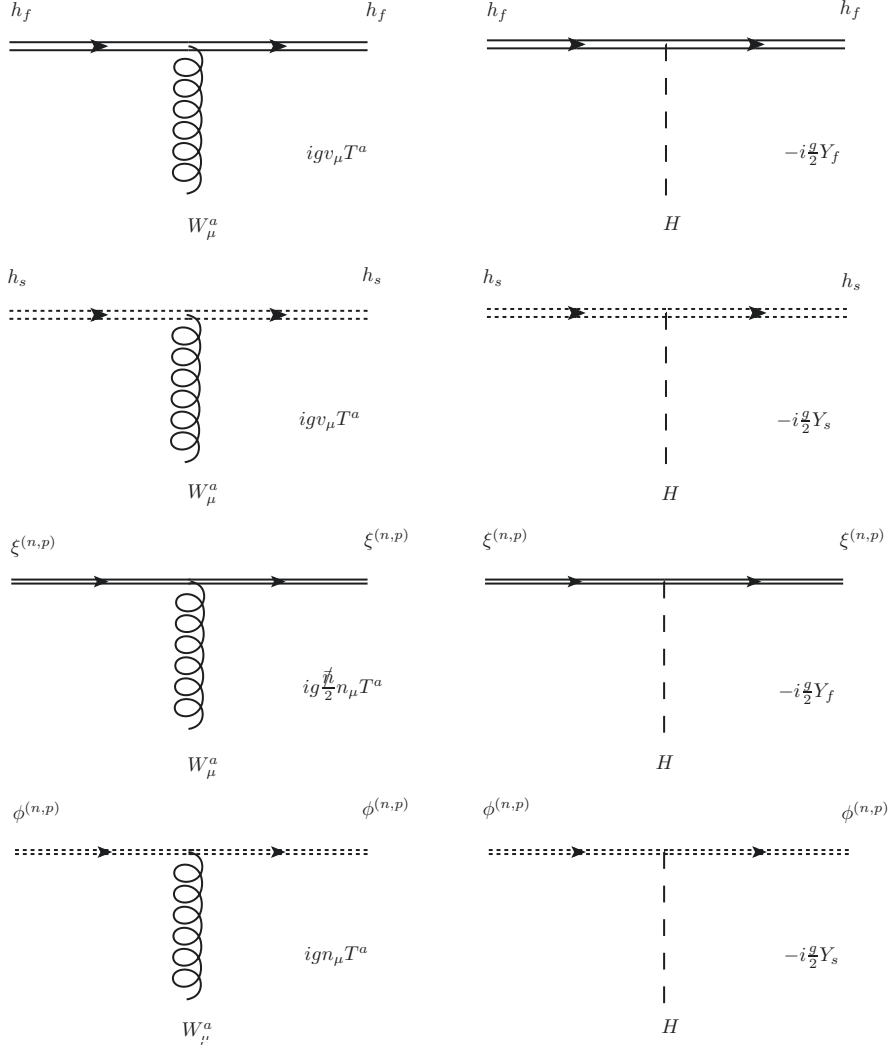
## A.2 Bosonic couplings to fermions/scalars



## A.3 Gauge boson coupling to Higgs and Goldstone bosons



## A.4 Effective theory couplings



In the effective theory vertices, solid (dashed) lines correspond to fermions (scalars) and widely (thinly) spaced lines correspond to heavy (co-linear) particles. As for the co-linear vertices we do not distinguish between soft/Wilson line couplings as they are identical to the order we are considering.

## B Mass Renormalisation

### B.1 Matching at $\mu \sim m_{1,2}$ :

$$\begin{aligned}
\Delta B_1^{(2)} = & -\frac{3C_A m_+^2 \mathcal{L}_{m_1}^3 C_F^3}{2m_-^2} + 3C_A \mathcal{L}_{m_1}^3 C_F^3 - \frac{3C_A m_-^2 \mathcal{L}_{m_1}^3 C_F^3}{2m_+^2} - \frac{3C_A m_+^2 \mathcal{L}_{m_2}^3 C_F^3}{2m_-^2} + 3C_A \mathcal{L}_{m_2}^3 C_F^3 - \frac{3C_A m_-^2 \mathcal{L}_{m_2}^3 C_F^3}{2m_+^2} + \\
& \frac{C_A m_+^2 \mathcal{L}_{m_1}^2 C_F^3}{2m_-^2} + \frac{3C_A m_+^2 \mathcal{L}_{m_1}^2 C_F^3}{2\epsilon m_-^2} - 4C_A \mathcal{L}_{m_1}^2 C_F^3 + \frac{12C_A m_+ \mathcal{L}_{m_1}^2 C_F^3}{m_-} - \frac{3C_A \mathcal{L}_{m_1}^2 C_F^3}{\epsilon} + \frac{12C_A m_- \mathcal{L}_{m_1}^2 C_F^3}{m_+} + \\
& \frac{7C_A m_-^2 \mathcal{L}_{m_1}^2 C_F^3}{2m_+^2} + \frac{3C_A m_-^2 \mathcal{L}_{m_1}^2 C_F^3}{2\epsilon m_+^2} + \frac{C_A m_+^2 \mathcal{L}_{m_2}^2 C_F^3}{2m_-^2} + \frac{3C_A m_+^2 \mathcal{L}_{m_2}^2 C_F^3}{2\epsilon m_-^2} - 4C_A \mathcal{L}_{m_2}^2 C_F^3 - \frac{12C_A m_+ \mathcal{L}_{m_2}^2 C_F^3}{m_-} + \\
& \frac{3C_A m_+^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 C_F^3}{2m_-^2} - 3C_A \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 C_F^3 + \frac{3C_A m_-^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 C_F^3}{2m_+^2} - \frac{3C_A \mathcal{L}_{m_2}^2 C_F^3}{\epsilon} - \frac{12C_A m_- \mathcal{L}_{m_2}^2 C_F^3}{m_+} + \frac{7C_A m_-^2 \mathcal{L}_{m_2}^2 C_F^3}{2m_+^2} +
\end{aligned}$$



$$\begin{aligned}
& \frac{3C_A m_-^2 \mathcal{L}_{m_2}^2 C_F^3}{2\epsilon m_+^2} - 16C_A C_F^3 + 12C_A \mathcal{L}_{m_1} C_F^3 - \frac{4C_A m_+ \mathcal{L}_{m_1} C_F^3}{m_-} - \frac{12C_A m_+ \mathcal{L}_{m_1} C_F^3}{\epsilon m_-} - \frac{28C_A m_- \mathcal{L}_{m_1} C_F^3}{m_+} \\
& \frac{12C_A m_- \mathcal{L}_{m_1} C_F^3}{\epsilon m_+} + \frac{3C_A m_+^2 \mathcal{L}_{m_1}^2 \mathcal{L}_{m_2} C_F^3}{2m_-^2} - 3C_A \mathcal{L}_{m_1}^2 \mathcal{L}_{m_2} C_F^3 + \frac{3C_A m_-^2 \mathcal{L}_{m_1}^2 \mathcal{L}_{m_2} C_F^3}{2m_+^2} + 12C_A \mathcal{L}_{m_2} C_F^3 + \frac{4C_A m_+ \mathcal{L}_{m_2} C_F^3}{m_-} \\
& \frac{12C_A m_+ \mathcal{L}_{m_2} C_F^3}{\epsilon m_-} - \frac{C_A m_+^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{m_-^2} - \frac{3C_A m_+^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{\epsilon m_-^2} + 8C_A \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3 + \frac{6C_A \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{\epsilon} \\
& \frac{7C_A m_-^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{m_+^2} - \frac{3C_A m_-^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{\epsilon m_+^2} + \frac{28C_A m_- \mathcal{L}_{m_2} C_F^3}{m_+} + \frac{12C_A m_- \mathcal{L}_{m_2} C_F^3}{\epsilon m_+} - \frac{12C_A C_F^3}{\epsilon} - \frac{3m_+^2 \mathcal{L}_{m_1}^3 C_F^2}{4m_-^2} \\
& \frac{3m_-^2 \mathcal{L}_{m_1}^3 C_F^2}{4m_+^2} + \frac{3}{2} \mathcal{L}_{m_1}^3 C_F^2 - \frac{3m_+^2 \mathcal{L}_{m_2}^3 C_F^2}{4m_-^2} - \frac{3m_-^2 \mathcal{L}_{m_2}^3 C_F^2}{4m_+^2} + \frac{3}{2} \mathcal{L}_{m_2}^3 C_F^2 + \frac{3m_+^2 \mathcal{L}_{m_1}^2 C_F^2}{4\epsilon m_-^2} + \frac{m_+^2 \mathcal{L}_{m_1}^2 C_F^2}{4m_-^2} + \frac{6m_+ \mathcal{L}_{m_1}^2 C_F^2}{m_-} \\
& \frac{3\mathcal{L}_{m_1}^2 C_F^2}{2\epsilon} + \frac{6m_- \mathcal{L}_{m_1}^2 C_F^2}{m_+} + \frac{3m_-^2 \mathcal{L}_{m_1}^2 C_F^2}{4\epsilon m_+^2} + \frac{7m_-^2 \mathcal{L}_{m_1}^2 C_F^2}{4m_+^2} - 2\mathcal{L}_{m_1}^2 C_F^2 + \frac{3m_+^2 \mathcal{L}_{m_2}^2 C_F^2}{4\epsilon m_-^2} + \frac{m_+^2 \mathcal{L}_{m_2}^2 C_F^2}{4m_-^2} - \frac{6m_+ \mathcal{L}_{m_2}^2 C_F^2}{m_-} \\
& \frac{3m_+^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2}{4m_-^2} + \frac{3m_-^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2}{4m_+^2} - \frac{3}{2} \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2 - \frac{3\mathcal{L}_{m_2}^2 C_F^2}{2\epsilon} - \frac{6m_- \mathcal{L}_{m_2}^2 C_F^2}{m_+} + \frac{3m_-^2 \mathcal{L}_{m_2}^2 C_F^2}{4\epsilon m_+^2} \\
& \frac{7m_-^2 \mathcal{L}_{m_2}^2 C_F^2}{4m_+^2} - 2\mathcal{L}_{m_2}^2 C_F^2 - \frac{6m_+ \mathcal{L}_{m_1} C_F^2}{\epsilon m_-} - \frac{2m_+ \mathcal{L}_{m_1} C_F^2}{m_-} - \frac{6m_- \mathcal{L}_{m_1} C_F^2}{\epsilon m_+} - \frac{14m_- \mathcal{L}_{m_1} C_F^2}{m_+} + 6\mathcal{L}_{m_1} C_F^2 \\
& \frac{3m_+^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2}{4m_-^2} + \frac{3m_-^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2}{4m_+^2} - \frac{3}{2} \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2 + \frac{6m_+ \mathcal{L}_{m_2} C_F^2}{\epsilon m_-} + \frac{2m_+ \mathcal{L}_{m_2} C_F^2}{m_-} - \frac{3m_+^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2}{2\epsilon m_-^2} \\
& \frac{m_+^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2}{2m_-^2} + \frac{3\mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2}{\epsilon} - \frac{3m_-^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2}{2\epsilon m_+^2} - \frac{7m_-^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2}{2m_+^2} + 4\mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2 + \frac{6m_- \mathcal{L}_{m_2} C_F^2}{\epsilon m_+} \\
& \frac{14m_- \mathcal{L}_{m_2} C_F^2}{m_+} + 6\mathcal{L}_{m_2} C_F^2 - \frac{6C_F^2}{\epsilon} - 8C_F^2 - \frac{C_A m_+ \beta_0 \mathcal{L}_{m_1}^3 C_F}{12m_-} - \frac{C_A m_- \beta_0 \mathcal{L}_{m_1}^3 C_F}{12m_+} + \frac{C_A m_+ \beta_0 \mathcal{L}_{m_1}^3 C_F}{12m_-} \\
& \frac{C_A m_- \beta_0 \mathcal{L}_{m_1}^3 C_F}{12m_+} - \frac{1}{4} C_A \beta_0 \mathcal{L}_{m_1}^2 C_F - \frac{C_A m_+ \beta_0 \mathcal{L}_{m_1}^2 C_F}{4m_-} + \frac{C_A m_+ \beta_0 \mathcal{L}_{m_1}^2 C_F}{4\epsilon m_-} + \frac{C_A m_- \beta_0 \mathcal{L}_{m_1}^2 C_F}{4m_+} + \frac{C_A m_- \beta_0 \mathcal{L}_{m_1}^2 C_F}{4\epsilon m_+} \\
& \frac{1}{4} C_A \beta_0 \mathcal{L}_{m_2}^2 C_F + \frac{C_A m_+ \beta_0 \mathcal{L}_{m_2}^2 C_F}{4m_-} - \frac{C_A m_+ \beta_0 \mathcal{L}_{m_2}^2 C_F}{4\epsilon m_-} - \frac{C_A m_- \beta_0 \mathcal{L}_{m_2}^2 C_F}{4m_+} - \frac{C_A m_- \beta_0 \mathcal{L}_{m_2}^2 C_F}{4\epsilon m_+} - \frac{C_A \beta_0 C_F}{\epsilon^2} \\
& \frac{1}{12} C_A \pi^2 \beta_0 C_F + \frac{C_A m_+ \beta_0 \mathcal{L}_{m_1} C_F}{2\epsilon m_-} - \frac{C_A m_+ \beta_0 \mathcal{L}_{m_1} C_F}{2\epsilon^2 m_-} + \frac{C_A \beta_0 \mathcal{L}_{m_1} C_F}{2\epsilon} - \frac{2C_A m_- \beta_0 \mathcal{L}_{m_1} C_F}{m_+} - \frac{C_A m_- \beta_0 \mathcal{L}_{m_1} C_F}{2\epsilon m_+} \\
& \frac{C_A m_- \beta_0 \mathcal{L}_{m_1} C_F}{2\epsilon^2 m_+} - \frac{C_A m_+ \pi^2 \beta_0 \mathcal{L}_{m_1} C_F}{24m_-} - \frac{C_A m_- \pi^2 \beta_0 \mathcal{L}_{m_1} C_F}{24m_+} - \frac{C_A m_+ \beta_0 \mathcal{L}_{m_2} C_F}{2\epsilon m_-} + \frac{C_A m_+ \beta_0 \mathcal{L}_{m_2} C_F}{2\epsilon^2 m_-} \\
& \frac{C_A \beta_0 \mathcal{L}_{m_2} C_F}{2\epsilon} + \frac{2C_A m_- \beta_0 \mathcal{L}_{m_2} C_F}{m_+} + \frac{C_A m_- \beta_0 \mathcal{L}_{m_2} C_F}{2\epsilon m_+} + \frac{C_A m_- \beta_0 \mathcal{L}_{m_2} C_F}{2\epsilon^2 m_+} + \frac{C_A m_+ \pi^2 \beta_0 \mathcal{L}_{m_2} C_F}{24m_-} \\
& \frac{C_A m_- \pi^2 \beta_0 \mathcal{L}_{m_2} C_F}{24m_+} - \frac{3C_A C_F^2 Y_f^2 \mathcal{L}_{m_1}^4}{m_-^2 m_+^2} - \frac{3C_F Y_f^2 \mathcal{L}_{m_1}^4}{2m_-^2 m_+^2} + \frac{3C_A C_F^2 Y_f^2 \mathcal{L}_{m_2}^4}{m_-^2 m_+^2} + \frac{3C_F Y_f^2 \mathcal{L}_{m_2}^4}{2m_-^2 m_+^2} \\
& \frac{9C_A Y_f^2 \beta_0 \mathcal{L}_{m_2}^2 m_1^4}{32m_-^2 m_+^2} + \frac{3C_A C_F^2 m_2 Y_f^2 \mathcal{L}_{m_1}^3 m_1^3}{m_-^2 m_+^2} + \frac{3C_F m_2 Y_f^2 \mathcal{L}_{m_1}^3 m_1^3}{2m_-^2 m_+^2} + \frac{C_A m_2 Y_f^2 \beta_0 \mathcal{L}_{m_1}^3 m_1^3}{12m_-^2 m_+^2} + \frac{3C_A C_F^2 m_2 Y_f^2 \mathcal{L}_{m_2}^3 m_1^3}{m_-^2 m_+^2} \\
& \frac{3C_F m_2 Y_f^2 \mathcal{L}_{m_2}^3 m_1^3}{2m_-^2 m_+^2} - \frac{C_A m_2 Y_f^2 \beta_0 \mathcal{L}_{m_2}^3 m_1^3}{12m_-^2 m_+^2} - \frac{35C_A C_F^2 m_2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^3}{2m_-^2 m_+^2} - \frac{35C_F m_2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^3}{4m_-^2 m_+^2} - \frac{3C_A C_F^2 m_2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^3}{\epsilon m_-^2 m_+^2} \\
& \frac{3C_F m_2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^3}{2\epsilon m_-^2 m_+^2} + \frac{13C_A C_F^2 m_2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^3}{2m_-^2 m_+^2} + \frac{13C_F m_2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^3}{4m_-^2 m_+^2} - \frac{3C_A C_F^2 m_2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^3}{\epsilon m_-^2 m_+^2} - \frac{3C_F m_2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^3}{2\epsilon m_-^2 m_+^2} \\
& \frac{C_A m_2 Y_f^2 \beta_0 \mathcal{L}_{m_2}^2 m_1^3}{4m_-^2 m_+^2} + \frac{C_A m_2 Y_f^2 \beta_0 \mathcal{L}_{m_2}^2 m_1^3}{4\epsilon m_-^2 m_+^2} - \frac{3C_A^2 C_F m_2 Y_f^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 m_1^3}{2m_-^2 m_+^2} - \frac{3C_A^2 C_F m_2 Y_f^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 m_1^3}{2m_-^2 m_+^2} \\
& \frac{11C_A^2 C_F m_2 Y_f^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} m_1^3}{2m_-^2 m_+^2} + \frac{3C_A^2 C_F m_2 Y_f^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} m_1^3}{\epsilon m_-^2 m_+^2} + \frac{3C_A C_F^2 m_2^2 Y_f^2 \mathcal{L}_{m_1}^3 m_1^2}{m_-^2 m_+^2} + \frac{3C_F m_2^2 Y_f^2 \mathcal{L}_{m_1}^3 m_1^2}{2m_-^2 m_+^2} \\
& \frac{3C_A C_F^2 m_2^2 Y_f^2 \mathcal{L}_{m_2}^3 m_1^2}{m_-^2 m_+^2} + \frac{3C_F m_2^2 Y_f^2 \mathcal{L}_{m_2}^3 m_1^2}{2m_-^2 m_+^2} - \frac{13C_A C_F^2 m_2^2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^2}{m_-^2 m_+^2} - \frac{13C_F m_2^2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^2}{2m_-^2 m_+^2} - \frac{3C_A C_F^2 m_2^2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^2}{\epsilon m_-^2 m_+^2} \\
& \frac{3C_F m_2^2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^2}{2\epsilon m_-^2 m_+^2} - \frac{9C_A Y_f^2 \beta_0 \mathcal{L}_{m_1}^2 m_1^2}{32m_- m_+} - \frac{13C_A C_F^2 m_2^2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^2}{m_-^2 m_+^2} - \frac{13C_F m_2^2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^2}{2m_-^2 m_+^2} - \frac{3C_A C_F^2 m_2^2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^2}{\epsilon m_-^2 m_+^2} \\
& \frac{3C_F m_2^2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^2}{2\epsilon m_-^2 m_+^2} + \frac{9C_A m_2^2 Y_f^2 \beta_0 \mathcal{L}_{m_2}^2 m_1^2}{16m_-^2 m_+^2} - \frac{3C_A^2 C_F m_2^2 Y_f^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 m_1^2}{2m_-^2 m_+^2} + \frac{31C_A C_F^2 Y_f^2 \mathcal{L}_{m_1} m_1^2}{m_- m_+} + \frac{31C_F Y_f^2 \mathcal{L}_{m_1} m_1^2}{2m_- m_+} \\
& \frac{3C_A C_F^2 Y_f^2 \mathcal{L}_{m_1} m_1^2}{\epsilon m_- m_+} + \frac{3C_F Y_f^2 \mathcal{L}_{m_1} m_1^2}{2\epsilon m_- m_+} + \frac{5C_A Y_f^2 \beta_0 \mathcal{L}_{m_1} m_1^2}{16m_- m_+} + \frac{9C_A Y_f^2 \beta_0 \mathcal{L}_{m_1} m_1^2}{16\epsilon m_- m_+} + \frac{5C_A C_F^2 Y_f^2 \mathcal{L}_{m_2} m_1^2}{m_- m_+} + \frac{5C_F Y_f^2 \mathcal{L}_{m_2} m_1^2}{2m_- m_+} \\
& \frac{3C_A C_F^2 Y_f^2 \mathcal{L}_{m_2} m_1^2}{\epsilon m_- m_+} - \frac{3C_F Y_f^2 \mathcal{L}_{m_2} m_1^2}{2\epsilon m_- m_+} - \frac{3C_A^2 C_F m_2^2 Y_f^2 \mathcal{L}_{m_1}^2 \mathcal{L}_{m_2} m_1^2}{2m_-^2 m_+^2} + \frac{9C_A Y_f^2 \beta_0 \mathcal{L}_{m_2} m_1^2}{16m_- m_+} + \frac{9C_A Y_f^2 \beta_0 \mathcal{L}_{m_2} m_1^2}{16\epsilon m_- m_+} \\
& \frac{13C_A^2 C_F m_2^2 Y_f^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} m_1^2}{m_-^2 m_+^2} + \frac{3C_A^2 C_F m_2^2 Y_f^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} m_1^2}{\epsilon m_-^2 m_+^2} + \frac{3C_A C_F^2 m_2^3 Y_f^2 \mathcal{L}_{m_1}^3 m_1}{m_-^2 m_+^2} + \frac{3C_F m_2^3 Y_f^2 \mathcal{L}_{m_1}^3 m_1}{2m_-^2 m_+^2}
\end{aligned}$$

$$\begin{aligned}
& \frac{CA m_2^3 Y_f^2 \beta_0 C_{m_1}^3}{12 m_+^2 m_+^2} + \frac{3CA C_F^2 m_2^3 Y_f^2 L_{m_2}^3}{m_+^2 m_+^2} + \frac{3CF m_2^3 Y_f^2 L_{m_2}^3}{2 m_+^2 m_+^2} + \frac{CA m_2^3 Y_f^2 \beta_0 L_{m_2}^3}{12 m_+^2 m_+^2} + \frac{13CA C_F^2 m_2^3 Y_f^2 L_{m_1}^2}{2 m_+^2 m_+^2} + \\
& \frac{13CF m_2^3 Y_f^2 L_{m_1}^2}{4 m_+^2 m_+^2} - \frac{3CA C_F^2 m_2^3 Y_f^2 L_{m_1}^2}{\epsilon m_+^2 m_+^2} - \frac{3CF m_2^3 Y_f^2 L_{m_1}^2}{2 \epsilon m_+^2 m_+^2} - \frac{CA m_2 Y_f^2 \beta_0 L_{m_1}^2}{4 m_+ m_+} - \frac{CA m_2 Y_f^2 \beta_0 L_{m_1}^2}{4 \epsilon m_+ m_+} - \\
& \frac{35CA C_F^2 m_2^3 Y_f^2 L_{m_2}^2}{2 m_+^2 m_+^2} - \frac{35CF m_2^3 Y_f^2 L_{m_2}^2}{4 m_+^2 m_+^2} - \frac{3CA C_F^2 m_2^3 Y_f^2 L_{m_2}^2}{\epsilon m_+^2 m_+^2} - \frac{3CF m_2^3 Y_f^2 L_{m_2}^2}{2 \epsilon m_+^2 m_+^2} - \frac{CA m_2^3 Y_f^2 \beta_0 L_{m_2}^2}{4 m_+^2 m_+^2} - \\
& \frac{CA m_2^3 Y_f^2 \beta_0 L_{m_2}^2}{4 \epsilon m_+ m_+} - \frac{3CA C_F m_2^3 Y_f^2 L_{m_1} L_{m_2}^2}{2 m_+ m_+} + \frac{22CA C_F^2 m_2 Y_f^2 L_{m_1} m_1}{m_+ m_+} + \frac{11CF m_2 Y_f^2 L_{m_1} m_1}{m_+ m_+} + \\
& \frac{12CA C_F^2 m_2 Y_f^2 L_{m_1} m_1}{\epsilon m_+ m_+} + \frac{6CF m_2 Y_f^2 L_{m_1} m_1}{\epsilon m_+ m_+} + \frac{CA m_2 Y_f^2 \beta_0 L_{m_1} m_1}{m_+ m_+} + \frac{CA m_2 Y_f^2 \beta_0 L_{m_1} m_1}{2 \epsilon m_+ m_+} + \frac{CA m_2 Y_f^2 \beta_0 L_{m_1} m_1}{2 \epsilon^2 m_+ m_+} + \\
& \frac{CA m_2 \pi^2 Y_f^2 \beta_0 L_{m_1} m_1}{24 m_+ m_+} - \frac{22CA C_F^2 m_2 Y_f^2 L_{m_2} m_1}{m_+ m_+} - \frac{11CF m_2 Y_f^2 L_{m_2} m_1}{m_+ m_+} - \frac{12CA C_F^2 m_2 Y_f^2 L_{m_2} m_1}{\epsilon m_+ m_+} - \frac{6CF m_2 Y_f^2 L_{m_2} m_1}{\epsilon m_+ m_+} - \\
& \frac{3CA C_F m_2^3 Y_f^2 L_{m_1} L_{m_2} m_1}{2 m_+^2 m_+^2} - \frac{CA m_2 Y_f^2 \beta_0 L_{m_2} m_1}{m_+ m_+} - \frac{CA m_2 Y_f^2 \beta_0 L_{m_2} m_1}{2 \epsilon m_+ m_+} - \frac{CA m_2 Y_f^2 \beta_0 L_{m_2} m_1}{2 \epsilon^2 m_+ m_+} - \frac{CA m_2 \pi^2 Y_f^2 \beta_0 L_{m_2} m_1}{24 m_+ m_+} + \\
& \frac{11CA C_F m_2^3 Y_f^2 L_{m_1} L_{m_2} m_1}{2 m_+^2 m_+^2} + \frac{3CA C_F m_2^3 Y_f^2 L_{m_1} L_{m_2} m_1}{\epsilon m_+^2 m_+^2} - 42CA C_F^2 Y_f^2 - 21CF Y_f^2 - \frac{18CA C_F^2 Y_f^2}{\epsilon} - \frac{9CF Y_f^2}{\epsilon} + \\
& \frac{3CA C_F^2 m_2^4 Y_f^2 L_{m_1}^2}{m_+^2 m_+^2} + \frac{3CF m_2^4 Y_f^2 L_{m_1}^2}{2 m_+^2 m_+^2} + \frac{9CA m_2^2 Y_f^2 \beta_0 L_{m_1}^2}{32 m_+ m_+} - \frac{3CA C_F^2 m_2^4 Y_f^2 L_{m_2}^2}{m_+^2 m_+^2} - \frac{3CF m_2^4 Y_f^2 L_{m_2}^2}{2 m_+^2 m_+^2} - \frac{9CA m_2^4 Y_f^2 \beta_0 L_{m_2}^2}{32 m_+^2 m_+^2} - \\
& \frac{9}{8} CA Y_f^2 \beta_0 - \frac{7CA Y_f^2 \beta_0}{8 \epsilon} - \frac{9CA Y_f^2 \beta_0}{8 \epsilon^2} - \frac{3}{32} CA \pi^2 Y_f^2 \beta_0 - \frac{5CA C_F^2 m_2^3 Y_f^2 L_{m_1}}{m_+ m_+} - \frac{5CF m_2^3 Y_f^2 L_{m_1}}{2 m_+ m_+} + \frac{3CA C_F^2 m_2^3 Y_f^2 L_{m_1}}{\epsilon m_+ m_+} + \\
& \frac{3CF m_2^3 Y_f^2 L_{m_1}}{2 \epsilon m_+ m_+} - \frac{9CA m_2^3 Y_f^2 \beta_0 L_{m_1}}{16 m_+ m_+} - \frac{9CA m_2^3 Y_f^2 \beta_0 L_{m_1}}{16 \epsilon m_+ m_+} - \frac{31CA C_F^2 m_2^3 Y_f^2 L_{m_2}}{m_+ m_+} - \frac{31CF m_2^3 Y_f^2 L_{m_2}}{2 m_+ m_+} - \frac{3CA C_F^2 m_2^3 Y_f^2 L_{m_2}}{\epsilon m_+ m_+} - \\
& \frac{3CF m_2^3 Y_f^2 L_{m_2}}{2 \epsilon m_+ m_+} - \frac{5CA m_2^3 Y_f^2 \beta_0 L_{m_2}}{16 m_+ m_+} - \frac{9CA m_2^3 Y_f^2 \beta_0 L_{m_2}}{16 \epsilon m_+ m_+} + \frac{3m_1^3 m_2 Y_f^4 C_A^2 L_{m_1}^2}{16 \epsilon m_+^2 m_+^2} + \frac{3m_1^3 m_2 Y_f^4 C_A^2 L_{m_2}^2}{16 \epsilon m_+^2 m_+^2} - \\
& \frac{3m_1^3 m_2 Y_f^4 C_A^2 L_{m_1} L_{m_2}}{8 \epsilon m_+^2 m_+^2} - \frac{3m_1^3 m_2 Y_f^4 C_A^2 L_{m_1}^3}{16 m_+^2 m_+^2} - \frac{3m_1^3 m_2 Y_f^4 C_A^2 L_{m_2}^3}{16 m_+^2 m_+^2} + \frac{11m_1^3 m_2 Y_f^4 C_A^2 L_{m_1}^2}{8 m_+^2 m_+^2} + \frac{3m_1^3 m_2 Y_f^4 C_A^2 L_{m_1} L_{m_2}^2}{16 m_+^2 m_+^2} - \\
& \frac{m_1^3 m_2 Y_f^4 C_A^2 L_{m_2}^2}{8 m_+^2 m_+^2} + \frac{3m_1^3 m_2 Y_f^4 C_A^2 L_{m_1} L_{m_2}}{16 m_+^2 m_+^2} - \frac{5m_1^3 m_2 Y_f^4 C_A^2 L_{m_1} L_{m_2}}{4 m_+^2 m_+^2} - \frac{3m_1^3 m_2 Y_f^4 C_A^2 L_{m_1}^2}{16 m_+^2 m_+^2} - \frac{3m_1^3 m_2 Y_f^4 C_A^2 L_{m_2}^2}{16 m_+^2 m_+^2} + \\
& \frac{3m_1^2 m_2^2 Y_f^4 C_A^2 L_{m_1} L_{m_2}}{8 m_+^2 m_+^2} - \frac{27m_1^2 Y_f^4 C_A^2 L_{m_2}}{32 m_+ m_+} - \frac{15m_1^2 Y_f^4 C_A^2 L_{m_1}}{32 m_+ m_+} + \frac{3m_1 m_2^3 Y_f^4 C_A^2 L_{m_1}^2}{16 \epsilon m_+^2 m_+^2} + \frac{3m_1 m_2^3 Y_f^4 C_A^2 L_{m_2}^2}{16 \epsilon m_+^2 m_+^2} - \\
& \frac{3m_1 m_2^3 Y_f^4 C_A^2 L_{m_1} L_{m_2}}{8 \epsilon m_+^2 m_+^2} - \frac{3m_1 m_2^3 Y_f^4 C_A^2 L_{m_1}^3}{16 m_+^2 m_+^2} - \frac{3m_1 m_2^3 Y_f^4 C_A^2 L_{m_2}^3}{16 m_+^2 m_+^2} - \frac{m_1 m_2^3 Y_f^4 C_A^2 L_{m_1}^2}{8 m_+^2 m_+^2} + \frac{3m_1 m_2^3 Y_f^4 C_A^2 L_{m_1} L_{m_2}^2}{16 m_+^2 m_+^2} + \\
& \frac{11m_1 m_2^3 Y_f^4 C_A^2 L_{m_2}^2}{8 m_+^2 m_+^2} + \frac{3m_1 m_2^3 Y_f^4 C_A^2 L_{m_1} L_{m_2}}{16 m_+^2 m_+^2} - \frac{5m_1 m_2^3 Y_f^4 C_A^2 L_{m_1} L_{m_2}}{4 m_+^2 m_+^2} + \frac{27m_1^2 Y_f^4 C_A^2 L_{m_1}}{32 m_+ m_+} - \frac{3m_1 m_2 Y_f^4 C_A^2 L_{m_1}}{4 \epsilon m_+ m_+} + \\
& \frac{3m_1 m_2 Y_f^4 C_A^2 L_{m_2}}{4 \epsilon m_+ m_+} - \frac{5m_1 m_2 Y_f^4 C_A^2 L_{m_1}}{2 m_+ m_+} + \frac{5m_1 m_2 Y_f^4 C_A^2 L_{m_2}}{2 m_+ m_+} + \frac{15m_2^2 Y_f^4 C_A^2 L_{m_2}}{32 m_+ m_+} + \frac{21Y_f^4 C_A^2}{32 \epsilon} + \frac{29}{16} Y_f^4 C_A^2 \quad (B.1)
\end{aligned}$$

$$\begin{aligned}
\Delta B_2^{(2)} &= 3CA L_{m_1}^3 C_F^3 - \frac{3CA m_+^2 L_{m_1}^3 C_F^3}{m_+^2} + 3CA L_{m_2}^3 C_F^3 - \frac{3CA m_+^2 L_{m_2}^3 C_F^3}{m_+^2} + \frac{3CA m_+^2 L_{m_1}^2 C_F^3}{4 m_+^2} - \frac{17}{2} CA L_{m_1}^2 C_F^3 - \frac{3CA L_{m_1}^2 C_F^3}{\epsilon} + \\
& \frac{24CA m_+ L_{m_1}^2 C_F^3}{m_+} + \frac{31CA m_+^2 L_{m_1}^2 C_F^3}{4 m_+^2} + \frac{3CA m_+^2 L_{m_1}^2 C_F^3}{\epsilon m_+^2} + \frac{3CA m_+^2 L_{m_2}^2 C_F^3}{4 m_+^2} - \frac{17}{2} CA L_{m_2}^2 C_F^3 - 3CA L_{m_1} L_{m_2}^2 C_F^3 + \\
& \frac{3CA m_+^2 L_{m_1} L_{m_2}^2 C_F^3}{m_+^2} - \frac{3CA L_{m_2}^2 C_F^3}{\epsilon} - \frac{24CA m_+ L_{m_2}^2 C_F^3}{m_+} + \frac{31CA m_+^2 L_{m_2}^2 C_F^3}{4 m_+^2} + \\
& \frac{3CA m_+^2 L_{m_2}^2 C_F^3}{\epsilon m_+^2} + 68CA C_F^3 - 24CA L_{m_1} C_F^3 - \frac{6CA m_+ L_{m_1} C_F^3}{m_+} - \frac{62CA m_+ L_{m_1} C_F^3}{m_+} - \\
& \frac{24CA m_+ L_{m_1} C_F^3}{\epsilon m_+} - 3CA L_{m_1}^2 L_{m_2} C_F^3 + \frac{3CA m_+^2 L_{m_1}^2 L_{m_2} C_F^3}{m_+^2} - 24CA L_{m_2} C_F^3 + \frac{6CA m_+ L_{m_2} C_F^3}{m_+} - \\
& \frac{3CA m_+^2 L_{m_1} L_{m_2} C_F^3}{2 m_+^2} + 17CA L_{m_1} L_{m_2} C_F^3 + \frac{6CA L_{m_1} L_{m_2} C_F^3}{\epsilon} - \frac{31CA m_+^2 L_{m_1} L_{m_2} C_F^3}{2 m_+^2} - \frac{6CA m_+^2 L_{m_1} L_{m_2} C_F^3}{\epsilon m_+^2} + \\
& \frac{62CA m_+ L_{m_2} C_F^3}{m_+} + \frac{24CA m_+ L_{m_2} C_F^3}{\epsilon m_+} + \frac{24CA C_F^3}{\epsilon} - \frac{3m_+^2 L_{m_1}^3 C_F^2}{2 m_+^2} + \frac{3}{2} L_{m_1}^3 C_F^2 - \frac{3m_+^2 L_{m_2}^3 C_F^2}{2 m_+^2} + \frac{3}{2} L_{m_2}^3 C_F^2 + \\
& \frac{3m_+^2 L_{m_1}^2 C_F^2}{8 m_+^2} - \frac{3L_{m_1}^2 C_F^2}{2 \epsilon} + \frac{12m_+ L_{m_1}^2 C_F^2}{m_+} + \frac{3m_+^2 L_{m_1}^2 C_F^2}{2 \epsilon m_+^2} + \frac{31m_+^2 L_{m_1}^2 C_F^2}{8 m_+^2} - \frac{17}{4} L_{m_1}^2 C_F^2 + \frac{3m_+^2 L_{m_2}^2 C_F^2}{8 m_+^2} + \\
& \frac{3m_+^2 L_{m_1} L_{m_2} C_F^2}{2 m_+^2} - \frac{3}{2} L_{m_1} L_{m_2} C_F^2 - \frac{3L_{m_2}^2 C_F^2}{2 \epsilon} - \frac{12m_+ L_{m_2}^2 C_F^2}{m_+} + \frac{3m_+^2 L_{m_2}^2 C_F^2}{2 \epsilon m_+^2} + \frac{31m_+^2 L_{m_2}^2 C_F^2}{8 m_+^2} - \frac{17}{4} L_{m_2}^2 C_F^2 -
\end{aligned}$$

$$\begin{aligned}
& \frac{3m_+ \mathcal{L}_{m_1} C_F^2}{m_-} - \frac{12m_- \mathcal{L}_{m_1} C_F^2}{\varepsilon m_+} - \frac{31m_- \mathcal{L}_{m_1} C_F^2}{m_+} - 12\mathcal{L}_{m_1} C_F^2 + \frac{3m_-^2 \mathcal{L}_{m_1}^2 \mathcal{L}_{m_2} C_F^2}{2m_+^2} - \frac{3}{2} \mathcal{L}_{m_1}^2 \mathcal{L}_{m_2} C_F^2 + \frac{3m_+ \mathcal{L}_{m_2} C_F^2}{m_-} - \\
& \frac{3m_+^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2}{4m_-^2} + \frac{3\mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2}{\varepsilon} - \frac{3m_-^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2}{\varepsilon m_+^2} - \frac{31m_-^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2}{4m_+^2} + \frac{17}{2} \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^2 + \frac{12m_- \mathcal{L}_{m_2} C_F^2}{\varepsilon m_+} + \\
& \frac{31m_- \mathcal{L}_{m_2} C_F^2}{m_+} - 12\mathcal{L}_{m_2} C_F^2 + \frac{12C_F^2}{\varepsilon} + 34C_F^2 - \frac{C_A m_- \beta_0 \mathcal{L}_{m_1}^3 C_F}{6m_+} + \frac{C_A m_- \beta_0 \mathcal{L}_{m_2}^3 C_F}{6m_+} + \frac{1}{2} C_A \beta_0 \mathcal{L}_{m_1}^2 C_F + \\
& \frac{C_A m_+ \beta_0 \mathcal{L}_{m_1}^2 C_F}{8m_-} + \frac{5C_A m_- \beta_0 \mathcal{L}_{m_1}^2 C_F}{8m_+} + \frac{C_A m_- \beta_0 \mathcal{L}_{m_1}^2 C_F}{2\varepsilon m_+} + \frac{1}{2} C_A \beta_0 \mathcal{L}_{m_2}^2 C_F - \frac{C_A m_+ \beta_0 \mathcal{L}_{m_2}^2 C_F}{8m_-} - \frac{5C_A m_- \beta_0 \mathcal{L}_{m_2}^2 C_F}{8m_+} - \\
& \frac{C_A m_- \beta_0 \mathcal{L}_{m_2}^2 C_F}{2\varepsilon m_+} + 5C_A \beta_0 C_F + \frac{3C_A \beta_0 C_F}{\varepsilon} + \frac{2C_A m_- \beta_0 C_F}{m_+} + \frac{2C_A \beta_0 C_F}{\varepsilon^2} + \frac{1}{6} C_A \pi^2 \beta_0 C_F - C_A \beta_0 \mathcal{L}_{m_1} C_F - \\
& \frac{C_A m_+ \beta_0 \mathcal{L}_{m_1} C_F}{4m_-} - \frac{C_A m_+ \beta_0 \mathcal{L}_{m_1} C_F}{4\varepsilon m_-} - \frac{C_A \beta_0 \mathcal{L}_{m_1} C_F}{\varepsilon} - \frac{13C_A m_- \beta_0 \mathcal{L}_{m_1} C_F}{4m_+} - \frac{5C_A m_- \beta_0 \mathcal{L}_{m_1} C_F}{4\varepsilon m_+} - \frac{C_A m_- \beta_0 \mathcal{L}_{m_1} C_F}{\varepsilon^2 m_+} - \\
& \frac{C_A m_-^2 \beta_0 \mathcal{L}_{m_1} C_F}{2m_+^2} - \frac{C_A m_- \pi^2 \beta_0 \mathcal{L}_{m_1} C_F}{12m_+} - 2C_A \beta_0 \mathcal{L}_{m_2} C_F + \frac{C_A m_+ \beta_0 \mathcal{L}_{m_2} C_F}{4m_-} + \frac{C_A m_+ \beta_0 \mathcal{L}_{m_2} C_F}{4\varepsilon m_-} - \frac{C_A \beta_0 \mathcal{L}_{m_2} C_F}{\varepsilon} + \\
& \frac{13C_A m_- \beta_0 \mathcal{L}_{m_2} C_F}{4m_+} + \frac{5C_A m_- \beta_0 \mathcal{L}_{m_2} C_F}{4\varepsilon m_+} + \frac{C_A m_- \beta_0 \mathcal{L}_{m_2} C_F}{\varepsilon^2 m_+} + \frac{C_A m_-^2 \beta_0 \mathcal{L}_{m_2} C_F}{2m_+^2} + \frac{C_A m_- \pi^2 \beta_0 \mathcal{L}_{m_2} C_F}{12m_+} + \\
& \frac{3C_A C_F^2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^4}{m_- m_+^3} + \frac{3C_F Y_f^2 \mathcal{L}_{m_2}^2 m_1^4}{2m_- m_+^3} - \frac{9C_A Y_f^2 \beta_0 \mathcal{L}_{m_2}^2 m_1^4}{32m_- m_+^3} - \frac{3C_A C_F^2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^3}{m_+^3} - \frac{3C_F Y_f^2 \mathcal{L}_{m_1}^2 m_1^3}{2m_+^3} - \frac{11C_A Y_f^2 \beta_0 \mathcal{L}_{m_2}^2 m_1^3}{32m_- m_+^2} + \\
& \frac{5C_A C_F^2 m_2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^3}{m_- m_+^3} + \frac{5C_F m_2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^3}{2m_- m_+^3} + \frac{3C_A C_F^2 m_2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^3}{2\varepsilon m_- m_+^3} + \frac{3C_F m_2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^3}{4\varepsilon m_- m_+^3} - \frac{9C_A m_2 Y_f^2 \beta_0 \mathcal{L}_{m_2}^2 m_1^3}{16m_- m_+^3} + \\
& \frac{31C_A C_F^2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^3}{m_- m_+^2} + \frac{31C_F Y_f^2 \mathcal{L}_{m_1}^2 m_1^3}{2m_- m_+^2} + \frac{3C_A C_F^2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^3}{\varepsilon m_- m_+^2} + \frac{3C_F Y_f^2 \mathcal{L}_{m_1}^2 m_1^3}{2\varepsilon m_- m_+^2} + \frac{9C_A Y_f^2 \beta_0 \mathcal{L}_{m_1}^2 m_1^3}{16m_- m_+^2} + \\
& \frac{11C_A Y_f^2 \beta_0 \mathcal{L}_{m_1}^2 m_1^3}{16\varepsilon m_- m_+^2} + \frac{5C_A C_F^2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^3}{m_- m_+^2} + \frac{5C_F Y_f^2 \mathcal{L}_{m_2}^2 m_1^3}{2m_- m_+^2} - \frac{3C_A C_F^2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^3}{\varepsilon m_- m_+^2} - \frac{3C_F Y_f^2 \mathcal{L}_{m_2}^2 m_1^3}{2\varepsilon m_- m_+^2} + \frac{9C_A Y_f^2 \beta_0 \mathcal{L}_{m_2}^2 m_1^3}{16m_- m_+^2} + \\
& \frac{9C_A Y_f^2 \beta_0 \mathcal{L}_{m_2}^2 m_1^3}{16\varepsilon m_- m_+^2} + \frac{2C_A C_F^2 m_2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^2}{m_+^3} + \frac{C_F m_2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^2}{m_+^3} + \frac{3C_A C_F^2 m_2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^2}{2\varepsilon m_+^3} + \frac{3C_F m_2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^2}{4\varepsilon m_+^3} - \\
& \frac{11C_A m_2 Y_f^2 \beta_0 \mathcal{L}_{m_1}^2 m_1^2}{32m_- m_+^2} - \frac{3C_A C_F^2 m_2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^2}{m_- m_+^3} - \frac{3C_F m_2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^2}{2m_- m_+^3} - \frac{3C_A m_2^2 Y_f^2 \beta_0 \mathcal{L}_{m_2}^2 m_1^2}{16m_- m_+^3} + \frac{14C_A C_F^2 m_2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^2}{m_- m_+^2} + \\
& \frac{7C_F m_2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^2}{m_- m_+^2} - \frac{3C_A C_F^2 m_2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^2}{\varepsilon m_- m_+^2} - \frac{3C_F m_2 Y_f^2 \mathcal{L}_{m_1}^2 m_1^2}{2\varepsilon m_- m_+^2} + \frac{17C_A m_2 Y_f^2 \beta_0 \mathcal{L}_{m_1}^2 m_1^2}{16m_- m_+^2} + \frac{11C_A m_2 Y_f^2 \beta_0 \mathcal{L}_{m_1}^2 m_1^2}{16\varepsilon m_- m_+^2} + \\
& \frac{22C_A C_F^2 m_2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^2}{m_- m_+^2} + \frac{11C_F m_2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^2}{m_- m_+^2} + \frac{3C_A C_F^2 m_2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^2}{\varepsilon m_- m_+^2} + \frac{3C_F m_2 Y_f^2 \mathcal{L}_{m_2}^2 m_1^2}{2\varepsilon m_- m_+^2} + \frac{9C_A m_2 Y_f^2 \beta_0 \mathcal{L}_{m_2}^2 m_1^2}{16m_- m_+^2} + \\
& \frac{9C_A m_2 Y_f^2 \beta_0 \mathcal{L}_{m_2}^2 m_1^2}{16\varepsilon m_- m_+^2} - \frac{5C_A^2 C_F m_2 Y_f^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} m_1^2}{m_+^3} - \frac{3C_A^2 C_F m_2 Y_f^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} m_1^2}{2\varepsilon m_+^3} - \frac{3C_A^2 C_F m_2 Y_f^2 \mathcal{L}_{m_1}^3 m_1}{4m_+^2} - \\
& \frac{3C_A^2 C_F m_2 Y_f^2 \mathcal{L}_{m_2}^3 m_1}{4m_+^2} - \frac{45C_A C_F^2 Y_f^2 m_1}{m_+} - \frac{45C_F Y_f^2 m_1}{2m_+} - \frac{18C_A C_F^2 Y_f^2 m_1}{\varepsilon m_+} - \frac{9C_F Y_f^2 m_1}{\varepsilon m_+} + \frac{11C_A C_F^2 m_2^2 Y_f^2 \mathcal{L}_{m_1}^2 m_1}{m_+^3} + \\
& \frac{11C_F m_2^2 Y_f^2 \mathcal{L}_{m_1}^2 m_1}{2m_+^3} + \frac{3C_A C_F^2 m_2^2 Y_f^2 \mathcal{L}_{m_1}^2 m_1}{2\varepsilon m_+^3} + \frac{3C_F m_2^2 Y_f^2 \mathcal{L}_{m_1}^2 m_1}{4\varepsilon m_+^3} + \frac{17C_A m_2^2 Y_f^2 \beta_0 \mathcal{L}_{m_1}^2 m_1}{32m_- m_+^2} - \frac{8C_A C_F^2 m_2^3 Y_f^2 \mathcal{L}_{m_2}^2 m_1}{m_- m_+^3} - \\
& \frac{4C_F m_2^3 Y_f^2 \mathcal{L}_{m_2}^2 m_1}{m_- m_+^3} - \frac{3C_A C_F^2 m_2^3 Y_f^2 \mathcal{L}_{m_2}^2 m_1}{2\varepsilon m_- m_+^3} - \frac{3C_F m_2^3 Y_f^2 \mathcal{L}_{m_2}^2 m_1}{4\varepsilon m_- m_+^3} + \frac{7C_A m_2^3 Y_f^2 \beta_0 \mathcal{L}_{m_2}^2 m_1}{16m_- m_+^3} + \frac{3C_A^2 C_F m_2 Y_f^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 m_1}{4m_+^2} - \\
& \frac{2C_A Y_f^2 \beta_0 m_1}{m_+} - \frac{9C_A Y_f^2 \beta_0 m_1}{8\varepsilon m_+} - \frac{5C_A Y_f^2 \beta_0 m_1}{4\varepsilon^2 m_+} - \frac{5C_A \pi^2 Y_f^2 \beta_0 m_1}{48m_+} - \frac{40C_A C_F^2 m_2^2 Y_f^2 \mathcal{L}_{m_1} m_1}{m_- m_+^2} - \frac{20C_F m_2^2 Y_f^2 \mathcal{L}_{m_1} m_1}{m_- m_+^2} - \\
& \frac{3C_A C_F^2 m_2^2 Y_f^2 \mathcal{L}_{m_1} m_1}{\varepsilon m_- m_+^2} - \frac{3C_F m_2^2 Y_f^2 \mathcal{L}_{m_1} m_1}{2\varepsilon m_- m_+^2} - \frac{25C_A m_2^2 Y_f^2 \beta_0 \mathcal{L}_{m_1} m_1}{16m_- m_+^2} - \frac{17C_A m_2^2 Y_f^2 \beta_0 \mathcal{L}_{m_1} m_1}{16\varepsilon m_- m_+^2} + \frac{4C_A C_F^2 m_2^3 Y_f^2 \mathcal{L}_{m_2} m_1}{m_- m_+^2} + \\
& \frac{2C_F m_2^3 Y_f^2 \mathcal{L}_{m_2} m_1}{m_- m_+^2} + \frac{3C_A C_F^2 m_2^3 Y_f^2 \mathcal{L}_{m_2} m_1}{\varepsilon m_- m_+^2} + \frac{3C_F m_2^3 Y_f^2 \mathcal{L}_{m_2} m_1}{2\varepsilon m_- m_+^2} + \frac{3C_A^2 C_F m_2 Y_f^2 \mathcal{L}_{m_1}^2 \mathcal{L}_{m_2} m_1}{4m_+^2} + \frac{7C_A m_2^2 Y_f^2 \beta_0 \mathcal{L}_{m_2} m_1}{16m_- m_+^2} - \\
& \frac{3C_A m_2^2 Y_f^2 \beta_0 \mathcal{L}_{m_2} m_1}{16\varepsilon m_- m_+^2} - \frac{8C_A^2 C_F m_2^2 Y_f^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} m_1}{m_+^3} - \frac{3C_A^2 C_F m_2^2 Y_f^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} m_1}{2\varepsilon m_+^3} - \frac{51C_A C_F^2 m_2 Y_f^2}{m_+} - \frac{51C_F m_2 Y_f^2}{2m_+} - \\
& \frac{18C_A C_F^2 m_2 Y_f^2}{\varepsilon m_+} - \frac{9C_F m_2 Y_f^2}{\varepsilon m_+} + \frac{3C_A C_F^2 m_2^3 Y_f^2 \mathcal{L}_{m_1}^2 m_1}{m_+^3} + \frac{3C_F m_2^3 Y_f^2 \mathcal{L}_{m_1}^2 m_1}{2m_+^3} + \frac{9C_A m_2^3 Y_f^2 \beta_0 \mathcal{L}_{m_1}^2 m_1}{32m_- m_+^2} + \frac{3C_A C_F^2 m_2^4 Y_f^2 \mathcal{L}_{m_2}^2 m_1}{m_- m_+^3} + \\
& \frac{3C_F m_2^4 Y_f^2 \mathcal{L}_{m_2}^2 m_1}{2m_- m_+^3} + \frac{11C_A m_2^4 Y_f^2 \beta_0 \mathcal{L}_{m_2}^2 m_1}{32m_- m_+^3} - \frac{7C_A m_2 Y_f^2 \beta_0}{2m_+} - \frac{13C_A m_2 Y_f^2 \beta_0}{8\varepsilon m_+} - \frac{5C_A m_2 Y_f^2 \beta_0}{4\varepsilon^2 m_+} - \frac{5C_A m_2 \pi^2 Y_f^2 \beta_0}{48m_+} - \\
& \frac{5C_A C_F^2 m_2^3 Y_f^2 \mathcal{L}_{m_1} m_1}{m_- m_+^2} - \frac{5C_F m_2^3 Y_f^2 \mathcal{L}_{m_1} m_1}{2m_- m_+^2} + \frac{3C_A C_F^2 m_2^3 Y_f^2 \mathcal{L}_{m_1} m_1}{\varepsilon m_- m_+^2} + \frac{3C_F m_2^3 Y_f^2 \mathcal{L}_{m_1} m_1}{2\varepsilon m_- m_+^2} - \frac{9C_A m_2^3 Y_f^2 \beta_0 \mathcal{L}_{m_1} m_1}{16m_- m_+^2} - \frac{9C_A m_2^3 Y_f^2 \beta_0 \mathcal{L}_{m_1} m_1}{16\varepsilon m_- m_+^2} - \\
& \frac{31C_A C_F^2 m_2^3 Y_f^2 \mathcal{L}_{m_2} m_1}{m_- m_+^2} - \frac{31C_F m_2^3 Y_f^2 \mathcal{L}_{m_2} m_1}{2m_- m_+^2} - \frac{3C_A C_F^2 m_2^3 Y_f^2 \mathcal{L}_{m_2} m_1}{\varepsilon m_- m_+^2} - \frac{3C_F m_2^3 Y_f^2 \mathcal{L}_{m_2} m_1}{2\varepsilon m_- m_+^2} - \frac{17C_A m_2^3 Y_f^2 \beta_0 \mathcal{L}_{m_2} m_1}{16m_- m_+^2} -
\end{aligned}$$

$$\begin{aligned}
& \frac{11C_A m_2^3 Y_f^2 \beta_0 \mathcal{L} m_2}{16\epsilon m_- m_+^2} - \frac{27m_1^5 Y_f^4 C_A^2 \mathcal{L} m_2}{32m_-^2 m_+^3} - \frac{33m_1^5 Y_f^4 C_A^2 \mathcal{L} m_1}{32m_-^2 m_+^3} - \frac{39m_1^4 m_2 Y_f^4 C_A^2 \mathcal{L} m_1}{32m_-^2 m_+^3} - \frac{21m_1^4 m_2 Y_f^4 C_A^2 \mathcal{L} m_2}{32m_-^2 m_+^3} - \\
& \frac{9m_1^3 m_2^2 Y_f^4 C_A^2 \mathcal{L} m_1}{32m_-^2 m_+^3} - \frac{9m_1^3 m_2^2 Y_f^4 C_A^2 \mathcal{L} m_2}{32m_-^2 m_+^3} + \frac{45m_1^3 m_2^2 Y_f^4 C_A^2 \mathcal{L} m_1}{16m_-^2 m_+^3} + \frac{15m_1^3 m_2^2 Y_f^4 C_A^2 \mathcal{L} m_2}{16m_-^2 m_+^3} + \frac{9m_1^3 m_2^2 Y_f^4 C_A^2 \mathcal{L} m_1 \mathcal{L} m_2}{16m_-^2 m_+^3} + \\
& \frac{9m_1^2 m_2^3 Y_f^4 C_A^2 \mathcal{L} m_1}{32m_-^2 m_+^3} + \frac{9m_1^2 m_2^3 Y_f^4 C_A^2 \mathcal{L} m_2}{32m_-^2 m_+^3} + \frac{33m_1^2 m_2^3 Y_f^4 C_A^2 \mathcal{L} m_1}{16m_-^2 m_+^3} + \frac{27m_1^2 m_2^3 Y_f^4 C_A^2 \mathcal{L} m_2}{16m_-^2 m_+^3} - \frac{9m_1^2 m_2^3 Y_f^4 C_A^2 \mathcal{L} m_1 \mathcal{L} m_2}{16m_-^2 m_+^3} - \\
& \frac{27m_1^5 Y_f^4 C_A^2 \mathcal{L} m_1}{32m_-^2 m_+^3} - \frac{3m_1 m_2^4 Y_f^4 C_A^2 \mathcal{L} m_1}{16m_-^2 m_+^3} - \frac{3m_1 m_2^4 Y_f^4 C_A^2 \mathcal{L} m_2}{16m_-^2 m_+^3} - \frac{57m_1 m_2^4 Y_f^4 C_A^2 \mathcal{L} m_1}{32m_-^2 m_+^3} - \frac{3m_1 m_2^4 Y_f^4 C_A^2 \mathcal{L} m_2}{32m_-^2 m_+^3} + \\
& \frac{3m_1 m_2^4 Y_f^4 C_A^2 \mathcal{L} m_1 \mathcal{L} m_2}{8m_-^2 m_+^3} - \frac{33m_1^5 Y_f^4 C_A^2 \mathcal{L} m_2}{32m_-^2 m_+^3} + \frac{15m_1 Y_f^4 C_A^2}{16\epsilon m_+} + \frac{97m_1 Y_f^4 C_A^2}{32m_+} + \frac{15m_2 Y_f^4 C_A^2}{16\epsilon m_+} + \frac{109m_2 Y_f^4 C_A^2}{32m_+} \quad (B.2)
\end{aligned}$$

$$\begin{aligned}
\Delta B_3^{(2)} = & -\frac{3C_A m_+^2 \mathcal{L} m_1 C_F^3}{2m_-^2} + 3C_A \mathcal{L} m_1 C_F^3 - \frac{3C_A m_-^2 \mathcal{L} m_1 C_F^3}{2m_+^2} - \frac{3C_A m_+^2 \mathcal{L} m_2 C_F^3}{2m_-^2} + 3C_A \mathcal{L} m_2 C_F^3 - \frac{3C_A m_-^2 \mathcal{L} m_2 C_F^3}{2m_+^2} + \\
& \frac{7C_A m_+^2 \mathcal{L} m_1 C_F^3}{2m_-^2} + \frac{3C_A m_+^2 \mathcal{L} m_1 C_F^3}{2\epsilon m_-^2} - 10C_A \mathcal{L} m_1 C_F^3 + \frac{12C_A m_+ \mathcal{L} m_1 C_F^3}{m_-} - \frac{3C_A \mathcal{L} m_1 C_F^3}{\epsilon} + \frac{12C_A m_- \mathcal{L} m_1 C_F^3}{m_+} + \\
& \frac{13C_A m_-^2 \mathcal{L} m_1 C_F^3}{2m_+^2} + \frac{3C_A m_-^2 \mathcal{L} m_1 C_F^3}{2\epsilon m_+^2} + \frac{7C_A m_+^2 \mathcal{L} m_2 C_F^3}{2m_-^2} + \frac{3C_A m_+^2 \mathcal{L} m_2 C_F^3}{2\epsilon m_-^2} - 10C_A \mathcal{L} m_2 C_F^3 - \frac{12C_A m_+ \mathcal{L} m_2 C_F^3}{m_-} + \\
& \frac{3C_A m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^3}{2m_-^2} - 3C_A \mathcal{L} m_1 \mathcal{L} m_2 C_F^3 + \frac{3C_A m_-^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^3}{2m_+^2} - \frac{3C_A \mathcal{L} m_2 C_F^3}{\epsilon} - \frac{12C_A m_- \mathcal{L} m_2 C_F^3}{m_+} + \\
& \frac{13C_A m_-^2 \mathcal{L} m_2 C_F^3}{2m_+^2} + \frac{3C_A m_-^2 \mathcal{L} m_2 C_F^3}{2\epsilon m_+^2} + 72C_A C_F^3 - 36C_A \mathcal{L} m_1 C_F^3 - \frac{28C_A m_+ \mathcal{L} m_1 C_F^3}{m_-} - \frac{12C_A m_+ \mathcal{L} m_1 C_F^3}{\epsilon m_-} - \\
& \frac{52C_A m_- \mathcal{L} m_1 C_F^3}{m_+} - \frac{12C_A m_- \mathcal{L} m_1 C_F^3}{\epsilon m_+} + \frac{3C_A m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^3}{2m_-^2} - 3C_A \mathcal{L} m_1 \mathcal{L} m_2 C_F^3 + \\
& \frac{3C_A m_-^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^3}{2m_+^2} - 36C_A \mathcal{L} m_2 C_F^3 + \frac{28C_A m_+ \mathcal{L} m_2 C_F^3}{m_-} + \frac{12C_A m_+ \mathcal{L} m_2 C_F^3}{\epsilon m_-} - \frac{7C_A m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^3}{m_-^2} - \\
& \frac{3C_A m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^3}{\epsilon m_-^2} + 20C_A \mathcal{L} m_1 \mathcal{L} m_2 C_F^3 + \frac{6C_A \mathcal{L} m_1 \mathcal{L} m_2 C_F^3}{\epsilon} - \frac{13C_A m_-^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^3}{m_+^2} - \frac{3C_A m_-^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^3}{\epsilon m_+^2} + \\
& \frac{52C_A m_- \mathcal{L} m_2 C_F^3}{m_+} + \frac{12C_A m_- \mathcal{L} m_2 C_F^3}{\epsilon m_+} + \frac{36C_A C_F^3}{\epsilon} - \frac{3m_+^2 \mathcal{L} m_1 C_F^2}{4m_-^2} - \frac{3m_-^2 \mathcal{L} m_1 C_F^2}{4m_+^2} + \frac{3}{2} \mathcal{L} m_1 C_F^2 - \frac{3m_+^2 \mathcal{L} m_2 C_F^2}{4m_-^2} - \\
& \frac{3m_-^2 \mathcal{L} m_2 C_F^2}{4m_+^2} + \frac{3}{2} \mathcal{L} m_2 C_F^2 + \frac{3m_+^2 \mathcal{L} m_1 C_F^2}{4\epsilon m_-^2} + \frac{7m_+^2 \mathcal{L} m_1 C_F^2}{4m_-^2} + \frac{6m_+ \mathcal{L} m_1 C_F^2}{m_-} - \frac{3\mathcal{L} m_1 C_F^2}{2\epsilon} + \frac{6m_- \mathcal{L} m_1 C_F^2}{m_+} + \frac{3m_-^2 \mathcal{L} m_1 C_F^2}{4\epsilon m_+^2} + \\
& \frac{13m_-^2 \mathcal{L} m_1 C_F^2}{4m_+^2} - 5\mathcal{L} m_1 C_F^2 + \frac{3m_+^2 \mathcal{L} m_2 C_F^2}{4\epsilon m_-^2} + \frac{7m_+^2 \mathcal{L} m_2 C_F^2}{4m_-^2} - \frac{6m_+ \mathcal{L} m_2 C_F^2}{m_-} + \frac{3m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^2}{4m_-^2} + \frac{3m_-^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^2}{4m_+^2} - \\
& \frac{3}{2} \mathcal{L} m_1 \mathcal{L} m_2 C_F^2 - \frac{3\mathcal{L} m_2 C_F^2}{2\epsilon} - \frac{6m_- \mathcal{L} m_2 C_F^2}{m_+} + \frac{3m_-^2 \mathcal{L} m_2 C_F^2}{4\epsilon m_+^2} + \frac{13m_-^2 \mathcal{L} m_2 C_F^2}{4m_+^2} - 5\mathcal{L} m_2 C_F^2 - \frac{6m_+ \mathcal{L} m_1 C_F^2}{\epsilon m_-} - \frac{14m_+ \mathcal{L} m_1 C_F^2}{m_-} - \\
& \frac{6m_- \mathcal{L} m_1 C_F^2}{\epsilon m_+} - \frac{26m_- \mathcal{L} m_1 C_F^2}{m_+} - 18\mathcal{L} m_1 C_F^2 + \frac{3m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^2}{4m_-^2} + \frac{3m_-^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^2}{4m_+^2} - \frac{3}{2} \mathcal{L} m_1 \mathcal{L} m_2 C_F^2 + \frac{6m_+ \mathcal{L} m_2 C_F^2}{\epsilon m_-} + \\
& \frac{14m_+ \mathcal{L} m_2 C_F^2}{m_-} - \frac{3m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^2}{2\epsilon m_-^2} - \frac{7m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^2}{2m_-^2} + \frac{3\mathcal{L} m_1 \mathcal{L} m_2 C_F^2}{\epsilon} - \frac{3m_-^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^2}{2\epsilon m_+^2} - \\
& \frac{13m_-^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^2}{2m_+^2} + 10\mathcal{L} m_1 \mathcal{L} m_2 C_F^2 + \frac{6m_- \mathcal{L} m_2 C_F^2}{\epsilon m_+} + \frac{26m_- \mathcal{L} m_2 C_F^2}{m_+} - 18\mathcal{L} m_2 C_F^2 + \frac{18C_F^2}{\epsilon} + 36C_F^2 - \frac{C_A m_+ \beta_0 \mathcal{L} m_1 C_F}{12m_-} - \\
& \frac{C_A m_- \beta_0 \mathcal{L} m_1 C_F}{12m_+} + \frac{C_A m_+ \beta_0 \mathcal{L} m_2 C_F}{12m_-} + \frac{C_A m_- \beta_0 \mathcal{L} m_2 C_F}{12m_+} + \frac{3}{4} C_A \beta_0 \mathcal{L} m_1 C_F + \frac{C_A m_+ \beta_0 \mathcal{L} m_1 C_F}{4m_-} + \frac{C_A m_+ \beta_0 \mathcal{L} m_2 C_F}{4\epsilon m_-} + \\
& \frac{3C_A m_- \beta_0 \mathcal{L} m_1 C_F}{4m_+} + \frac{C_A m_- \beta_0 \mathcal{L} m_1 C_F}{4\epsilon m_+} + \frac{3}{4} C_A \beta_0 \mathcal{L} m_2 C_F - \frac{C_A m_+ \beta_0 \mathcal{L} m_2 C_F}{4m_-} - \frac{C_A m_+ \beta_0 \mathcal{L} m_2 C_F}{4\epsilon m_-} - \frac{3C_A m_- \beta_0 \mathcal{L} m_2 C_F}{4m_+} - \\
& \frac{C_A m_- \beta_0 \mathcal{L} m_2 C_F}{4\epsilon m_+} + \frac{2C_A \beta_0 C_F}{\epsilon} + \frac{3C_A \beta_0 C_F}{\epsilon^2} + \frac{1}{4} C_A \pi^2 \beta_0 C_F - C_A \beta_0 \mathcal{L} m_1 C_F - \frac{C_A m_+ \beta_0 \mathcal{L} m_1 C_F}{2m_-} - \frac{C_A m_+ \beta_0 \mathcal{L} m_1 C_F}{2\epsilon m_-} - \\
& \frac{C_A m_+ \beta_0 \mathcal{L} m_1 C_F}{2\epsilon^2 m_-} - \frac{3C_A \beta_0 \mathcal{L} m_1 C_F}{2\epsilon} - \frac{5C_A m_- \beta_0 \mathcal{L} m_1 C_F}{2m_+} - \frac{3C_A m_- \beta_0 \mathcal{L} m_1 C_F}{2\epsilon m_+} - \frac{C_A m_- \beta_0 \mathcal{L} m_1 C_F}{2\epsilon^2 m_+} - \\
& \frac{C_A m_+ \pi^2 \beta_0 \mathcal{L} m_1 C_F}{24m_-} - \frac{C_A m_- \pi^2 \beta_0 \mathcal{L} m_1 C_F}{24m_+} - C_A \beta_0 \mathcal{L} m_2 C_F + \frac{C_A m_+ \beta_0 \mathcal{L} m_2 C_F}{2m_-} + \frac{C_A m_+ \beta_0 \mathcal{L} m_2 C_F}{2\epsilon m_-} + \\
& \frac{C_A m_+ \beta_0 \mathcal{L} m_2 C_F}{2\epsilon^2 m_-} - \frac{3C_A \beta_0 \mathcal{L} m_2 C_F}{2\epsilon} + \frac{5C_A m_- \beta_0 \mathcal{L} m_2 C_F}{2m_+} + \frac{3C_A m_- \beta_0 \mathcal{L} m_2 C_F}{2\epsilon m_+} + \frac{C_A m_- \beta_0 \mathcal{L} m_2 C_F}{2\epsilon^2 m_+} + \\
& \frac{C_A m_+ \pi^2 \beta_0 \mathcal{L} m_2 C_F}{24m_-} + \frac{C_A m_- \pi^2 \beta_0 \mathcal{L} m_2 C_F}{24m_+} - \frac{3C_A C_F^2 Y_f^2 \mathcal{L} m_1 m_1^4}{m_-^2 m_+^2} - \frac{3C_F Y_f^2 \mathcal{L} m_1 m_1^4}{2m_-^2 m_+^2} + \frac{3C_A C_F^2 Y_f^2 \mathcal{L} m_2 m_1^4}{m_-^2 m_+^2} +
\end{aligned}$$



$$\frac{27m_2^2 Y_f^4 C_A^2 \mathcal{L} m_2}{32m_- m_+} + \frac{27Y_f^4 C_A^2}{32\varepsilon} + \frac{21}{8} Y_f^4 C_A^2 \quad (\text{B.3})$$

$$\begin{aligned} \Delta B_4^{(2)} = & -\frac{3C_A m_+^2 \mathcal{L}^3 m_1 C_F^3}{4m_-^2} + \frac{3}{2} C_A \mathcal{L}^3 m_1 C_F^3 - \frac{3C_A m_-^2 \mathcal{L}^3 m_1 C_F^3}{4m_+^2} - \frac{3C_A m_+^2 \mathcal{L}^3 m_2 C_F^3}{4m_-^2} + \frac{3}{2} C_A \mathcal{L}^3 m_2 C_F^3 - \frac{3C_A m_-^2 \mathcal{L}^3 m_2 C_F^3}{4m_+^2} + \\ & \frac{17C_A m_+^2 \mathcal{L}^2 m_1 C_F^3}{8m_-^2} + \frac{3C_A m_+^2 \mathcal{L}^2 m_1 C_F^3}{4\varepsilon m_-^2} - \frac{17}{4} C_A \mathcal{L}^2 m_1 C_F^3 + \frac{6C_A m_+ \mathcal{L}^2 m_1 C_F^3}{m_-} - \frac{3C_A \mathcal{L}^2 m_1 C_F^3}{2\varepsilon} + \frac{6C_A m_- \mathcal{L}^2 m_1 C_F^3}{m_+} + \\ & \frac{17C_A m_-^2 \mathcal{L}^2 m_1 C_F^3}{8m_+^2} + \frac{3C_A m_-^2 \mathcal{L}^2 m_1 C_F^3}{4\varepsilon m_+^2} + \frac{17C_A m_+^2 \mathcal{L}^2 m_2 C_F^3}{8m_-^2} + \frac{3C_A m_+^2 \mathcal{L}^2 m_2 C_F^3}{4\varepsilon m_-^2} - \frac{17}{4} C_A \mathcal{L}^2 m_2 C_F^3 - \frac{6C_A m_+ \mathcal{L}^2 m_2 C_F^3}{m_-} + \\ & \frac{3C_A m_-^2 \mathcal{L}^2 m_2 C_F^3}{4m_+^2} - \frac{3}{2} C_A \mathcal{L} m_1 \mathcal{L}^2 m_2 C_F^3 + \frac{3C_A m_-^2 \mathcal{L} m_1 \mathcal{L}^2 m_2 C_F^3}{4m_+^2} - \frac{3C_A \mathcal{L}^2 m_2 C_F^3}{2\varepsilon} - \frac{6C_A m_- \mathcal{L}^2 m_2 C_F^3}{m_+} + \frac{17C_A m_+^2 \mathcal{L}^2 m_2 C_F^3}{8m_-^2} + \\ & \frac{3C_A m_-^2 \mathcal{L}^2 m_2 C_F^3}{4\varepsilon m_+^2} - 4C_A C_F^3 - 6C_A \mathcal{L} m_1 C_F^3 - \frac{17C_A m_+ \mathcal{L} m_1 C_F^3}{m_-} - \frac{6C_A m_+ \mathcal{L} m_1 C_F^3}{\varepsilon m_-} - \frac{17C_A m_- \mathcal{L} m_1 C_F^3}{m_+} - \\ & \frac{6C_A m_- \mathcal{L} m_1 C_F^3}{\varepsilon m_+} + \frac{3C_A m_+^2 \mathcal{L}^2 m_1 \mathcal{L} m_2 C_F^3}{4m_-^2} - \frac{3}{2} C_A \mathcal{L}^2 m_1 \mathcal{L} m_2 C_F^3 + \frac{3C_A m_-^2 \mathcal{L}^2 m_1 \mathcal{L} m_2 C_F^3}{4m_+^2} - 6C_A \mathcal{L} m_2 C_F^3 + \frac{17C_A m_+ \mathcal{L} m_2 C_F^3}{m_-} + \\ & \frac{6C_A m_- \mathcal{L} m_2 C_F^3}{\varepsilon m_-} - \frac{17C_A m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^3}{4m_-^2} - \frac{3C_A m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^3}{2\varepsilon m_-^2} + \frac{17}{2} C_A \mathcal{L} m_1 \mathcal{L} m_2 C_F^3 + \frac{3C_A \mathcal{L} m_1 \mathcal{L} m_2 C_F^3}{\varepsilon} - \\ & \frac{17C_A m_-^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^3}{4m_+^2} - \frac{3C_A m_-^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^3}{2\varepsilon m_+^2} + \frac{17C_A m_- \mathcal{L} m_2 C_F^3}{m_+} + \frac{6C_A m_- \mathcal{L} m_2 C_F^3}{\varepsilon m_+} + \frac{6C_A C_F^3}{\varepsilon} - \frac{3m_+^2 \mathcal{L}^3 m_1 C_F^2}{8m_-^2} - \\ & \frac{3m_-^2 \mathcal{L}^3 m_1 C_F^2}{8m_+^2} + \frac{3}{4} \mathcal{L}^3 m_1 C_F^2 - \frac{3m_+^2 \mathcal{L}^3 m_2 C_F^2}{8m_-^2} - \frac{3m_-^2 \mathcal{L}^3 m_2 C_F^2}{8m_+^2} + \frac{3}{4} \mathcal{L}^3 m_2 C_F^2 + \frac{3m_+^2 \mathcal{L}^2 m_1 C_F^2}{8\varepsilon m_-^2} + \frac{17m_+^2 \mathcal{L}^2 m_1 C_F^2}{16m_-^2} + \frac{3m_+ \mathcal{L}^2 m_1 C_F^2}{m_-} + \\ & \frac{3\mathcal{L}^2 m_1 C_F^2}{4\varepsilon} + \frac{3m_- \mathcal{L}^2 m_1 C_F^2}{m_+} + \frac{3m_-^2 \mathcal{L}^2 m_1 C_F^2}{8\varepsilon m_+^2} + \frac{17m_-^2 \mathcal{L}^2 m_1 C_F^2}{16m_+^2} - \frac{17}{8} \mathcal{L}^2 m_1 C_F^2 + \frac{3m_+^2 \mathcal{L}^2 m_2 C_F^2}{8\varepsilon m_-^2} + \frac{17m_+^2 \mathcal{L}^2 m_2 C_F^2}{16m_-^2} - \\ & \frac{3m_+ \mathcal{L}^2 m_2 C_F^2}{m_-} + \frac{3m_+^2 \mathcal{L} m_1 \mathcal{L}^2 m_2 C_F^2}{8m_-^2} + \frac{3m_-^2 \mathcal{L} m_1 \mathcal{L}^2 m_2 C_F^2}{8m_+^2} - \frac{3}{4} \mathcal{L} m_1 \mathcal{L}^2 m_2 C_F^2 - \frac{3\mathcal{L}^2 m_2 C_F^2}{4\varepsilon} - \frac{3m_- \mathcal{L}^2 m_2 C_F^2}{m_+} + \frac{3m_-^2 \mathcal{L}^2 m_2 C_F^2}{8\varepsilon m_+^2} + \\ & \frac{17m_-^2 \mathcal{L}^2 m_2 C_F^2}{16m_+^2} - \frac{17}{8} \mathcal{L}^2 m_2 C_F^2 - \frac{3m_+ \mathcal{L} m_1 C_F^2}{\varepsilon m_-} - \frac{17m_+ \mathcal{L} m_1 C_F^2}{2m_-} - \frac{3m_- \mathcal{L} m_1 C_F^2}{\varepsilon m_+} - \frac{17m_- \mathcal{L} m_1 C_F^2}{2m_+} - 3\mathcal{L} m_1 C_F^2 + \\ & \frac{3m_+^2 \mathcal{L}^2 m_1 \mathcal{L} m_2 C_F^2}{8m_-^2} + \frac{3m_-^2 \mathcal{L}^2 m_1 \mathcal{L} m_2 C_F^2}{8m_+^2} - \frac{3}{4} \mathcal{L}^2 m_1 \mathcal{L} m_2 C_F^2 + \frac{3m_+ \mathcal{L} m_2 C_F^2}{\varepsilon m_-} + \frac{17m_+ \mathcal{L} m_2 C_F^2}{2m_-} - \frac{3m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^2}{4\varepsilon m_-^2} - \\ & \frac{17m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^2}{8m_-^2} + \frac{3\mathcal{L} m_1 \mathcal{L} m_2 C_F^2}{2\varepsilon} - \frac{3m_- \mathcal{L} m_1 \mathcal{L} m_2 C_F^2}{4\varepsilon m_+^2} - \frac{17m_-^2 \mathcal{L} m_1 \mathcal{L} m_2 C_F^2}{8m_+^2} + \frac{17}{4} \mathcal{L} m_1 \mathcal{L} m_2 C_F^2 + \frac{3m_- \mathcal{L} m_2 C_F^2}{\varepsilon m_+} + \\ & \frac{17m_- \mathcal{L} m_2 C_F^2}{2m_+} - 3\mathcal{L} m_2 C_F^2 + \frac{3C_F^2}{\varepsilon} - 2C_F^2 - \frac{C_A m_+ \beta_0 \mathcal{L}^3 m_1 C_F}{12m_-} - \frac{C_A m_- \beta_0 \mathcal{L}^3 m_1 C_F}{12m_+} + \frac{C_A m_+ \beta_0 \mathcal{L}^3 m_2 C_F}{12m_-} + \\ & \frac{C_A m_- \beta_0 \mathcal{L}^3 m_2 C_F}{12m_+} + \frac{1}{4} C_A \beta_0 \mathcal{L}^2 m_1 C_F + \frac{C_A m_+ \beta_0 \mathcal{L}^2 m_1 C_F}{8m_-} + \frac{C_A m_+ \beta_0 \mathcal{L}^2 m_1 C_F}{4\varepsilon m_-} + \frac{C_A m_- \beta_0 \mathcal{L}^2 m_1 C_F}{8m_+} + \frac{C_A m_- \beta_0 \mathcal{L}^2 m_1 C_F}{4\varepsilon m_+} + \\ & \frac{1}{4} C_A \beta_0 \mathcal{L}^2 m_2 C_F - \frac{C_A m_+ \beta_0 \mathcal{L}^2 m_2 C_F}{8m_-} - \frac{C_A m_+ \beta_0 \mathcal{L}^2 m_2 C_F}{4\varepsilon m_-} - \frac{C_A m_- \beta_0 \mathcal{L}^2 m_2 C_F}{8m_+} - \frac{C_A m_- \beta_0 \mathcal{L}^2 m_2 C_F}{4\varepsilon m_+} - 7C_A \beta_0 C_F - \\ & \frac{3C_A \beta_0 C_F}{\varepsilon} + \frac{C_A \beta_0 C_F}{\varepsilon^2} + \frac{1}{12} C_A \pi^2 \beta_0 C_F + \frac{3}{2} C_A \beta_0 \mathcal{L} m_1 C_F - \frac{C_A m_+ \beta_0 \mathcal{L} m_1 C_F}{4m_-} - \frac{C_A m_+ \beta_0 \mathcal{L} m_1 C_F}{4\varepsilon m_-} - \frac{C_A m_+ \beta_0 \mathcal{L} m_1 C_F}{2\varepsilon^2 m_-} - \\ & \frac{C_A \beta_0 \mathcal{L} m_1 C_F}{2\varepsilon} - \frac{C_A m_- \beta_0 \mathcal{L} m_1 C_F}{4m_+} - \frac{C_A m_- \beta_0 \mathcal{L} m_1 C_F}{4\varepsilon m_+} - \frac{C_A m_- \beta_0 \mathcal{L} m_1 C_F}{2\varepsilon^2 m_+} - \frac{C_A m_+ \pi^2 \beta_0 \mathcal{L} m_1 C_F}{24m_-} - \\ & \frac{C_A m_- \pi^2 \beta_0 \mathcal{L} m_1 C_F}{24m_+} + \frac{3}{2} C_A \beta_0 \mathcal{L} m_2 C_F + \frac{C_A m_+ \beta_0 \mathcal{L} m_2 C_F}{4m_-} + \frac{C_A m_+ \beta_0 \mathcal{L} m_2 C_F}{4\varepsilon m_-} + \frac{C_A m_+ \beta_0 \mathcal{L} m_2 C_F}{2\varepsilon^2 m_-} - \frac{C_A \beta_0 \mathcal{L} m_2 C_F}{2\varepsilon} + \\ & \frac{C_A m_- \beta_0 \mathcal{L} m_2 C_F}{4m_+} + \frac{C_A m_- \beta_0 \mathcal{L} m_2 C_F}{4\varepsilon m_+} + \frac{C_A m_- \beta_0 \mathcal{L} m_2 C_F}{2\varepsilon^2 m_+} + \frac{C_A m_+ \pi^2 \beta_0 \mathcal{L} m_2 C_F}{24m_-} + \frac{C_A m_- \pi^2 \beta_0 \mathcal{L} m_2 C_F}{24m_+} + \\ & \frac{C_A C_F^2 Y_s^2 \mathcal{L}^2 m_2 m_1^4}{m_2^2 m_-^2 m_+^2} + \frac{C_F Y_s^2 \mathcal{L}^2 m_2 m_1^4}{2m_2^2 m_-^2 m_+^2} - \frac{C_A Y_s^2 \beta_0 \mathcal{L}^2 m_2 m_1^4}{16m_2^2 m_-^2 m_+^2} + \frac{17C_A C_F^2 Y_s^2 \mathcal{L}^3 m_1 m_1^2}{24m_-^2 m_+^2} + \frac{17C_F Y_s^2 \mathcal{L}^3 m_1 m_1^2}{48m_-^2 m_+^2} + \frac{C_A Y_s^2 \beta_0 \mathcal{L}^3 m_1 m_1^2}{48m_-^2 m_+^2} + \\ & \frac{25C_A C_F^2 Y_s^2 \mathcal{L}^3 m_2 m_1^2}{24m_-^2 m_+^2} + \frac{25C_F Y_s^2 \mathcal{L}^3 m_2 m_1^2}{48m_-^2 m_+^2} - \frac{C_A Y_s^2 \beta_0 \mathcal{L}^3 m_2 m_1^2}{48m_-^2 m_+^2} - \frac{31C_A C_F^2 Y_s^2 \mathcal{L}^2 m_1 m_1^2}{8m_-^2 m_+^2} - \frac{31C_F Y_s^2 \mathcal{L}^2 m_1 m_1^2}{16m_-^2 m_+^2} - \\ & \frac{5C_A C_F^2 Y_s^2 \mathcal{L}^2 m_1 m_1^2}{8\varepsilon m_-^2 m_+^2} - \frac{5C_F Y_s^2 \mathcal{L}^2 m_1 m_1^2}{16\varepsilon m_-^2 m_+^2} - \frac{11C_A C_F^2 Y_s^2 \mathcal{L}^2 m_2 m_1^2}{8m_-^2 m_+^2} - \frac{11C_F Y_s^2 \mathcal{L}^2 m_2 m_1^2}{16m_-^2 m_+^2} - \frac{9C_A C_F^2 Y_s^2 \mathcal{L}^2 m_2 m_1^2}{8\varepsilon m_-^2 m_+^2} - \frac{9C_F Y_s^2 \mathcal{L}^2 m_2 m_1^2}{16\varepsilon m_-^2 m_+^2} + \\ & \frac{C_A Y_s^2 \beta_0 \mathcal{L}^2 m_2 m_1^2}{8m_-^2 m_+^2} + \frac{C_A Y_s^2 \beta_0 \mathcal{L}^2 m_2 m_1^2}{16\varepsilon m_-^2 m_+^2} - \frac{7C_A^2 C_F Y_s^2 \mathcal{L} m_1 \mathcal{L}^2 m_2 m_1^2}{16m_-^2 m_+^2} - \frac{C_A C_F^2 Y_s^2 \mathcal{L} m_2 m_1^2}{2m_2^2 m_- m_+} - \frac{C_F Y_s^2 \mathcal{L} m_2 m_1^2}{4m_2^2 m_- m_+} - \\ & \frac{C_A C_F^2 Y_s^2 \mathcal{L} m_2 m_1^2}{\varepsilon m_2^2 m_- m_+} - \frac{C_F Y_s^2 \mathcal{L} m_2 m_1^2}{2\varepsilon m_2^2 m_- m_+} - \frac{7C_A^2 C_F Y_s^2 \mathcal{L}^2 m_1 \mathcal{L} m_2 m_1^2}{16m_-^2 m_+^2} + \frac{C_A Y_s^2 \beta_0 \mathcal{L} m_2 m_1^2}{8m_2^2 m_- m_+} + \frac{C_A Y_s^2 \beta_0 \mathcal{L} m_2 m_1^2}{8\varepsilon m_2^2 m_- m_+} \end{aligned}$$

$$\begin{aligned}
& \frac{17C_A^2 C_F Y_s^2 L_{m_1} L_{m_2} m_+^2}{8m_-^2 m_+^2} + \frac{7C_A^2 C_F Y_s^2 L_{m_1} L_{m_2} m_+^2}{8\epsilon m_-^2 m_+^2} + \frac{25C_A C_F^2 m_2^2 Y_s^2 L_{m_1}^3}{24m_-^2 m_+^2} + \frac{25C_F m_2^2 Y_s^2 L_{m_1}^3}{48m_-^2 m_+^2} - \frac{C_A m_2^2 Y_s^2 \beta_0 L_{m_1}^3}{48m_-^2 m_+^2} + \\
& \frac{17C_A C_F^2 m_2^2 Y_s^2 L_{m_2}^3}{24m_-^2 m_+^2} + \frac{17C_F m_2^2 Y_s^2 L_{m_2}^3}{48m_-^2 m_+^2} + \frac{C_A m_2^2 Y_s^2 \beta_0 L_{m_2}^3}{48m_-^2 m_+^2} + \frac{C_A C_F^2 Y_s^2}{2m_+^2} + \frac{C_F Y_s^2}{4m_+^2} + \frac{C_A C_F^2 Y_s^2}{4\epsilon m_+^2} + \frac{C_F Y_s^2}{8\epsilon m_+^2} + \frac{C_A C_F^2 Y_s^2}{2\epsilon^2 m_+^2} + \\
& \frac{C_F Y_s^2}{4\epsilon^2 m_+^2} + \frac{C_A C_F^2 \pi^2 Y_s^2}{12m_+^2} + \frac{C_F \pi^2 Y_s^2}{24m_+^2} - \frac{11C_A C_F^2 m_2^2 Y_s^2 L_{m_1}^2}{8m_-^2 m_+^2} - \frac{11C_F m_2^2 Y_s^2 L_{m_1}^2}{16m_-^2 m_+^2} - \frac{9C_A C_F^2 m_2^2 Y_s^2 L_{m_1}^2}{8\epsilon m_-^2 m_+^2} - \frac{9C_F m_2^2 Y_s^2 L_{m_1}^2}{16\epsilon m_-^2 m_+^2} - \\
& \frac{C_A Y_s^2 \beta_0 L_{m_1}^2}{16m_- m_+} - \frac{C_A Y_s^2 \beta_0 L_{m_1}^2}{16\epsilon m_- m_+} - \frac{31C_A C_F^2 m_2^2 Y_s^2 L_{m_2}^2}{8m_-^2 m_+^2} - \frac{31C_F m_2^2 Y_s^2 L_{m_2}^2}{16m_-^2 m_+^2} - \frac{5C_A C_F^2 m_2^2 Y_s^2 L_{m_2}^2}{8\epsilon m_-^2 m_+^2} - \frac{5C_F m_2^2 Y_s^2 L_{m_2}^2}{16\epsilon m_-^2 m_+^2} - \\
& \frac{C_A m_2^2 Y_s^2 \beta_0 L_{m_2}^2}{16m_-^2 m_+^2} - \frac{C_A m_2^2 Y_s^2 \beta_0 L_{m_2}^2}{16\epsilon m_-^2 m_+^2} - \frac{7C_A^2 C_F m_2^2 Y_s^2 L_{m_1} L_{m_2}^2}{16m_-^2 m_+^2} - \frac{C_A Y_s^2 \beta_0}{4m_+^2} - \frac{C_A Y_s^2 \beta_0}{8\epsilon m_+^2} - \frac{C_A Y_s^2 \beta_0}{8\epsilon^2 m_+^2} - \frac{C_A \pi^2 Y_s^2 \beta_0}{96m_+^2} + \\
& \frac{25C_A C_F^2 Y_s^2 L_{m_1}}{4m_- m_+} + \frac{25C_F Y_s^2 L_{m_1}}{8m_- m_+} + \frac{7C_A C_F^2 Y_s^2 L_{m_1}}{4\epsilon m_- m_+} + \frac{7C_F Y_s^2 L_{m_1}}{8\epsilon m_- m_+} - \frac{C_A C_F^2 Y_s^2 L_{m_1}}{4\epsilon^2 m_- m_+} - \frac{C_F Y_s^2 L_{m_1}}{8\epsilon^2 m_- m_+} - \frac{C_A C_F^2 \pi^2 Y_s^2 L_{m_1}}{24m_- m_+} - \\
& \frac{C_F \pi^2 Y_s^2 L_{m_1}}{48m_- m_+} + \frac{C_A Y_s^2 \beta_0 L_{m_1}}{8m_- m_+} + \frac{C_A Y_s^2 \beta_0 L_{m_1}}{8\epsilon m_- m_+} + \frac{C_A Y_s^2 \beta_0 L_{m_1}}{8\epsilon^2 m_- m_+} + \frac{C_A \pi^2 Y_s^2 \beta_0 L_{m_1}}{96m_- m_+} - \frac{25C_A C_F^2 Y_s^2 L_{m_2}}{4m_- m_+} - \frac{25C_F Y_s^2 L_{m_2}}{8m_- m_+} - \\
& \frac{7C_A C_F^2 Y_s^2 L_{m_2}}{4\epsilon m_- m_+} - \frac{7C_F Y_s^2 L_{m_2}}{8\epsilon m_- m_+} + \frac{C_A C_F^2 Y_s^2 L_{m_2}}{4\epsilon^2 m_- m_+} + \frac{C_F Y_s^2 L_{m_2}}{8\epsilon^2 m_- m_+} + \frac{C_A C_F^2 \pi^2 Y_s^2 L_{m_2}}{24m_- m_+} + \frac{C_F \pi^2 Y_s^2 L_{m_2}}{48m_- m_+} - \\
& \frac{7C_A^2 C_F m_2^2 Y_s^2 L_{m_1} L_{m_2}}{16m_-^2 m_+^2} - \frac{C_A Y_s^2 \beta_0 L_{m_2}}{8m_- m_+} - \frac{C_A Y_s^2 \beta_0 L_{m_2}}{8\epsilon m_- m_+} - \frac{C_A Y_s^2 \beta_0 L_{m_2}}{8\epsilon^2 m_- m_+} - \frac{C_A \pi^2 Y_s^2 \beta_0 L_{m_2}}{96m_- m_+} + \frac{17C_A^2 C_F m_2^2 Y_s^2 L_{m_1} L_{m_2}}{8m_-^2 m_+^2} + \\
& \frac{7C_A^2 C_F m_2^2 Y_s^2 L_{m_1} L_{m_2}}{8\epsilon m_-^2 m_+^2} + \frac{C_A C_F^2 Y_s^2}{2m_+^2} + \frac{C_F Y_s^2}{4m_+^2} + \frac{C_A C_F^2 Y_s^2}{4\epsilon m_+^2} + \frac{C_F Y_s^2}{8\epsilon m_+^2} + \frac{C_A C_F^2 Y_s^2}{2\epsilon^2 m_+^2} + \frac{C_F Y_s^2}{4\epsilon^2 m_+^2} + \frac{C_A C_F^2 \pi^2 Y_s^2}{12m_+^2} + \frac{C_F \pi^2 Y_s^2}{24m_+^2} + \\
& \frac{C_A C_F^2 m_2^2 Y_s^2 L_{m_1}^2}{m_-^2 m_+^2 m_1^2} + \frac{C_F m_2^2 Y_s^2 L_{m_1}^2}{2m_-^2 m_+^2 m_1^2} + \frac{C_A m_2^2 Y_s^2 \beta_0 L_{m_1}^2}{16m_- m_+ m_1^2} - \frac{C_A Y_s^2 \beta_0}{4m_+^2} - \frac{C_A Y_s^2 \beta_0}{8\epsilon m_+^2} - \frac{C_A Y_s^2 \beta_0}{8\epsilon^2 m_+^2} - \frac{C_A \pi^2 Y_s^2 \beta_0}{96m_+^2} + \\
& \frac{C_A C_F^2 m_2^2 Y_s^2 L_{m_1}}{2m_- m_+ m_1^2} + \frac{C_F m_2^2 Y_s^2 L_{m_1}}{4m_- m_+ m_1^2} + \frac{C_A C_F^2 m_2^2 Y_s^2 L_{m_1}}{\epsilon m_- m_+ m_1^2} + \frac{C_F m_2^2 Y_s^2 L_{m_1}}{2\epsilon m_- m_+ m_1^2} - \frac{C_A m_2^2 Y_s^2 \beta_0 L_{m_1}}{8m_- m_+ m_1^2} - \frac{C_A m_2^2 Y_s^2 \beta_0 L_{m_1}}{8\epsilon m_- m_+ m_1^2} + \\
& \frac{m_2^2 Y_s^4 C_A^2 L_{m_1}}{16m_+^2 m_- m_+} - \frac{m_2^2 Y_s^4 C_A^2 L_{m_2}}{16m_+^2 m_- m_+} - \frac{m_2^2 Y_s^4 C_A^2 L_{m_1}^2}{16m_+^2 m_- m_+^2} - \frac{m_2^2 Y_s^4 C_A^2 L_{m_2}^2}{16m_+^2 m_- m_+^2} - \frac{Y_s^4 C_A^2 L_{m_1}}{16\epsilon m_+^2 m_- m_+} - \frac{3Y_s^4 C_A^2 L_{m_1}}{16m_+^2 m_- m_+} - \frac{Y_s^4 C_A^2 L_{m_1} L_{m_2}}{16\epsilon m_-^2 m_+^2} + \\
& \frac{Y_s^4 C_A^2 L_{m_1} L_{m_2}^2}{32m_-^2 m_+^2} + \frac{Y_s^4 C_A^2 L_{m_1}^2 L_{m_2}}{32m_-^2 m_+^2} - \frac{Y_s^4 C_A^2 L_{m_1} L_{m_2}}{8m_-^2 m_+^2} + \frac{Y_s^4 C_A^2 L_{m_1}}{32\epsilon m_-^2 m_+^2} - \frac{Y_s^4 C_A^2 L_{m_1}^3}{32m_-^2 m_+^2} + \frac{Y_s^4 C_A^2 L_{m_1}^2}{8m_-^2 m_+^2} + \frac{Y_s^4 C_A^2 L_{m_2}}{16\epsilon m_-^2 m_+^2} + \\
& \frac{3Y_s^4 C_A^2 L_{m_2}}{16m_-^2 m_+^2} + \frac{Y_s^4 C_A^2 L_{m_2}^2}{32\epsilon m_-^2 m_+^2} - \frac{Y_s^4 C_A^2 L_{m_2}^3}{32m_-^2 m_+^2} + \frac{Y_s^4 C_A^2 L_{m_2}^2}{8m_-^2 m_+^2} + \frac{Y_s^4 C_A^2}{32\epsilon m_+^4} + \frac{3Y_s^4 C_A^2}{32m_+^4} - \frac{Y_s^4 C_A^2}{16\epsilon m_+^2 m_2^2} - \frac{Y_s^4 C_A^2}{32\epsilon^2 m_+^2 m_2^2} - \frac{Y_s^4 C_A^2}{8m_+^2 m_2^2} - \\
& \frac{\pi^2 Y_s^4 C_A^2}{192m_+^2 m_2^2} + \frac{Y_s^4 C_A^2}{32\epsilon m_+^4} + \frac{3Y_s^4 C_A^2}{32m_+^4}
\end{aligned} \tag{B.4}$$

$$\begin{aligned}
\Delta B_5^{(2)} = & -\frac{10C_A C_F^3 L_{m_1}^2 m_-^5}{(m_1^2 + m_2^2)^2 m_+} - \frac{5C_F^2 L_{m_1}^2 m_-^5}{(m_1^2 + m_2^2)^2 m_+} - \frac{12C_A C_F^3 L_{m_1}^2 m_-^5}{\epsilon (m_1^2 + m_2^2)^2 m_+} - \frac{6C_F^2 L_{m_1}^2 m_-^5}{\epsilon (m_1^2 + m_2^2)^2 m_+} - \frac{9C_A C_F \beta_0 L_{m_2}^2 m_-^5}{8(m_1^2 + m_2^2)^2 m_+} - \\
& \frac{C_A C_F \beta_0 L_{m_2}^2 m_-^5}{\epsilon (m_1^2 + m_2^2)^2 m_+} + \frac{C_A C_F^3 L_{m_2}^2 m_-^5}{(m_1^2 + m_2^2)^2 m_+} + \frac{C_F^2 L_{m_2}^2 m_-^5}{2(m_1^2 + m_2^2)^2 m_+} + \frac{12C_A C_F^3 L_{m_2}^2 m_-^5}{\epsilon (m_1^2 + m_2^2)^2 m_+} + \frac{6C_F^2 L_{m_2}^2 m_-^5}{\epsilon (m_1^2 + m_2^2)^2 m_+} - \\
& \frac{3C_A C_F^3 L_{m_1} L_{m_2} m_-^5}{(m_1^2 + m_2^2)^2 m_+} - \frac{3C_F^2 L_{m_1} L_{m_2} m_-^5}{2(m_1^2 + m_2^2)^2 m_+} - \frac{5C_A C_F^3 L_{m_1}^2 m_-^4}{2(m_1^2 + m_2^2)^2} - \frac{5C_F^2 L_{m_1}^2 m_-^4}{4(m_1^2 + m_2^2)^2} - \frac{12C_A C_F^3 L_{m_1}^2 m_-^4}{\epsilon (m_1^2 + m_2^2)^2} - \frac{6C_F^2 L_{m_1}^2 m_-^4}{\epsilon (m_1^2 + m_2^2)^2} + \\
& \frac{4C_A C_F \beta_0 L_{m_2}^2 m_-^4}{(m_1^2 + m_2^2)^2} + \frac{43C_A C_F^3 L_{m_2}^2 m_-^4}{2(m_1^2 + m_2^2)^2} + \frac{43C_F^2 L_{m_2}^2 m_-^4}{4(m_1^2 + m_2^2)^2} - \frac{12C_A C_F^3 L_{m_2}^2 m_-^4}{\epsilon (m_1^2 + m_2^2)^2} - \frac{6C_F^2 L_{m_2}^2 m_-^4}{\epsilon (m_1^2 + m_2^2)^2} + \frac{77C_A C_F^3 L_{m_1} L_{m_2} m_-^4}{(m_1^2 + m_2^2)^2} + \\
& \frac{77C_F^2 L_{m_1} L_{m_2} m_-^4}{2(m_1^2 + m_2^2)^2} + \frac{24C_A C_F^3 L_{m_1} L_{m_2} m_-^4}{\epsilon (m_1^2 + m_2^2)^2} + \frac{12C_F^2 L_{m_1} L_{m_2} m_-^4}{\epsilon (m_1^2 + m_2^2)^2} + \frac{66C_A C_F^3 m_+ L_{m_1}^2 m_-^3}{(m_1^2 + m_2^2)^2} + \frac{33C_F^2 m_+ L_{m_1}^2 m_-^3}{(m_1^2 + m_2^2)^2} - \\
& \frac{36C_A C_F^3 m_+ L_{m_1}^2 m_-^3}{\epsilon (m_1^2 + m_2^2)^2} - \frac{18C_F^2 m_+ L_{m_1}^2 m_-^3}{\epsilon (m_1^2 + m_2^2)^2} + \frac{9C_A C_F \beta_0 L_{m_1}^2 m_-^3}{16(m_1^2 + m_2^2) m_+} + \frac{C_A C_F \beta_0 L_{m_1}^2 m_-^3}{2\epsilon (m_1^2 + m_2^2) m_+} - \frac{105C_A C_F^3 m_+ L_{m_2}^2 m_-^3}{(m_1^2 + m_2^2)^2} - \\
& \frac{105C_F^2 m_+ L_{m_2}^2 m_-^3}{2(m_1^2 + m_2^2)^2} + \frac{36C_A C_F^3 m_+ L_{m_2}^2 m_-^3}{\epsilon (m_1^2 + m_2^2)^2} + \frac{18C_F^2 m_+ L_{m_2}^2 m_-^3}{\epsilon (m_1^2 + m_2^2)^2} - \frac{55C_A C_F m_+ \beta_0 L_{m_2}^2 m_-^3}{8(m_1^2 + m_2^2)^2} - \frac{3C_A C_F m_+ \beta_0 L_{m_2}^2 m_-^3}{\epsilon (m_1^2 + m_2^2)^2} - \\
& \frac{19C_A C_F \beta_0 L_{m_1} m_-^3}{16(m_1^2 + m_2^2) m_+} - \frac{9C_A C_F \beta_0 L_{m_1} m_-^3}{8\epsilon (m_1^2 + m_2^2) m_+} - \frac{C_A C_F \beta_0 L_{m_1} m_-^3}{\epsilon^2 (m_1^2 + m_2^2) m_+} - \frac{C_A C_F \pi^2 \beta_0 L_{m_1} m_-^3}{12(m_1^2 + m_2^2) m_+} + \frac{5C_A C_F^3 L_{m_1} m_-^3}{8(m_1^2 + m_2^2) m_+} + \\
& \frac{5C_F^2 L_{m_1} m_-^3}{16(m_1^2 + m_2^2) m_+} + \frac{23C_A C_F^3 L_{m_1} m_-^3}{4\epsilon (m_1^2 + m_2^2) m_+} + \frac{23C_F^2 L_{m_1} m_-^3}{8\epsilon (m_1^2 + m_2^2) m_+} + \frac{6C_A C_F^3 L_{m_1} m_-^3}{\epsilon^2 (m_1^2 + m_2^2) m_+} + \frac{3C_F^2 L_{m_1} m_-^3}{\epsilon^2 (m_1^2 + m_2^2) m_+} + \\
& \frac{C_A C_F^3 \pi^2 L_{m_1} m_-^3}{(m_1^2 + m_2^2) m_+} + \frac{C_F^2 \pi^2 L_{m_1} m_-^3}{2(m_1^2 + m_2^2) m_+} + \frac{19C_A C_F \beta_0 L_{m_2} m_-^3}{16(m_1^2 + m_2^2) m_+} + \frac{9C_A C_F \beta_0 L_{m_2} m_-^3}{8\epsilon (m_1^2 + m_2^2) m_+} + \frac{C_A C_F \beta_0 L_{m_2} m_-^3}{\epsilon^2 (m_1^2 + m_2^2) m_+} + \\
& \frac{C_A C_F \pi^2 \beta_0 L_{m_2} m_-^3}{12(m_1^2 + m_2^2) m_+} + \frac{3C_A C_F^3 m_+ L_{m_1} L_{m_2} m_-^3}{(m_1^2 + m_2^2)^2} + \frac{3C_F^2 m_+ L_{m_1} L_{m_2} m_-^3}{2(m_1^2 + m_2^2)^2} + \frac{155C_A C_F^3 L_{m_2} m_-^3}{8(m_1^2 + m_2^2) m_+} + \frac{155C_F^2 L_{m_2} m_-^3}{16(m_1^2 + m_2^2) m_+} +
\end{aligned}$$

$$\begin{aligned}
& \frac{C_A C_F^3 \mathcal{L} m_2 m_-^3}{4\epsilon (m_1^2 + m_2^2) m_+} + \frac{C_F^2 \mathcal{L} m_2 m_-^3}{8\epsilon (m_1^2 + m_2^2) m_+} - \frac{6C_A C_F^3 \mathcal{L} m_2 m_-^3}{\epsilon^2 (m_1^2 + m_2^2) m_+} - \frac{3C_F^2 \mathcal{L} m_2 m_-^3}{\epsilon^2 (m_1^2 + m_2^2) m_+} - \frac{C_A C_F^3 \pi^2 \mathcal{L} m_2 m_-^3}{(m_1^2 + m_2^2) m_+} - \frac{C_F^2 \pi^2 \mathcal{L} m_2 m_-^3}{2(m_1^2 + m_2^2) m_+} + \\
& \frac{37C_A C_F^3 m_+^2 \mathcal{L} m_1 m_-^2}{(m_1^2 + m_2^2)^2} + \frac{37C_F^2 m_+^2 \mathcal{L} m_1 m_-^2}{2(m_1^2 + m_2^2)^2} - \frac{24C_A C_F^3 m_+^2 \mathcal{L} m_1 m_-^2}{\epsilon (m_1^2 + m_2^2)^2} - \frac{12C_F^2 m_+^2 \mathcal{L} m_1 m_-^2}{\epsilon (m_1^2 + m_2^2)^2} + \frac{2C_A C_F \beta_0 \mathcal{L} m_1 m_-^2}{m_1^2 + m_2^2} + \\
& \frac{85C_A C_F^3 m_+^2 \mathcal{L} m_2 m_-^2}{(m_1^2 + m_2^2)^2} + \frac{85C_F^2 m_+^2 \mathcal{L} m_2 m_-^2}{2(m_1^2 + m_2^2)^2} - \frac{24C_A C_F^3 m_+^2 \mathcal{L} m_2 m_-^2}{\epsilon (m_1^2 + m_2^2)^2} - \frac{12C_F^2 m_+^2 \mathcal{L} m_2 m_-^2}{\epsilon (m_1^2 + m_2^2)^2} + \frac{8C_A C_F m_+^2 \beta_0 \mathcal{L} m_2 m_-^2}{(m_1^2 + m_2^2)^2} - \\
& \frac{9C_A C_F \beta_0 \mathcal{L} m_1 m_-^2}{4(m_1^2 + m_2^2)} - \frac{4C_A C_F \beta_0 \mathcal{L} m_1 m_-^2}{\epsilon (m_1^2 + m_2^2)} - \frac{111C_A C_F^3 \mathcal{L} m_1 m_-^2}{m_1^2 + m_2^2} - \frac{111C_F^2 \mathcal{L} m_1 m_-^2}{2(m_1^2 + m_2^2)} - \frac{18C_A C_F^3 \mathcal{L} m_1 m_-^2}{\epsilon (m_1^2 + m_2^2)} - \frac{9C_F^2 \mathcal{L} m_1 m_-^2}{\epsilon (m_1^2 + m_2^2)} - \\
& \frac{9C_A C_F \beta_0 \mathcal{L} m_2 m_-^2}{4(m_1^2 + m_2^2)} - \frac{4C_A C_F \beta_0 \mathcal{L} m_2 m_-^2}{\epsilon (m_1^2 + m_2^2)} + \frac{70C_A C_F^3 m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 m_-^2}{(m_1^2 + m_2^2)^2} + \frac{35C_F^2 m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 m_-^2}{(m_1^2 + m_2^2)^2} + \\
& \frac{48C_A C_F^3 m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 m_-^2}{\epsilon (m_1^2 + m_2^2)^2} + \frac{24C_F^2 m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 m_-^2}{\epsilon (m_1^2 + m_2^2)^2} - \frac{151C_A C_F^3 \mathcal{L} m_2 m_-^2}{m_1^2 + m_2^2} - \frac{151C_F^2 \mathcal{L} m_2 m_-^2}{2(m_1^2 + m_2^2)} - \frac{30C_A C_F^3 \mathcal{L} m_2 m_-^2}{\epsilon (m_1^2 + m_2^2)} - \\
& \frac{15C_F^2 \mathcal{L} m_2 m_-^2}{\epsilon (m_1^2 + m_2^2)} - \frac{C_A C_F \beta_0 \mathcal{L} m_1 m_-}{12m_+} + \frac{2C_A C_F^3 \mathcal{L} m_1 m_-}{m_+} + \frac{C_F^2 \mathcal{L} m_1 m_-}{m_+} + \frac{C_A C_F \beta_0 \mathcal{L} m_2 m_-}{12m_+} + \frac{2C_A C_F^3 \mathcal{L} m_2 m_-}{m_+} - \\
& \frac{C_F^2 \mathcal{L} m_2 m_-}{m_+} + \frac{66C_A C_F^3 m_+^2 \mathcal{L} m_1 m_-}{(m_1^2 + m_2^2)^2} + \frac{33C_F^2 m_+^2 \mathcal{L} m_1 m_-}{(m_1^2 + m_2^2)^2} - \frac{36C_A C_F^3 m_+^2 \mathcal{L} m_1 m_-}{\epsilon (m_1^2 + m_2^2)^2} - \frac{18C_F^2 m_+^2 \mathcal{L} m_1 m_-}{\epsilon (m_1^2 + m_2^2)^2} + \\
& \frac{23C_A C_F m_+ \beta_0 \mathcal{L} m_1 m_-}{8(m_1^2 + m_2^2)} + \frac{C_A C_F m_+ \beta_0 \mathcal{L} m_1 m_-}{\epsilon (m_1^2 + m_2^2)} - \frac{105C_A C_F^3 m_+^2 \mathcal{L} m_2 m_-}{(m_1^2 + m_2^2)^2} - \frac{105C_F^2 m_+^2 \mathcal{L} m_2 m_-}{2(m_1^2 + m_2^2)^2} + \frac{36C_A C_F^3 m_+^2 \mathcal{L} m_2 m_-}{\epsilon (m_1^2 + m_2^2)^2} + \\
& \frac{18C_F^2 m_+^2 \mathcal{L} m_2 m_-}{\epsilon (m_1^2 + m_2^2)^2} - \frac{55C_A C_F m_+^2 \beta_0 \mathcal{L} m_2 m_-}{8(m_1^2 + m_2^2)^2} - \frac{3C_A C_F m_+^2 \beta_0 \mathcal{L} m_2 m_-}{\epsilon (m_1^2 + m_2^2)^2} - \frac{861C_A C_F^3 m_+ \mathcal{L} m_1 m_-}{4(m_1^2 + m_2^2)} - \frac{861C_F^2 m_+ \mathcal{L} m_1 m_-}{8(m_1^2 + m_2^2)} - \\
& \frac{79C_A C_F^3 m_+ \mathcal{L} m_1 m_-}{2\epsilon (m_1^2 + m_2^2)} - \frac{79C_F^2 m_+ \mathcal{L} m_1 m_-}{4\epsilon (m_1^2 + m_2^2)} + \frac{12C_A C_F^3 m_+ \mathcal{L} m_1 m_-}{\epsilon^2 (m_1^2 + m_2^2)} + \frac{6C_F^2 m_+ \mathcal{L} m_1 m_-}{\epsilon^2 (m_1^2 + m_2^2)} - \frac{81C_A C_F m_+ \beta_0 \mathcal{L} m_1 m_-}{8(m_1^2 + m_2^2)} - \\
& \frac{23C_A C_F m_+ \beta_0 \mathcal{L} m_1 m_-}{4\epsilon (m_1^2 + m_2^2)} - \frac{2C_A C_F m_+ \beta_0 \mathcal{L} m_1 m_-}{\epsilon^2 (m_1^2 + m_2^2)} - \frac{C_A C_F m_+ \pi^2 \beta_0 \mathcal{L} m_1 m_-}{6(m_1^2 + m_2^2)} + \frac{2C_A C_F^3 m_+ \pi^2 \mathcal{L} m_1 m_-}{m_1^2 + m_2^2} + \\
& \frac{C_F^2 m_+ \pi^2 \mathcal{L} m_1 m_-}{m_1^2 + m_2^2} + \frac{1021C_A C_F^3 m_+ \mathcal{L} m_2 m_-}{4(m_1^2 + m_2^2)} + \frac{1021C_F^2 m_+ \mathcal{L} m_2 m_-}{8(m_1^2 + m_2^2)} + \frac{103C_A C_F^3 m_+ \mathcal{L} m_2 m_-}{2\epsilon (m_1^2 + m_2^2)} + \frac{103C_F^2 m_+ \mathcal{L} m_2 m_-}{4\epsilon (m_1^2 + m_2^2)} - \\
& \frac{12C_A C_F^3 m_+ \mathcal{L} m_2 m_-}{\epsilon^2 (m_1^2 + m_2^2)} - \frac{6C_F^2 m_+ \mathcal{L} m_2 m_-}{\epsilon^2 (m_1^2 + m_2^2)} + \frac{81C_A C_F m_+ \beta_0 \mathcal{L} m_2 m_-}{8(m_1^2 + m_2^2)} + \frac{23C_A C_F m_+ \beta_0 \mathcal{L} m_2 m_-}{4\epsilon (m_1^2 + m_2^2)} + \frac{2C_A C_F m_+ \beta_0 \mathcal{L} m_2 m_-}{\epsilon^2 (m_1^2 + m_2^2)} + \\
& \frac{C_A C_F m_+ \pi^2 \beta_0 \mathcal{L} m_2 m_-}{6(m_1^2 + m_2^2)} + \frac{3C_A C_F^3 m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 m_-}{(m_1^2 + m_2^2)^2} + \frac{3C_F^2 m_+^2 \mathcal{L} m_1 \mathcal{L} m_2 m_-}{2(m_1^2 + m_2^2)^2} - \frac{2C_A C_F^3 m_+ \pi^2 \mathcal{L} m_2 m_-}{m_1^2 + m_2^2} - \\
& \frac{C_F^2 m_+ \pi^2 \mathcal{L} m_2 m_-}{m_1^2 + m_2^2} - \frac{25C_A C_F^3 m_-}{2m_+} - \frac{25C_F^2 m_-}{4m_+} - \frac{5C_A C_F^3 m_-}{\epsilon m_+} - \frac{5C_F^2 m_-}{2\epsilon m_+} - \frac{3C_A C_F^3 m_-}{2\epsilon^2 m_+} - \frac{3C_F^2 m_-}{4\epsilon^2 m_+} - \frac{C_A C_F^3 \pi^2 m_-}{4m_+} - \\
& \frac{C_F^2 \pi^2 m_-}{8m_+} + \frac{701}{4} C_A C_F^3 + 3C_A C_F^3 \mathcal{L} m_1 + \frac{3}{2} C_F^2 \mathcal{L} m_1 + 3C_A C_F^3 \mathcal{L} m_2 + \frac{3}{2} C_F^2 \mathcal{L} m_2 + \frac{701C_F^2}{8} - \frac{5C_A C_F^3 m_+^4 \mathcal{L} m_1}{2(m_1^2 + m_2^2)^2} - \\
& \frac{5C_F^2 m_+^4 \mathcal{L} m_1}{4(m_1^2 + m_2^2)^2} - \frac{12C_A C_F^3 m_+^4 \mathcal{L} m_1}{\epsilon (m_1^2 + m_2^2)^2} - \frac{6C_F^2 m_+^4 \mathcal{L} m_1}{\epsilon (m_1^2 + m_2^2)^2} + \frac{2C_A C_F m_+^2 \beta_0 \mathcal{L} m_1}{m_1^2 + m_2^2} + \frac{43C_A C_F^3 m_+^4 \mathcal{L} m_2}{2(m_1^2 + m_2^2)^2} + \frac{43C_F^2 m_+^4 \mathcal{L} m_2}{4(m_1^2 + m_2^2)^2} - \\
& \frac{12C_A C_F^3 m_+^4 \mathcal{L} m_2}{\epsilon (m_1^2 + m_2^2)^2} - \frac{6C_F^2 m_+^4 \mathcal{L} m_2}{\epsilon (m_1^2 + m_2^2)^2} + \frac{4C_A C_F m_+^2 \beta_0 \mathcal{L} m_2}{(m_1^2 + m_2^2)^2} - 3C_A C_F^3 \mathcal{L} m_1 \mathcal{L} m_2 - \frac{3}{2} C_F^2 \mathcal{L} m_1 \mathcal{L} m_2 + \frac{19}{8} C_A C_F \beta_0 + \frac{9C_A C_F \beta_0}{4\epsilon} + \\
& \frac{4C_A C_F \beta_0}{\epsilon^2} + \frac{1}{3} C_A C_F \pi^2 \beta_0 - \frac{111C_A C_F^3 m_+^2 \mathcal{L} m_1}{m_1^2 + m_2^2} - \frac{111C_F^2 m_+^2 \mathcal{L} m_1}{2(m_1^2 + m_2^2)} - \frac{18C_A C_F^3 m_+^2 \mathcal{L} m_1}{\epsilon (m_1^2 + m_2^2)} - \frac{9C_F^2 m_+^2 \mathcal{L} m_1}{\epsilon (m_1^2 + m_2^2)} - \\
& \frac{9C_A C_F m_+^2 \beta_0 \mathcal{L} m_1}{4(m_1^2 + m_2^2)} - \frac{4C_A C_F m_+^2 \beta_0 \mathcal{L} m_1}{\epsilon (m_1^2 + m_2^2)} - \frac{151C_A C_F^3 m_+^2 \mathcal{L} m_2}{m_1^2 + m_2^2} - \frac{151C_F^2 m_+^2 \mathcal{L} m_2}{2(m_1^2 + m_2^2)} - \frac{30C_A C_F^3 m_+^2 \mathcal{L} m_2}{\epsilon (m_1^2 + m_2^2)} - \\
& \frac{15C_F^2 m_+^2 \mathcal{L} m_2}{\epsilon (m_1^2 + m_2^2)} - 3C_A C_F^3 \mathcal{L} m_1 \mathcal{L} m_2 - \frac{3}{2} C_F^2 \mathcal{L} m_1 \mathcal{L} m_2 - \frac{9C_A C_F m_+^2 \beta_0 \mathcal{L} m_2}{4(m_1^2 + m_2^2)} - \frac{4C_A C_F m_+^2 \beta_0 \mathcal{L} m_2}{\epsilon (m_1^2 + m_2^2)} + \frac{77C_A C_F^3 m_+^4 \mathcal{L} m_1 \mathcal{L} m_2}{(m_1^2 + m_2^2)^2} + \\
& \frac{77C_F^2 m_+^4 \mathcal{L} m_1 \mathcal{L} m_2}{2(m_1^2 + m_2^2)^2} + \frac{24C_A C_F^3 m_+^4 \mathcal{L} m_1 \mathcal{L} m_2}{\epsilon (m_1^2 + m_2^2)^2} + \frac{12C_F^2 m_+^4 \mathcal{L} m_1 \mathcal{L} m_2}{\epsilon (m_1^2 + m_2^2)^2} + \frac{131C_A C_F^3}{2\epsilon} + \frac{131C_F^2}{4\epsilon} + \frac{12C_A C_F^3}{\epsilon^2} + \\
& \frac{6C_F^2}{\epsilon^2} + 2C_A C_F^3 \pi^2 + C_F^2 \pi^2 + \frac{2C_A C_F^3 m_+ \mathcal{L} m_1}{m_-} + \frac{C_F^2 m_+ \mathcal{L} m_1}{m_-} - \frac{C_A C_F m_+ \beta_0 \mathcal{L} m_1}{12m_-} - \frac{2C_A C_F^3 m_+ \mathcal{L} m_2}{m_-} - \frac{C_F^2 m_+ \mathcal{L} m_2}{m_-} + \\
& \frac{C_A C_F m_+ \beta_0 \mathcal{L} m_2}{12m_-} - \frac{10C_A C_F^3 m_+^5 \mathcal{L} m_1}{(m_1^2 + m_2^2)^2 m_-} - \frac{5C_F^2 m_+^5 \mathcal{L} m_1}{(m_1^2 + m_2^2)^2 m_-} - \frac{12C_A C_F^3 m_+^5 \mathcal{L} m_1}{\epsilon (m_1^2 + m_2^2)^2 m_-} - \frac{6C_F^2 m_+^5 \mathcal{L} m_1}{\epsilon (m_1^2 + m_2^2)^2 m_-} + \\
& \frac{9C_A C_F m_+^3 \beta_0 \mathcal{L} m_1}{16(m_1^2 + m_2^2) m_-} + \frac{C_A C_F m_+^3 \beta_0 \mathcal{L} m_1}{2\epsilon (m_1^2 + m_2^2) m_-} + \frac{C_A C_F^3 m_+^5 \mathcal{L} m_2}{(m_1^2 + m_2^2)^2 m_-} + \frac{C_F^2 m_+^5 \mathcal{L} m_2}{2(m_1^2 + m_2^2)^2 m_-} + \frac{12C_A C_F^3 m_+^5 \mathcal{L} m_2}{\epsilon (m_1^2 + m_2^2)^2 m_-} + \\
& \frac{6C_F^2 m_+^5 \mathcal{L} m_2}{\epsilon (m_1^2 + m_2^2)^2 m_-} - \frac{9C_A C_F m_+^5 \beta_0 \mathcal{L} m_2}{8(m_1^2 + m_2^2)^2 m_-} - \frac{C_A C_F m_+^5 \beta_0 \mathcal{L} m_2}{\epsilon (m_1^2 + m_2^2)^2 m_-} - \frac{25C_A C_F^3 m_+}{2m_-} - \frac{25C_F^2 m_+}{4m_-} - \frac{5C_A C_F^3 m_+}{\epsilon m_-} - \frac{5C_F^2 m_+}{2\epsilon m_-} -
\end{aligned}$$



$$\begin{aligned}
& \frac{3C_A C_F^3 m_+}{2\epsilon^2 m_-} - \frac{3C_F^2 m_+}{4\epsilon^2 m_-} + \frac{5C_A C_F^3 m_+^3 \mathcal{L} m_1}{8(m_1^2 + m_2^2) m_-} + \frac{5C_F^2 m_+^3 \mathcal{L} m_1}{16(m_1^2 + m_2^2) m_-} + \frac{23C_A C_F^3 m_+^3 \mathcal{L} m_1}{4\epsilon(m_1^2 + m_2^2) m_-} + \frac{23C_F^2 m_+^3 \mathcal{L} m_1}{8\epsilon(m_1^2 + m_2^2) m_-} + \\
& \frac{6C_A C_F^3 m_+^3 \mathcal{L} m_1}{\epsilon^2(m_1^2 + m_2^2) m_-} + \frac{3C_F^2 m_+^3 \mathcal{L} m_1}{\epsilon^2(m_1^2 + m_2^2) m_-} - \frac{19C_A C_F m_+^3 \beta_0 \mathcal{L} m_1}{16(m_1^2 + m_2^2) m_-} - \frac{9C_A C_F m_+^3 \beta_0 \mathcal{L} m_1}{8\epsilon(m_1^2 + m_2^2) m_-} - \frac{C_A C_F m_+^3 \beta_0 \mathcal{L} m_1}{\epsilon^2(m_1^2 + m_2^2) m_-} - \\
& \frac{C_A C_F m_+^3 \pi^2 \beta_0 \mathcal{L} m_1}{12(m_1^2 + m_2^2) m_-} + \frac{C_A C_F^3 m_+^3 \pi^2 \mathcal{L} m_1}{(m_1^2 + m_2^2) m_-} + \frac{C_F^2 m_+^3 \pi^2 \mathcal{L} m_1}{2(m_1^2 + m_2^2) m_-} + \frac{155C_A C_F^3 m_+^3 \mathcal{L} m_2}{8(m_1^2 + m_2^2) m_-} + \frac{155C_F^2 m_+^3 \mathcal{L} m_2}{16(m_1^2 + m_2^2) m_-} + \\
& \frac{C_A C_F^3 m_+^3 \mathcal{L} m_2}{4\epsilon(m_1^2 + m_2^2) m_-} + \frac{C_F^2 m_+^3 \mathcal{L} m_2}{8\epsilon(m_1^2 + m_2^2) m_-} - \frac{6C_A C_F^3 m_+^3 \mathcal{L} m_2}{\epsilon^2(m_1^2 + m_2^2) m_-} - \frac{3C_F^2 m_+^3 \mathcal{L} m_2}{\epsilon^2(m_1^2 + m_2^2) m_-} + \frac{19C_A C_F m_+^3 \beta_0 \mathcal{L} m_2}{16(m_1^2 + m_2^2) m_-} + \\
& \frac{9C_A C_F m_+^3 \beta_0 \mathcal{L} m_2}{8\epsilon(m_1^2 + m_2^2) m_-} + \frac{C_A C_F m_+^3 \beta_0 \mathcal{L} m_2}{\epsilon^2(m_1^2 + m_2^2) m_-} + \frac{C_A C_F m_+^3 \pi^2 \beta_0 \mathcal{L} m_2}{12(m_1^2 + m_2^2) m_-} - \frac{3C_A C_F^3 m_+^5 \mathcal{L} m_1 \mathcal{L} m_2}{(m_1^2 + m_2^2)^2 m_-} - \frac{3C_F^2 m_+^5 \mathcal{L} m_1 \mathcal{L} m_2}{2(m_1^2 + m_2^2)^2 m_-} - \\
& \frac{C_A C_F^3 m_+^3 \pi^2 \mathcal{L} m_2}{(m_1^2 + m_2^2) m_-} - \frac{C_F^2 m_+^3 \pi^2 \mathcal{L} m_2}{2(m_1^2 + m_2^2) m_-} - \frac{C_A C_F^3 m_+ \pi^2}{4m_-} - \frac{C_F^2 m_+ \pi^2}{8m_-} + \frac{C_A C_F^2 Y_s^2 \mathcal{L} m_2^6 m_1^6}{m_2^2(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_F Y_s^2 \mathcal{L} m_2^6 m_1^6}{2m_2^2(m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{C_A Y_s^2 \beta_0 \mathcal{L} m_2^6 m_1^6}{16m_2^2(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A C_F^2 Y_s^2 \mathcal{L} m_2^6 m_1^6}{2m_2^2(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_F Y_s^2 \mathcal{L} m_2^6 m_1^6}{4m_2^2(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A C_F^2 Y_s^2 \mathcal{L} m_2^6 m_1^6}{\epsilon m_2^2(m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{C_F Y_s^2 \mathcal{L} m_2^6 m_1^6}{2\epsilon m_2^2(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A Y_s^2 \beta_0 \mathcal{L} m_2^6 m_1^6}{8\epsilon m_2^2(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A Y_s^2 \beta_0 \mathcal{L} m_2^6 m_1^6}{8\epsilon m_2^2(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A C_F^2 Y_s^2 m_1^4}{2m_2^2(m_1^2 + m_2^2)^2 m_- m_+} + \\
& \frac{C_F Y_s^2 m_1^4}{4m_2^2(m_1^2 + m_2^2) m_- m_+} + \frac{C_A C_F^2 Y_s^2 m_1^4}{4\epsilon m_2^2(m_1^2 + m_2^2) m_- m_+} + \frac{C_F Y_s^2 m_1^4}{8\epsilon m_2^2(m_1^2 + m_2^2) m_- m_+} + \frac{C_A C_F^2 Y_s^2 m_1^4}{2\epsilon^2 m_2^2(m_1^2 + m_2^2) m_- m_+} + \\
& \frac{C_F Y_s^2 m_1^4}{4\epsilon^2 m_2^2(m_1^2 + m_2^2) m_- m_+} + \frac{C_A C_F^2 \pi^2 Y_s^2 m_1^4}{12m_2^2(m_1^2 + m_2^2) m_- m_+} + \frac{C_F \pi^2 Y_s^2 m_1^4}{24m_2^2(m_1^2 + m_2^2) m_- m_+} - \frac{3C_A C_F^2 Y_s^2 \mathcal{L} m_1^4 m_1^4}{8(m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{3C_F Y_s^2 \mathcal{L} m_1^4 m_1^4}{16(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A C_F^2 Y_s^2 \mathcal{L} m_1^4 m_1^4}{\epsilon(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_F Y_s^2 \mathcal{L} m_1^4 m_1^4}{2\epsilon(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A Y_s^2 \beta_0 \mathcal{L} m_1^4 m_1^4}{8(m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{C_A Y_s^2 \beta_0 \mathcal{L} m_1^4 m_1^4}{16\epsilon(m_1^2 + m_2^2)^2 m_- m_+} + \frac{17C_A C_F^2 Y_s^2 \mathcal{L} m_2^4 m_1^4}{8(m_1^2 + m_2^2)^2 m_- m_+} + \frac{17C_F Y_s^2 \mathcal{L} m_2^4 m_1^4}{16(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A C_F^2 Y_s^2 \mathcal{L} m_2^4 m_1^4}{\epsilon(m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{C_F Y_s^2 \mathcal{L} m_2^4 m_1^4}{2\epsilon(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A Y_s^2 \beta_0 \mathcal{L} m_2^4 m_1^4}{16\epsilon(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A Y_s^2 \beta_0 m_1^4}{4m_2^2(m_1^2 + m_2^2) m_- m_+} - \frac{C_A Y_s^2 \beta_0 m_1^4}{8\epsilon m_2^2(m_1^2 + m_2^2) m_- m_+} - \\
& \frac{C_A Y_s^2 \beta_0 m_1^4}{8\epsilon^2 m_2^2(m_1^2 + m_2^2) m_- m_+} - \frac{C_A \pi^2 Y_s^2 \beta_0 m_1^4}{96m_2^2(m_1^2 + m_2^2) m_- m_+} + \frac{49C_A C_F^2 Y_s^2 \mathcal{L} m_1^4 m_1^4}{8(m_1^2 + m_2^2)^2 m_- m_+} + \frac{49C_F Y_s^2 \mathcal{L} m_1^4 m_1^4}{16(m_1^2 + m_2^2)^2 m_- m_+} + \\
& \frac{C_A C_F^2 Y_s^2 \mathcal{L} m_1^4 m_1^4}{4\epsilon(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_F Y_s^2 \mathcal{L} m_1^4 m_1^4}{8\epsilon(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A C_F^2 Y_s^2 \mathcal{L} m_1^4 m_1^4}{\epsilon^2(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_F Y_s^2 \mathcal{L} m_1^4 m_1^4}{2\epsilon^2(m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{C_A C_F^2 \pi^2 Y_s^2 \mathcal{L} m_1^4 m_1^4}{6(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_F \pi^2 Y_s^2 \mathcal{L} m_1^4 m_1^4}{12(m_1^2 + m_2^2)^2 m_- m_+} - \frac{57C_A C_F^2 Y_s^2 \mathcal{L} m_2^4 m_1^4}{8(m_1^2 + m_2^2)^2 m_- m_+} - \frac{57C_F Y_s^2 \mathcal{L} m_2^4 m_1^4}{16(m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{9C_A C_F^2 Y_s^2 \mathcal{L} m_2^4 m_1^4}{4\epsilon(m_1^2 + m_2^2)^2 m_- m_+} - \frac{9C_F Y_s^2 \mathcal{L} m_2^4 m_1^4}{8\epsilon(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A C_F^2 Y_s^2 \mathcal{L} m_2^4 m_1^4}{\epsilon^2(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_F Y_s^2 \mathcal{L} m_2^4 m_1^4}{2\epsilon^2(m_1^2 + m_2^2)^2 m_- m_+} + \\
& \frac{C_A C_F^2 \pi^2 Y_s^2 \mathcal{L} m_2^4 m_1^4}{6(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_F \pi^2 Y_s^2 \mathcal{L} m_2^4 m_1^4}{12(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A Y_s^2 \beta_0 \mathcal{L} m_2^4 m_1^4}{8(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A Y_s^2 \beta_0 \mathcal{L} m_2^4 m_1^4}{8\epsilon^2(m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{C_A \pi^2 Y_s^2 \beta_0 \mathcal{L} m_2^4 m_1^4}{96(m_1^2 + m_2^2)^2 m_- m_+} + \frac{15C_A C_F^2 Y_s^2 m_1^2}{4(m_1^2 + m_2^2) m_- m_+} + \frac{15C_F Y_s^2 m_1^2}{8(m_1^2 + m_2^2) m_- m_+} + \frac{C_A C_F^2 Y_s^2 m_1^2}{4\epsilon(m_1^2 + m_2^2) m_- m_+} + \frac{C_F Y_s^2 m_1^2}{8\epsilon(m_1^2 + m_2^2) m_- m_+} + \\
& \frac{C_A C_F^2 Y_s^2 m_1^2}{2\epsilon^2(m_1^2 + m_2^2) m_- m_+} + \frac{C_F Y_s^2 m_1^2}{4\epsilon^2(m_1^2 + m_2^2) m_- m_+} + \frac{C_A C_F^2 \pi^2 Y_s^2 m_1^2}{12(m_1^2 + m_2^2) m_- m_+} + \frac{C_F \pi^2 Y_s^2 m_1^2}{24(m_1^2 + m_2^2) m_- m_+} + \\
& \frac{9C_A C_F^2 m_2^2 Y_s^2 \mathcal{L} m_1^2 m_1^2}{4(m_1^2 + m_2^2)^2 m_- m_+} + \frac{9C_F m_2^2 Y_s^2 \mathcal{L} m_1^2 m_1^2}{8(m_1^2 + m_2^2)^2 m_- m_+} + \frac{2C_A C_F^2 m_2^2 Y_s^2 \mathcal{L} m_1^2 m_1^2}{\epsilon(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_F m_2^2 Y_s^2 \mathcal{L} m_1^2 m_1^2}{\epsilon(m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{3C_A m_2^2 Y_s^2 \beta_0 \mathcal{L} m_1^2 m_1^2}{16(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A m_2^2 Y_s^2 \beta_0 \mathcal{L} m_1^2 m_1^2}{8\epsilon(m_1^2 + m_2^2)^2 m_- m_+} - \frac{3C_A C_F^2 m_2^2 Y_s^2 \mathcal{L} m_2^2 m_1^2}{4(m_1^2 + m_2^2)^2 m_- m_+} - \frac{3C_F m_2^2 Y_s^2 \mathcal{L} m_2^2 m_1^2}{8(m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{2C_A C_F^2 \pi^2 Y_s^2 \mathcal{L} m_2^2 m_1^2}{\epsilon(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_F \pi^2 Y_s^2 \mathcal{L} m_2^2 m_1^2}{\epsilon(m_1^2 + m_2^2)^2 m_- m_+} + \frac{3C_A m_2^2 Y_s^2 \beta_0 \mathcal{L} m_2^2 m_1^2}{16(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A m_2^2 Y_s^2 \beta_0 \mathcal{L} m_2^2 m_1^2}{8\epsilon(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A Y_s^2 \beta_0 m_1^2}{2(m_1^2 + m_2^2) m_- m_+} + \\
& \frac{C_A Y_s^2 \beta_0 m_1^2}{8\epsilon(m_1^2 + m_2^2) m_- m_+} - \frac{C_A Y_s^2 \beta_0 m_1^2}{8\epsilon^2(m_1^2 + m_2^2) m_- m_+} - \frac{C_A \pi^2 Y_s^2 \beta_0 m_1^2}{96(m_1^2 + m_2^2) m_- m_+} - \frac{C_A C_F^2 m_2^2 Y_s^2 \mathcal{L} m_1^2 m_1^2}{2(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_F m_2^2 Y_s^2 \mathcal{L} m_1^2 m_1^2}{4(m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{5C_A C_F^2 m_2^2 Y_s^2 \mathcal{L} m_1^2 m_1^2}{2\epsilon(m_1^2 + m_2^2)^2 m_- m_+} - \frac{5C_F m_2^2 Y_s^2 \mathcal{L} m_1^2 m_1^2}{4\epsilon(m_1^2 + m_2^2)^2 m_- m_+} - \frac{2C_A C_F^2 m_2^2 Y_s^2 \mathcal{L} m_1^2 m_1^2}{\epsilon^2(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_F m_2^2 Y_s^2 \mathcal{L} m_1^2 m_1^2}{\epsilon^2(m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{C_A C_F^2 m_2^2 \pi^2 Y_s^2 \mathcal{L} m_1^2 m_1^2}{3(m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_F m_2^2 \pi^2 Y_s^2 \mathcal{L} m_1^2 m_1^2}{6(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A Y_s^2 \beta_0 \mathcal{L} m_1^2 m_1^2}{8(m_1^2 + m_2^2) m_- m_+} + \frac{C_A Y_s^2 \beta_0 \mathcal{L} m_1^2 m_1^2}{4\epsilon(m_1^2 + m_2^2) m_- m_+} + \frac{C_A Y_s^2 \beta_0 \mathcal{L} m_1^2 m_1^2}{8\epsilon^2(m_1^2 + m_2^2) m_- m_+} + \\
& \frac{C_A \pi^2 Y_s^2 \beta_0 \mathcal{L} m_1^2 m_1^2}{96(m_1^2 + m_2^2) m_- m_+} - \frac{11C_A C_F^2 m_2^2 Y_s^2 \mathcal{L} m_2^2 m_1^2}{2(m_1^2 + m_2^2)^2 m_- m_+} - \frac{11C_F m_2^2 Y_s^2 \mathcal{L} m_2^2 m_1^2}{4(m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A C_F^2 m_2^2 Y_s^2 \mathcal{L} m_2^2 m_1^2}{2\epsilon(m_1^2 + m_2^2)^2 m_- m_+} +
\end{aligned}$$

$$\begin{aligned}
& \frac{C_F m_2^2 Y_s^2 \mathcal{L} m_2 m_1^2}{4\epsilon (m_1^2 + m_2^2)^2 m_- m_+} + \frac{2C_A C_F^2 m_2^2 Y_s^2 \mathcal{L} m_2 m_1^2}{\epsilon^2 (m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_F m_2^2 Y_s^2 \mathcal{L} m_2 m_1^2}{\epsilon^2 (m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A C_F^2 m_2^2 \pi^2 Y_s^2 \mathcal{L} m_2 m_1^2}{3 (m_1^2 + m_2^2)^2 m_- m_+} + \\
& \frac{C_F m_2^2 \pi^2 Y_s^2 \mathcal{L} m_2 m_1^2}{6 (m_1^2 + m_2^2)^2 m_- m_+} - \frac{3C_A m_2^2 Y_s^2 \beta_0 \mathcal{L} m_2 m_1^2}{8 (m_1^2 + m_2^2)^2 m_- m_+} - \frac{3C_A m_2^2 Y_s^2 \beta_0 \mathcal{L} m_2 m_1^2}{8\epsilon (m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A m_2^2 Y_s^2 \beta_0 \mathcal{L} m_2 m_1^2}{4\epsilon^2 (m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{C_A m_2^2 \pi^2 Y_s^2 \beta_0 \mathcal{L} m_2 m_1^2}{48 (m_1^2 + m_2^2)^2 m_- m_+} - \frac{2C_A C_F^2 Y_s^2 \mathcal{L} m_1^3}{3m_- m_+} - \frac{C_F Y_s^2 \mathcal{L} m_1^3}{3m_- m_+} + \frac{C_A Y_s^2 \beta_0 \mathcal{L} m_1^3}{48m_- m_+} + \frac{2C_A C_F^2 Y_s^2 \mathcal{L} m_2^3}{3m_- m_+} + \frac{C_F Y_s^2 \mathcal{L} m_2^3}{3m_- m_+} - \\
& \frac{C_A Y_s^2 \beta_0 \mathcal{L} m_2^3}{48m_- m_+} + \frac{13C_A C_F^2 m_2^2 Y_s^2}{4 (m_1^2 + m_2^2) m_- m_+} + \frac{13C_F m_2^2 Y_s^2}{8 (m_1^2 + m_2^2) m_- m_+} + \frac{11C_A C_F^2 m_2^2 Y_s^2}{4\epsilon (m_1^2 + m_2^2) m_- m_+} + \frac{11C_F m_2^2 Y_s^2}{8\epsilon (m_1^2 + m_2^2) m_- m_+} + \\
& \frac{C_A C_F^2 m_2^2 Y_s^2}{2\epsilon^2 (m_1^2 + m_2^2) m_- m_+} + \frac{C_F m_2^2 Y_s^2}{4\epsilon^2 (m_1^2 + m_2^2) m_- m_+} + \frac{C_A C_F^2 m_2^2 \pi^2 Y_s^2}{12 (m_1^2 + m_2^2) m_- m_+} + \frac{C_F m_2^2 \pi^2 Y_s^2}{24 (m_1^2 + m_2^2) m_- m_+} - \\
& \frac{15C_A C_F^2 m_2^4 Y_s^2 \mathcal{L} m_1}{8 (m_1^2 + m_2^2)^2 m_- m_+} - \frac{15C_F m_2^4 Y_s^2 \mathcal{L} m_1}{16 (m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A C_F^2 m_2^4 Y_s^2 \mathcal{L} m_1}{\epsilon (m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_F m_2^4 Y_s^2 \mathcal{L} m_1}{2\epsilon (m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{C_A m_2^4 Y_s^2 \beta_0 \mathcal{L} m_1}{16\epsilon (m_1^2 + m_2^2)^2 m_- m_+} + \frac{13C_A C_F^2 m_2^4 Y_s^2 \mathcal{L} m_2}{8 (m_1^2 + m_2^2)^2 m_- m_+} + \frac{13C_F m_2^4 Y_s^2 \mathcal{L} m_2}{16 (m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A C_F^2 m_2^4 Y_s^2 \mathcal{L} m_2}{\epsilon (m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{C_F m_2^4 Y_s^2 \mathcal{L} m_2}{2\epsilon (m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A m_2^4 Y_s^2 \beta_0 \mathcal{L} m_2}{8 (m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A m_2^4 Y_s^2 \beta_0 \mathcal{L} m_2}{16\epsilon (m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A m_2^4 Y_s^2 \beta_0}{2 (m_1^2 + m_2^2) m_- m_+} - \\
& \frac{C_A m_2^2 Y_s^2 \beta_0}{8\epsilon (m_1^2 + m_2^2) m_- m_+} + \frac{C_A m_2^2 Y_s^2 \beta_0}{8\epsilon^2 (m_1^2 + m_2^2) m_- m_+} + \frac{C_A m_2^2 \pi^2 Y_s^2 \beta_0}{96 (m_1^2 + m_2^2) m_- m_+} + \frac{45C_A C_F^2 m_2^4 Y_s^2 \mathcal{L} m_1}{8 (m_1^2 + m_2^2)^2 m_- m_+} + \\
& \frac{45C_F m_2^4 Y_s^2 \mathcal{L} m_1}{16 (m_1^2 + m_2^2)^2 m_- m_+} + \frac{7C_A C_F^2 m_2^4 Y_s^2 \mathcal{L} m_1}{4\epsilon (m_1^2 + m_2^2)^2 m_- m_+} + \frac{7C_F m_2^4 Y_s^2 \mathcal{L} m_1}{8\epsilon (m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A C_F^2 m_2^4 Y_s^2 \mathcal{L} m_1}{\epsilon^2 (m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{C_F m_2^4 Y_s^2 \mathcal{L} m_1}{2\epsilon^2 (m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A C_F^2 m_2^4 \pi^2 Y_s^2 \mathcal{L} m_1}{6 (m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_F m_2^4 \pi^2 Y_s^2 \mathcal{L} m_1}{12 (m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A m_2^2 Y_s^2 \beta_0 \mathcal{L} m_1}{4 (m_1^2 + m_2^2) m_- m_+} + \\
& \frac{C_A m_2^2 Y_s^2 \beta_0 \mathcal{L} m_1}{8\epsilon (m_1^2 + m_2^2) m_- m_+} + \frac{C_A m_2^2 Y_s^2 \beta_0 \mathcal{L} m_1}{8\epsilon^2 (m_1^2 + m_2^2) m_- m_+} + \frac{C_A m_2^2 \pi^2 Y_s^2 \beta_0 \mathcal{L} m_1}{96 (m_1^2 + m_2^2) m_- m_+} - \frac{85C_A C_F^2 m_2^4 Y_s^2 \mathcal{L} m_2}{8 (m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{85C_F m_2^4 Y_s^2 \mathcal{L} m_2}{16 (m_1^2 + m_2^2)^2 m_- m_+} - \frac{7C_A C_F^2 m_2^4 Y_s^2 \mathcal{L} m_2}{4\epsilon (m_1^2 + m_2^2)^2 m_- m_+} - \frac{7C_F m_2^4 Y_s^2 \mathcal{L} m_2}{8\epsilon (m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A C_F^2 m_2^4 Y_s^2 \mathcal{L} m_2}{\epsilon^2 (m_1^2 + m_2^2)^2 m_- m_+} + \\
& \frac{C_F m_2^4 Y_s^2 \mathcal{L} m_2}{2\epsilon^2 (m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A C_F^2 m_2^4 \pi^2 Y_s^2 \mathcal{L} m_2}{6 (m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_F m_2^4 \pi^2 Y_s^2 \mathcal{L} m_2}{12 (m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A m_2^4 Y_s^2 \beta_0 \mathcal{L} m_2}{8 (m_1^2 + m_2^2)^2 m_- m_+} - \\
& \frac{C_A m_2^4 Y_s^2 \beta_0 \mathcal{L} m_2}{4\epsilon (m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A m_2^4 Y_s^2 \beta_0 \mathcal{L} m_2}{8\epsilon^2 (m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A m_2^4 \pi^2 Y_s^2 \beta_0 \mathcal{L} m_2}{96 (m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A^2 C_F Y_s^2 \mathcal{L} m_1 \mathcal{L} m_2}{8m_- m_+} - \\
& \frac{C_A C_F^2 m_2^4 Y_s^2}{2 (m_1^2 + m_2^2) m_- m_+ m_1^2} - \frac{C_F m_2^4 Y_s^2}{4 (m_1^2 + m_2^2) m_- m_+ m_1^2} - \frac{C_A C_F^2 m_2^4 Y_s^2}{4\epsilon (m_1^2 + m_2^2) m_- m_+ m_1^2} - \frac{C_F m_2^4 Y_s^2}{8\epsilon (m_1^2 + m_2^2) m_- m_+ m_1^2} - \\
& \frac{C_A C_F^2 m_2^4 Y_s^2}{2\epsilon^2 (m_1^2 + m_2^2) m_- m_+ m_1^2} - \frac{C_F m_2^4 Y_s^2}{4\epsilon^2 (m_1^2 + m_2^2) m_- m_+ m_1^2} - \frac{C_A C_F^2 m_2^4 \pi^2 Y_s^2}{12 (m_1^2 + m_2^2) m_- m_+ m_1^2} - \frac{C_F m_2^4 \pi^2 Y_s^2}{24 (m_1^2 + m_2^2) m_- m_+ m_1^2} + \\
& \frac{C_A C_F^2 m_2^6 Y_s^2 \mathcal{L} m_1}{(m_1^2 + m_2^2)^2 m_- m_+ m_1^2} - \frac{C_F m_2^6 Y_s^2 \mathcal{L} m_1}{2 (m_1^2 + m_2^2)^2 m_- m_+ m_1^2} + \frac{C_A m_2^6 Y_s^2 \beta_0 \mathcal{L} m_1}{16 (m_1^2 + m_2^2)^2 m_- m_+ m_1^2} + \frac{C_A m_2^4 Y_s^2 \beta_0}{4 (m_1^2 + m_2^2) m_- m_+ m_1^2} + \\
& \frac{C_A m_2^4 Y_s^2 \beta_0}{8\epsilon (m_1^2 + m_2^2) m_- m_+ m_1^2} + \frac{C_A m_2^4 Y_s^2 \beta_0}{8\epsilon^2 (m_1^2 + m_2^2) m_- m_+ m_1^2} + \frac{C_A m_2^4 \pi^2 Y_s^2 \beta_0}{96 (m_1^2 + m_2^2) m_- m_+ m_1^2} + \frac{C_A C_F^2 m_2^6 Y_s^2 \mathcal{L} m_1}{2 (m_1^2 + m_2^2)^2 m_- m_+ m_1^2} + \\
& \frac{C_F m_2^6 Y_s^2 \mathcal{L} m_1}{4 (m_1^2 + m_2^2)^2 m_- m_+ m_1^2} + \frac{C_A C_F^2 m_2^6 Y_s^2 \mathcal{L} m_1}{\epsilon (m_1^2 + m_2^2)^2 m_- m_+ m_1^2} + \frac{C_F m_2^6 Y_s^2 \mathcal{L} m_1}{2\epsilon (m_1^2 + m_2^2)^2 m_- m_+ m_1^2} - \frac{C_A m_2^4 Y_s^2 \beta_0 \mathcal{L} m_1}{8 (m_1^2 + m_2^2) m_- m_+ m_1^2} - \\
& \frac{C_A m_2^4 Y_s^2 \beta_0 \mathcal{L} m_1}{8\epsilon (m_1^2 + m_2^2) m_- m_+ m_1^2} - \frac{C_A^2 Y_s^4 \mathcal{L} m_2 m_1^6}{16m_2^4 (m_1^2 + m_2^2)^2 m_- m_+} + \frac{3C_A^2 Y_s^4 m_1^4}{32m_2^4 (m_1^2 + m_2^2)^2} + \frac{C_A^2 Y_s^4 m_1^4}{32\epsilon m_2^4 (m_1^2 + m_2^2)^2} + \\
& \frac{C_A^2 Y_s^4 \mathcal{L} m_2 m_1^4}{8m_2^2 (m_1^2 + m_2^2)^2 m_- m_+} + \frac{C_A^2 Y_s^4 \mathcal{L} m_2 m_1^4}{16\epsilon m_2^2 (m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A^2 Y_s^4 m_1^2}{16m_2^2 (m_1^2 + m_2^2)^2} - \frac{C_A^2 Y_s^4 m_1^2}{32\epsilon m_2^2 (m_1^2 + m_2^2)^2} - \frac{C_A^2 Y_s^4 m_1^2}{32\epsilon^2 m_2^2 (m_1^2 + m_2^2)^2} - \\
& \frac{C_A^2 \pi^2 Y_s^4 m_1^2}{192m_2^2 (m_1^2 + m_2^2)^2} - \frac{C_A^2 Y_s^4 \mathcal{L} m_2 m_1^2}{16m_2^2 (m_1^2 + m_2^2) m_- m_+} + \frac{C_A^2 Y_s^4 \mathcal{L} m_1 m_1^2}{16 (m_1^2 + m_2^2)^2 m_- m_+} + \frac{3C_A^2 Y_s^4 \mathcal{L} m_2 m_1^2}{8 (m_1^2 + m_2^2)^2 m_- m_+} + \\
& \frac{C_A^2 Y_s^4 \mathcal{L} m_2 m_1^2}{16\epsilon (m_1^2 + m_2^2)^2 m_- m_+} - \frac{15C_A^2 Y_s^4}{16 (m_1^2 + m_2^2)^2} - \frac{C_A^2 Y_s^4}{4\epsilon (m_1^2 + m_2^2)^2} - \frac{C_A^2 Y_s^4}{16\epsilon^2 (m_1^2 + m_2^2)^2} - \frac{C_A^2 \pi^2 Y_s^4}{96 (m_1^2 + m_2^2)^2} - \\
& \frac{3C_A^2 m_2^2 Y_s^4 \mathcal{L} m_1}{8 (m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A^2 m_2^2 Y_s^4 \mathcal{L} m_1}{16\epsilon (m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A^2 m_2^2 Y_s^4 \mathcal{L} m_2}{16 (m_1^2 + m_2^2)^2 m_- m_+} - \frac{C_A^2 m_2^2 Y_s^4}{16 (m_1^2 + m_2^2)^2 m_1^2} - \frac{C_A^2 m_2^2 Y_s^4}{32\epsilon (m_1^2 + m_2^2)^2 m_1^2} - \\
& \frac{C_A^2 m_2^2 Y_s^4}{32\epsilon^2 (m_1^2 + m_2^2)^2 m_1^2} - \frac{C_A^2 m_2^2 \pi^2 Y_s^4}{192 (m_1^2 + m_2^2)^2 m_1^2} + \frac{C_A^2 m_2^2 Y_s^4 \mathcal{L} m_1}{16 (m_1^2 + m_2^2) m_- m_+ m_1^2} - \frac{C_A^2 m_2^4 Y_s^4 \mathcal{L} m_1}{8 (m_1^2 + m_2^2)^2 m_- m_+ m_1^2} - \\
& \frac{C_A^2 m_2^4 Y_s^4 \mathcal{L} m_1}{16\epsilon (m_1^2 + m_2^2)^2 m_- m_+ m_1^2} + \frac{3C_A^2 m_2^4 Y_s^4}{32 (m_1^2 + m_2^2)^2 m_1^4} + \frac{C_A^2 m_2^4 Y_s^4}{32\epsilon (m_1^2 + m_2^2)^2 m_1^4} + \frac{C_A^2 m_2^6 Y_s^4 \mathcal{L} m_1}{16 (m_1^2 + m_2^2)^2 m_- m_+ m_1^4}
\end{aligned} \tag{B.5}$$

$$\begin{aligned}
\Delta B_6^{(2)} = & -\frac{7C_A m_+^2 \mathcal{L}_{m_1}^3 C_F^3}{4m_-^2} + 3C_A \mathcal{L}_{m_1}^3 C_F^3 + \frac{C_A m_+ \mathcal{L}_{m_1}^3 C_F^3}{4m_-} + \frac{C_A m_- \mathcal{L}_{m_1}^3 C_F^3}{4m_+} - \frac{7C_A m_-^2 \mathcal{L}_{m_1}^3 C_F^3}{4m_+^2} + \\
& \frac{C_A m_+^2 \mathcal{L}_{m_2}^3 C_F^3}{4m_-^2} + 3C_A \mathcal{L}_{m_2}^3 C_F^3 - \frac{7C_A m_+ \mathcal{L}_{m_2}^3 C_F^3}{4m_-} - \frac{7C_A m_- \mathcal{L}_{m_2}^3 C_F^3}{4m_+} + \frac{C_A m_-^2 \mathcal{L}_{m_2}^3 C_F^3}{4m_+^2} + \frac{33C_A m_+^2 \mathcal{L}_{m_1}^2 C_F^3}{8m_-^2} + \\
& \frac{9C_A m_+^2 \mathcal{L}_{m_1}^2 C_F^3}{4\epsilon m_-^2} - \frac{11}{2} C_A \mathcal{L}_{m_1}^2 C_F^3 + \frac{41C_A m_+ \mathcal{L}_{m_1}^2 C_F^3}{4m_-} - \frac{3C_A m_+ \mathcal{L}_{m_1}^2 C_F^3}{4\epsilon m_-} - \frac{3C_A \mathcal{L}_{m_1}^2 C_F^3}{\epsilon} + \frac{19C_A m_- \mathcal{L}_{m_1}^2 C_F^3}{2m_+} - \\
& \frac{3C_A m_- \mathcal{L}_{m_1}^2 C_F^3}{4\epsilon m_+} + \frac{63C_A m_-^2 \mathcal{L}_{m_1}^2 C_F^3}{8m_+^2} + \frac{9C_A m_-^2 \mathcal{L}_{m_1}^2 C_F^3}{4\epsilon m_+^2} - \frac{9C_A m_+^3 \mathcal{L}_{m_1}^2 C_F^3}{4m_+^3} - \frac{11C_A m_+^2 \mathcal{L}_{m_2}^2 C_F^3}{8m_-^2} - \frac{3C_A m_+^2 \mathcal{L}_{m_2}^2 C_F^3}{4\epsilon m_-^2} - \\
& \frac{23}{2} C_A \mathcal{L}_{m_2}^2 C_F^3 + \frac{9C_A m_+ \mathcal{L}_{m_2}^2 C_F^3}{4m_-} + \frac{9C_A m_+ \mathcal{L}_{m_2}^2 C_F^3}{4\epsilon m_-} + \frac{3C_A m_+^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 C_F^3}{4m_-^2} - 3C_A \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 C_F^3 + \frac{3C_A m_+ \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 C_F^3}{4m_-} + \\
& \frac{3C_A m_- \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 C_F^3}{4m_+} + \frac{3C_A m_-^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 C_F^3}{4m_+^2} - \frac{3C_A \mathcal{L}_{m_2}^2 C_F^3}{\epsilon} - \frac{3C_A m_- \mathcal{L}_{m_2}^2 C_F^3}{2m_+} + \frac{9C_A m_- \mathcal{L}_{m_2}^2 C_F^3}{4\epsilon m_+} - \frac{5C_A m_-^2 \mathcal{L}_{m_2}^2 C_F^3}{8m_+^2} - \\
& \frac{3C_A m_-^2 \mathcal{L}_{m_2}^2 C_F^3}{4\epsilon m_+^2} + \frac{3C_A m_+^3 \mathcal{L}_{m_2}^2 C_F^3}{4m_+^3} + 36C_A C_F^3 + \frac{31C_A m_+ C_F^3}{2m_-} + \frac{15C_A m_+ C_F^3}{2\epsilon m_-} + \frac{9C_A m_+ C_F^3}{2\epsilon^2 m_-} - \frac{23C_A m_+^2 \mathcal{L}_{m_1} C_F^3}{4m_-^2} - \\
& \frac{11C_A m_+^2 \mathcal{L}_{m_1} C_F^3}{4\epsilon m_-^2} - \frac{3C_A m_+^2 \mathcal{L}_{m_1} C_F^3}{2\epsilon^2 m_-^2} - 34C_A \mathcal{L}_{m_1} C_F^3 - \frac{29C_A m_+ \mathcal{L}_{m_1} C_F^3}{2m_-} - \frac{17C_A m_+ \mathcal{L}_{m_1} C_F^3}{2\epsilon m_-} + \frac{3C_A m_+ \mathcal{L}_{m_1} C_F^3}{2\epsilon^2 m_-} - \\
& \frac{3C_A \mathcal{L}_{m_1} C_F^3}{\epsilon} - \frac{7C_A m_- \mathcal{L}_{m_1} C_F^3}{m_+} - \frac{4C_A m_- \mathcal{L}_{m_1} C_F^3}{\epsilon m_+} + \frac{3C_A m_- \mathcal{L}_{m_1} C_F^3}{2\epsilon^2 m_+} - \frac{53C_A m_-^2 \mathcal{L}_{m_1} C_F^3}{4m_+^2} - \frac{29C_A m_-^2 \mathcal{L}_{m_1} C_F^3}{4\epsilon m_+^2} - \\
& \frac{3C_A m_-^2 \mathcal{L}_{m_1} C_F^3}{2\epsilon^2 m_+^2} + \frac{13C_A m_+^3 \mathcal{L}_{m_1} C_F^3}{2m_+^3} + \frac{3C_A m_+^3 \mathcal{L}_{m_1} C_F^3}{2\epsilon m_+^3} - \frac{C_A m_+^2 \pi^2 \mathcal{L}_{m_1} C_F^3}{4m_-^2} + \frac{C_A m_+ \pi^2 \mathcal{L}_{m_1} C_F^3}{4m_-} + \frac{C_A m_- \pi^2 \mathcal{L}_{m_1} C_F^3}{4m_+} - \\
& \frac{C_A m_-^2 \pi^2 \mathcal{L}_{m_1} C_F^3}{4m_+^2} + \frac{23C_A m_+^2 \mathcal{L}_{m_2} C_F^3}{4m_-^2} + \frac{11C_A m_+^2 \mathcal{L}_{m_2} C_F^3}{4\epsilon m_-^2} + \frac{3C_A m_+^2 \mathcal{L}_{m_2} C_F^3}{2\epsilon^2 m_-^2} + \\
& \frac{3C_A m_+^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{4m_-^2} - 3C_A \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3 + \frac{3C_A m_+ \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{4m_-} + \frac{3C_A m_- \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{4m_+} + \\
& \frac{3C_A m_-^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{4m_+^2} - 2C_A \mathcal{L}_{m_2} C_F^3 - \frac{C_A m_+ \mathcal{L}_{m_2} C_F^3}{2m_-} - \frac{C_A m_+ \mathcal{L}_{m_2} C_F^3}{2\epsilon m_-} - \frac{3C_A m_+ \mathcal{L}_{m_2} C_F^3}{2\epsilon^2 m_-} - \frac{11C_A m_+^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{4m_-^2} - \\
& \frac{3C_A m_+^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{2\epsilon m_-^2} + 17C_A \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3 - \frac{7C_A m_+ \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{2m_-} - \frac{3C_A m_+ \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{2\epsilon m_-} + \frac{6C_A \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{\epsilon} - \\
& \frac{11C_A m_- \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{m_+} - \frac{3C_A m_- \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{2\epsilon m_+} - \frac{5C_A m_-^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{4m_+^2} - \frac{3C_A m_-^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{2\epsilon m_+^2} + \frac{3C_A m_+^3 \mathcal{L}_{m_1} \mathcal{L}_{m_2} C_F^3}{2m_+^3} + \\
& \frac{3C_A \mathcal{L}_{m_2} C_F^3}{\epsilon} + \frac{32C_A m_- \mathcal{L}_{m_2} C_F^3}{m_+} + \frac{7C_A m_- \mathcal{L}_{m_2} C_F^3}{\epsilon m_+} - \frac{3C_A m_- \mathcal{L}_{m_2} C_F^3}{2\epsilon^2 m_+} - \frac{3C_A m_-^2 \mathcal{L}_{m_2} C_F^3}{4m_+^2} + \frac{5C_A m_-^2 \mathcal{L}_{m_2} C_F^3}{4\epsilon m_+^2} + \\
& \frac{3C_A m_-^2 \mathcal{L}_{m_2} C_F^3}{2\epsilon^2 m_+^2} - \frac{13C_A m_+^3 \mathcal{L}_{m_2} C_F^3}{2m_+^3} - \frac{3C_A m_+^3 \mathcal{L}_{m_2} C_F^3}{2\epsilon m_+^3} + \frac{C_A m_+^2 \pi^2 \mathcal{L}_{m_2} C_F^3}{4m_-^2} - \frac{C_A m_+ \pi^2 \mathcal{L}_{m_2} C_F^3}{4m_-} - \frac{C_A m_- \pi^2 \mathcal{L}_{m_2} C_F^3}{4m_+} - \\
& \frac{C_A m_-^2 \pi^2 \mathcal{L}_{m_2} C_F^3}{4m_+^2} + \frac{18C_A C_F^3}{\epsilon} - \frac{105C_A m_- C_F^3}{2m_+} - \frac{25C_A m_- C_F^3}{2\epsilon m_+} - \frac{3C_A m_- C_F^3}{2\epsilon^2 m_+} + \frac{15C_A m_-^2 C_F^3}{m_+^2} + \frac{7C_A m_-^2 C_F^3}{\epsilon m_+^2} + \\
& \frac{3C_A m_-^2 C_F^3}{\epsilon^2 m_+^2} + \frac{3C_A m_+ \pi^2 C_F^3}{4m_-} - \frac{C_A m_- \pi^2 C_F^3}{4m_+} + \frac{C_A m_-^2 \pi^2 C_F^3}{2m_+^2} - \frac{7m_+^2 \mathcal{L}_{m_1}^3 C_F^2}{8m_-^2} + \frac{m_+ \mathcal{L}_{m_1}^3 C_F^2}{8m_-} + \frac{m_- \mathcal{L}_{m_1}^3 C_F^2}{8m_+} - \\
& \frac{7m_+^2 \mathcal{L}_{m_1}^3 C_F^2}{8m_-^2} + \frac{3}{2} \mathcal{L}_{m_1}^3 C_F^2 + \frac{m_+^2 \mathcal{L}_{m_2}^3 C_F^2}{8m_-^2} - \frac{7m_+ \mathcal{L}_{m_2}^3 C_F^2}{8m_-} - \frac{7m_- \mathcal{L}_{m_2}^3 C_F^2}{8m_+} + \frac{m_-^2 \mathcal{L}_{m_2}^3 C_F^2}{8m_+^2} + \frac{3}{2} \mathcal{L}_{m_2}^3 C_F^2 + \frac{9m_+^2 \mathcal{L}_{m_1}^2 C_F^2}{8\epsilon m_-^2} + \\
& \frac{33m_+^2 \mathcal{L}_{m_1}^2 C_F^2}{16m_-^2} - \frac{3m_+ \mathcal{L}_{m_1}^2 C_F^2}{8\epsilon m_-} + \frac{41m_+ \mathcal{L}_{m_1}^2 C_F^2}{8m_-} - \frac{3\mathcal{L}_{m_1}^2 C_F^2}{2\epsilon} - \frac{3m_- \mathcal{L}_{m_1}^2 C_F^2}{8\epsilon m_+} + \frac{19m_- \mathcal{L}_{m_1}^2 C_F^2}{4m_+} + \frac{9m_-^2 \mathcal{L}_{m_1}^2 C_F^2}{8\epsilon m_+^2} + \\
& \frac{63m_-^2 \mathcal{L}_{m_1}^2 C_F^2}{16m_+^2} - \frac{9m_-^3 \mathcal{L}_{m_1}^2 C_F^2}{8m_+^3} - \frac{11}{4} \mathcal{L}_{m_1}^2 C_F^2 - \frac{3m_+^2 \mathcal{L}_{m_2}^2 C_F^2}{8\epsilon m_-^2} - \frac{11m_+^2 \mathcal{L}_{m_2}^2 C_F^2}{16m_-^2} + \frac{9m_+ \mathcal{L}_{m_2}^2 C_F^2}{8\epsilon m_-} + \frac{9m_+ \mathcal{L}_{m_2}^2 C_F^2}{8m_-} + \\
& \frac{3m_+^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 C_F^2}{8m_-^2} + \frac{3m_+ \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 C_F^2}{8m_-} + \frac{3m_- \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 C_F^2}{8m_+} + \frac{3m_-^2 \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 C_F^2}{8m_+^2} - \frac{3}{2} \mathcal{L}_{m_1} \mathcal{L}_{m_2}^2 C_F^2 - \frac{3\mathcal{L}_{m_2}^2 C_F^2}{2\epsilon} + \\
& \frac{9m_- \mathcal{L}_{m_2}^2 C_F^2}{8\epsilon m_+} - \frac{3m_- \mathcal{L}_{m_2}^2 C_F^2}{4m_+} - \frac{3m_-^2 \mathcal{L}_{m_2}^2 C_F^2}{8\epsilon m_+^2} - \frac{5m_-^2 \mathcal{L}_{m_2}^2 C_F^2}{16m_+^2} + \frac{3m_+^3 \mathcal{L}_{m_2}^2 C_F^2}{8m_+^3} - \frac{23}{4} \mathcal{L}_{m_2}^2 C_F^2 + \frac{15m_+ C_F^2}{4\epsilon m_-} + \frac{9m_+ C_F^2}{4\epsilon^2 m_-} + \\
& \frac{31m_+ C_F^2}{4m_-} - \frac{11m_+^2 \mathcal{L}_{m_1} C_F^2}{8\epsilon m_-^2} - \frac{3m_+^2 \mathcal{L}_{m_1} C_F^2}{4\epsilon^2 m_-^2} - \frac{23m_+^2 \mathcal{L}_{m_1} C_F^2}{8m_-^2} - \frac{17m_+ \mathcal{L}_{m_1} C_F^2}{4\epsilon m_-} + \frac{3m_+ \mathcal{L}_{m_1} C_F^2}{4\epsilon^2 m_-} - \frac{29m_+ \mathcal{L}_{m_1} C_F^2}{4m_-} - \\
& \frac{3\mathcal{L}_{m_1} C_F^2}{2\epsilon} - \frac{2m_- \mathcal{L}_{m_1} C_F^2}{\epsilon m_+} + \frac{3m_- \mathcal{L}_{m_1} C_F^2}{4\epsilon^2 m_+} - \frac{7m_- \mathcal{L}_{m_1} C_F^2}{2m_+} - \frac{29m_-^2 \mathcal{L}_{m_1} C_F^2}{8\epsilon m_+^2} - \frac{3m_-^2 \mathcal{L}_{m_1} C_F^2}{4\epsilon^2 m_+^2} - \frac{53m_-^2 \mathcal{L}_{m_1} C_F^2}{8m_+^2} + \\
& \frac{3m_-^3 \mathcal{L}_{m_1} C_F^2}{4\epsilon m_+^3} + \frac{13m_-^3 \mathcal{L}_{m_1} C_F^2}{4m_+^3} - \frac{m_+^2 \pi^2 \mathcal{L}_{m_1} C_F^2}{8m_-^2} + \frac{m_+ \pi^2 \mathcal{L}_{m_1} C_F^2}{8m_-} + \frac{m_- \pi^2 \mathcal{L}_{m_1} C_F^2}{8m_+} - \frac{m_-^2 \pi^2 \mathcal{L}_{m_1} C_F^2}{8m_+^2} - 17\mathcal{L}_{m_1} C_F^2 +
\end{aligned}$$

$$\begin{aligned}
& \frac{11m_+^2 \mathcal{L}m_2 C_F^2}{8\epsilon m_-^2} + \frac{3m_+^2 \mathcal{L}m_2 C_F^2}{4\epsilon^2 m_-^2} + \frac{23m_+^2 \mathcal{L}m_2 C_F^2}{8m_-^2} + \frac{3m_+^2 \mathcal{L}_m^2 \mathcal{L}m_2 C_F^2}{8m_-^2} + \frac{3m_+ \mathcal{L}_m^2 \mathcal{L}m_2 C_F^2}{8m_-} + \frac{3m_- \mathcal{L}_m^2 \mathcal{L}m_2 C_F^2}{8m_+} + \\
& \frac{3m_-^2 \mathcal{L}_m^2 \mathcal{L}m_2 C_F^2}{8m_+^2} - \frac{3}{2} \mathcal{L}_m^2 \mathcal{L}m_2 C_F^2 - \frac{m_+ \mathcal{L}m_2 C_F^2}{4\epsilon m_-} - \frac{3m_+ \mathcal{L}m_2 C_F^2}{4\epsilon^2 m_-} - \frac{m_+ \mathcal{L}m_2 C_F^2}{4m_-} - \frac{3m_+^2 \mathcal{L}m_1 \mathcal{L}m_2 C_F^2}{4\epsilon m_-^2} - \\
& \frac{11m_+^2 \mathcal{L}m_1 \mathcal{L}m_2 C_F^2}{8m_-^2} - \frac{3m_+ \mathcal{L}m_1 \mathcal{L}m_2 C_F^2}{4\epsilon m_-} - \frac{7m_+ \mathcal{L}m_1 \mathcal{L}m_2 C_F^2}{4m_-} + \frac{3\mathcal{L}m_1 \mathcal{L}m_2 C_F^2}{\epsilon} - \frac{3m_- \mathcal{L}m_1 \mathcal{L}m_2 C_F^2}{4\epsilon m_+} - \frac{11m_- \mathcal{L}m_1 \mathcal{L}m_2 C_F^2}{2m_+} - \\
& \frac{3m_-^2 \mathcal{L}m_1 \mathcal{L}m_2 C_F^2}{4\epsilon m_+^2} - \frac{5m_-^2 \mathcal{L}m_1 \mathcal{L}m_2 C_F^2}{8m_+^2} + \frac{3m_-^3 \mathcal{L}m_1 \mathcal{L}m_2 C_F^2}{4m_+^3} + \frac{17}{2} \mathcal{L}m_1 \mathcal{L}m_2 C_F^2 + \frac{3\mathcal{L}m_2 C_F^2}{2\epsilon} + \frac{7m_- \mathcal{L}m_2 C_F^2}{2\epsilon m_+} - \\
& \frac{3m_- \mathcal{L}m_2 C_F^2}{4\epsilon^2 m_+} + \frac{16m_- \mathcal{L}m_2 C_F^2}{m_+} + \frac{5m_-^2 \mathcal{L}m_2 C_F^2}{8\epsilon m_+^2} + \frac{3m_-^2 \mathcal{L}m_2 C_F^2}{4\epsilon^2 m_+^2} - \frac{3m_-^2 \mathcal{L}m_2 C_F^2}{8m_+^2} - \frac{3m_-^3 \mathcal{L}m_2 C_F^2}{4\epsilon m_+^3} - \frac{13m_-^3 \mathcal{L}m_2 C_F^2}{4m_+^3} + \\
& \frac{m_+^2 \pi^2 \mathcal{L}m_2 C_F^2}{8m_-^2} - \frac{m_+ \pi^2 \mathcal{L}m_2 C_F^2}{8m_-} - \frac{m_- \pi^2 \mathcal{L}m_2 C_F^2}{8m_+} + \frac{m_-^2 \pi^2 \mathcal{L}m_2 C_F^2}{8m_+^2} - \mathcal{L}m_2 C_F^2 + \frac{9C_F^2}{\epsilon} - \frac{25m_- C_F^2}{4\epsilon m_+} - \frac{3m_- C_F^2}{4\epsilon^2 m_+} - \\
& \frac{105m_- C_F^2}{4m_+} + \frac{7m_-^2 C_F^2}{2\epsilon m_+^2} + \frac{3m_-^2 C_F^2}{2\epsilon^2 m_+^2} + \frac{15m_-^2 C_F^2}{2m_+^2} + \frac{3m_+ \pi^2 C_F^2}{8m_-} - \frac{m_- \pi^2 C_F^2}{8m_+} + \frac{m_-^2 \pi^2 C_F^2}{4m_+^2} + 18C_F^2 - \frac{CA^{m_+ \beta_0} \mathcal{L}_m^3 C_F}{12m_-} - \\
& \frac{CA^{m_- \beta_0} \mathcal{L}_m^3 C_F}{12m_+} + \frac{CA^{m_+ \beta_0} \mathcal{L}_m^3 C_F}{12m_-} + \frac{CA^{m_- \beta_0} \mathcal{L}_m^3 C_F}{12m_+} + \frac{3}{8} CA \beta_0 \mathcal{L}_m^2 C_F + \frac{CA^{m_+ \beta_0} \mathcal{L}_m^2 C_F}{8m_-} + \frac{CA^{m_+ \beta_0} \mathcal{L}_m^2 C_F}{4\epsilon m_-} + \\
& \frac{3CA^{m_- \beta_0} \mathcal{L}_m^2 C_F}{8m_+} + \frac{CA^{m_- \beta_0} \mathcal{L}_m^2 C_F}{4\epsilon m_+} - \frac{CA^{m_- \beta_0} \mathcal{L}_m^2 C_F}{8m_+^2} + \frac{3}{8} CA \beta_0 \mathcal{L}_m^2 C_F - \frac{CA^{m_+ \beta_0} \mathcal{L}_m^2 C_F}{8m_-} - \frac{CA^{m_+ \beta_0} \mathcal{L}_m^2 C_F}{4\epsilon m_-} - \\
& \frac{3CA^{m_- \beta_0} \mathcal{L}_m^2 C_F}{8m_+} - \frac{CA^{m_- \beta_0} \mathcal{L}_m^2 C_F}{4\epsilon m_+} + \frac{CA^{m_- \beta_0} \mathcal{L}_m^2 C_F}{8m_+^2} - 3CA \beta_0 C_F - \frac{CA \beta_0 C_F}{\epsilon} - \frac{3CA^{m_- \beta_0} C_F}{m_+} - \frac{CA^{m_- \beta_0} C_F}{\epsilon m_+} + \\
& \frac{3CA \beta_0 C_F}{2\epsilon^2} + \frac{1}{8} CA \pi^2 \beta_0 C_F + \frac{1}{4} CA \beta_0 \mathcal{L}_m C_F - \frac{CA^{m_+ \beta_0} \mathcal{L}_m C_F}{4m_-} - \frac{CA^{m_+ \beta_0} \mathcal{L}_m C_F}{4\epsilon m_-} - \frac{CA^{m_+ \beta_0} \mathcal{L}_m C_F}{2\epsilon^2 m_-} - \\
& \frac{3CA \beta_0 \mathcal{L}_m C_F}{4\epsilon} - \frac{3CA^{m_- \beta_0} \mathcal{L}_m C_F}{4m_+} - \frac{3CA^{m_- \beta_0} \mathcal{L}_m C_F}{4\epsilon m_+} - \frac{CA^{m_- \beta_0} \mathcal{L}_m C_F}{2\epsilon^2 m_+} + \frac{3CA^{m_- \beta_0} \mathcal{L}_m C_F}{4m_+^2} + \\
& \frac{CA^{m_- \beta_0} \mathcal{L}_m C_F}{4\epsilon m_+^2} - \frac{CA^{m_+ \pi^2 \beta_0} \mathcal{L}_m C_F}{24m_-} - \frac{CA^{m_- \pi^2 \beta_0} \mathcal{L}_m C_F}{24m_+} + \frac{3}{4} CA \beta_0 \mathcal{L}_m C_F + \frac{CA^{m_+ \beta_0} \mathcal{L}_m C_F}{4m_-} + \\
& \frac{CA^{m_+ \beta_0} \mathcal{L}_m C_F}{4\epsilon m_-} + \frac{CA^{m_+ \beta_0} \mathcal{L}_m C_F}{2\epsilon^2 m_-} - \frac{3CA \beta_0 \mathcal{L}_m C_F}{4\epsilon} + \frac{7CA^{m_- \beta_0} \mathcal{L}_m C_F}{4m_+} + \frac{3CA^{m_- \beta_0} \mathcal{L}_m C_F}{4\epsilon m_+} + \frac{CA^{m_- \beta_0} \mathcal{L}_m C_F}{2\epsilon^2 m_+} - \\
& \frac{3CA^{m_- \beta_0} \mathcal{L}_m C_F}{4m_+^2} - \frac{CA^{m_- \beta_0} \mathcal{L}_m C_F}{4\epsilon m_+^2} + \frac{CA^{m_+ \pi^2 \beta_0} \mathcal{L}_m C_F}{24m_-} + \frac{CA^{m_- \pi^2 \beta_0} \mathcal{L}_m C_F}{24m_+} - \frac{3CA C_F^2 Y_f^2 \mathcal{L}_m^2 m_1^5}{m_-^2 m_+^3} - \\
& \frac{3C_F Y_f^2 \mathcal{L}_m^2 m_1^5}{2m_-^2 m_+^3} + \frac{25CA C_F^2 Y_f^2 \mathcal{L}_m m_1^5}{m_-^2 m_+^3} + \frac{25C_F Y_f^2 \mathcal{L}_m m_1^5}{2m_-^2 m_+^3} + \frac{3CA C_F^2 Y_f^2 \mathcal{L}_m m_1^5}{\epsilon m_-^2 m_+^3} + \frac{3C_F Y_f^2 \mathcal{L}_m m_1^5}{2\epsilon m_-^2 m_+^3} - \\
& \frac{3CA C_F^2 m_2 Y_f^2 \mathcal{L}_m^2 m_1^4}{m_-^2 m_+^3} - \frac{3C_F m_2 Y_f^2 \mathcal{L}_m^2 m_1^4}{2m_-^2 m_+^3} + \frac{28CA C_F^2 m_2 Y_f^2 \mathcal{L}_m m_1^4}{m_-^2 m_+^3} + \frac{14C_F m_2 Y_f^2 \mathcal{L}_m m_1^4}{m_-^2 m_+^3} + \frac{3CA C_F^2 m_2 Y_f^2 \mathcal{L}_m m_1^4}{\epsilon m_-^2 m_+^3} + \\
& \frac{3C_F m_2 Y_f^2 \mathcal{L}_m m_1^4}{2\epsilon m_-^2 m_+^3} - \frac{57CA C_F^2 Y_f^2 m_1^3}{2m_- m_+^2} - \frac{57C_F Y_f^2 m_1^3}{4m_- m_+^2} - \frac{25CA C_F^2 Y_f^2 m_1^3}{2\epsilon m_- m_+^2} - \frac{25C_F Y_f^2 m_1^3}{4\epsilon m_- m_+^2} - \frac{3CA C_F^2 Y_f^2 m_1^3}{2\epsilon^2 m_- m_+^2} - \frac{3C_F Y_f^2 m_1^3}{4\epsilon^2 m_- m_+^2} - \\
& \frac{CA C_F^2 \pi^2 Y_f^2 m_1^3}{4m_- m_+^2} - \frac{C_F \pi^2 Y_f^2 m_1^3}{8m_- m_+^2} - \frac{21CA C_F^2 m_2^2 Y_f^2 \mathcal{L}_m^2 m_1^3}{2m_-^2 m_+^3} - \frac{21C_F m_2^2 Y_f^2 \mathcal{L}_m^2 m_1^3}{4m_-^2 m_+^3} - \frac{9CA C_F^2 m_2^2 Y_f^2 \mathcal{L}_m^2 m_1^3}{2\epsilon m_-^2 m_+^3} - \\
& \frac{9C_F m_2^2 Y_f^2 \mathcal{L}_m^2 m_1^3}{4\epsilon m_-^2 m_+^3} + \frac{7C_A^2 C_F m_2^2 Y_f^2 \mathcal{L}_m^2 m_1^3}{4m_-^2 m_+^3} + \frac{3C_A^2 C_F m_2^2 Y_f^2 \mathcal{L}_m^2 m_1^3}{4\epsilon m_-^2 m_+^3} - \frac{17CA Y_f^2 \beta_0 m_1^3}{16m_- m_+^2} - \frac{9CA Y_f^2 \beta_0 m_1^3}{16\epsilon m_- m_+^2} - \frac{9CA Y_f^2 \beta_0 m_1^3}{16\epsilon^2 m_- m_+^2} - \\
& \frac{3CA \pi^2 Y_f^2 \beta_0 m_1^3}{64m_- m_+^2} - \frac{27CA C_F^2 m_2^2 Y_f^2 \mathcal{L}_m m_1^3}{m_-^2 m_+^3} - \frac{27C_F m_2^2 Y_f^2 \mathcal{L}_m m_1^3}{2m_-^2 m_+^3} + \frac{7CA C_F^2 m_2^2 Y_f^2 \mathcal{L}_m m_1^3}{\epsilon m_-^2 m_+^3} + \frac{7C_F m_2^2 Y_f^2 \mathcal{L}_m m_1^3}{2\epsilon m_-^2 m_+^3} + \\
& \frac{3CA C_F^2 m_2^2 Y_f^2 \mathcal{L}_m m_1^3}{\epsilon^2 m_-^2 m_+^3} + \frac{3C_F m_2^2 Y_f^2 \mathcal{L}_m m_1^3}{2\epsilon^2 m_-^2 m_+^3} + \frac{CA C_F^2 m_2^2 \pi^2 Y_f^2 \mathcal{L}_m m_1^3}{2m_-^2 m_+^3} + \frac{C_F m_2^2 \pi^2 Y_f^2 \mathcal{L}_m m_1^3}{4m_-^2 m_+^3} - \frac{6C_A^2 C_F m_2^2 Y_f^2 \mathcal{L}_m^2 m_1^3}{m_-^2 m_+^3} - \\
& \frac{7C_A^2 C_F m_2^2 Y_f^2 \mathcal{L}_m^2 m_1^3}{2\epsilon m_-^2 m_+^3} - \frac{3C_A^2 C_F m_2^2 Y_f^2 \mathcal{L}_m^2 m_1^3}{2\epsilon^2 m_-^2 m_+^3} - \frac{C_A^2 C_F m_2^2 \pi^2 Y_f^2 \mathcal{L}_m^2 m_1^3}{4m_-^2 m_+^3} + \frac{7C_A^2 C_F m_2^2 Y_f^2 \mathcal{L}_m^2 m_1^3}{2m_-^2 m_+^3} + \\
& \frac{3C_A^2 C_F m_2^2 Y_f^2 \mathcal{L}_m \mathcal{L}_m^2 m_1^3}{2\epsilon m_-^2 m_+^3} + \frac{7C_A^2 C_F m_2^2 Y_f^2 \mathcal{L}_m^3 m_1^2}{4m_-^2 m_+^2} - \frac{C_A^2 C_F m_2^2 Y_f^2 \mathcal{L}_m^3 m_1^2}{4m_-^2 m_+^2} - \frac{79CA C_F^2 m_2 Y_f^2 m_1^2}{2m_- m_+^2} - \frac{79C_F m_2 Y_f^2 m_1^2}{4m_- m_+^2} - \\
& \frac{14CA C_F^2 m_2 Y_f^2 m_1^2}{\epsilon m_- m_+^2} - \frac{7C_F m_2 Y_f^2 m_1^2}{\epsilon m_- m_+^2} - \frac{3CA C_F^2 m_2 Y_f^2 m_1^2}{2\epsilon^2 m_- m_+^2} - \frac{3C_F m_2 Y_f^2 m_1^2}{4\epsilon^2 m_- m_+^2} - \frac{CA C_F^2 m_2 \pi^2 Y_f^2 m_1^2}{4m_- m_+^2} - \frac{C_F m_2 \pi^2 Y_f^2 m_1^2}{8m_- m_+^2} - \\
& \frac{69CA C_F^2 m_2^3 Y_f^2 \mathcal{L}_m^2 m_1^2}{4m_-^2 m_+^3} - \frac{69C_F m_2^3 Y_f^2 \mathcal{L}_m^2 m_1^2}{8m_-^2 m_+^3} - \frac{9CA C_F^2 m_2^3 Y_f^2 \mathcal{L}_m^2 m_1^2}{2\epsilon m_-^2 m_+^3} - \frac{9C_F m_2^3 Y_f^2 \mathcal{L}_m^2 m_1^2}{4\epsilon m_-^2 m_+^3} + \frac{23C_A^2 C_F m_2^3 Y_f^2 \mathcal{L}_m^2 m_1^2}{8m_-^2 m_+^3} + \\
& \frac{3C_A^2 C_F m_2^3 Y_f^2 \mathcal{L}_m^2 m_1^2}{4\epsilon m_-^2 m_+^3} - \frac{3C_A^2 C_F m_2^3 Y_f^2 \mathcal{L}_m \mathcal{L}_m^2 m_1^2}{4m_-^2 m_+^2} - \frac{17CA m_2 Y_f^2 \beta_0 m_1^2}{16m_- m_+^2} - \frac{9CA m_2 Y_f^2 \beta_0 m_1^2}{16\epsilon m_- m_+^2} - \frac{9CA m_2 Y_f^2 \beta_0 m_1^2}{16\epsilon^2 m_- m_+^2} - \\
& \frac{3CA m_2 \pi^2 Y_f^2 \beta_0 m_1^2}{64m_- m_+^2} + \frac{3CA C_F^2 m_2^3 Y_f^2 \mathcal{L}_m m_1^2}{m_-^2 m_+^3} + \frac{3C_F m_2^3 Y_f^2 \mathcal{L}_m m_1^2}{2m_-^2 m_+^3} + \frac{23CA C_F^2 m_2^3 Y_f^2 \mathcal{L}_m m_1^2}{2\epsilon m_-^2 m_+^3} + \frac{23C_F m_2^3 Y_f^2 \mathcal{L}_m m_1^2}{4\epsilon m_-^2 m_+^3} +
\end{aligned}$$

$$\begin{aligned}
& \frac{3C_A C_F^2 m_3^3 Y_f^2 L_{m_1} m_1^2}{\varepsilon^2 m_-^2 m_+^3} + \frac{3C_F m_3^3 Y_f^2 L_{m_1} m_1^2}{2\varepsilon^2 m_-^2 m_+^3} + \frac{C_A C_F^2 m_3^3 \pi^2 Y_f^2 L_{m_1} m_1^2}{2m_-^2 m_+^3} + \frac{C_F m_3^3 \pi^2 Y_f^2 L_{m_1} m_1^2}{4m_-^2 m_+^3} - \frac{18C_A^2 C_F m_3^3 Y_f^2 L_{m_2} m_1^2}{m_-^2 m_+^3} \\
& \frac{23C_A^2 C_F m_3^3 Y_f^2 L_{m_2} m_1^2}{4\varepsilon m_-^2 m_+^3} - \frac{3C_A^2 C_F m_3^3 Y_f^2 L_{m_2} m_1^2}{2\varepsilon^2 m_-^2 m_+^3} - \frac{C_A^2 C_F m_3^3 \pi^2 Y_f^2 L_{m_2} m_1^2}{4m_-^2 m_+^3} - \frac{3C_A^2 C_F m_3^2 Y_f^2 L_{m_1}^2 L_{m_2} m_1^2}{4m_-^2 m_+^2} + \\
& \frac{23C_A^2 C_F m_3^3 Y_f^2 L_{m_1} L_{m_2} m_1^2}{4m_-^2 m_+^3} + \frac{3C_A^2 C_F m_3^3 Y_f^2 L_{m_1} L_{m_2} m_1^2}{2\varepsilon m_-^2 m_+^3} + \frac{49C_A C_F^2 m_3^2 Y_f^2 m_1}{2m_- m_+^2} + \frac{49C_F m_3^2 Y_f^2 m_1}{4m_- m_+^2} + \frac{7C_A C_F^2 m_3^2 Y_f^2 m_1}{\varepsilon m_- m_+^2} + \\
& \frac{7C_F m_3^2 Y_f^2 m_1}{2\varepsilon m_- m_+^2} - \frac{3C_A C_F^2 m_3^2 Y_f^2 m_1}{2\varepsilon^2 m_- m_+^2} - \frac{3C_F m_3^2 Y_f^2 m_1}{4\varepsilon^2 m_- m_+^2} - \frac{C_A C_F^2 m_3^2 \pi^2 Y_f^2 m_1}{4m_- m_+^2} - \frac{C_F m_3^2 \pi^2 Y_f^2 m_1}{8m_- m_+^2} + \frac{21C_A C_F^2 m_3^2 Y_f^2 L_{m_1} m_1}{4m_-^2 m_+^3} \\
& \frac{21C_F m_3^2 Y_f^2 L_{m_1} m_1}{8m_-^2 m_+^3} - \frac{3C_A^2 C_F m_3^2 Y_f^2 L_{m_2} m_1}{8m_-^2 m_+^3} + \frac{17C_A m_3^2 Y_f^2 \beta_0 m_1}{16m_- m_+^2} + \frac{9C_A m_3^2 Y_f^2 \beta_0 m_1}{16\varepsilon m_- m_+^2} + \frac{9C_A m_3^2 Y_f^2 \beta_0 m_1}{16\varepsilon^2 m_- m_+^2} + \\
& \frac{3C_A m_3^2 \pi^2 Y_f^2 \beta_0 m_1}{64m_- m_+^2} + \frac{6C_A C_F^2 m_3^2 Y_f^2 L_{m_1} m_1}{m_-^2 m_+^3} + \frac{3C_F m_3^2 Y_f^2 L_{m_1} m_1}{m_-^2 m_+^3} - \frac{9C_A C_F^2 m_3^2 Y_f^2 L_{m_1} m_1}{2\varepsilon m_-^2 m_+^3} - \frac{9C_F m_3^2 Y_f^2 L_{m_1} m_1}{4\varepsilon m_-^2 m_+^3} + \\
& \frac{4C_A^2 C_F m_3^2 Y_f^2 L_{m_2} m_1}{m_-^2 m_+^3} + \frac{3C_A^2 C_F m_3^2 Y_f^2 L_{m_2} m_1}{4\varepsilon m_-^2 m_+^3} - \frac{3C_A^2 C_F m_3^2 Y_f^2 L_{m_1} L_{m_2} m_1}{4m_-^2 m_+^3} + \frac{7C_A C_F^2 m_3^3 Y_f^2}{2m_- m_+^2} + \frac{7C_F m_3^3 Y_f^2}{4m_- m_+^2} + \\
& \frac{5C_A C_F^2 m_3^3 Y_f^2}{2\varepsilon m_- m_+^2} + \frac{5C_F m_3^3 Y_f^2}{4\varepsilon m_- m_+^2} - \frac{3C_A C_F^2 m_3^3 Y_f^2}{2\varepsilon^2 m_- m_+^2} - \frac{3C_F m_3^3 Y_f^2}{4\varepsilon^2 m_- m_+^2} - \frac{C_A C_F^2 m_3^3 \pi^2 Y_f^2}{4m_- m_+^2} - \frac{C_F m_3^3 \pi^2 Y_f^2}{8m_- m_+^2} + \frac{3C_A C_F^2 m_3^2 Y_f^2 L_{m_1}^2}{m_-^2 m_+^3} + \\
& \frac{3C_F m_3^2 Y_f^2 L_{m_1}^2}{2m_-^2 m_+^3} - \frac{9}{32} C_A Y_f^2 \beta_0 L_{m_1}^2 + \frac{17C_A m_3^3 Y_f^2 \beta_0}{16m_- m_+^2} + \frac{9C_A m_3^3 Y_f^2 \beta_0}{16\varepsilon m_- m_+^2} + \frac{9C_A m_3^3 Y_f^2 \beta_0}{16\varepsilon^2 m_- m_+^2} + \frac{3C_A m_3^3 \pi^2 Y_f^2 \beta_0}{64m_- m_+^2} + \\
& \frac{5C_A C_F^2 m_3^2 Y_f^2 L_{m_1}}{m_-^2 m_+^3} + \frac{5C_F m_3^2 Y_f^2 L_{m_1}}{2m_-^2 m_+^3} - \frac{3C_A C_F^2 m_3^2 Y_f^2 L_{m_1}}{\varepsilon m_-^2 m_+^3} - \frac{3C_F m_3^2 Y_f^2 L_{m_1}}{2\varepsilon m_-^2 m_+^3} + \frac{9}{16} C_A Y_f^2 \beta_0 L_{m_1} + \frac{9C_A Y_f^2 \beta_0 L_{m_1}}{16\varepsilon} + \\
& \frac{C_A C_F^2 Y_s^2 m_1^3}{2m_-^2 m_- m_+^2} + \frac{C_F Y_s^2 m_1^3}{4m_-^2 m_- m_+^2} + \frac{C_A C_F^2 Y_s^2 m_1^3}{4\varepsilon m_-^2 m_- m_+^2} + \frac{C_F Y_s^2 m_1^3}{8\varepsilon m_-^2 m_- m_+^2} + \frac{C_A C_F^2 Y_s^2 m_1^3}{2\varepsilon^2 m_-^2 m_- m_+^2} + \frac{C_F Y_s^2 m_1^3}{4\varepsilon^2 m_-^2 m_- m_+^2} + \frac{C_A C_F^2 \pi^2 Y_s^2 m_1^3}{12m_-^2 m_- m_+^2} + \\
& \frac{C_F \pi^2 Y_s^2 m_1^3}{24m_-^2 m_- m_+^2} + \frac{C_A C_F^2 Y_s^2 L_{m_2}^2 m_1^3}{m_-^2 m_- m_+^2} + \frac{C_F Y_s^2 L_{m_2}^2 m_1^3}{2m_-^2 m_- m_+^2} - \frac{C_A Y_s^2 \beta_0 L_{m_2}^2 m_1^3}{16m_-^2 m_- m_+^2} - \frac{C_A Y_s^2 \beta_0 m_1^3}{4m_-^2 m_- m_+^2} - \frac{C_A Y_s^2 \beta_0 m_1^3}{8\varepsilon m_-^2 m_- m_+^2} - \frac{C_A Y_s^2 \beta_0 m_1^3}{8\varepsilon^2 m_-^2 m_- m_+^2} + \\
& \frac{C_A \pi^2 Y_s^2 \beta_0 m_1^3}{96m_-^2 m_- m_+^2} - \frac{C_A C_F^2 Y_s^2 L_{m_2} m_1^3}{2m_-^2 m_- m_+^2} - \frac{C_F Y_s^2 L_{m_2} m_1^3}{4m_-^2 m_- m_+^2} - \frac{C_A C_F^2 Y_s^2 L_{m_2} m_1^3}{\varepsilon m_-^2 m_- m_+^2} - \frac{C_F Y_s^2 L_{m_2} m_1^3}{2\varepsilon m_-^2 m_- m_+^2} + \frac{C_A Y_s^2 \beta_0 L_{m_2} m_1^3}{8m_-^2 m_- m_+^2} + \\
& \frac{C_A Y_s^2 \beta_0 L_{m_2} m_1^3}{8\varepsilon m_-^2 m_- m_+^2} + \frac{C_A C_F^2 Y_s^2 m_1^2}{2m_- m_- m_+^2} + \frac{C_F Y_s^2 m_1^2}{4m_- m_- m_+^2} + \frac{C_A C_F^2 Y_s^2 m_1^2}{4\varepsilon m_- m_- m_+^2} + \frac{C_F Y_s^2 m_1^2}{8\varepsilon m_- m_- m_+^2} + \frac{C_A C_F^2 Y_s^2 m_1^2}{2\varepsilon^2 m_- m_- m_+^2} + \frac{C_F Y_s^2 m_1^2}{4\varepsilon^2 m_- m_- m_+^2} + \\
& \frac{C_A C_F^2 \pi^2 Y_s^2 m_1^2}{12m_- m_- m_+^2} + \frac{C_F \pi^2 Y_s^2 m_1^2}{24m_- m_- m_+^2} + \frac{C_A C_F^2 Y_s^2 L_{m_2}^2 m_1^2}{m_- m_- m_+^2} + \frac{C_F Y_s^2 L_{m_2}^2 m_1^2}{2m_- m_- m_+^2} - \frac{C_A Y_s^2 \beta_0 L_{m_2}^2 m_1^2}{16m_- m_- m_+^2} - \frac{C_A Y_s^2 \beta_0 m_1^2}{4m_- m_- m_+^2} - \frac{C_A Y_s^2 \beta_0 m_1^2}{8\varepsilon m_- m_- m_+^2} + \\
& \frac{C_A Y_s^2 \beta_0 m_1^2}{8\varepsilon^2 m_- m_- m_+^2} - \frac{C_A \pi^2 Y_s^2 \beta_0 m_1^2}{96m_- m_- m_+^2} - \frac{C_A C_F^2 Y_s^2 L_{m_2} m_1^2}{2m_- m_- m_+^2} - \frac{C_F Y_s^2 L_{m_2} m_1^2}{4m_- m_- m_+^2} - \frac{C_A C_F^2 Y_s^2 L_{m_2} m_1^2}{\varepsilon m_- m_- m_+^2} - \frac{C_F Y_s^2 L_{m_2} m_1^2}{2\varepsilon m_- m_- m_+^2} + \\
& \frac{C_A Y_s^2 \beta_0 L_{m_2} m_1^2}{8m_- m_- m_+^2} + \frac{C_A Y_s^2 \beta_0 L_{m_2} m_1^2}{8\varepsilon m_- m_- m_+^2} - \frac{21C_A C_F^2 Y_s^2 m_1}{2m_- m_+^2} - \frac{21C_F Y_s^2 m_1}{4m_- m_+^2} - \frac{13C_A C_F^2 Y_s^2 m_1}{4\varepsilon m_- m_+^2} - \frac{13C_F Y_s^2 m_1}{8\varepsilon m_- m_+^2} - \frac{C_A C_F^2 Y_s^2 m_1}{\varepsilon^2 m_- m_+^2} \\
& \frac{C_F Y_s^2 m_1}{2\varepsilon^2 m_- m_+^2} - \frac{C_A C_F^2 \pi^2 Y_s^2 m_1}{6m_- m_+^2} - \frac{C_F \pi^2 Y_s^2 m_1}{12m_- m_+^2} + \frac{C_A^2 C_F Y_s^2 L_{m_1}^2 m_1}{8\varepsilon m_- m_+^2} - \frac{2C_A C_F^2 Y_s^2 L_{m_2}^2 m_1}{m_- m_+^2} - \frac{C_F Y_s^2 L_{m_2}^2 m_1}{m_- m_+^2} \\
& \frac{3C_A C_F^2 Y_s^2 L_{m_2}^2 m_1}{4\varepsilon m_- m_+^2} - \frac{3C_F Y_s^2 L_{m_2}^2 m_1}{8\varepsilon m_- m_+^2} + \frac{C_A Y_s^2 \beta_0 L_{m_2}^2 m_1}{16m_- m_+^2} + \frac{C_A Y_s^2 \beta_0 m_1}{4m_- m_+^2} + \frac{C_A Y_s^2 \beta_0 m_1}{8\varepsilon m_- m_+^2} + \frac{C_A Y_s^2 \beta_0 m_1}{8\varepsilon^2 m_- m_+^2} + \frac{C_A \pi^2 Y_s^2 \beta_0 m_1}{96m_- m_+^2} + \\
& \frac{C_A^2 C_F Y_s^2 L_{m_1} m_1}{m_- m_+^2} - \frac{C_A^2 C_F Y_s^2 L_{m_1} m_1}{4\varepsilon^2 m_- m_+^2} - \frac{C_A^2 C_F \pi^2 Y_s^2 L_{m_1} m_1}{24m_- m_+^2} + \frac{9C_A C_F^2 Y_s^2 L_{m_2} m_1}{2m_- m_+^2} + \frac{9C_F Y_s^2 L_{m_2} m_1}{4m_- m_+^2} + \\
& \frac{2C_A C_F^2 Y_s^2 L_{m_2} m_1}{\varepsilon m_- m_+^2} + \frac{C_F Y_s^2 L_{m_2} m_1}{\varepsilon m_- m_+^2} + \frac{C_A C_F^2 Y_s^2 L_{m_2} m_1}{2\varepsilon^2 m_- m_+^2} + \frac{C_F Y_s^2 L_{m_2} m_1}{4\varepsilon^2 m_- m_+^2} + \frac{C_A C_F^2 \pi^2 Y_s^2 L_{m_2} m_1}{12m_- m_+^2} + \frac{C_F \pi^2 Y_s^2 L_{m_2} m_1}{24m_- m_+^2} \\
& \frac{C_A Y_s^2 \beta_0 L_{m_2} m_1}{8m_- m_+^2} - \frac{C_A Y_s^2 \beta_0 L_{m_2} m_1}{8\varepsilon m_- m_+^2} + \frac{C_A^2 C_F Y_s^2 L_{m_1} L_{m_2} m_1}{4\varepsilon m_- m_+^2} - \frac{C_A^2 C_F Y_s^2 L_{m_1}^3}{24m_- m_+} + \frac{7C_A^2 C_F Y_s^2 L_{m_2}^3}{24m_- m_+} + \frac{27C_A C_F^2 m_2 Y_s^2}{2m_- m_+^2} + \\
& \frac{27C_F m_2 Y_s^2}{4m_- m_+^2} + \frac{19C_A C_F^2 m_2 Y_s^2}{4\varepsilon m_- m_+^2} + \frac{19C_F m_2 Y_s^2}{8\varepsilon m_- m_+^2} + \frac{C_A C_F^2 m_2 Y_s^2}{\varepsilon^2 m_- m_+^2} + \frac{C_F m_2 Y_s^2}{2\varepsilon^2 m_- m_+^2} + \frac{C_A C_F^2 m_2 \pi^2 Y_s^2}{6m_- m_+^2} + \frac{C_F m_2 \pi^2 Y_s^2}{12m_- m_+^2} + \\
& \frac{C_A^2 C_F m_2 Y_s^2 L_{m_1}^2}{2m_- m_+^2} + \frac{C_A^2 C_F m_2 Y_s^2 L_{m_1}^2}{8\varepsilon m_- m_+^2} - \frac{C_A C_F^2 m_2 Y_s^2 L_{m_2}^2}{m_- m_+^2} - \frac{C_F m_2 Y_s^2 L_{m_2}^2}{2m_- m_+^2} - \frac{3C_A C_F^2 m_2 Y_s^2 L_{m_2}^2}{4\varepsilon m_- m_+^2} - \frac{3C_F m_2 Y_s^2 L_{m_2}^2}{8\varepsilon m_- m_+^2} + \\
& \frac{C_A m_2 Y_s^2 \beta_0 L_{m_2}^2}{16m_- m_+^2} - \frac{C_A^2 C_F Y_s^2 L_{m_1} L_{m_2}^2}{8m_- m_+} + \frac{C_A m_2 Y_s^2 \beta_0}{4m_- m_+^2} + \frac{C_A m_2 Y_s^2 \beta_0}{8\varepsilon m_- m_+^2} + \frac{C_A m_2 Y_s^2 \beta_0}{8\varepsilon^2 m_- m_+^2} + \frac{C_A m_2 \pi^2 Y_s^2 \beta_0}{96m_- m_+^2} \\
& \frac{3C_A^2 C_F m_2 Y_s^2 L_{m_1}}{m_- m_+^2} - \frac{C_A^2 C_F m_2 Y_s^2 L_{m_1}}{\varepsilon m_- m_+^2} - \frac{C_A^2 C_F m_2 Y_s^2 L_{m_1}}{4\varepsilon^2 m_- m_+^2} - \frac{C_A^2 C_F m_2 \pi^2 Y_s^2 L_{m_1}}{24m_- m_+^2} - \frac{7C_A C_F^2 m_2 Y_s^2 L_{m_2}}{2m_- m_+^2} \\
& \frac{7C_F m_2 Y_s^2 L_{m_2}}{4m_- m_+^2} + \frac{C_A C_F^2 m_2 Y_s^2 L_{m_2}}{2\varepsilon^2 m_- m_+^2} + \frac{C_F m_2 Y_s^2 L_{m_2}}{4\varepsilon^2 m_- m_+^2} + \frac{C_A C_F^2 m_2 \pi^2 Y_s^2 L_{m_2}}{12m_- m_+^2} + \frac{C_F m_2 \pi^2 Y_s^2 L_{m_2}}{24m_- m_+^2} - \frac{C_A^2 C_F Y_s^2 L_{m_1}^2 L_{m_2}}{8m_- m_+} \\
& \frac{C_A m_2 Y_s^2 \beta_0 L_{m_2}}{8m_- m_+^2} - \frac{C_A m_2 Y_s^2 \beta_0 L_{m_2}}{8\varepsilon m_- m_+^2} + \frac{C_A^2 C_F m_2 Y_s^2 L_{m_1} L_{m_2}}{m_- m_+^2} + \frac{C_A^2 C_F m_2 Y_s^2 L_{m_1} L_{m_2}}{4\varepsilon m_- m_+^2} - \frac{9C_A C_F^2 Y_s^2 L_{m_1}^2 m_1^4}{m_-^2 m_+^3}
\end{aligned}$$

$$\begin{aligned}
& \frac{9C_F Y_f Y_s \mathcal{L}_2^2 m_1^4}{2m_2^2 m_3^4} - \frac{9C_A C_F^2 Y_f Y_s \mathcal{L}_2^2 m_1^4}{4\epsilon m_2^2 m_3^4} - \frac{9C_F Y_f Y_s \mathcal{L}_2^2 m_1^4}{8\epsilon m_2^2 m_3^4} + \frac{4C_A C_F^2 Y_f Y_s \mathcal{L}_2^2 m_1^4}{m_2^2 m_3^4} + \frac{2C_F Y_f Y_s \mathcal{L}_2^2 m_1^4}{m_2^2 m_3^4} + \\
& \frac{3C_A C_F^2 Y_f Y_s \mathcal{L}_2^2 m_1^4}{4\epsilon m_2^2 m_3^4} + \frac{3C_F Y_f Y_s \mathcal{L}_2^2 m_1^4}{8\epsilon m_2^2 m_3^4} + \frac{C_A Y_f Y_s \beta_0 \mathcal{L}_2^2 m_1^4}{8\epsilon m_2^2 m_3^4} + \frac{9C_A C_F^2 Y_f Y_s \mathcal{L}_2 m_1^4}{m_2^2 m_3^4} + \frac{9C_F Y_f Y_s \mathcal{L}_2 m_1^4}{2m_2^2 m_3^4} + \\
& \frac{8C_A C_F^2 Y_f Y_s \mathcal{L}_2 m_1^4}{\epsilon m_2^2 m_3^4} + \frac{4C_F Y_f Y_s \mathcal{L}_2 m_1^4}{\epsilon m_2^2 m_3^4} + \frac{3C_A C_F^2 Y_f Y_s \mathcal{L}_2 m_1^4}{2\epsilon^2 m_2^2 m_3^4} + \frac{3C_F Y_f Y_s \mathcal{L}_2 m_1^4}{4\epsilon^2 m_2^2 m_3^4} + \frac{C_A C_F^2 \pi^2 Y_f Y_s \mathcal{L}_2 m_1^4}{4m_2^2 m_3^4} + \\
& \frac{C_F \pi^2 Y_f Y_s \mathcal{L}_2 m_1^4}{8m_2^2 m_3^4} - \frac{11C_A C_F^2 Y_f Y_s \mathcal{L}_2 m_2^4}{m_2^2 m_3^4} - \frac{11C_F Y_f Y_s \mathcal{L}_2 m_2^4}{2m_2^2 m_3^4} - \frac{5C_A C_F^2 Y_f Y_s \mathcal{L}_2 m_2^4}{\epsilon m_2^2 m_3^4} - \frac{5C_F Y_f Y_s \mathcal{L}_2 m_2^4}{2\epsilon m_2^2 m_3^4} - \\
& \frac{3C_A C_F^2 Y_f Y_s \mathcal{L}_2 m_2^4}{2\epsilon^2 m_2^2 m_3^4} - \frac{3C_F Y_f Y_s \mathcal{L}_2 m_2^4}{4\epsilon^2 m_2^2 m_3^4} - \frac{C_A C_F^2 \pi^2 Y_f Y_s \mathcal{L}_2 m_2^4}{4m_2^2 m_3^4} - \frac{C_F \pi^2 Y_f Y_s \mathcal{L}_2 m_2^4}{8m_2^2 m_3^4} - \frac{C_A Y_f Y_s \beta_0 \mathcal{L}_2 m_2^4}{4\epsilon^2 m_2^2 m_3^4} - \\
& \frac{C_A \pi^2 Y_f Y_s \beta_0 \mathcal{L}_2 m_2^4}{48m_2^2 m_3^4} + \frac{C_A^2 C_F Y_f Y_s \mathcal{L}_2 m_2^4}{m_2^2 m_3^4} + \frac{3C_A^2 C_F Y_f Y_s \mathcal{L}_2 m_2^4}{4\epsilon m_2^2 m_3^4} + \frac{7C_A C_F^2 Y_f Y_s \mathcal{L}_2^3 m_1^3}{4m_2^2 m_3^4} + \\
& \frac{7C_F Y_f Y_s \mathcal{L}_2^3 m_1^3}{8m_2^2 m_3^4} + \frac{C_A Y_f Y_s \beta_0 \mathcal{L}_2^3 m_1^3}{24m_2^2 m_3^4} - \frac{C_A C_F^2 Y_f Y_s \mathcal{L}_2^3 m_1^3}{4m_2^2 m_3^4} - \frac{C_F Y_f Y_s \mathcal{L}_2^3 m_1^3}{8m_2^2 m_3^4} - \frac{C_A Y_f Y_s \beta_0 \mathcal{L}_2^3 m_1^3}{24m_2^2 m_3^4} - \\
& \frac{27C_A C_F^2 m_2 Y_f Y_s \mathcal{L}_2^2 m_1^3}{2m_2^2 m_3^4} - \frac{27C_F m_2 Y_f Y_s \mathcal{L}_2^2 m_1^3}{4m_2^2 m_3^4} - \frac{9C_A C_F^2 m_2 Y_f Y_s \mathcal{L}_2^2 m_1^3}{4\epsilon m_2^2 m_3^4} - \frac{9C_F m_2 Y_f Y_s \mathcal{L}_2^2 m_1^3}{8\epsilon m_2^2 m_3^4} + \\
& \frac{11C_A C_F^2 m_2 Y_f Y_s \mathcal{L}_2^2 m_1^3}{2m_2^2 m_3^4} + \frac{11C_F m_2 Y_f Y_s \mathcal{L}_2^2 m_1^3}{4m_2^2 m_3^4} + \frac{3C_A C_F^2 m_2 Y_f Y_s \mathcal{L}_2^2 m_1^3}{4\epsilon m_2^2 m_3^4} + \frac{3C_F m_2 Y_f Y_s \mathcal{L}_2^2 m_1^3}{8\epsilon m_2^2 m_3^4} + \\
& \frac{C_A m_2 Y_f Y_s \beta_0 \mathcal{L}_2^2 m_1^3}{8m_2^2 m_3^4} + \frac{C_A m_2 Y_f Y_s \beta_0 \mathcal{L}_2^2 m_1^3}{8\epsilon m_2^2 m_3^4} - \frac{3C_A^2 C_F Y_f Y_s \mathcal{L}_2 m_1^3}{8m_2^2 m_3^4} + \frac{31C_A C_F^2 m_2 Y_f Y_s \mathcal{L}_2 m_1^3}{m_2^2 m_3^4} + \\
& \frac{31C_F m_2 Y_f Y_s \mathcal{L}_2 m_1^3}{2m_2^2 m_3^4} + \frac{11C_A C_F^2 m_2 Y_f Y_s \mathcal{L}_2 m_1^3}{\epsilon m_2^2 m_3^4} + \frac{11C_F m_2 Y_f Y_s \mathcal{L}_2 m_1^3}{2\epsilon m_2^2 m_3^4} + \frac{3C_A C_F^2 m_2 Y_f Y_s \mathcal{L}_2 m_1^3}{2\epsilon^2 m_2^2 m_3^4} + \\
& \frac{3C_F m_2 Y_f Y_s \mathcal{L}_2 m_1^3}{4\epsilon^2 m_2^2 m_3^4} + \frac{C_A C_F^2 m_2 \pi^2 Y_f Y_s \mathcal{L}_2 m_1^3}{4m_2^2 m_3^4} + \frac{C_F m_2 \pi^2 Y_f Y_s \mathcal{L}_2 m_1^3}{8m_2^2 m_3^4} - \frac{3C_A^2 C_F Y_f Y_s \mathcal{L}_2^2 m_1 m_2^3}{8m_2^2 m_3^4} - \\
& \frac{24C_A C_F^2 m_2 Y_f Y_s \mathcal{L}_2 m_2^3}{m_2^2 m_3^4} - \frac{12C_F m_2 Y_f Y_s \mathcal{L}_2 m_2^3}{m_2^2 m_3^4} - \frac{8C_A C_F^2 m_2 Y_f Y_s \mathcal{L}_2 m_2^3}{\epsilon m_2^2 m_3^4} - \frac{4C_F m_2 Y_f Y_s \mathcal{L}_2 m_2^3}{\epsilon m_2^2 m_3^4} - \\
& \frac{3C_A C_F^2 m_2 Y_f Y_s \mathcal{L}_2 m_2^3}{2\epsilon^2 m_2^2 m_3^4} - \frac{3C_F m_2 Y_f Y_s \mathcal{L}_2 m_2^3}{4\epsilon^2 m_2^2 m_3^4} - \frac{C_A C_F^2 m_2 \pi^2 Y_f Y_s \mathcal{L}_2 m_2^3}{4m_2^2 m_3^4} - \frac{C_F m_2 \pi^2 Y_f Y_s \mathcal{L}_2 m_2^3}{8m_2^2 m_3^4} - \\
& \frac{C_A m_2 Y_f Y_s \beta_0 \mathcal{L}_2 m_2^3}{2m_2^2 m_3^4} - \frac{C_A m_2 Y_f Y_s \beta_0 \mathcal{L}_2 m_2^3}{4\epsilon m_2^2 m_3^4} - \frac{C_A m_2 Y_f Y_s \beta_0 \mathcal{L}_2 m_2^3}{4\epsilon^2 m_2^2 m_3^4} - \frac{C_A m_2 \pi^2 Y_f Y_s \beta_0 \mathcal{L}_2 m_2^3}{48m_2^2 m_3^4} + \\
& \frac{5C_A^2 C_F m_2 Y_f Y_s \mathcal{L}_2 m_1^3}{2m_2^2 m_3^4} + \frac{3C_A^2 C_F m_2 Y_f Y_s \mathcal{L}_2 m_1^3}{4\epsilon m_2^2 m_3^4} + \frac{21C_A C_F^2 m_2^2 Y_f Y_s \mathcal{L}_2^2 m_1^2}{4m_2^2 m_3^4} + \frac{21C_F m_2^2 Y_f Y_s \mathcal{L}_2^2 m_1^2}{8m_2^2 m_3^4} - \\
& \frac{9C_A C_F^2 m_2^2 Y_f Y_s \mathcal{L}_2^2 m_1^2}{4\epsilon m_2^2 m_3^4} - \frac{9C_F m_2^2 Y_f Y_s \mathcal{L}_2^2 m_1^2}{8\epsilon m_2^2 m_3^4} - \frac{C_A Y_f Y_s \beta_0 \mathcal{L}_2^2 m_1^2}{8\epsilon m_2^2 m_3^4} - \frac{23C_A C_F^2 m_2^2 Y_f Y_s \mathcal{L}_2^2 m_1^2}{4m_2^2 m_3^4} - \\
& \frac{23C_F m_2^2 Y_f Y_s \mathcal{L}_2^2 m_1^2}{8m_2^2 m_3^4} + \frac{3C_A C_F^2 m_2^2 Y_f Y_s \mathcal{L}_2^2 m_1^2}{4\epsilon m_2^2 m_3^4} + \frac{3C_F m_2^2 Y_f Y_s \mathcal{L}_2^2 m_1^2}{8\epsilon m_2^2 m_3^4} - \frac{C_A m_2^2 Y_f Y_s \beta_0 \mathcal{L}_2^2 m_1^2}{8\epsilon m_2^2 m_3^4} + \frac{9C_A C_F^2 Y_f Y_s m_2^2}{m_2^2 m_3^4} + \\
& \frac{9C_F Y_f Y_s m_2^2}{2m_2^2 m_3^4} + \frac{C_A C_F^2 Y_f Y_s m_2^2}{\epsilon m_2^2 m_3^4} + \frac{C_F Y_f Y_s m_2^2}{2\epsilon m_2^2 m_3^4} - \frac{3C_A C_F^2 Y_f Y_s m_2^2}{2\epsilon^2 m_2^2 m_3^4} - \frac{3C_F Y_f Y_s m_2^2}{4\epsilon^2 m_2^2 m_3^4} - \frac{C_A C_F^2 \pi^2 Y_f Y_s m_2^2}{4m_2^2 m_3^4} - \frac{C_F \pi^2 Y_f Y_s m_2^2}{8m_2^2 m_3^4} + \\
& \frac{3C_A Y_f Y_s \beta_0 m_2^2}{4m_2^2 m_3^4} + \frac{C_A Y_f Y_s \beta_0 m_2^2}{4\epsilon m_2^2 m_3^4} - \frac{21C_A C_F^2 m_2^2 Y_f Y_s \mathcal{L}_2 m_1^2}{2m_2^2 m_3^4} - \frac{21C_F m_2^2 Y_f Y_s \mathcal{L}_2 m_1^2}{4m_2^2 m_3^4} - \frac{11C_A C_F^2 m_2^2 Y_f Y_s \mathcal{L}_2 m_1^2}{2\epsilon m_2^2 m_3^4} - \\
& \frac{11C_F m_2^2 Y_f Y_s \mathcal{L}_2 m_1^2}{4\epsilon m_2^2 m_3^4} + \frac{3C_A C_F^2 m_2^2 Y_f Y_s \mathcal{L}_2 m_1^2}{2\epsilon^2 m_2^2 m_3^4} + \frac{3C_F m_2^2 Y_f Y_s \mathcal{L}_2 m_1^2}{4\epsilon^2 m_2^2 m_3^4} + \frac{C_A C_F^2 m_2^2 \pi^2 Y_f Y_s \mathcal{L}_2 m_1^2}{4m_2^2 m_3^4} + \\
& \frac{C_F m_2^2 \pi^2 Y_f Y_s \mathcal{L}_2 m_1^2}{8m_2^2 m_3^4} - \frac{C_A Y_f Y_s \beta_0 \mathcal{L}_2 m_1^2}{4m_2^2 m_3^4} + \frac{C_A Y_f Y_s \beta_0 \mathcal{L}_2 m_1^2}{4\epsilon^2 m_2^2 m_3^4} + \frac{C_A \pi^2 Y_f Y_s \beta_0 \mathcal{L}_2 m_1^2}{48m_2^2 m_3^4} + \frac{45C_A C_F^2 m_2^2 Y_f Y_s \mathcal{L}_2 m_2^2}{2m_2^2 m_3^4} + \\
& \frac{45C_F m_2^2 Y_f Y_s \mathcal{L}_2 m_2^2}{4m_2^2 m_3^4} + \frac{11C_A C_F^2 m_2^2 Y_f Y_s \mathcal{L}_2 m_2^2}{2\epsilon m_2^2 m_3^4} + \frac{11C_F m_2^2 Y_f Y_s \mathcal{L}_2 m_2^2}{4\epsilon m_2^2 m_3^4} - \frac{3C_A C_F^2 m_2^2 Y_f Y_s \mathcal{L}_2 m_2^2}{2\epsilon^2 m_2^2 m_3^4} - \\
& \frac{3C_F m_2^2 Y_f Y_s \mathcal{L}_2 m_2^2}{4\epsilon^2 m_2^2 m_3^4} - \frac{C_A C_F^2 m_2^2 \pi^2 Y_f Y_s \mathcal{L}_2 m_2^2}{4m_2^2 m_3^4} - \frac{C_F m_2^2 \pi^2 Y_f Y_s \mathcal{L}_2 m_2^2}{4m_2^2 m_3^4} + \\
& \frac{C_A m_2^2 Y_f Y_s \beta_0 \mathcal{L}_2 m_2^2}{4\epsilon^2 m_2^2 m_3^4} + \frac{C_A m_2^2 \pi^2 Y_f Y_s \beta_0 \mathcal{L}_2 m_2^2}{48m_2^2 m_3^4} + \frac{C_A^2 C_F m_2^2 Y_f Y_s \mathcal{L}_2 m_1^2}{4m_2^2 m_3^4} + \frac{3C_A^2 C_F m_2^2 Y_f Y_s \mathcal{L}_2 m_1^2}{4\epsilon m_2^2 m_3^4} + \\
& \frac{7C_A C_F^2 m_2^2 Y_f Y_s \mathcal{L}_2^3 m_1}{4m_2^2 m_3^4} + \frac{7C_F m_2^2 Y_f Y_s \mathcal{L}_2^3 m_1}{8m_2^2 m_3^4} - \frac{C_A m_2^2 Y_f Y_s \beta_0 \mathcal{L}_2^3 m_1}{24m_2^2 m_3^4} - \frac{C_A C_F^2 m_2^2 Y_f Y_s \mathcal{L}_2^3 m_1}{4m_2^2 m_3^4} - \\
& \frac{C_F m_2^2 Y_f Y_s \mathcal{L}_2^3 m_1}{8m_2^2 m_3^4} + \frac{C_A m_2^2 Y_f Y_s \beta_0 \mathcal{L}_2^3 m_1}{24m_2^2 m_3^4} + \frac{3C_A C_F^2 m_2^2 Y_f Y_s \mathcal{L}_2^2 m_1}{4m_2^2 m_3^4} + \frac{3C_F m_2^2 Y_f Y_s \mathcal{L}_2^2 m_1}{8m_2^2 m_3^4} - \\
& \frac{9C_A C_F^2 m_2^2 Y_f Y_s \mathcal{L}_2^2 m_1}{4\epsilon m_2^2 m_3^4} - \frac{9C_F m_2^2 Y_f Y_s \mathcal{L}_2^2 m_1}{8\epsilon m_2^2 m_3^4} - \frac{C_A m_2^2 Y_f Y_s \beta_0 \mathcal{L}_2^2 m_1}{8m_2^2 m_3^4} - \frac{C_A m_2^2 Y_f Y_s \beta_0 \mathcal{L}_2^2 m_1}{8\epsilon m_2^2 m_3^4} -
\end{aligned}$$

$$\begin{aligned}
& \frac{5C_A C_F^2 m_2^3 Y_f Y_s \mathcal{L}_2^2 m_1}{4m^2 m_+^3} - \frac{5C_F m_2^3 Y_f Y_s \mathcal{L}_2^2 m_1}{8m^2 m_+^3} + \frac{3C_A C_F^2 m_2^3 Y_f Y_s \mathcal{L}_2^2 m_1}{4\epsilon m^2 m_+^3} + \frac{3C_F m_2^3 Y_f Y_s \mathcal{L}_2^2 m_1}{8\epsilon m^2 m_+^3} - \\
& \frac{C_A m_2^3 Y_f Y_s \beta_0 \mathcal{L}_2^2 m_1}{8m^2 m_+^3} - \frac{C_A m_2^3 Y_f Y_s \beta_0 \mathcal{L}_2^2 m_1}{8\epsilon m^2 m_+^3} - \frac{3C_A^2 C_F m_2^3 Y_f Y_s \mathcal{L}_2 m_1 \mathcal{L}_2^2 m_1}{8m^2 m_+^2} - \frac{21C_A C_F^2 m_2 Y_f Y_s m_1}{2m_- m_+^2} - \\
& \frac{21C_F m_2 Y_f Y_s m_1}{4m_- m_+^2} - \frac{7C_A C_F^2 m_2 Y_f Y_s m_1}{2\epsilon m_- m_+^2} - \frac{7C_F m_2 Y_f Y_s m_1}{4\epsilon m_- m_+^2} - \frac{3C_A C_F^2 m_2 Y_f Y_s m_1}{2\epsilon^2 m_- m_+^2} - \frac{3C_F m_2 Y_f Y_s m_1}{4\epsilon^2 m_- m_+^2} - \\
& \frac{C_A C_F^2 m_2 \pi^2 Y_f Y_s m_1}{4m_- m_+^2} - \frac{C_F m_2 \pi^2 Y_f Y_s m_1}{8m_- m_+^2} - \frac{13C_A C_F^2 m_2^3 Y_f Y_s \mathcal{L}_2 m_1 m_1}{2m_-^2 m_+^3} - \frac{13C_F m_2^3 Y_f Y_s \mathcal{L}_2 m_1 m_1}{4m_-^2 m_+^3} - \frac{5C_A C_F^2 m_2^3 Y_f Y_s \mathcal{L}_2 m_1 m_1}{2\epsilon m_-^2 m_+^3} - \\
& \frac{5C_F m_2^3 Y_f Y_s \mathcal{L}_2 m_1 m_1}{4\epsilon m_-^2 m_+^3} + \frac{3C_A C_F^2 m_2^3 Y_f Y_s \mathcal{L}_2 m_1 m_1}{2\epsilon^2 m_-^2 m_+^3} + \frac{3C_F m_2^3 Y_f Y_s \mathcal{L}_2 m_1 m_1}{4\epsilon^2 m_-^2 m_+^3} + \frac{C_A C_F^2 m_2^3 \pi^2 Y_f Y_s \mathcal{L}_2 m_1 m_1}{4m_-^2 m_+^3} + \\
& \frac{C_F m_2^3 \pi^2 Y_f Y_s \mathcal{L}_2 m_1 m_1}{8m_-^2 m_+^3} + \frac{C_A m_2 Y_f Y_s \beta_0 \mathcal{L}_2 m_1 m_1}{2m_- m_+^2} + \frac{C_A m_2 Y_f Y_s \beta_0 \mathcal{L}_2 m_1 m_1}{4\epsilon m_- m_+^2} + \frac{C_A m_2 Y_f Y_s \beta_0 \mathcal{L}_2 m_1 m_1}{4\epsilon^2 m_- m_+^2} + \\
& \frac{C_A m_2 \pi^2 Y_f Y_s \beta_0 \mathcal{L}_2 m_1 m_1}{48m_- m_+^2} - \frac{3C_A^2 C_F m_2^3 Y_f Y_s \mathcal{L}_2 m_1 \mathcal{L}_2 m_2 m_1}{8m_-^2 m_+^2} - \frac{C_A C_F^2 m_2^3 Y_f Y_s \mathcal{L}_2 m_2 m_1}{2m_-^2 m_+^3} - \frac{C_F m_2^3 Y_f Y_s \mathcal{L}_2 m_2 m_1}{4m_-^2 m_+^3} - \\
& \frac{C_A C_F^2 m_2^3 Y_f Y_s \mathcal{L}_2 m_2 m_1}{2\epsilon m_-^2 m_+^3} - \frac{C_F m_2^3 Y_f Y_s \mathcal{L}_2 m_2 m_1}{4\epsilon m_-^2 m_+^3} - \frac{3C_A C_F^2 m_2^3 Y_f Y_s \mathcal{L}_2 m_2 m_1}{2\epsilon^2 m_-^2 m_+^3} - \frac{3C_F m_2^3 Y_f Y_s \mathcal{L}_2 m_2 m_1}{4\epsilon^2 m_-^2 m_+^3} - \\
& \frac{C_A C_F^2 m_2^3 \pi^2 Y_f Y_s \mathcal{L}_2 m_2 m_1}{4m_-^2 m_+^3} - \frac{C_F m_2^3 \pi^2 Y_f Y_s \mathcal{L}_2 m_2 m_1}{8m_-^2 m_+^3} + \frac{C_A m_2^3 Y_f Y_s \beta_0 \mathcal{L}_2 m_2 m_1}{2m_-^2 m_+^3} + \frac{C_A m_2^3 Y_f Y_s \beta_0 \mathcal{L}_2 m_2 m_1}{4\epsilon m_-^2 m_+^3} + \\
& \frac{C_A m_2^3 Y_f Y_s \beta_0 \mathcal{L}_2 m_2 m_1}{4\epsilon^2 m_-^2 m_+^3} + \frac{C_A m_2^3 \pi^2 Y_f Y_s \beta_0 \mathcal{L}_2 m_2 m_1}{48m_-^2 m_+^3} + \frac{7C_A^2 C_F m_2^3 Y_f Y_s \mathcal{L}_2 m_1 \mathcal{L}_2 m_2 m_1}{4m_-^2 m_+^3} + \frac{3C_A^2 C_F m_2^3 Y_f Y_s \mathcal{L}_2 m_1 \mathcal{L}_2 m_2 m_1}{4\epsilon m_-^2 m_+^3} + \\
& \frac{3C_A C_F^2 m_2^4 Y_f Y_s \mathcal{L}_2^2 m_1}{m_-^2 m_+^3} + \frac{3C_F m_2^4 Y_f Y_s \mathcal{L}_2^2 m_1}{2m_-^2 m_+^3} - \frac{14C_A C_F^2 m_2^3 Y_f Y_s}{m_- m_+^2} - \frac{7C_F m_2^3 Y_f Y_s}{m_- m_+^2} - \frac{5C_A C_F^2 m_2^3 Y_f Y_s}{\epsilon m_- m_+^2} - \frac{5C_F m_2^3 Y_f Y_s}{2\epsilon m_- m_+^2} - \\
& \frac{3C_A C_F^2 m_2^3 Y_f Y_s}{2\epsilon^2 m_- m_+^2} - \frac{3C_F m_2^3 Y_f Y_s}{4\epsilon^2 m_- m_+^2} - \frac{C_A C_F^2 m_2^3 \pi^2 Y_f Y_s}{4m_- m_+^2} - \frac{C_F m_2^3 \pi^2 Y_f Y_s}{8m_- m_+^2} - \frac{3C_A m_2^3 Y_f Y_s \beta_0}{4m_- m_+^2} - \frac{C_A m_2^3 Y_f Y_s \beta_0}{4\epsilon m_- m_+^2} - \\
& \frac{10C_A C_F^2 m_2^4 Y_f Y_s \mathcal{L}_2 m_2}{m_-^2 m_+^3} - \frac{5C_F m_2^4 Y_f Y_s \mathcal{L}_2 m_2}{\epsilon m_-^2 m_+^3} - \frac{3C_A C_F^2 m_2^4 Y_f Y_s \mathcal{L}_2 m_2}{2\epsilon m_-^2 m_+^3} - \frac{3C_F m_2^4 Y_f Y_s \mathcal{L}_2 m_2}{2\epsilon m_-^2 m_+^3} - \frac{C_A m_2^4 Y_f Y_s \beta_0 \mathcal{L}_2 m_2}{4m_-^2 m_+^3} + \\
& \frac{m_1^2 Y_f Y_s^3 C_A^2 \mathcal{L}_2 m_2}{8\epsilon m_2^2 m_- m_+^2} - \frac{m_1^2 Y_f Y_s^3 C_A^2 \mathcal{L}_2 m_2}{8m_2^2 m_- m_+^2} + \frac{m_1^2 Y_f Y_s^3 C_A^2 \mathcal{L}_2 m_2}{8\epsilon m_2 m_- m_+^2} + \frac{m_1 Y_f Y_s^3 C_A^2 \mathcal{L}_2 m_2}{8\epsilon m_2 m_- m_+^2} - \frac{m_1 Y_f Y_s^3 C_A^2 \mathcal{L}_2 m_2}{8m_2 m_- m_+^2} + \frac{3m_1 Y_f Y_s^3 C_A^2 \mathcal{L}_2 m_2}{8m_2 m_- m_+^2} - \\
& \frac{Y_f Y_s^3 C_A^2 \mathcal{L}_2 m_2}{8m_- m_+^2} + \frac{Y_f Y_s^3 C_A^2 \mathcal{L}_2 m_2}{8m_- m_+^2} - \frac{3Y_f Y_s^3 C_A^2 \mathcal{L}_2 m_2}{8m_- m_+^2} - \frac{m_1^2 Y_f Y_s^3 C_A^2}{8\epsilon m_2^2 m_- m_+^2} - \frac{m_1^2 Y_f Y_s^3 C_A^2}{16\epsilon^2 m_2^2 m_- m_+^2} - \frac{m_1^2 Y_f Y_s^3 C_A^2}{4m_2^2 m_- m_+^2} - \frac{\pi^2 Y_f Y_s^3 C_A^2}{96m_2^2 m_- m_+^2} - \\
& \frac{3m_1 Y_f Y_s^3 C_A^2}{16\epsilon m_2 m_- m_+^2} - \frac{m_1 Y_f Y_s^3 C_A^2}{16\epsilon^2 m_2 m_- m_+^2} - \frac{m_1 Y_f Y_s^3 C_A^2}{2m_2 m_- m_+^2} - \frac{\pi^2 m_1 Y_f Y_s^3 C_A^2}{96m_2 m_- m_+^2} + \frac{3Y_f Y_s^3 C_A^2}{16\epsilon m_- m_+^2} + \frac{Y_f Y_s^3 C_A^2}{16\epsilon^2 m_- m_+^2} + \frac{Y_f Y_s^3 C_A^2}{2m_- m_+^2} + \frac{\pi^2 Y_f Y_s^3 C_A^2}{96m_- m_+^2} + \\
& \frac{45C_A^2 Y_f^3 Y_s \mathcal{L}_2^2 m_1^4}{64m^2 m_+^3} + \frac{9C_A^2 Y_f^3 Y_s \mathcal{L}_2^2 m_1^4}{64\epsilon m^2 m_+^3} - \frac{7C_A^2 Y_f^3 Y_s \mathcal{L}_2^2 m_1^4}{64m^2 m_+^3} - \frac{3C_A^2 Y_f^3 Y_s \mathcal{L}_2^2 m_1^4}{64\epsilon m^2 m_+^3} - \frac{37C_A^2 Y_f^3 Y_s \mathcal{L}_2 m_1^4}{32m^2 m_+^3} - \\
& \frac{19C_A^2 Y_f^3 Y_s \mathcal{L}_2 m_1^4}{32\epsilon m^2 m_+^3} - \frac{3C_A^2 Y_f^3 Y_s \mathcal{L}_2 m_1^4}{32\epsilon^2 m^2 m_+^3} - \frac{C_A^2 \pi^2 Y_f^3 Y_s \mathcal{L}_2 m_1^4}{64m^2 m_+^3} + \frac{15C_A^2 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1^4}{32m^2 m_+^3} + \frac{7C_A^2 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1^4}{32\epsilon m^2 m_+^3} + \\
& \frac{3C_A^2 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1^4}{32\epsilon^2 m^2 m_+^3} + \frac{C_A^2 \pi^2 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1^4}{64m^2 m_+^3} - \frac{7C_A^2 Y_f^3 Y_s \mathcal{L}_2 m_1 \mathcal{L}_2 m_2 m_1^4}{32m^2 m_+^3} - \frac{3C_A^2 Y_f^3 Y_s \mathcal{L}_2 m_1 \mathcal{L}_2 m_2 m_1^4}{32\epsilon m^2 m_+^3} - \frac{7C_A^2 Y_f^3 Y_s \mathcal{L}_2^3 m_1^4}{64m^2 m_+^2} + \\
& \frac{C_A^2 Y_f^3 Y_s \mathcal{L}_2^3 m_1^4}{64m^2 m_+^2} + \frac{63C_A^2 m_2 Y_f^3 Y_s \mathcal{L}_2^2 m_1^3}{64m^2 m_+^3} + \frac{9C_A^2 m_2 Y_f^3 Y_s \mathcal{L}_2^2 m_1^3}{64\epsilon m^2 m_+^3} - \frac{13C_A^2 m_2 Y_f^3 Y_s \mathcal{L}_2^2 m_1^3}{64m^2 m_+^3} - \frac{3C_A^2 m_2 Y_f^3 Y_s \mathcal{L}_2^2 m_1^3}{64\epsilon m^2 m_+^3} + \\
& \frac{3C_A^2 Y_f^3 Y_s \mathcal{L}_2 m_1 \mathcal{L}_2^2 m_1^3}{64m^2 m_+^2} - \frac{87C_A^2 m_2 Y_f^3 Y_s \mathcal{L}_2 m_1^3}{32m^2 m_+^3} - \frac{25C_A^2 m_2 Y_f^3 Y_s \mathcal{L}_2 m_1^3}{32\epsilon m^2 m_+^3} - \frac{3C_A^2 m_2 Y_f^3 Y_s \mathcal{L}_2 m_1^3}{32\epsilon^2 m^2 m_+^3} - \\
& \frac{C_A^2 m_2 \pi^2 Y_f^3 Y_s \mathcal{L}_2 m_1^3}{64m^2 m_+^3} + \frac{3C_A^2 Y_f^3 Y_s \mathcal{L}_2^2 m_1 \mathcal{L}_2 m_2 m_1^3}{64m^2 m_+^2} + \frac{41C_A^2 m_2 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1^3}{32m^2 m_+^3} + \frac{13C_A^2 m_2 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1^3}{32\epsilon m^2 m_+^3} + \\
& \frac{3C_A^2 m_2 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1^3}{32\epsilon^2 m^2 m_+^3} + \frac{C_A^2 m_2 \pi^2 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1^3}{64m^2 m_+^3} - \frac{13C_A^2 m_2 Y_f^3 Y_s \mathcal{L}_2 m_1 \mathcal{L}_2 m_2 m_1^3}{32m^2 m_+^3} - \frac{3C_A^2 m_2 Y_f^3 Y_s \mathcal{L}_2 m_1 \mathcal{L}_2 m_2 m_1^3}{32\epsilon m^2 m_+^3} - \\
& \frac{3C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_2 m_1^2 m_1^2}{16m^2 m_+^3} + \frac{9C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_2 m_1^2 m_1^2}{64\epsilon m^2 m_+^3} - \frac{C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1^2}{16m^2 m_+^3} - \frac{3C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1^2}{64\epsilon m^2 m_+^3} + \frac{C_A^2 Y_f^3 Y_s m_1^2}{4m_- m_+^2} + \\
& \frac{11C_A^2 Y_f^3 Y_s m_1^2}{32\epsilon m_- m_+^2} + \frac{3C_A^2 Y_f^3 Y_s m_1^2}{16\epsilon^2 m_- m_+^2} + \frac{C_A^2 \pi^2 Y_f^3 Y_s m_1^2}{32m_- m_+^2} + \frac{29C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_2 m_1^2 m_1^2}{32m^2 m_+^3} + \frac{C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_2 m_1^2 m_1^2}{4\epsilon m^2 m_+^3} - \frac{3C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_2 m_1^2 m_1^2}{32\epsilon^2 m^2 m_+^3} - \\
& \frac{C_A^2 m_2^2 \pi^2 Y_f^3 Y_s \mathcal{L}_2 m_1^2 m_1^2}{64m^2 m_+^3} - \frac{7C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1^2}{32m^2 m_+^3} + \frac{C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1^2}{8\epsilon m^2 m_+^3} + \frac{3C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1^2}{32\epsilon^2 m^2 m_+^3} + \frac{C_A^2 m_2^2 \pi^2 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1^2}{64m^2 m_+^3} - \\
& \frac{C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_2 m_1 \mathcal{L}_2 m_2 m_1^2}{8m^2 m_+^3} - \frac{3C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_2 m_1 \mathcal{L}_2 m_2 m_1^2}{32\epsilon m^2 m_+^3} - \frac{7C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_2^3 m_1^2}{64m^2 m_+^2} + \frac{C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_2^3 m_1^2}{64m^2 m_+^2} + \\
& \frac{3C_A^2 m_2^3 Y_f^3 Y_s \mathcal{L}_2 m_1 m_1}{32m^2 m_+^3} + \frac{9C_A^2 m_2^3 Y_f^3 Y_s \mathcal{L}_2 m_1 m_1}{64\epsilon m^2 m_+^3} - \frac{5C_A^2 m_2^3 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1}{32m^2 m_+^3} - \frac{3C_A^2 m_2^3 Y_f^3 Y_s \mathcal{L}_2 m_2 m_1}{64\epsilon m^2 m_+^3} +
\end{aligned}$$

$$\begin{aligned}
& \frac{3C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_{m_1} c_{m_2}^2 m_1}{64m_-^2 m_+^2} + \frac{9C_A^2 m_2 Y_f^3 Y_s m_1}{4m_- m_+^2} + \frac{23C_A^2 m_2 Y_f^3 Y_s m_1}{32\epsilon m_- m_+^2} + \frac{3C_A^2 m_2 Y_f^3 Y_s m_1}{16\epsilon^2 m_- m_+^2} + \frac{C_A^2 m_2 \pi^2 Y_f^3 Y_s m_1}{32m_- m_+^2} + \\
& \frac{15C_A^2 m_2^3 Y_f^3 Y_s \mathcal{L}_{m_1} m_1}{32m_-^2 m_+^3} + \frac{C_A^2 m_2^3 Y_f^3 Y_s \mathcal{L}_{m_1} m_1}{16\epsilon m_-^2 m_+^3} - \frac{3C_A^2 m_2^3 Y_f^3 Y_s \mathcal{L}_{m_1} m_1}{32\epsilon^2 m_-^2 m_+^3} - \frac{C_A^2 m_2^3 \pi^2 Y_f^3 Y_s \mathcal{L}_{m_1} m_1}{64m_-^2 m_+^3} + \\
& \frac{3C_A^2 m_2^2 Y_f^3 Y_s \mathcal{L}_{m_1} \mathcal{L}_{m_2} m_1}{64m_-^2 m_+^2} + \frac{31C_A^2 m_2^3 Y_f^3 Y_s \mathcal{L}_{m_2} m_1}{32m_-^2 m_+^3} + \frac{5C_A^2 m_2^3 Y_f^3 Y_s \mathcal{L}_{m_2} m_1}{16\epsilon m_-^2 m_+^3} + \frac{3C_A^2 m_2^3 \pi^2 Y_f^3 Y_s \mathcal{L}_{m_2} m_1}{32\epsilon^2 m_-^2 m_+^3} + \\
& \frac{C_A^2 m_2^3 \pi^2 Y_f^3 Y_s \mathcal{L}_{m_2} m_1}{64m_-^2 m_+^3} - \frac{5C_A^2 m_2^3 Y_f^3 Y_s \mathcal{L}_{m_1} \mathcal{L}_{m_2} m_1}{16m_-^2 m_+^3} - \frac{3C_A^2 m_2^3 Y_f^3 Y_s \mathcal{L}_{m_1} \mathcal{L}_{m_2} m_1}{32\epsilon m_-^2 m_+^3} - \frac{27}{32} Y_f^4 C_A^2 \mathcal{L}_{m_1} + \frac{27Y_f^4 C_A^2}{64\epsilon} + \frac{45}{32} Y_f^4 C_A^2 - \\
& \frac{Y_s^4 C_A^2 \mathcal{L}_{m_2}}{16m_+^4} + \frac{Y_s^4 C_A^2}{32\epsilon m_+^4} + \frac{3Y_s^4 C_A^2}{32m_+^4}
\end{aligned} \tag{B.6}$$

## B.2 Matching at $\mu \sim M$ :

The one-loop bosonic mass renormalisation contributions are non-trivial as they involve two mass scales,  $M_H$  and  $M_W$ , and are required to  $\mathcal{O}(\epsilon^2)$  to provide the correct contributions at two-loop order.

$$\begin{aligned}
\frac{\delta M_W^2}{\alpha M_W^2} = & -\frac{1}{48}\epsilon \log^2(r(r-s))r^6 - \frac{1}{48}\epsilon \log^2(r(r+s))r^6 + \frac{1}{48}\epsilon \log^2(u)r^6 + \frac{1}{48}\epsilon \log^2(v)r^6 - \frac{1}{9}\epsilon \log(r^2)r^6 + \frac{1}{24}\epsilon \mathcal{L}_{M_W} \log(r^2)r^6 - \\
& \frac{1}{24}\log(r^2)r^6 + \frac{1}{24}\epsilon \log(r^2)\log(r(r-s))r^6 + \frac{1}{24}\epsilon \log(r^2)\log(r(r+s))r^6 - \frac{1}{24}\epsilon \log\left(\frac{s-r}{2s}\right)\log(r(r+s))r^6 - \\
& \frac{1}{24}\epsilon \log(r(r-s))\log\left(\frac{r+s}{2s}\right)r^6 + \frac{1}{24}\epsilon \log\left(-\frac{u}{2rs}\right)\log(v)r^6 + \frac{1}{24}\epsilon \log(u)\log\left(\frac{v}{2rs}\right)r^6 - \frac{1}{24}\epsilon \text{Li}_2\left(\frac{r+s}{2s}\right)r^6 + \\
& \frac{1}{24}\epsilon \text{Li}_2\left(-\frac{u}{2rs}\right)r^6 + \frac{1}{24}\epsilon \text{Li}_2\left(\frac{v}{2rs}\right)r^6 + \frac{1}{48}\epsilon s \log^2(r(r-s))r^5 - \frac{1}{48}\epsilon s \log^2(r(r+s))r^5 - \frac{1}{48}\epsilon s \log^2(u)r^5 + \frac{1}{48}\epsilon s \log^2(v)r^5 - \\
& \frac{1}{24}\epsilon s \log(r^2)\log(r(r-s))r^5 - \frac{1}{9}\epsilon s \log(s-r)r^5 - \frac{1}{24}\epsilon s \log(s-r)r^5 + \frac{1}{24}\epsilon s \mathcal{L}_{M_W} \log(s-r)r^5 + \frac{1}{9}\epsilon s \log(r+s)r^5 + \\
& \frac{1}{24}\epsilon s \log(r+s)r^5 - \frac{1}{24}\epsilon s \mathcal{L}_{M_W} \log(r+s)r^5 + \frac{1}{24}\epsilon s \log(r^2)\log(r(r+s))r^5 - \frac{1}{24}\epsilon s \log\left(\frac{s-r}{2s}\right)\log(r(r+s))r^5 + \\
& \frac{1}{24}\epsilon s \log(r(r-s))\log\left(\frac{r+s}{2s}\right)r^5 + \frac{1}{144}\epsilon s \log(64)\log(u)r^5 + \frac{1}{9}\epsilon s \log(-v)r^5 + \frac{1}{24}\epsilon s \log(-v)r^5 - \frac{1}{24}\epsilon s \mathcal{L}_{M_W} \log(-v)r^5 + \\
& \frac{1}{24}\epsilon s \log\left(-\frac{u}{2rs}\right)\log(v)r^5 - \frac{1}{9}\epsilon s \log\left(\frac{v}{rs}\right)r^5 - \frac{1}{24}\epsilon s \log\left(\frac{v}{rs}\right)r^5 + \frac{1}{24}\epsilon s \mathcal{L}_{M_W} \log\left(\frac{v}{rs}\right)r^5 - \frac{1}{24}\epsilon s \log(u)\log\left(\frac{v}{rs}\right)r^5 - \\
& \frac{1}{24}\epsilon s \text{Li}_2\left(\frac{r+s}{2s}\right)r^5 - \frac{1}{24}\epsilon s \text{Li}_2\left(-\frac{u}{2rs}\right)r^5 + \frac{1}{24}\epsilon s \text{Li}_2\left(\frac{v}{2rs}\right)r^5 - \frac{1}{24}\epsilon \mathcal{L}_{M_H}^2 r^4 + \frac{1}{24}\epsilon \mathcal{L}_{M_W}^2 r^4 + \frac{1}{24}\epsilon \log^2(r^2)r^4 + \\
& \frac{1}{8}\epsilon \log^2(r(r-s))r^4 + \frac{1}{8}\epsilon \log^2(r(r+s))r^4 - \frac{1}{8}\epsilon \log^2(u)r^4 - \frac{1}{8}\epsilon \log^2(v)r^4 + \frac{11\epsilon r^4}{36} + \frac{5}{36}\epsilon \mathcal{L}_{M_H} r^4 + \frac{1}{12}\mathcal{L}_{M_H} r^4 - \frac{2}{9}\epsilon \mathcal{L}_{M_W} r^4 - \\
& \frac{1}{12}\mathcal{L}_{M_W} r^4 + \frac{4}{9}\epsilon \log(r^2)r^4 - \frac{1}{6}\epsilon \mathcal{L}_{M_W} \log(r^2)r^4 + \frac{1}{6}\log(r^2)r^4 - \frac{1}{4}\epsilon \log(r^2)\log(r(r-s))r^4 - \frac{1}{4}\epsilon \log(r^2)\log(r(r+s))r^4 + \\
& \frac{1}{4}\epsilon \log\left(\frac{s-r}{2s}\right)\log(r(r+s))r^4 + \frac{1}{4}\epsilon \log(r(r-s))\log\left(\frac{r+s}{2s}\right)r^4 + \frac{1}{12}\epsilon \log(2)\log(u)r^4 - \frac{1}{4}\epsilon \log\left(-\frac{u}{2rs}\right)\log(v)r^4 - \\
& \frac{1}{6}\epsilon \log(u)\log\left(\frac{v}{2rs}\right)r^4 - \frac{1}{12}\epsilon \log(u)\log\left(\frac{v}{rs}\right)r^4 - \frac{1}{24}\epsilon u \text{Li}_2\left(\frac{s-r}{2s}\right)r^4 + \frac{1}{4}\epsilon \text{Li}_2\left(\frac{r+s}{2s}\right)r^4 - \frac{1}{4}\epsilon \text{Li}_2\left(-\frac{u}{2rs}\right)r^4 - \\
& \frac{1}{4}\epsilon \text{Li}_2\left(\frac{v}{2rs}\right)r^4 + \frac{1}{144}\epsilon \pi^2 r^4 + \frac{r^4}{12} - \frac{1}{12}\epsilon s \log^2(r(r-s))r^3 + \frac{1}{12}\epsilon s \log^2(r(r+s))r^3 + \frac{1}{12}\epsilon s \log^2(u)r^3 - \frac{1}{12}\epsilon s \log^2(v)r^3 + \\
& \frac{1}{6}\epsilon s \log(r^2)\log(r(r-s))r^3 + \frac{4}{9}\epsilon s \log(s-r)r^3 + \frac{1}{6}\epsilon s \log(s-r)r^3 - \frac{1}{6}\epsilon s \mathcal{L}_{M_W} \log(s-r)r^3 - \frac{4}{9}\epsilon s \log(r+s)r^3 - \\
& \frac{1}{6}\epsilon s \log(r+s)r^3 + \frac{1}{6}\epsilon s \mathcal{L}_{M_W} \log(r+s)r^3 - \frac{1}{6}\epsilon s \log(r^2)\log(r(r+s))r^3 + \frac{1}{6}\epsilon s \log\left(\frac{s-r}{2s}\right)\log(r(r+s))r^3 - \\
& \frac{1}{6}\epsilon s \log(r(r-s))\log\left(\frac{r+s}{2s}\right)r^3 - \frac{1}{36}\epsilon s \log(64)\log(u)r^3 - \frac{4}{9}\epsilon s \log(-v)r^3 - \frac{1}{6}\epsilon s \log(-v)r^3 + \frac{1}{6}\epsilon s \mathcal{L}_{M_W} \log(-v)r^3 - \\
& \frac{1}{6}\epsilon s \log\left(-\frac{u}{2rs}\right)\log(v)r^3 + \frac{4}{9}\epsilon s \log\left(\frac{v}{rs}\right)r^3 + \frac{1}{6}\epsilon s \log\left(\frac{v}{rs}\right)r^3 - \frac{1}{6}\epsilon s \mathcal{L}_{M_W} \log\left(\frac{v}{rs}\right)r^3 + \frac{1}{6}\epsilon s \log(u)\log\left(\frac{v}{rs}\right)r^3 + \\
& \frac{1}{6}\epsilon s \text{Li}_2\left(\frac{r+s}{2s}\right)r^3 + \frac{1}{6}\epsilon s \text{Li}_2\left(-\frac{u}{2rs}\right)r^3 - \frac{1}{6}\epsilon s \text{Li}_2\left(\frac{v}{2rs}\right)r^3 - \frac{1}{8}\epsilon \mathcal{L}_{M_W}^2 r^2 - \frac{1}{6}\epsilon \log^2(r^2)r^2 - \frac{5}{12}\epsilon \log^2(r(r-s))r^2 - \\
& \frac{5}{12}\epsilon \log^2(r(r+s))r^2 + \frac{5}{12}\epsilon \log^2(u)r^2 + \frac{5}{12}\epsilon \log^2(v)r^2 - \frac{7\epsilon r^2}{4} + \frac{3}{4}\epsilon \mathcal{L}_{M_W} r^2 + \frac{1}{4}\mathcal{L}_{M_W} r^2 - \epsilon \log(r^2)r^2 + \\
& \frac{1}{2}\epsilon \mathcal{L}_{M_W} \log(r^2)r^2 - \frac{1}{2}\log(r^2)r^2 + \frac{5}{6}\epsilon \log(r^2)\log(r(r-s))r^2 + \frac{5}{6}\epsilon \log(r^2)\log(r(r+s))r^2 - \\
& \frac{5}{6}\epsilon \log\left(\frac{s-r}{2s}\right)\log(r(r+s))r^2 - \frac{5}{6}\epsilon \log(r(r-s))\log\left(\frac{r+s}{2s}\right)r^2 - \frac{1}{3}\epsilon \log(2)\log(u)r^2 + \frac{5}{6}\epsilon \log\left(-\frac{u}{2rs}\right)\log(v)r^2 + \\
& \frac{1}{2}\epsilon \log(u)\log\left(\frac{v}{2rs}\right)r^2 + \frac{1}{3}\epsilon \log(u)\log\left(\frac{v}{rs}\right)r^2 + \frac{1}{6}\epsilon u \text{Li}_2\left(\frac{s-r}{2s}\right)r^2 - \frac{5}{6}\epsilon \text{Li}_2\left(\frac{r+s}{2s}\right)r^2 + \frac{5}{6}\epsilon \text{Li}_2\left(-\frac{u}{2rs}\right)r^2 + \\
& \frac{5}{6}\epsilon \text{Li}_2\left(\frac{v}{2rs}\right)r^2 - \frac{r^2}{4\epsilon} - \frac{7}{144}\epsilon \pi^2 r^2 - \frac{3r^2}{4} + \frac{1}{4}\epsilon s \log^2(r(r-s))r - \frac{1}{4}\epsilon s \log^2(r(r+s))r - \frac{1}{4}\epsilon s \log^2(u)r + \frac{1}{4}\epsilon s \log^2(v)r - \\
& \frac{1}{2}\epsilon s \log(r^2)\log(r(r-s))r - \epsilon s \log(s-r)r - \frac{1}{2}\epsilon s \log(s-r)r + \frac{1}{2}\epsilon s \mathcal{L}_{M_W} \log(s-r)r + \epsilon s \log(r+s)r + \frac{1}{2}\epsilon s \log(r+s)r -
\end{aligned}$$



$$\begin{aligned}
& \frac{1}{2}\varepsilon s \mathcal{L}M_W \log(r+s)r + \frac{1}{2}\varepsilon s \log(r^2) \log(r(r+s))r - \frac{1}{2}\varepsilon s \log\left(\frac{s-r}{2s}\right) \log(r(r+s))r + \frac{1}{2}\varepsilon s \log(r(r-s)) \log\left(\frac{r+s}{2s}\right)r + \\
& \frac{1}{12}\varepsilon s \log(64) \log(u)r + \varepsilon s \log(-v)r + \frac{1}{2}s \log(-v)r - \frac{1}{2}\varepsilon s \mathcal{L}M_W \log(-v)r + \frac{1}{2}\varepsilon s \log\left(-\frac{u}{2rs}\right) \log(v)r - \varepsilon s \log\left(\frac{v}{s}\right)r - \\
& \frac{1}{2}s \log\left(\frac{v}{s}\right)r + \frac{1}{2}\varepsilon s \mathcal{L}M_W \log\left(\frac{v}{s}\right)r - \frac{1}{2}\varepsilon s \log(u) \log\left(\frac{v}{rs}\right)r - \frac{1}{2}\varepsilon s \text{Li}_2\left(\frac{r+s}{2s}\right)r - \frac{1}{2}\varepsilon s \text{Li}_2\left(-\frac{u}{2rs}\right)r + \frac{1}{2}\varepsilon s \text{Li}_2\left(\frac{v}{2rs}\right)r - \\
& \frac{47}{16}C_A \varepsilon \mathcal{L}M_W^2 + \frac{5}{12}\varepsilon \mathcal{L}M_W^2 + \frac{8}{3}\varepsilon n_f T_f \mathcal{L}M_W^2 + \frac{1}{2}\varepsilon \log^2(r^2) + \frac{1}{2}\varepsilon \log^2(r(r-s)) + \frac{1}{2}\varepsilon \log^2(r(r+s)) - \frac{1}{2}\varepsilon \log^2(u) - \frac{1}{2}\varepsilon \log^2(v) - \\
& \frac{115C_A}{12} - \frac{623C_A \varepsilon}{36} + \frac{197\varepsilon}{54} + \frac{112}{27}\varepsilon n_f T_f + \frac{4n_f T_f}{3\varepsilon} + \frac{20n_f T_f}{9} - \frac{7}{9}\varepsilon n_f \pi^2 T_f - \frac{20}{9}i\varepsilon n_f \pi T_f - \frac{4}{3}i n_f \pi T_f + \frac{47}{8}C_A \mathcal{L}M_W + \\
& \frac{115}{12}C_A \varepsilon \mathcal{L}M_W - \frac{31}{18}\varepsilon \mathcal{L}M_W - \frac{11}{8}\sqrt{3}C_A \varepsilon \pi \mathcal{L}M_W - \frac{5\mathcal{L}M_W \varepsilon \pi}{6} - \frac{40}{9}\varepsilon n_f T_f \mathcal{L}M_W - \frac{8}{3}n_f T_f \mathcal{L}M_W + \frac{8}{3}i\varepsilon n_f \pi T_f \mathcal{L}M_W - \\
& \varepsilon \log(r^2) \log(r(r-s)) - \varepsilon \log(r^2) \log(r(r+s)) + \varepsilon \log\left(\frac{s-r}{2s}\right) \log(r(r+s)) + \varepsilon \log(r(r-s)) \log\left(\frac{r+s}{2s}\right) + \frac{20}{9}\varepsilon n_f T_f \log(\mu^2) - \\
& \frac{8}{3}\varepsilon n_f T_f \mathcal{L}M_W \log(\mu^2) + \varepsilon \log(2) \log(u) - \varepsilon \log\left(-\frac{u}{2rs}\right) \log(v) - \varepsilon \log(u) \log\left(\frac{v}{rs}\right) - \frac{1}{2}\varepsilon u \text{Li}_2\left(\frac{s-r}{2s}\right) + \varepsilon \text{Li}_2\left(\frac{r+s}{2s}\right) - \\
& \varepsilon \text{Li}_2\left(-\frac{u}{2rs}\right) - \varepsilon \text{Li}_2\left(\frac{v}{2rs}\right) - \frac{47C_A}{8\varepsilon} + \frac{5}{6\varepsilon} - \frac{33}{8}i\sqrt{3}C_A \varepsilon \text{Li}_2\left(\frac{1}{6}(3+i\sqrt{3})\right) + \frac{33}{8}i\sqrt{3}C_A \varepsilon \text{Li}_2\left(\frac{1}{6}(3-i\sqrt{3})\right) - \\
& \frac{11}{16}\sqrt{3}C_A \varepsilon \pi \log(3) - \frac{5}{6}C_A \varepsilon \pi^2 + \frac{11\varepsilon \pi^2}{72} + \frac{11}{8}\sqrt{3}C_A \pi + \frac{8C_A \varepsilon \pi}{\sqrt{3}} + \frac{31}{18} \tag{B.7}
\end{aligned}$$

$$\begin{aligned}
\frac{\delta M_H^2}{aM_H^2} &= \frac{9}{64}\varepsilon \mathcal{L}M_H^2 r^2 + \frac{1}{128}C_A C_F \varepsilon \mathcal{L}M_W^2 r^2 + \frac{1}{256}C_A C_F \varepsilon \log^2(r(r-s))r^2 - \frac{1}{256}C_A C_F \varepsilon \log^2(r(s-r))r^2 \\
& - \frac{1}{256}C_A C_F \varepsilon \log^2(-r(r+s))r^2 + \frac{1}{256}C_A C_F \varepsilon \log^2(r(r+s))r^2 + \frac{1}{32}C_A C_F r^2 + \frac{1}{16}C_A C_F \varepsilon r^2 + \frac{9\varepsilon r^2}{8} - \frac{9}{16}\varepsilon \mathcal{L}M_H r^2 \\
& + \frac{3}{32}\sqrt{3}\varepsilon \pi \mathcal{L}M_H r^2 - \frac{9}{32}\mathcal{L}M_H r^2 - \frac{1}{64}C_A C_F \mathcal{L}M_W r^2 - \frac{1}{32}C_A C_F \varepsilon \mathcal{L}M_W r^2 + \frac{1}{128}C_A C_F \varepsilon \log(2) \log(r(s-r))r^2 \\
& - \frac{1}{128}C_A C_F \varepsilon \log\left(\frac{s-r}{2s}\right) \log(-r(r+s))r^2 + \frac{1}{128}C_A C_F \varepsilon \log\left(\frac{s-r}{2s}\right) \log(r(r+s))r^2 + \frac{1}{128}C_A C_F \varepsilon \log(r(r-s)) \log\left(\frac{r+s}{2s}\right)r^2 \\
& - \frac{1}{128}C_A C_F \varepsilon \log(r(s-r)) \log\left(\frac{r+s}{s}\right)r^2 + \frac{C_A C_F r^2}{64\varepsilon} + \frac{9r^2}{32\varepsilon} + \frac{9}{32}i\sqrt{3}\varepsilon \text{Li}_2\left(\frac{1}{6}(3+i\sqrt{3})\right)r^2 - \frac{9}{32}i\sqrt{3}\varepsilon \text{Li}_2\left(\frac{1}{6}(3-i\sqrt{3})\right)r^2 \\
& + \frac{3}{64}\sqrt{3}\varepsilon \pi \log(3)r^2 + \frac{1}{384}C_A C_F \varepsilon \pi^2 r^2 + \frac{3}{64}\varepsilon \pi^2 r^2 - \frac{3}{16}\sqrt{3}\varepsilon \pi r^2 - \frac{3}{32}\sqrt{3}\pi r^2 + \frac{9r^2}{16} + \frac{1}{256}C_A C_F \varepsilon s \log^2(r(r-s))r \\
& + \frac{1}{256}C_A C_F \varepsilon s \log^2(r(s-r))r - \frac{1}{256}C_A C_F \varepsilon s \log^2(-r(r+s))r - \frac{1}{256}C_A C_F \varepsilon s \log^2(r(r+s))r - \frac{1}{64}C_A C_F s \log(s-r)r \\
& - \frac{1}{32}C_A C_F \varepsilon s \log(s-r)r + \frac{1}{64}C_A C_F \varepsilon s \mathcal{L}M_W \log(s-r)r + \frac{1}{64}C_A C_F s \log(r+s)r + \frac{1}{32}C_A C_F \varepsilon s \log(r+s)r \\
& - \frac{1}{64}C_A C_F \varepsilon s \mathcal{L}M_W \log(r+s)r - \frac{1}{128}C_A C_F \varepsilon s \log\left(\frac{s-r}{2s}\right) \log(-r(r+s))r - \frac{1}{128}C_A C_F \varepsilon s \log\left(\frac{s-r}{2s}\right) \log(r(r+s))r \\
& + \frac{1}{128}C_A C_F \varepsilon s \log(r(r-s)) \log\left(\frac{r+s}{2s}\right)r + \frac{1}{128}C_A C_F \varepsilon s \log(r(s-r)) \log\left(\frac{r+s}{2s}\right)r + \frac{1}{64}C_A C_F \varepsilon s \text{Li}_2\left(\frac{s-r}{2s}\right)r \\
& - \frac{1}{64}C_A C_F \varepsilon s \text{Li}_2\left(\frac{r+s}{2s}\right)r - \frac{1}{8}C_A C_F \varepsilon \mathcal{L}M_W - \frac{1}{16}C_A C_F \varepsilon \log^2(r(r-s)) + \frac{1}{16}C_A C_F \varepsilon \log^2(r(s-r)) \\
& + \frac{1}{16}C_A C_F \varepsilon \log^2(-r(r+s)) - \frac{1}{16}C_A C_F \varepsilon \log^2(r(r+s)) - \frac{C_A C_F}{2} - C_A C_F \varepsilon + \frac{1}{4}C_A C_F \mathcal{L}M_W + \frac{1}{2}C_A C_F \varepsilon \mathcal{L}M_W \\
& - \frac{1}{8}C_A C_F \varepsilon \log(2) \log(r(s-r)) + \frac{1}{8}C_A C_F \varepsilon \log\left(\frac{s-r}{2s}\right) \log(-r(r+s)) - \frac{1}{8}C_A C_F \varepsilon \log\left(\frac{s-r}{2s}\right) \log(r(r+s)) \\
& - \frac{1}{8}C_A C_F \varepsilon \log(r(r-s)) \log\left(\frac{r+s}{2s}\right) + \frac{1}{8}C_A C_F \varepsilon \log(r(s-r)) \log\left(\frac{r+s}{s}\right) - \frac{C_A C_F}{4\varepsilon} - \frac{1}{24}C_A C_F \varepsilon \pi^2 - \frac{C_A C_F \varepsilon s \log^2(r(r-s))}{16r} \\
& - \frac{C_A C_F \varepsilon s \log^2(r(s-r))}{16r} + \frac{C_A C_F \varepsilon s \log^2(-r(r+s))}{16r} + \frac{C_A C_F \varepsilon s \log^2(r(r+s))}{16r} + \frac{C_A C_F s \log(s-r)}{4r} + \frac{C_A C_F \varepsilon s \log(s-r)}{2r} \\
& - \frac{C_A C_F \varepsilon s \mathcal{L}M_W \log(s-r)}{4r} - \frac{C_A C_F s \log(r+s)}{4r} - \frac{C_A C_F \varepsilon s \log(r+s)}{2r} + \frac{C_A C_F \varepsilon s \mathcal{L}M_W \log(r+s)}{4r} \\
& + \frac{C_A C_F \varepsilon s \log\left(\frac{s-r}{2s}\right) \log(-r(r+s))}{8r} + \frac{C_A C_F \varepsilon s \log\left(\frac{s-r}{2s}\right) \log(r(r+s))}{8r} - \frac{C_A C_F \varepsilon s \log(r(r-s)) \log\left(\frac{r+s}{2s}\right)}{8r} \\
& - \frac{C_A C_F \varepsilon s \log(r(s-r)) \log\left(\frac{r+s}{2s}\right)}{8r} - \frac{C_A C_F \varepsilon s \text{Li}_2\left(\frac{s-r}{2s}\right)}{4r} + \frac{C_A C_F \varepsilon s \text{Li}_2\left(\frac{r+s}{2s}\right)}{4r} + \frac{3C_A C_F \varepsilon \mathcal{L}M_W^2}{8r^2} + \frac{7C_A C_F \varepsilon \log^2(r(r-s))}{32r^2} \\
& - \frac{7C_A C_F \varepsilon \log^2(r(s-r))}{32r^2} - \frac{7C_A C_F \varepsilon \log^2(-r(r+s))}{32r^2} + \frac{7C_A C_F \varepsilon \log^2(r(r+s))}{32r^2} + \frac{9C_A C_F}{8r^2} + \frac{19C_A C_F \varepsilon}{8r^2} - \frac{3C_A C_F \mathcal{L}M_W}{4r^2} \\
& - \frac{9C_A C_F \varepsilon \mathcal{L}M_W}{8r^2} + \frac{7C_A C_F \varepsilon \log(2) \log(r(s-r))}{16r^2} - \frac{7C_A C_F \varepsilon \log\left(\frac{s-r}{2s}\right) \log(-r(r+s))}{16r^2} + \frac{7C_A C_F \varepsilon \log\left(\frac{s-r}{2s}\right) \log(r(r+s))}{16r^2} \\
& + \frac{7C_A C_F \varepsilon \log(r(r-s)) \log\left(\frac{r+s}{2s}\right)}{16r^2} - \frac{7C_A C_F \varepsilon \log(r(s-r)) \log\left(\frac{r+s}{s}\right)}{16r^2} + \frac{3C_A C_F}{4\varepsilon r^2} + \frac{13C_A C_F \varepsilon \pi^2}{96r^2} + \frac{7C_A C_F \varepsilon s \log^2(r(r-s))}{32r^3} \\
& + \frac{7C_A C_F \varepsilon s \log^2(r(s-r))}{32r^3} - \frac{7C_A C_F \varepsilon s \log^2(-r(r+s))}{32r^3} - \frac{7C_A C_F \varepsilon s \log^2(r(r+s))}{32r^3} - \frac{7C_A C_F s \log(s-r)}{8r^3} \\
& - \frac{5C_A C_F \varepsilon s \log(s-r)}{4r^3} + \frac{7C_A C_F \varepsilon s \mathcal{L}M_W \log(s-r)}{8r^3} + \frac{7C_A C_F s \log(r+s)}{8r^3} + \frac{5C_A C_F \varepsilon s \log(r+s)}{4r^3}
\end{aligned}$$

$$\begin{aligned}
& -\frac{7C_A C_F \varepsilon s \mathcal{L}_{M_W} \log(r+s)}{8r^3} - \frac{7C_A C_F \varepsilon s \log\left(\frac{s-r}{2s}\right) \log(-r(r+s))}{16r^3} - \frac{7C_A C_F \varepsilon s \log\left(\frac{s-r}{2s}\right) \log(r(r+s))}{16r^3} \\
& + \frac{7C_A C_F \varepsilon s \log(r(r-s)) \log\left(\frac{r+s}{2s}\right)}{16r^3} + \frac{7C_A C_F \varepsilon s \log(r(s-r)) \log\left(\frac{r+s}{2s}\right)}{16r^3} + \frac{7C_A C_F \varepsilon s \text{Li}_2\left(\frac{s-r}{2s}\right)}{8r^3} - \frac{7C_A C_F \varepsilon s \text{Li}_2\left(\frac{r+s}{2s}\right)}{8r^3}
\end{aligned} \tag{B.8}$$

where  $r = M_H/M_W$ ,  $s = \sqrt{r^2 - 4}$ ,  $u = r^2 - rs - 2$  and  $v = r^2 + rs - 2$  and  $\text{Li}_2(z)$  is the dilogarithm function,

$$\text{Li}_2(z) = \sum_{k=1}^{\infty} \frac{z^k}{k^2} : \quad z \in \mathbb{C}. \tag{B.9}$$

These corrections are provided for two-loop order matching contributions for the heavy-heavy graphs at the scale  $\mu \sim M_{W,H}$ , we provide these contributions here.

$$\begin{aligned}
\Delta U_1^{(2)} = & -\frac{1}{12} C_A C_F \log(r^2) r^6 + \frac{1}{12} C_A C_F w \log(r^2) r(w) r^6 - \frac{1}{12} C_A C_F s \log(s-r) r^5 + \frac{1}{12} C_A C_F s \log(r+s) r^5 \\
& + \frac{1}{12} C_A C_F s \log(-v) r^5 - \frac{1}{12} C_A C_F s \log\left(\frac{v}{s}\right) r^5 + \frac{1}{12} C_A C_F s w \log(s-r) r(w) r^5 - \frac{1}{12} C_A C_F s w \log(r+s) r(w) r^5 \\
& - \frac{1}{12} C_A C_F s w \log(-v) r(w) r^5 + \frac{1}{12} C_A C_F s w \log\left(\frac{v}{s}\right) r(w) r^5 + \frac{1}{6} C_A C_F r^4 + \frac{1}{6} C_A C_F \mathcal{L}_{M_H} r^4 - \frac{1}{6} C_A C_F \mathcal{L}_{M_W} r^4 \\
& + \frac{1}{3} C_A C_F \log(r^2) r^4 - \frac{1}{6} C_A C_F w r(w) r^4 - \frac{1}{6} C_A C_F w \mathcal{L}_{M_H} r(w) r^4 + \frac{1}{6} C_A C_F w \mathcal{L}_{M_W} r(w) r^4 - \frac{1}{3} C_A C_F w \log(r^2) r(w) r^4 \\
& + \frac{1}{3} C_A C_F s \log(s-r) r^3 - \frac{1}{3} C_A C_F s \log(r+s) r^3 - \frac{1}{3} C_A C_F s \log(-v) r^3 + \frac{1}{3} C_A C_F s \log\left(\frac{v}{s}\right) r^3 - \frac{1}{3} C_A C_F s w \log(s-r) r(w) r^3 \\
& + \frac{1}{3} C_A C_F s w \log(r+s) r(w) r^3 + \frac{1}{3} C_A C_F s w \log(-v) r(w) r^3 - \frac{1}{3} C_A C_F s w \log\left(\frac{v}{s}\right) r(w) r^3 - \frac{1}{16} C_A C_F^2 Y_f^2 r^2 - \frac{9}{16} C_A Y_f^2 r^2 \\
& - \frac{1}{32} C_F Y_f^2 r^2 - \frac{C_A C_F^2 Y_f^2 r^2}{64\varepsilon} - \frac{9C_A Y_f^2 r^2}{64\varepsilon} - \frac{C_F Y_f^2 r^2}{128\varepsilon} - \frac{5}{2} C_A C_F r^2 + \frac{1}{64} C_A C_F^2 Y_f^2 \mathcal{L}_{M_H} r^2 + \frac{9}{32} C_A Y_f^2 \mathcal{L}_{M_H} r^2 \\
& + \frac{1}{128} C_F Y_f^2 \mathcal{L}_{M_H} r^2 + \frac{1}{64} C_A C_F^2 Y_f^2 \mathcal{L}_{M_W} r^2 + \frac{1}{128} C_F Y_f^2 \mathcal{L}_{M_W} r^2 + C_A C_F \mathcal{L}_{M_W} r^2 - C_A C_F \log(r^2) r^2 \\
& + \frac{1}{32} C_A C_F^2 Y_f^2 r(w) r^2 + \frac{9}{32} C_A Y_f^2 r(w) r^2 + \frac{1}{64} C_F Y_f^2 r(w) r^2 + \frac{C_A C_F^2 Y_f^2 r(w) r^2}{64\varepsilon} + \frac{9C_A Y_f^2 r(w) r^2}{64\varepsilon} + \frac{C_F Y_f^2 r(w) r^2}{128\varepsilon} \\
& + \frac{3}{2} C_A C_F w r(w) r^2 + \frac{C_A C_F w r(w) r^2}{2\varepsilon} - \frac{1}{64} C_A C_F^2 Y_f^2 \mathcal{L}_{M_H} r(w) r^2 - \frac{9}{32} C_A Y_f^2 \mathcal{L}_{M_H} r(w) r^2 - \frac{1}{128} C_F Y_f^2 \mathcal{L}_{M_H} r(w) r^2 \\
& - \frac{1}{64} C_A C_F^2 Y_f^2 \mathcal{L}_{M_W} r(w) r^2 - \frac{1}{128} C_F Y_f^2 \mathcal{L}_{M_W} r(w) r^2 - C_A C_F w \mathcal{L}_{M_W} r(w) r^2 + C_A C_F w \log(r^2) r(w) r^2 \\
& - \frac{3}{64} \sqrt{3} C_A Y_f^2 \pi r(w) r^2 - \frac{C_A C_F r^2}{2\varepsilon} + \frac{3}{64} \sqrt{3} C_A Y_f^2 \pi r^2 + \frac{1}{64} C_A C_F^2 Y_f^2 s \log(s-r) r \\
& + \frac{1}{128} C_F Y_f^2 s \log(s-r) r - C_A C_F s \log(s-r) r - \frac{1}{64} C_A C_F^2 Y_f^2 s \log(r+s) r - \frac{1}{128} C_F Y_f^2 s \log(r+s) r \\
& + C_A C_F s \log(r+s) r + C_A C_F s \log(-v) r - C_A C_F s \log\left(\frac{v}{s}\right) r - \frac{1}{64} C_A C_F^2 Y_f^2 s \log(s-r) r(w) r - \frac{1}{128} C_F Y_f^2 s \log(s-r) r(w) r \\
& + C_A C_F s w \log(s-r) r(w) r + \frac{1}{64} C_A C_F^2 Y_f^2 s \log(r+s) r(w) r + \frac{1}{128} C_F Y_f^2 s \log(r+s) r(w) r \\
& - C_A C_F s w \log(r+s) r(w) r - C_A C_F s w \log(-v) r(w) r + C_A C_F s w \log\left(\frac{v}{s}\right) r(w) r - \frac{256}{3} C_A C_F^2 + C_A C_F^2 Y_f^2 + \frac{C_F Y_f^2}{2} \\
& + \frac{C_A C_F^2 Y_f^2}{4\varepsilon} + \frac{C_F Y_f^2}{8\varepsilon} + \frac{1}{4} C_A Y_f^2 \beta_0 \mathcal{L}_{M_H}^2 - C_A C_F \beta_0 \mathcal{L}_{M_W}^2 + \frac{61C_A C_F}{9} - \frac{128C_F}{3} + \frac{88}{9} C_A C_F n_f T_f + \frac{8C_A C_F n_f T_f}{3\varepsilon} \\
& - \frac{8}{3} i C_A C_F n_f \pi T_f + \frac{C_A Y_f^2 \beta_0}{2\varepsilon^2} - \frac{2C_A C_F \beta_0}{\varepsilon^2} + \frac{1}{24} C_A Y_f^2 \pi^2 \beta_0 - \frac{1}{6} C_A C_F \pi^2 \beta_0 - \frac{1}{4} C_A C_F^2 Y_f^2 \mathcal{L}_{M_H} - \frac{1}{8} C_F Y_f^2 \mathcal{L}_{M_H} \\
& - C_A Y_f^2 \beta_0 \mathcal{L}_{M_H} - \frac{C_A Y_f^2 \beta_0 \mathcal{L}_{M_H}}{2\varepsilon} + 47C_A C_F^2 \mathcal{L}_{M_W} - \frac{1}{4} C_A C_F^2 Y_f^2 \mathcal{L}_{M_W} - \frac{1}{8} C_F Y_f^2 \mathcal{L}_{M_W} - \frac{10}{3} C_A C_F \mathcal{L}_{M_W} + \frac{47}{2} C_F \mathcal{L}_{M_W} \\
& - \frac{16}{3} C_A C_F n_f T_f \mathcal{L}_{M_W} + 4C_A C_F \beta_0 \mathcal{L}_{M_W} + \frac{2C_A C_F \beta_0 \mathcal{L}_{M_W}}{\varepsilon} - \frac{1}{2} C_A C_F^2 Y_f^2 r(w) - \frac{1}{4} C_F Y_f^2 r(w) - \frac{C_A C_F^2 Y_f^2 r(w)}{4\varepsilon} \\
& - \frac{C_F Y_f^2 r(w)}{8\varepsilon} - \frac{1}{4} C_A Y_f^2 \beta_0 \mathcal{L}_{M_H}^2 r(w) + C_A C_F w \beta_0 \mathcal{L}_{M_W}^2 r(w) + \frac{115}{3} C_A C_F^2 w r(w) - \frac{31}{9} C_A C_F w r(w) + \frac{115}{6} C_F w r(w) \\
& - \frac{40}{9} C_A C_F n_f T_f w r(w) - \frac{8C_A C_F n_f T_f w r(w)}{3\varepsilon} + \frac{8}{3} i C_A C_F n_f \pi T_f w r(w) + \frac{47C_A C_F^2 w r(w)}{2\varepsilon} - \frac{5C_A C_F w r(w)}{3\varepsilon} + \frac{47C_F w r(w)}{4\varepsilon} \\
& - \frac{11}{2} \sqrt{3} C_A C_F^2 \pi w r(w) - \frac{11}{4} \sqrt{3} C_F \pi w r(w) - \frac{C_A Y_f^2 \beta_0 r(w)}{2\varepsilon^2} + \frac{2C_A C_F w \beta_0 r(w)}{\varepsilon^2} + \frac{1}{6} C_A C_F \pi^2 w \beta_0 r(w) - \frac{1}{24} C_A Y_f^2 \pi^2 \beta_0 r(w) \\
& + \frac{1}{4} C_A C_F^2 Y_f^2 \mathcal{L}_{M_H} r(w) + \frac{1}{8} C_F Y_f^2 \mathcal{L}_{M_H} r(w) + \frac{C_A Y_f^2 \beta_0 \mathcal{L}_{M_H} r(w)}{2\varepsilon} + \frac{1}{4} C_A C_F^2 Y_f^2 \mathcal{L}_{M_W} r(w) + \frac{1}{8} C_F Y_f^2 \mathcal{L}_{M_W} r(w) \\
& - 47C_A C_F^2 w \mathcal{L}_{M_W} r(w) + \frac{10}{3} C_A C_F w \mathcal{L}_{M_W} r(w) - \frac{47}{2} C_F w \mathcal{L}_{M_W} r(w) + \frac{16}{3} C_A C_F n_f T_f w \mathcal{L}_{M_W} r(w) - \frac{2C_A C_F w \beta_0 \mathcal{L}_{M_W} r(w)}{\varepsilon} \\
& - \frac{47C_A C_F^2}{2\varepsilon} + \frac{5C_A C_F}{3\varepsilon} - \frac{47C_F}{4\varepsilon} + \frac{11}{2} \sqrt{3} C_A C_F^2 \pi + \frac{11}{4} \sqrt{3} C_F \pi - \frac{C_A C_F^2 Y_f^2 s \log(s-r)}{4r} - \frac{C_F Y_f^2 s \log(s-r)}{8r} \\
& + \frac{C_A C_F^2 Y_f^2 s \log(r+s)}{4r} + \frac{C_F Y_f^2 s \log(r+s)}{8r} + \frac{C_A C_F^2 Y_f^2 s \log(s-r) r(w)}{4r} - \frac{C_F Y_f^2 s \log(s-r) r(w)}{8r} - \frac{C_A C_F^2 Y_f^2 s \log(r+s) r(w)}{4r} \\
& - \frac{C_F Y_f^2 s \log(r+s) r(w)}{8r} - \frac{21C_A C_F^2 Y_f^2}{8r^2} - \frac{21C_F Y_f^2}{16r^2} - \frac{3C_A C_F^2 Y_f^2}{4\varepsilon r^2} - \frac{3C_F Y_f^2}{8\varepsilon r^2} + \frac{3C_A C_F^2 Y_f^2 \mathcal{L}_{M_H}}{4r^2} + \frac{3C_F Y_f^2 \mathcal{L}_{M_H}}{8r^2}
\end{aligned}$$

$$\begin{aligned}
& + \frac{3C_A C_F^2 Y_f^2 \mathcal{L}_{M_W}}{4r^2} + \frac{3C_F Y_f^2 \mathcal{L}_{M_W}}{8r^2} + \frac{9C_A C_F^2 Y_f^2 r(w)}{8r^2} + \frac{9C_F Y_f^2 r(w)}{16r^2} + \frac{3C_A C_F^2 Y_f^2 r(w)}{4\epsilon r^2} + \frac{3C_F Y_f^2 r(w)}{8\epsilon r^2} \\
& - \frac{3C_A C_F^2 Y_f^2 \mathcal{L}_{M_H} r(w)}{4r^2} - \frac{3C_F Y_f^2 \mathcal{L}_{M_H} r(w)}{8r^2} - \frac{3C_A C_F^2 Y_f^2 \mathcal{L}_{M_W} r(w)}{4r^2} - \frac{3C_F Y_f^2 \mathcal{L}_{M_W} r(w)}{8r^2} + \frac{7C_A C_F^2 Y_f^2 s \log(s-r)}{8r^3} \\
& + \frac{7C_F Y_f^2 s \log(s-r)}{16r^3} - \frac{7C_A C_F^2 Y_f^2 s \log(r+s)}{8r^3} - \frac{7C_F Y_f^2 s \log(r+s)}{16r^3} - \frac{7C_A C_F^2 Y_f^2 s \log(s-r)r(w)}{8r^3} - \frac{7C_F Y_f^2 s \log(s-r)r(w)}{16r^3} \\
& + \frac{7C_A C_F^2 Y_f^2 s \log(r+s)r(w)}{8r^3} + \frac{7C_F Y_f^2 s \log(r+s)r(w)}{16r^3} \tag{B.10}
\end{aligned}$$

$$\begin{aligned}
\Delta U_2^{(2)} = & -\frac{1}{12} C_A C_F \log(r^2) r^6 + \frac{1}{12} C_A C_F w \log(r^2) r(w) r^6 - \frac{1}{12} C_A C_F s \log(s-r) r^5 + \frac{1}{12} C_A C_F s \log(r+s) r^5 \\
& + \frac{1}{12} C_A C_F s \log(-v) r^5 - \frac{1}{12} C_A C_F s \log\left(\frac{v}{s}\right) r^5 + \frac{1}{12} C_A C_F s w \log(s-r) r(w) r^5 - \frac{1}{12} C_A C_F s w \log(r+s) r(w) r^5 \\
& - \frac{1}{12} C_A C_F s w \log(-v) r(w) r^5 + \frac{1}{12} C_A C_F s w \log\left(\frac{v}{s}\right) r(w) r^5 + \frac{1}{6} C_A C_F r^4 + \frac{1}{6} C_A C_F \mathcal{L}_{M_H} r^4 - \frac{1}{6} C_A C_F \mathcal{L}_{M_W} r^4 \\
& + \frac{1}{3} C_A C_F \log(r^2) r^4 - \frac{1}{6} C_A C_F w r(w) r^4 - \frac{1}{6} C_A C_F w \mathcal{L}_{M_H} r(w) r^4 + \frac{1}{6} C_A C_F w \mathcal{L}_{M_W} r(w) r^4 - \frac{1}{3} C_A C_F w \log(r^2) r(w) r^4 \\
& + \frac{1}{3} C_A C_F s \log(s-r) r^3 - \frac{1}{3} C_A C_F s \log(r+s) r^3 - \frac{1}{3} C_A C_F s \log(-v) r^3 + \frac{1}{3} C_A C_F s \log\left(\frac{v}{s}\right) r^3 - \frac{1}{3} C_A C_F s w \log(s-r) r(w) r^3 \\
& + \frac{1}{3} C_A C_F s w \log(r+s) r(w) r^3 + \frac{1}{3} C_A C_F s w \log(-v) r(w) r^3 - \frac{1}{3} C_A C_F s w \log\left(\frac{v}{s}\right) r(w) r^3 - \frac{1}{16} C_A C_F^2 Y_s^2 r^2 - \frac{9}{16} C_A Y_s^2 r^2 \\
& - \frac{1}{32} C_F Y_s^2 r^2 - \frac{C_A C_F^2 Y_s^2 r^2}{64\epsilon} - \frac{9C_A Y_s^2 r^2}{64\epsilon} - \frac{C_F Y_s^2 r^2}{128\epsilon} - \frac{5}{2} C_A C_F r^2 + \frac{1}{64} C_A C_F^2 Y_s^2 \mathcal{L}_{M_H} r^2 + \frac{9}{32} C_A Y_s^2 \mathcal{L}_{M_H} r^2 \\
& + \frac{1}{128} C_F Y_s^2 \mathcal{L}_{M_H} r^2 + \frac{1}{64} C_A C_F^2 Y_s^2 \mathcal{L}_{M_W} r^2 + \frac{1}{128} C_F Y_s^2 \mathcal{L}_{M_W} r^2 + C_A C_F \mathcal{L}_{M_W} r^2 - C_A C_F \log(r^2) r^2 \\
& + \frac{1}{32} C_A C_F^2 Y_s^2 r(w) r^2 + \frac{9}{32} C_A Y_s^2 r(w) r^2 + \frac{1}{64} C_F Y_s^2 r(w) r^2 + \frac{C_A C_F^2 Y_s^2 r(w) r^2}{64\epsilon} + \frac{9C_A Y_s^2 r(w) r^2}{64\epsilon} + \frac{C_F Y_s^2 r(w) r^2}{128\epsilon} \\
& + \frac{3}{2\epsilon} C_A C_F w r(w) r^2 + \frac{C_A C_F w r(w) r^2}{2\epsilon} - \frac{1}{64} C_A C_F^2 Y_s^2 \mathcal{L}_{M_H} r(w) r^2 - \frac{9}{32} C_A Y_s^2 \mathcal{L}_{M_H} r(w) r^2 - \frac{1}{128} C_F Y_s^2 \mathcal{L}_{M_H} r(w) r^2 \\
& - \frac{1}{64} C_A C_F^2 Y_s^2 \mathcal{L}_{M_W} r(w) r^2 - \frac{1}{128} C_F Y_s^2 \mathcal{L}_{M_W} r(w) r^2 - C_A C_F w \mathcal{L}_{M_W} r(w) r^2 + C_A C_F w \log(r^2) r(w) r^2 \\
& - \frac{3}{64} \sqrt{3} C_A Y_s^2 \pi r(w) r^2 - \frac{C_A C_F r^2}{2\epsilon} + \frac{3}{64} \sqrt{3} C_A Y_s^2 \pi r^2 + \frac{1}{64} C_A C_F^2 Y_s^2 s \log(s-r) r \\
& + \frac{1}{128} C_F Y_s^2 s \log(s-r) r - C_A C_F s \log(s-r) r - \frac{1}{64} C_A C_F^2 Y_s^2 s \log(r+s) r - \frac{1}{128} C_F Y_s^2 s \log(r+s) r + C_A C_F s \log(r+s) r \\
& + C_A C_F s \log(-v) r - C_A C_F s \log\left(\frac{v}{s}\right) r - \frac{1}{64} C_A C_F^2 Y_s^2 s \log(s-r) r(w) r \\
& - \frac{1}{128} C_F Y_s^2 s \log(s-r) r(w) r + C_A C_F s w \log(s-r) r(w) r + \frac{1}{64} C_A C_F^2 Y_s^2 s \log(r+s) r(w) r \\
& + \frac{1}{128} C_F Y_s^2 s \log(r+s) r(w) r - C_A C_F s w \log(r+s) r(w) r - C_A C_F s w \log(-v) r(w) r + C_A C_F s w \log\left(\frac{v}{s}\right) r(w) r \\
& - \frac{256}{3} C_A C_F^2 + C_A C_F^2 Y_s^2 + \frac{C_F Y_s^2}{2} + \frac{C_A C_F^2 Y_s^2}{4\epsilon} + \frac{C_F Y_s^2}{8\epsilon} + \frac{1}{4} C_A Y_s^2 \beta_0 \mathcal{L}_{M_H}^2 - C_A C_F \beta_0 \mathcal{L}_{M_W}^2 + \frac{61 C_A C_F}{9} - \frac{128 C_F}{3} \\
& + \frac{88}{9} C_A C_F n_f T_f + \frac{8 C_A C_F n_f T_f}{3\epsilon} - \frac{8}{3} i C_A C_F n_f \pi T_f + \frac{C_A Y_s^2 \beta_0}{2\epsilon^2} - \frac{2 C_A C_F \beta_0}{\epsilon^2} + \frac{1}{24} C_A Y_s^2 \pi^2 \beta_0 - \frac{1}{6} C_A C_F \pi^2 \beta_0 \\
& - \frac{1}{4} C_A C_F^2 Y_s^2 \mathcal{L}_{M_H} - \frac{1}{8} C_F Y_s^2 \mathcal{L}_{M_H} - C_A Y_s^2 \beta_0 \mathcal{L}_{M_H} - \frac{C_A Y_s^2 \beta_0 \mathcal{L}_{M_H}}{2\epsilon} + 47 C_A C_F^2 \mathcal{L}_{M_W} - \frac{1}{4} C_A C_F^2 Y_s^2 \mathcal{L}_{M_W} \\
& - \frac{1}{8} C_F Y_s^2 \mathcal{L}_{M_W} + \frac{10}{3} C_A C_F \mathcal{L}_{M_W} + \frac{47}{2} C_F \mathcal{L}_{M_W} - \frac{16}{3} C_A C_F n_f T_f \mathcal{L}_{M_W} + 4 C_A C_F \beta_0 \mathcal{L}_{M_W} + \frac{2 C_A C_F \beta_0 \mathcal{L}_{M_W}}{\epsilon} \\
& - \frac{1}{2} C_A C_F^2 Y_s^2 r(w) - \frac{1}{4} C_F Y_s^2 r(w) - \frac{C_A C_F^2 Y_s^2 r(w)}{4\epsilon} - \frac{C_F Y_s^2 r(w)}{8\epsilon} - \frac{1}{4} C_A Y_s^2 \beta_0 \mathcal{L}_{M_H}^2 r(w) + C_A C_F w \beta_0 \mathcal{L}_{M_W}^2 r(w) \\
& + \frac{115}{3} C_A C_F^2 w r(w) - \frac{31}{6} C_A C_F w r(w) + \frac{115}{6} C_F w r(w) - \frac{40}{9} C_A C_F n_f T_f w r(w) - \frac{8 C_A C_F n_f T_f w r(w)}{3\epsilon} + \frac{8}{3} i C_A C_F n_f \pi T_f w r(w) \\
& + \frac{47 C_A C_F^2 w r(w)}{2\epsilon} - \frac{5 C_A C_F w r(w)}{3\epsilon} + \frac{47 C_F w r(w)}{4\epsilon} - \frac{11}{2} \sqrt{3} C_A C_F^2 \pi w r(w) - \frac{11}{4} \sqrt{3} C_F \pi w r(w) - \frac{C_A Y_s^2 \beta_0 r(w)}{2\epsilon^2} \\
& + \frac{2 C_A C_F w \beta_0 r(w)}{\epsilon^2} + \frac{1}{6} C_A C_F \pi^2 w \beta_0 r(w) - \frac{1}{24} C_A Y_s^2 \pi^2 \beta_0 r(w) + \frac{1}{4} C_A C_F^2 Y_s^2 \mathcal{L}_{M_H} r(w) + \frac{1}{8} C_F Y_s^2 \mathcal{L}_{M_H} r(w) \\
& + \frac{C_A Y_s^2 \beta_0 \mathcal{L}_{M_H} r(w)}{2\epsilon} + \frac{1}{4} C_A C_F^2 Y_s^2 \mathcal{L}_{M_W} r(w) + \frac{1}{8} C_F Y_s^2 \mathcal{L}_{M_W} r(w) - 47 C_A C_F^2 w \mathcal{L}_{M_W} r(w) + \frac{10}{3} C_A C_F w \mathcal{L}_{M_W} r(w) \\
& - \frac{47}{2} C_F w \mathcal{L}_{M_W} r(w) + \frac{16}{3} C_A C_F n_f T_f w \mathcal{L}_{M_W} r(w) - \frac{2 C_A C_F w \beta_0 \mathcal{L}_{M_W} r(w)}{\epsilon} - \frac{47 C_A C_F^2}{2\epsilon} + \frac{5 C_A C_F}{3\epsilon} - \frac{47 C_F}{4\epsilon} + \frac{11}{2} \sqrt{3} C_A C_F^2 \pi \\
& + \frac{11}{4} \sqrt{3} C_F \pi - \frac{C_A C_F^2 Y_s^2 s \log(s-r)}{4r} - \frac{C_F Y_s^2 s \log(s-r)}{8r} + \frac{C_A C_F^2 Y_s^2 s \log(r+s)}{4r} + \frac{C_F Y_s^2 s \log(r+s)}{8r} \\
& + \frac{C_A C_F^2 Y_s^2 s \log(s-r) r(w)}{4r} + \frac{C_F Y_s^2 s \log(s-r) r(w)}{8r} - \frac{C_A C_F^2 Y_s^2 s \log(r+s) r(w)}{4r} - \frac{C_F Y_s^2 s \log(r+s) r(w)}{8r} - \frac{21 C_A C_F^2 Y_s^2}{8r^2} \\
& - \frac{21 C_F Y_s^2}{16r^2} - \frac{3 C_A C_F^2 Y_s^2}{4\epsilon r^2} - \frac{3 C_F Y_s^2}{8\epsilon r^2} + \frac{3 C_A C_F^2 Y_s^2 \mathcal{L}_{M_H}}{4r^2} + \frac{3 C_F Y_s^2 \mathcal{L}_{M_H}}{8r^2} + \frac{3 C_A C_F^2 Y_s^2 \mathcal{L}_{M_W}}{4r^2} + \frac{3 C_F Y_s^2 \mathcal{L}_{M_W}}{8r^2} \\
& + \frac{9 C_A C_F^2 Y_s^2 r(w)}{8r^2} + \frac{9 C_F Y_s^2 r(w)}{16r^2} + \frac{3 C_A C_F^2 Y_s^2 r(w)}{4\epsilon r^2} + \frac{3 C_F Y_s^2 r(w)}{8\epsilon r^2} - \frac{3 C_A C_F^2 Y_s^2 \mathcal{L}_{M_H} r(w)}{4r^2} - \frac{3 C_F Y_s^2 \mathcal{L}_{M_H} r(w)}{8r^2}
\end{aligned}$$

$$\begin{aligned}
& - \frac{3C_A C_F^2 Y_s^2 \mathcal{L}_{M_W} r(w)}{4r^2} - \frac{3C_F Y_s^2 \mathcal{L}_{M_W} r(w)}{8r^2} + \frac{7C_A C_F^2 Y_s^2 s \log(s-r)}{8r^3} + \frac{7C_F Y_s^2 s \log(s-r)}{16r^3} - \frac{7C_A C_F^2 Y_s^2 s \log(r+s)}{8r^3} \\
& - \frac{7C_F Y_s^2 s \log(r+s)}{16r^3} - \frac{7C_A C_F^2 Y_s^2 s \log(s-r)r(w)}{8r^3} - \frac{7C_F Y_s^2 s \log(s-r)r(w)}{16r^3} + \frac{7C_A C_F^2 Y_s^2 s \log(r+s)r(w)}{8r^3} \\
& + \frac{7C_F Y_s^2 s \log(r+s)r(w)}{16r^3}
\end{aligned} \tag{B.11}$$

$$\begin{aligned}
\Delta U_3^{(2)} = & - \frac{1}{12} C_A C_F \log(r^2) r^6 + \frac{1}{12} C_A C_F w \log(r^2) r(w) r^6 - \frac{1}{12} C_A C_F s \log(s-r) r^5 + \frac{1}{12} C_A C_F s \log(r+s) r^5 \\
& + \frac{1}{12} C_A C_F s \log(-v) r^5 - \frac{1}{12} C_A C_F s \log\left(\frac{v}{s}\right) r^5 + \frac{1}{12} C_A C_F s w \log(s-r) r(w) r^5 - \frac{1}{12} C_A C_F s w \log(r+s) r(w) r^5 \\
& - \frac{1}{12} C_A C_F s w \log(-v) r(w) r^5 + \frac{1}{12} C_A C_F s w \log\left(\frac{v}{s}\right) r(w) r^5 + \frac{1}{6} C_A C_F r^4 + \frac{1}{6} C_A C_F \mathcal{L}_{M_H} r^4 - \frac{1}{6} C_A C_F \mathcal{L}_{M_W} r^4 \\
& + \frac{1}{3} C_A C_F \log(r^2) r^4 - \frac{1}{6} C_A C_F w r(w) r^4 - \frac{1}{6} C_A C_F w \mathcal{L}_{M_H} r(w) r^4 + \frac{1}{6} C_A C_F w \mathcal{L}_{M_W} r(w) r^4 - \frac{1}{3} C_A C_F w \log(r^2) r(w) r^4 \\
& + \frac{1}{3} C_A C_F s \log(s-r) r^3 - \frac{1}{3} C_A C_F s \log(r+s) r^3 - \frac{1}{3} C_A C_F s \log(-v) r^3 + \frac{1}{3} C_A C_F s \log\left(\frac{v}{s}\right) r^3 - \frac{1}{3} C_A C_F s w \log(s-r) r(w) r^3 \\
& + \frac{1}{3} C_A C_F s w \log(r+s) r(w) r^3 + \frac{1}{3} C_A C_F s w \log(-v) r(w) r^3 - \frac{1}{3} C_A C_F s w \log\left(\frac{v}{s}\right) r(w) r^3 - \frac{1}{32} C_A C_F^2 Y_f^2 r^2 - \frac{9}{32} C_A Y_f^2 r^2 \\
& - \frac{1}{64} C_F Y_f^2 r^2 - \frac{C_A C_F^2 Y_f^2 r^2}{128\epsilon} - \frac{9C_A Y_f^2 r^2}{128\epsilon} - \frac{C_F Y_f^2 r^2}{256\epsilon} - \frac{1}{32} C_A C_F^2 Y_s^2 r^2 - \frac{9}{32} C_A Y_s^2 r^2 - \frac{1}{64} C_F Y_s^2 r^2 - \frac{C_A C_F^2 Y_s^2 r^2}{128\epsilon} \\
& - \frac{9C_A Y_s^2 r^2}{128\epsilon} - \frac{C_F Y_s^2 r^2}{256\epsilon} - \frac{5}{2} C_A C_F r^2 + \frac{1}{128} C_A C_F^2 Y_f^2 \mathcal{L}_{M_H} r^2 + \frac{9}{64} C_A Y_f^2 \mathcal{L}_{M_H} r^2 + \frac{1}{256} C_F Y_f^2 \mathcal{L}_{M_H} r^2 \\
& + \frac{1}{128} C_A C_F^2 Y_s^2 \mathcal{L}_{M_H} r^2 + \frac{9}{64} C_A Y_s^2 \mathcal{L}_{M_H} r^2 + \frac{1}{256} C_F Y_s^2 \mathcal{L}_{M_H} r^2 + \frac{1}{128} C_A C_F^2 Y_f^2 \mathcal{L}_{M_W} r^2 + \frac{1}{256} C_F Y_f^2 \mathcal{L}_{M_W} r^2 \\
& + \frac{1}{128} C_A C_F^2 Y_s^2 \mathcal{L}_{M_W} r^2 + \frac{1}{256} C_F Y_s^2 \mathcal{L}_{M_W} r^2 + C_A C_F \mathcal{L}_{M_W} r^2 - C_A C_F \log(r^2) r^2 + \frac{1}{32} C_A C_F^2 Y_f Y_s r(w) r^2 \\
& + \frac{9}{32} C_A Y_f Y_s r(w) r^2 + \frac{1}{64} C_F Y_f Y_s r(w) r^2 + \frac{C_A C_F^2 Y_f Y_s r(w) r^2}{64\epsilon} + \frac{9C_A Y_f Y_s r(w) r^2}{64\epsilon} + \frac{C_F Y_f Y_s r(w) r^2}{128\epsilon} + \frac{3}{2} C_A C_F w r(w) r^2 \\
& + \frac{C_A C_F w r(w) r^2}{2\epsilon} - \frac{1}{64} C_A C_F^2 Y_f Y_s \mathcal{L}_{M_H} r(w) r^2 - \frac{9}{32} C_A Y_f Y_s \mathcal{L}_{M_H} r(w) r^2 - \frac{1}{128} C_F Y_f Y_s \mathcal{L}_{M_H} r(w) r^2 \\
& - \frac{1}{64} C_A C_F^2 Y_f Y_s \mathcal{L}_{M_W} r(w) r^2 - \frac{1}{128} C_F Y_f Y_s \mathcal{L}_{M_W} r(w) r^2 - C_A C_F w \mathcal{L}_{M_W} r(w) r^2 + C_A C_F w \log(r^2) r(w) r^2 \\
& - \frac{3}{64} \sqrt{3} C_A Y_f Y_s \pi r(w) r^2 - \frac{C_A C_F r^2}{2\epsilon} + \frac{3}{128} \sqrt{3} C_A Y_f^2 \pi r^2 + \frac{3}{128} \sqrt{3} C_A Y_s^2 \pi r^2 + \frac{1}{128} C_A C_F^2 Y_f^2 s \log(s-r) r \\
& + \frac{1}{256} C_F Y_f^2 s \log(s-r) r + \frac{1}{128} C_A C_F^2 Y_s^2 s \log(s-r) r + \frac{1}{256} C_F Y_s^2 s \log(s-r) r - C_A C_F s \log(s-r) r \\
& - \frac{1}{128} C_A C_F^2 Y_f^2 s \log(r+s) r - \frac{1}{256} C_F Y_f^2 s \log(r+s) r - \frac{1}{128} C_A C_F^2 Y_s^2 s \log(r+s) r - \frac{1}{256} C_F Y_s^2 s \log(r+s) r \\
& + C_A C_F s \log(r+s) r + C_A C_F s \log(-v) r - C_A C_F s \log\left(\frac{v}{s}\right) r - \frac{1}{64} C_A C_F^2 Y_f Y_s s \log(s-r) r(w) r \\
& - \frac{1}{128} C_F Y_f Y_s s \log(s-r) r(w) r + C_A C_F s w \log(s-r) r(w) r + \frac{1}{64} C_A C_F^2 Y_f Y_s s \log(r+s) r(w) r \\
& + \frac{1}{128} C_F Y_f Y_s s \log(r+s) r(w) r - C_A C_F s w \log(r+s) r(w) r - C_A C_F s w \log(-v) r(w) r + C_A C_F s w \log\left(\frac{v}{s}\right) r(w) r - \frac{256}{3} C_A C_F^2 \\
& + \frac{1}{2} C_A C_F^2 Y_f^2 + \frac{C_F Y_f^2}{4} + \frac{C_A C_F^2 Y_f^2}{8\epsilon} + \frac{C_F Y_f^2}{16\epsilon} + \frac{1}{2} C_A C_F^2 Y_s^2 + \frac{C_F Y_s^2}{4} + \frac{C_A C_F^2 Y_s^2}{8\epsilon} + \frac{C_F Y_s^2}{16\epsilon} + \frac{1}{8} C_A Y_f^2 \beta_0 \mathcal{L}_{M_H}^2 \\
& + \frac{1}{8} C_A Y_s^2 \beta_0 \mathcal{L}_{M_H}^2 - C_A C_F \beta_0 \mathcal{L}_{M_W}^2 + \frac{61C_A C_F}{9} - \frac{128C_F}{3} + \frac{88}{9} C_A C_F n_f T_f + \frac{8C_A C_F n_f T_f}{3\epsilon} - \frac{8}{3} i C_A C_F n_f \pi T_f + \frac{C_A Y_f^2 \beta_0}{4\epsilon^2} \\
& + \frac{C_A Y_s^2 \beta_0}{4\epsilon^2} - \frac{2C_A C_F \beta_0}{\epsilon^2} + \frac{1}{48} C_A Y_f^2 \pi^2 \beta_0 + \frac{1}{48} C_A Y_s^2 \pi^2 \beta_0 - \frac{1}{6} C_A C_F \pi^2 \beta_0 - \frac{1}{8} C_A C_F^2 Y_f^2 \mathcal{L}_{M_H} - \frac{1}{16} C_F Y_f^2 \mathcal{L}_{M_H} \\
& - \frac{1}{8} C_A C_F^2 Y_s^2 \mathcal{L}_{M_H} - \frac{1}{16} C_F Y_s^2 \mathcal{L}_{M_H} - \frac{1}{2} C_A Y_f^2 \beta_0 \mathcal{L}_{M_H} - \frac{C_A Y_f^2 \beta_0 \mathcal{L}_{M_H}}{4\epsilon} - \frac{1}{2} C_A Y_s^2 \beta_0 \mathcal{L}_{M_H} - \frac{C_A Y_s^2 \beta_0 \mathcal{L}_{M_H}}{4\epsilon} \\
& + 47C_A C_F^2 \mathcal{L}_{M_W} - \frac{1}{8} C_A C_F^2 Y_f^2 \mathcal{L}_{M_W} - \frac{1}{16} C_F Y_f^2 \mathcal{L}_{M_W} - \frac{1}{8} C_A C_F^2 Y_s^2 \mathcal{L}_{M_W} - \frac{1}{16} C_F Y_s^2 \mathcal{L}_{M_W} - \frac{10}{3} C_A C_F \mathcal{L}_{M_W} \\
& + \frac{47}{2} C_F \mathcal{L}_{M_W} - \frac{16}{3} C_A C_F n_f T_f \mathcal{L}_{M_W} + 4C_A C_F \beta_0 \mathcal{L}_{M_W} + \frac{2C_A C_F \beta_0 \mathcal{L}_{M_W}}{\epsilon} - \frac{1}{4} C_A Y_f Y_s \beta_0 \mathcal{L}_{M_H}^2 r(w) \\
& + C_A C_F w \beta_0 \mathcal{L}_{M_W}^2 r(w) - \frac{1}{2} C_A C_F^2 Y_f Y_s r(w) - \frac{1}{4} C_F Y_f Y_s r(w) - \frac{C_A C_F^2 Y_f Y_s r(w)}{4\epsilon} - \frac{C_F Y_f Y_s r(w)}{8\epsilon} + \frac{115}{3} C_A C_F^2 w r(w) \\
& - \frac{31}{9} C_A C_F w r(w) + \frac{115}{6} C_F w r(w) - \frac{40}{9} C_A C_F n_f T_f w r(w) - \frac{8C_A C_F n_f T_f w r(w)}{3\epsilon} + \frac{8}{3} i C_A C_F n_f \pi T_f w r(w) + \frac{47C_A C_F^2 w r(w)}{2\epsilon} \\
& - \frac{5C_A C_F w r(w)}{3\epsilon} + \frac{47C_F w r(w)}{4\epsilon} - \frac{11}{2} \sqrt{3} C_A C_F^2 \pi w r(w) - \frac{11}{4} \sqrt{3} C_F \pi w r(w) - \frac{C_A Y_f Y_s \beta_0 r(w)}{2\epsilon^2} + \frac{2C_A C_F w \beta_0 r(w)}{\epsilon^2} \\
& + \frac{1}{6} C_A C_F \pi^2 w \beta_0 r(w) - \frac{1}{24} C_A Y_f Y_s \pi^2 \beta_0 r(w) + \frac{1}{4} C_A C_F^2 Y_f Y_s \mathcal{L}_{M_H} r(w) + \frac{1}{8} C_F Y_f Y_s \mathcal{L}_{M_H} r(w) + \frac{C_A Y_f Y_s \beta_0 \mathcal{L}_{M_H} r(w)}{2\epsilon} \\
& + \frac{1}{4} C_A C_F^2 Y_f Y_s \mathcal{L}_{M_W} r(w) + \frac{1}{8} C_F Y_f Y_s \mathcal{L}_{M_W} r(w) - 47C_A C_F^2 w \mathcal{L}_{M_W} r(w) + \frac{10}{3} C_A C_F w \mathcal{L}_{M_W} r(w) - \frac{47}{2} C_F w \mathcal{L}_{M_W} r(w) \\
& + \frac{16}{3} C_A C_F n_f T_f w \mathcal{L}_{M_W} r(w) - \frac{2C_A C_F w \beta_0 \mathcal{L}_{M_W} r(w)}{\epsilon} - \frac{47C_A C_F^2}{2\epsilon} + \frac{5C_A C_F}{3\epsilon} - \frac{47C_F}{4\epsilon} + \frac{11}{2} \sqrt{3} C_A C_F^2 \pi + \frac{11}{4} \sqrt{3} C_F \pi
\end{aligned}$$

$$\begin{aligned}
& - \frac{C_A C_F^2 Y_f^2 s \log(s-r)}{8r} - \frac{C_F Y_f^2 s \log(s-r)}{16r} - \frac{C_A C_F^2 Y_s^2 s \log(s-r)}{8r} - \frac{C_F Y_s^2 s \log(s-r)}{16r} + \frac{C_A C_F^2 Y_f^2 s \log(r+s)}{8r} \\
& + \frac{C_F Y_f^2 s \log(r+s)}{16r} + \frac{C_A C_F^2 Y_s^2 s \log(r+s)}{8r} + \frac{C_F Y_s^2 s \log(r+s)}{16r} + \frac{C_A C_F^2 Y_f Y_s s \log(s-r)r(w)}{4r} + \frac{C_F Y_f Y_s s \log(s-r)r(w)}{8r} \\
& - \frac{C_A C_F^2 Y_f Y_s s \log(r+s)r(w)}{4r} - \frac{C_F Y_f Y_s s \log(r+s)r(w)}{8r} - \frac{21 C_A C_F^2 Y_f^2}{16r^2} - \frac{21 C_F Y_f^2}{32r^2} - \frac{3 C_A C_F^2 Y_f^2}{8\epsilon r^2} - \frac{3 C_F Y_f^2}{16\epsilon r^2} - \frac{21 C_A C_F^2 Y_s^2}{16r^2} \\
& - \frac{21 C_F Y_s^2}{32r^2} - \frac{3 C_A C_F^2 Y_s^2}{8\epsilon r^2} - \frac{3 C_F Y_s^2}{16\epsilon r^2} + \frac{3 C_A C_F^2 Y_f^2 \mathcal{L}_{MH}}{8r^2} + \frac{3 C_F Y_f^2 \mathcal{L}_{MH}}{16r^2} + \frac{3 C_A C_F^2 Y_s^2 \mathcal{L}_{MH}}{8r^2} + \frac{3 C_F Y_s^2 \mathcal{L}_{MH}}{16r^2} \\
& + \frac{3 C_A C_F^2 Y_f^2 \mathcal{L}_{MW}}{8r^2} + \frac{3 C_F Y_f^2 \mathcal{L}_{MW}}{16r^2} + \frac{3 C_A C_F^2 Y_s^2 \mathcal{L}_{MW}}{8r^2} + \frac{3 C_F Y_s^2 \mathcal{L}_{MW}}{16r^2} + \frac{9 C_A C_F^2 Y_f Y_s r(w)}{8r^2} + \frac{9 C_F Y_f Y_s r(w)}{16r^2} \\
& + \frac{3 C_A C_F^2 Y_f Y_s r(w)}{4\epsilon r^2} + \frac{3 C_F Y_f Y_s r(w)}{8\epsilon r^2} - \frac{3 C_A C_F^2 Y_f Y_s \mathcal{L}_{MH} r(w)}{4r^2} - \frac{3 C_F Y_f Y_s \mathcal{L}_{MH} r(w)}{8r^2} - \frac{3 C_A C_F^2 Y_f Y_s \mathcal{L}_{MW} r(w)}{4r^2} \\
& - \frac{3 C_F Y_f Y_s \mathcal{L}_{MW} r(w)}{8r^2} + \frac{7 C_A C_F^2 Y_f^2 s \log(s-r)}{16r^3} + \frac{7 C_F Y_f^2 s \log(s-r)}{32r^3} + \frac{7 C_A C_F^2 Y_s^2 s \log(s-r)}{16r^3} + \frac{7 C_F Y_s^2 s \log(s-r)}{32r^3} \\
& - \frac{7 C_A C_F^2 Y_f^2 s \log(r+s)}{16r^3} - \frac{7 C_F Y_f^2 s \log(r+s)}{32r^3} - \frac{7 C_A C_F^2 Y_s^2 s \log(r+s)}{16r^3} - \frac{7 C_F Y_s^2 s \log(r+s)}{32r^3} - \frac{7 C_A C_F^2 Y_f Y_s s \log(s-r)r(w)}{8r^3} \\
& - \frac{7 C_F Y_f Y_s s \log(s-r)r(w)}{16r^3} + \frac{7 C_A C_F^2 Y_f Y_s s \log(r+s)r(w)}{8r^3} + \frac{7 C_F Y_f Y_s s \log(r+s)r(w)}{16r^3}
\end{aligned} \tag{B.12}$$

## C Vertex Corrections

### C.1 Matching at $\mu \sim Q$ :

$$\begin{aligned}
v_1^{(Q)} = & \frac{1}{96} C_A \mathcal{L}_Q^4 Y_f^4 - \frac{5}{12} C_A \mathcal{L}_Q^3 Y_f^4 - \frac{C_A \mathcal{L}_Q^3 Y_f^4}{48\epsilon} + \frac{105}{128} C_A \mathcal{L}_Q^2 Y_f^4 + \frac{5 C_A \mathcal{L}_Q^2 Y_f^4}{8\epsilon} + \frac{C_A \mathcal{L}_Q^2 Y_f^4}{32\epsilon^2} - \frac{2463 C_A Y_f^4}{512} - \frac{329}{256} C_A \mathcal{L}_Q Y_f^4 \\
& - \frac{105 C_A \mathcal{L}_Q Y_f^4}{128\epsilon} - \frac{5 C_A \mathcal{L}_Q Y_f^4}{8\epsilon^2} - \frac{C_A \mathcal{L}_Q Y_f^4}{32\epsilon^3} + \frac{3}{128} C_A \pi^2 \mathcal{L}_Q Y_f^4 + \frac{329 C_A Y_f^4}{512\epsilon} + \frac{105 C_A Y_f^4}{256\epsilon^2} + \frac{5 C_A Y_f^4}{16\epsilon^3} + \frac{C_A Y_f^4}{64\epsilon^4} - \frac{83}{192} C_A \zeta_3 Y_f^4 \\
& - \frac{5}{256} C_A \pi^2 Y_f^4 - \frac{3 C_A \pi^2 Y_f^4}{256\epsilon} + \frac{1}{4} C_A C_F \mathcal{L}_Q^4 Y_f^2 - \frac{17}{2} C_A C_F \mathcal{L}_Q^3 Y_f^2 - \frac{C_A C_F \mathcal{L}_Q^3 Y_f^2}{2\epsilon} + \frac{19}{4} C_A C_F^2 Y_f^2 + \frac{1}{2} C_A C_F^2 \mathcal{L}_Q^2 Y_f^2 \\
& + \frac{333}{16} C_A C_F \mathcal{L}_Q^2 Y_f^2 + \frac{1}{4} C_F \mathcal{L}_Q^2 Y_f^2 + \frac{51 C_A C_F \mathcal{L}_Q^2 Y_f^2}{4\epsilon} + \frac{3 C_A C_F \mathcal{L}_Q^2 Y_f^2}{4\epsilon^2} + \frac{1}{12} C_A C_F \pi^2 \mathcal{L}_Q^2 Y_f^2 - \frac{167}{2} C_A C_F Y_f^2 + \frac{19 C_F Y_f^2}{8} \\
& - \frac{5}{2} C_A C_F^2 \mathcal{L}_Q Y_f^2 + \frac{2501}{16} C_A C_F \mathcal{L}_Q Y_f^2 - \frac{5}{4} C_F \mathcal{L}_Q Y_f^2 - \frac{C_A C_F^2 \mathcal{L}_Q Y_f^2}{2\epsilon} - \frac{333 C_A C_F \mathcal{L}_Q Y_f^2}{16\epsilon} - \frac{C_F \mathcal{L}_Q Y_f^2}{4\epsilon} - \frac{51 C_A C_F \mathcal{L}_Q Y_f^2}{4\epsilon^2} \\
& - \frac{3 C_A C_F \mathcal{L}_Q Y_f^2}{4\epsilon^3} - \frac{9}{8} C_A C_F \pi^2 \mathcal{L}_Q Y_f^2 - \frac{C_A C_F \pi^2 \mathcal{L}_Q Y_f^2}{12\epsilon} + \frac{5 C_A C_F^2 Y_f^2}{4\epsilon} - \frac{2501 C_A C_F Y_f^2}{32\epsilon} + \frac{5 C_F Y_f^2}{8\epsilon} + \frac{C_A C_F^2 Y_f^2}{4\epsilon^2} \\
& + \frac{333 C_A C_F Y_f^2}{32\epsilon^2} + \frac{C_F Y_f^2}{8\epsilon^2} + \frac{51 C_A C_F Y_f^2}{8\epsilon^3} + \frac{3 C_A C_F Y_f^2}{8\epsilon^4} + \frac{83}{12} C_A C_F \zeta_3 Y_f^2 + \frac{1}{24} C_A C_F^2 \pi^2 Y_f^2 + \frac{47}{48} C_A C_F \pi^2 Y_f^2 + \frac{1}{48} C_F \pi^2 Y_f^2 \\
& + \frac{9 C_A C_F \pi^2 Y_f^2}{16\epsilon} + \frac{C_A C_F \pi^2 Y_f^2}{24\epsilon^2} - 6 C_A C_F^2 \mathcal{L}_Q^4 + 2 C_F \mathcal{L}_Q^4 - \frac{157}{6} C_A C_F^2 \mathcal{L}_Q^3 + \frac{1}{18} C_A C_F \mathcal{L}_Q^3 - \frac{133}{12} C_F \mathcal{L}_Q^3 - \frac{2}{9} C_A C_F n_f T_f \mathcal{L}_Q^3 \\
& + \frac{12 C_A C_F^2 \mathcal{L}_Q^3}{\epsilon} - \frac{4 C_F \mathcal{L}_Q^3}{\epsilon} - \frac{78349}{216} C_A C_F^2 + \frac{650}{3} C_A C_F^2 \mathcal{L}_Q^2 - \frac{2}{9} C_A C_F \mathcal{L}_Q^2 + \frac{263}{6} C_F \mathcal{L}_Q^2 - \frac{1}{9} C_A C_F n_f T_f \mathcal{L}_Q^2 + \frac{C_A C_F n_f T_f \mathcal{L}_Q^2}{3\epsilon} \\
& + \frac{157 C_A C_F^2 \mathcal{L}_Q^2}{4\epsilon} - \frac{C_A C_F \mathcal{L}_Q^2}{12\epsilon} + \frac{133 C_F \mathcal{L}_Q^2}{8\epsilon} - \frac{18 C_A C_F^2 \mathcal{L}_Q^2}{\epsilon^2} + \frac{6 C_F \mathcal{L}_Q^2}{\epsilon^2} + C_A C_F^2 \pi^2 \mathcal{L}_Q^2 - \frac{5}{3} C_F \pi^2 \mathcal{L}_Q^2 + \frac{1129 C_A C_F}{648} \\
& + \frac{362507 C_F}{432} - \frac{1615}{162} C_A C_F n_f T_f + \frac{8 C_A C_F n_f T_f}{27\epsilon} - \frac{C_A C_F n_f T_f}{18\epsilon^2} + \frac{C_A C_F n_f T_f}{6\epsilon^3} - \frac{4}{27} C_A C_F n_f \pi^2 T_f + \frac{C_A C_F n_f \pi^2 T_f}{36\epsilon} \\
& - \frac{5759}{18} C_A C_F^2 \mathcal{L}_Q + \frac{35}{54} C_A C_F \mathcal{L}_Q - \frac{31147}{72} C_F \mathcal{L}_Q - \frac{16}{27} C_A C_F n_f T_f \mathcal{L}_Q + \frac{C_A C_F n_f T_f \mathcal{L}_Q}{9\epsilon} - \frac{C_A C_F n_f T_f \mathcal{L}_Q}{3\epsilon^2} \\
& - \frac{1}{18} C_A C_F n_f \pi^2 T_f \mathcal{L}_Q - \frac{650 C_A C_F^2 \mathcal{L}_Q}{3\epsilon} + \frac{2 C_A C_F \mathcal{L}_Q}{9\epsilon} - \frac{263 C_F \mathcal{L}_Q}{6\epsilon} - \frac{157 C_A C_F^2 \mathcal{L}_Q}{4\epsilon^2} + \frac{C_A C_F \mathcal{L}_Q}{12\epsilon^2} - \frac{133 C_F \mathcal{L}_Q}{8\epsilon^2} \\
& + \frac{18 C_A C_F^2 \mathcal{L}_Q}{\epsilon^3} - \frac{6 C_F \mathcal{L}_Q}{\epsilon^3} + \frac{17}{8} C_A C_F^2 \pi^2 \mathcal{L}_Q + \frac{1}{72} C_A C_F \pi^2 \mathcal{L}_Q + \frac{93}{16} C_F \pi^2 \mathcal{L}_Q - \frac{C_A C_F^2 \pi^2 \mathcal{L}_Q}{\epsilon} + \frac{5 C_F \pi^2 \mathcal{L}_Q}{3\epsilon} + \frac{5759 C_A C_F^2}{36\epsilon} \\
& - \frac{35 C_A C_F}{108\epsilon} + \frac{31147 C_F}{144\epsilon} + \frac{325 C_A C_F^2}{3\epsilon^2} - \frac{C_A C_F}{9\epsilon^2} + \frac{263 C_F}{12\epsilon^2} + \frac{157 C_A C_F^2}{8\epsilon^3} - \frac{C_A C_F}{24\epsilon^3} + \frac{133 C_F}{16\epsilon^3} - \frac{9 C_A C_F^2}{\epsilon^4} + \frac{3 C_F}{\epsilon^4} - \frac{415 C_F \zeta_3}{6} \\
& + \frac{166}{3} C_F \mathcal{L}_Q \zeta_3 - \frac{83 C_F \zeta_3}{3\epsilon} - \frac{59 C_F \pi^4}{120} + \frac{389}{144} C_A C_F^2 \pi^2 + \frac{7}{432} C_A C_F \pi^2 - \frac{3835 C_F \pi^2}{288} - \frac{17 C_A C_F^2 \pi^2}{16\epsilon} - \frac{C_A C_F \pi^2}{144\epsilon} - \frac{93 C_F \pi^2}{32\epsilon} \\
& + \frac{C_A C_F^2 \pi^2}{2\epsilon^2} - \frac{5 C_F \pi^2}{6\epsilon^2}
\end{aligned} \tag{C.1}$$

$$\begin{aligned}
V_2^{(Q)} = & \frac{1}{96} C_A C_Q^4 Y_f^4 - \frac{31}{48} C_A C_Q^3 Y_f^4 - \frac{C_A C_Q^3 Y_f^4}{48\epsilon} + \frac{67}{64} C_A C_Q^2 Y_f^4 + \frac{31 C_A C_Q^2 Y_f^4}{32\epsilon} + \frac{C_A C_Q^2 Y_f^4}{32\epsilon^2} + \frac{1}{192} C_A \pi^2 C_Q^2 Y_f^4 + \frac{823 C_A Y_f^4}{512} \\
& + \frac{2595}{256} C_A C_Q Y_f^4 - \frac{67 C_A C_Q Y_f^4}{64\epsilon} - \frac{31 C_A C_Q Y_f^4}{32\epsilon^2} - \frac{C_A C_Q Y_f^4}{32\epsilon^3} - \frac{17}{384} C_A \pi^2 C_Q Y_f^4 - \frac{C_A \pi^2 C_Q Y_f^4}{192\epsilon} - \frac{2595 C_A Y_f^4}{512\epsilon} + \frac{67 C_A Y_f^4}{128\epsilon^2} \\
& + \frac{31 C_A Y_f^4}{64\epsilon^3} + \frac{C_A Y_f^4}{64\epsilon^4} + \frac{83}{192} C_A \zeta_3 Y_f^4 + \frac{55}{768} C_A \pi^2 Y_f^4 + \frac{17 C_A \pi^2 Y_f^4}{768\epsilon} + \frac{C_A \pi^2 Y_f^4}{384\epsilon^2} - \frac{2}{3} C_A C_F C_Q^4 Y_f^2 + \frac{91}{9} C_A C_F C_Q^3 Y_f^2 \\
& + \frac{4 C_A C_F C_Q^3 Y_f^2}{3\epsilon} - \frac{37}{32} C_A C_F^2 Y_f^2 - \frac{1}{4} C_A C_F^2 C_Q^2 Y_f^2 + \frac{1493}{144} C_A C_F C_Q^2 Y_f^2 - \frac{1}{8} C_F C_Q^2 Y_f^2 - \frac{91 C_A C_F C_Q^2 Y_f^2}{6\epsilon} - \frac{2 C_A C_F C_Q^2 Y_f^2}{\epsilon^2} \\
& - \frac{1}{8} C_A C_F \pi^2 C_Q^2 Y_f^2 + \frac{1244593 C_A C_F Y_f^2}{5184} - \frac{37 C_F Y_f^2}{64} + \frac{7}{8} C_A C_F^2 C_Q Y_f^2 - \frac{91409}{864} C_A C_F C_Q Y_f^2 + \frac{7}{16} C_F C_Q Y_f^2 + \frac{C_A C_F^2 C_Q Y_f^2}{4\epsilon} \\
& - \frac{1493 C_A C_F C_Q Y_f^2}{144\epsilon} + \frac{C_F C_Q Y_f^2}{8\epsilon} + \frac{91 C_A C_F C_Q Y_f^2}{6\epsilon^2} + \frac{2 C_A C_F C_Q Y_f^2}{\epsilon^3} + \frac{163}{144} C_A C_F \pi^2 C_Q Y_f^2 + \frac{C_A C_F \pi^2 C_Q Y_f^2}{8\epsilon} - \frac{7 C_A C_F^2 Y_f^2}{16\epsilon} \\
& + \frac{91409 C_A C_F Y_f^2}{1728\epsilon} - \frac{7 C_F Y_f^2}{32\epsilon} - \frac{C_A C_F^2 Y_f^2}{8\epsilon^2} + \frac{1493 C_A C_F Y_f^2}{288\epsilon^2} - \frac{C_F Y_f^2}{16\epsilon^2} - \frac{91 C_A C_F Y_f^2}{12\epsilon^3} - \frac{C_A C_F Y_f^2}{\epsilon^4} - \frac{415}{24} C_A C_F \zeta_3 Y_f^2 \\
& - \frac{1}{48} C_A C_F^2 \pi^2 Y_f^2 - \frac{3041}{864} C_A C_F \pi^2 Y_f^2 - \frac{1}{96} C_F \pi^2 Y_f^2 - \frac{163 C_A C_F \pi^2 Y_f^2}{288\epsilon} - \frac{C_A C_F \pi^2 Y_f^2}{16\epsilon^2} - \frac{20}{3} C_A C_F^2 C_Q^4 - 2 C_F C_Q^4 \\
& - \frac{1783}{18} C_A C_F^2 C_Q^3 + \frac{1}{18} C_A C_F C_Q^3 - \frac{269}{12} C_F C_Q^3 - \frac{2}{9} C_A C_F n_f T_f C_Q^3 + \frac{40 C_A C_F^2 C_Q^3}{3\epsilon} + \frac{4 C_F C_Q^3}{\epsilon} - \frac{239309 C_A C_F^2}{5184} \\
& + \frac{43897}{72} C_A C_F^2 C_Q^2 - \frac{25}{72} C_A C_F C_Q^2 + \frac{20131}{48} C_F C_Q^2 + \frac{17}{9} C_A C_F n_f T_f C_Q^2 + \frac{C_A C_F n_f T_f C_Q^2}{3\epsilon} + \frac{1783 C_A C_F^2 C_Q^2}{12\epsilon} - \frac{C_A C_F C_Q^2}{12\epsilon} \\
& + \frac{269 C_F C_Q^2}{8\epsilon} - \frac{20 C_A C_F^2 C_Q^2}{\epsilon^2} - \frac{6 C_F C_Q^2}{\epsilon^2} + C_A C_F^2 \pi^2 C_Q^2 - \frac{5}{2} C_F \pi^2 C_Q^2 - \frac{283 C_A C_F}{5184} + \frac{8829121 C_F}{3456} + \frac{2683}{324} C_A C_F n_f T_f \\
& + \frac{259 C_A C_F n_f T_f}{54\epsilon} + \frac{17 C_A C_F n_f T_f}{18\epsilon^2} + \frac{C_A C_F n_f T_f}{6\epsilon^3} + \frac{1}{54} C_A C_F n_f \pi^2 T_f + \frac{C_A C_F n_f \pi^2 T_f}{36\epsilon} + \frac{201881}{432} C_A C_F^2 C_Q \\
& + \frac{631}{432} C_A C_F C_Q - \frac{283861}{288} C_F C_Q - \frac{259}{27} C_A C_F n_f T_f C_Q - \frac{17 C_A C_F n_f T_f C_Q}{9\epsilon} - \frac{C_A C_F n_f T_f C_Q}{3\epsilon^2} - \frac{1}{18} C_A C_F n_f \pi^2 T_f C_Q \\
& - \frac{43897 C_A C_F^2 C_Q}{72\epsilon} + \frac{25 C_A C_F C_Q}{72\epsilon} - \frac{20131 C_F C_Q}{48\epsilon} - \frac{1783 C_A C_F^2 C_Q}{12\epsilon^2} + \frac{C_A C_F C_Q}{12\epsilon^2} - \frac{269 C_F C_Q}{8\epsilon^2} + \frac{20 C_A C_F^2 C_Q}{\epsilon^3} + \frac{6 C_F C_Q}{\epsilon^3} \\
& - \frac{427}{72} C_A C_F^2 \pi^2 C_Q + \frac{1}{72} C_A C_F \pi^2 C_Q + \frac{535}{48} C_F \pi^2 C_Q - \frac{C_A C_F^2 \pi^2 C_Q}{\epsilon} + \frac{5 C_F \pi^2 C_Q}{2\epsilon} - \frac{201881 C_A C_F^2}{864\epsilon} - \frac{631 C_A C_F}{864\epsilon} + \frac{283861 C_F}{576\epsilon} \\
& + \frac{43897 C_A C_F^2}{144\epsilon^2} - \frac{25 C_A C_F}{144\epsilon^2} + \frac{20131 C_F}{96\epsilon^2} + \frac{1783 C_A C_F^2}{24\epsilon^3} - \frac{C_A C_F}{24\epsilon^3} + \frac{269 C_F}{16\epsilon^3} - \frac{10 C_A C_F^2}{\epsilon^4} - \frac{3 C_F}{\epsilon^4} - \frac{581 C_F \zeta_3}{3} + \frac{166}{3} C_F C_Q \zeta_3 \\
& - \frac{83 C_F \zeta_3}{3\epsilon} - \frac{59 C_F \pi^4}{120} + \frac{1999}{864} C_A C_F^2 \pi^2 + \frac{5}{864} C_A C_F \pi^2 - \frac{20507 C_F \pi^2}{576} + \frac{427 C_A C_F^2 \pi^2}{144\epsilon} - \frac{C_A C_F \pi^2}{144\epsilon} - \frac{535 C_F \pi^2}{96\epsilon} + \frac{C_A C_F^2 \pi^2}{2\epsilon^2} \\
& - \frac{5 C_F \pi^2}{4\epsilon^2} \tag{C.2}
\end{aligned}$$

$$\begin{aligned}
V_3^{(Q)} = & \frac{1}{192} C_A C_Q^4 Y_f^4 - \frac{9}{32} C_A C_Q^3 Y_f^4 - \frac{C_A C_Q^3 Y_f^4}{96\epsilon} + \frac{201}{256} C_A C_Q^2 Y_f^4 + \frac{27 C_A C_Q^2 Y_f^4}{64\epsilon} + \frac{C_A C_Q^2 Y_f^4}{64\epsilon^2} - \frac{1075 C_A Y_f^4}{512} - \frac{99}{128} C_A C_Q Y_f^4 \\
& - \frac{201 C_A C_Q Y_f^4}{256\epsilon} - \frac{27 C_A C_Q Y_f^4}{64\epsilon^2} - \frac{C_A C_Q Y_f^4}{64\epsilon^3} + \frac{1}{48} C_A \pi^2 C_Q Y_f^4 + \frac{99 C_A Y_f^4}{256\epsilon} + \frac{201 C_A Y_f^4}{512\epsilon^2} + \frac{27 C_A Y_f^4}{128\epsilon^3} - \frac{83}{192} C_A \zeta_3 Y_f^4 \\
& - \frac{1}{24} C_A \pi^2 Y_f^4 - \frac{C_A \pi^2 Y_f^4}{96\epsilon} + \frac{3}{8} C_A C_F C_Q^4 Y_f^2 - \frac{33}{4} C_A C_F C_Q^3 Y_f^2 - \frac{3 C_A C_F C_Q^3 Y_f^2}{4\epsilon} + \frac{19}{4} C_A C_F^2 Y_f^2 + \frac{1}{2} C_A C_F^2 C_Q^2 Y_f^2 \\
& - \frac{209}{32} C_A C_F C_Q^2 Y_f^2 + \frac{1}{4} C_F C_Q^2 Y_f^2 + \frac{99 C_A C_F C_Q^2 Y_f^2}{8\epsilon} + \frac{9 C_A C_F C_Q^2 Y_f^2}{8\epsilon^2} + \frac{1}{8} C_A C_F \pi^2 C_Q^2 Y_f^2 - \frac{8079}{64} C_A C_F Y_f^2 + \frac{19 C_F Y_f^2}{8} \\
& - \frac{5}{2} C_A C_F^2 C_Q Y_f^2 + \frac{809}{4} C_A C_F C_Q Y_f^2 - \frac{5}{4} C_F C_Q Y_f^2 - \frac{C_A C_F^2 C_Q Y_f^2}{2\epsilon} + \frac{209 C_A C_F C_Q Y_f^2}{32\epsilon} - \frac{C_F C_Q Y_f^2}{4\epsilon} - \frac{99 C_A C_F C_Q Y_f^2}{8\epsilon^2} \\
& - \frac{9 C_A C_F C_Q Y_f^2}{8\epsilon^3} - 2 C_A C_F \pi^2 C_Q Y_f^2 - \frac{C_A C_F \pi^2 C_Q Y_f^2}{8\epsilon} + \frac{5 C_A C_F^2 Y_f^2}{4\epsilon} - \frac{809 C_A C_F Y_f^2}{8\epsilon} + \frac{5 C_F Y_f^2}{8\epsilon} + \frac{C_A C_F^2 Y_f^2}{4\epsilon^2} - \frac{209 C_A C_F Y_f^2}{64\epsilon^2} \\
& + \frac{C_F Y_f^2}{8\epsilon^2} + \frac{99 C_A C_F Y_f^2}{16\epsilon^3} + \frac{9 C_A C_F Y_f^2}{16\epsilon^4} + \frac{581}{24} C_A C_F \zeta_3 Y_f^2 + \frac{1}{24} C_A C_F^2 \pi^2 Y_f^2 + \frac{101}{24} C_A C_F \pi^2 Y_f^2 + \frac{1}{48} C_F \pi^2 Y_f^2 + \frac{C_A C_F \pi^2 Y_f^2}{\epsilon} \\
& + \frac{C_A C_F \pi^2 Y_f^2}{16\epsilon^2} - \frac{19}{3} C_A C_F^2 C_Q^4 - 3 C_F C_Q^4 - \frac{665}{6} C_A C_F^2 C_Q^3 + \frac{1}{18} C_A C_F C_Q^3 - \frac{149}{12} C_F C_Q^3 - \frac{2}{9} C_A C_F n_f T_f C_Q^3 + \frac{38 C_A C_F^2 C_Q^3}{3\epsilon} \\
& + \frac{6 C_F C_Q^3}{\epsilon} - \frac{5689}{54} C_A C_F^2 + \frac{6377}{12} C_A C_F^2 C_Q^2 - \frac{2}{9} C_A C_F C_Q^2 + \frac{1144}{3} C_F C_Q^2 - \frac{1}{9} C_A C_F n_f T_f C_Q^2 + \frac{C_A C_F n_f T_f C_Q^2}{3\epsilon} \\
& + \frac{665 C_A C_F^2 C_Q^2}{4\epsilon} - \frac{C_A C_F C_Q^2}{12\epsilon} + \frac{149 C_F C_Q^2}{8\epsilon} - \frac{19 C_A C_F^2 C_Q^2}{\epsilon^2} - \frac{9 C_F C_Q^2}{\epsilon^2} + C_A C_F^2 \pi^2 C_Q^2 - \frac{17}{6} C_F \pi^2 C_Q^2 + \frac{1129 C_A C_F}{648} \\
& + \frac{267173 C_F}{108} - \frac{3235}{162} C_A C_F n_f T_f - \frac{46 C_A C_F n_f T_f}{27\epsilon} - \frac{C_A C_F n_f T_f}{18\epsilon^2} + \frac{C_A C_F n_f T_f}{6\epsilon^3} - \frac{4}{27} C_A C_F n_f \pi^2 T_f + \frac{C_A C_F n_f \pi^2 T_f}{36\epsilon} \\
& + \frac{7088}{9} C_A C_F^2 C_Q + \frac{35}{54} C_A C_F C_Q - \frac{71539}{72} C_F C_Q + \frac{92}{27} C_A C_F n_f T_f C_Q + \frac{C_A C_F n_f T_f C_Q}{9\epsilon} - \frac{C_A C_F n_f T_f C_Q}{3\epsilon^2} \\
& - \frac{1}{18} C_A C_F n_f \pi^2 T_f C_Q - \frac{6377 C_A C_F^2 C_Q}{12\epsilon} + \frac{2 C_A C_F C_Q}{9\epsilon} - \frac{1144 C_F C_Q}{3\epsilon} - \frac{665 C_A C_F^2 C_Q}{4\epsilon^2} + \frac{C_A C_F C_Q}{12\epsilon^2} - \frac{149 C_F C_Q}{8\epsilon^2}
\end{aligned}$$

$$\begin{aligned}
& + \frac{19C_A C_F^2 \mathcal{L}Q}{\epsilon^3} + \frac{9C_F \mathcal{L}Q}{\epsilon^3} - \frac{55}{8} C_A C_F^2 \pi^2 \mathcal{L}Q + \frac{1}{72} C_A C_F \pi^2 \mathcal{L}Q + \frac{205}{16} C_F \pi^2 \mathcal{L}Q - \frac{C_A C_F^2 \pi^2 \mathcal{L}Q}{\epsilon} + \frac{17C_F \pi^2 \mathcal{L}Q}{6\epsilon} - \frac{3544C_A C_F^2}{9\epsilon} \\
& - \frac{35C_A C_F}{108\epsilon} + \frac{71539C_F}{144\epsilon} + \frac{6377C_A C_F^2}{24\epsilon^2} - \frac{C_A C_F}{9\epsilon^2} + \frac{572C_F}{3\epsilon^2} + \frac{665C_A C_F^2}{8\epsilon^3} - \frac{C_A C_F}{24\epsilon^3} + \frac{149C_F}{16\epsilon^3} - \frac{19C_A C_F^2}{2\epsilon^4} - \frac{9C_F}{2\epsilon^4} - \frac{664C_F \zeta_3}{3} \\
& + \frac{166}{3} C_F \mathcal{L}Q \zeta_3 - \frac{83C_F \zeta_3}{3\epsilon} - \frac{59C_F \pi^4}{120} + \frac{413}{144} C_A C_F^2 \pi^2 + \frac{7}{432} C_A C_F \pi^2 - \frac{11347C_F \pi^2}{288} + \frac{55C_A C_F^2 \pi^2}{16\epsilon} - \frac{C_A C_F \pi^2}{144\epsilon} - \frac{205C_F \pi^2}{32\epsilon} \\
& + \frac{C_A C_F^2 \pi^2}{2\epsilon^2} - \frac{17C_F \pi^2}{12\epsilon^2} \tag{C.3}
\end{aligned}$$

$$\begin{aligned}
V_4^{(Q)} = & - \frac{4n_f C_A C_F T_f \mathcal{L}Q^2}{3\epsilon} + \frac{32n_f C_A C_F T_f \mathcal{L}Q}{9\epsilon} + \frac{4n_f C_A C_F T_f \mathcal{L}Q}{3\epsilon^2} + \frac{8}{9} n_f C_A C_F T_f \mathcal{L}Q^3 - \frac{32}{9} n_f C_A C_F T_f \mathcal{L}Q^2 + \frac{280}{27} n_f C_A C_F T_f \mathcal{L}Q \\
& + \frac{2}{9} \pi^2 n_f C_A C_F T_f \mathcal{L}Q - \frac{140n_f C_A C_F T_f}{27\epsilon} - \frac{16n_f C_A C_F T_f}{9\epsilon^2} - \frac{2n_f C_A C_F T_f}{3\epsilon^3} - \frac{\pi^2 n_f C_A C_F T_f}{9\epsilon} + \frac{2258}{81} n_f C_A C_F T_f \\
& + \frac{7}{27} \pi^2 n_f C_A C_F T_f + \frac{1579C_A C_F^2}{72\epsilon} - \frac{2047C_A C_F}{216\epsilon} + \frac{677C_A C_F^2}{24\epsilon^2} - \frac{155C_A C_F}{72\epsilon^2} + \frac{69C_A C_F^2}{8\epsilon^3} - \frac{C_A C_F}{24\epsilon^3} - \frac{8C_A C_F^2}{\epsilon^4} - \frac{59\pi^2 C_A C_F^2}{48\epsilon} \\
& - \frac{\pi^2 C_A C_F}{144\epsilon} + \frac{\pi^2 C_A C_F^2}{2\epsilon^2} + \frac{32C_A C_F^2 \mathcal{L}Q^3}{3\epsilon} + \frac{69C_A C_F^2 \mathcal{L}Q^2}{4\epsilon} - \frac{C_A C_F \mathcal{L}Q^2}{12\epsilon} - \frac{16C_A C_F^2 \mathcal{L}Q^2}{\epsilon^2} - \frac{677C_A C_F^2 \mathcal{L}Q}{12\epsilon} + \frac{155C_A C_F \mathcal{L}Q}{36\epsilon} \\
& - \frac{69C_A C_F^2 \mathcal{L}Q}{4\epsilon^2} + \frac{C_A C_F \mathcal{L}Q}{12\epsilon^2} + \frac{16C_A C_F^2 \mathcal{L}Q}{\epsilon^3} - \frac{\pi^2 C_A C_F^2 \mathcal{L}Q}{\epsilon} - \frac{16}{3} C_A C_F^2 \mathcal{L}Q^4 - \frac{23}{2} C_A C_F^2 \mathcal{L}Q^3 + \frac{1}{18} C_A C_F \mathcal{L}Q^3 + \frac{677}{12} C_A C_F^2 \mathcal{L}Q^2 \\
& - \frac{155}{36} C_A C_F \mathcal{L}Q^2 + \pi^2 C_A C_F^2 \mathcal{L}Q^2 - \frac{1579}{36} C_A C_F^2 \mathcal{L}Q + \frac{2047}{108} C_A C_F \mathcal{L}Q + \frac{59}{24} \pi^2 C_A C_F^2 \mathcal{L}Q + \frac{1}{72} \pi^2 C_A C_F \mathcal{L}Q - \frac{3823}{54} C_A C_F^2 \\
& - \frac{5141C_A C_F}{162} + \frac{29}{36} \pi^2 C_A C_F^2 - \frac{35}{108} \pi^2 C_A C_F + \frac{20767C_F}{144\epsilon} - \frac{1201C_F}{48\epsilon^2} + \frac{253C_F}{16\epsilon^3} + \frac{6C_F}{\epsilon^4} - \frac{81\pi^2 C_F}{32\epsilon} - \frac{2\pi^2 C_F}{3\epsilon^2} - \frac{8C_F \mathcal{L}Q^3}{\epsilon} \\
& + \frac{253C_F \mathcal{L}Q^2}{8\epsilon} + \frac{12C_F \mathcal{L}Q^2}{\epsilon^2} + \frac{1201C_F \mathcal{L}Q}{24\epsilon} - \frac{253C_F \mathcal{L}Q}{8\epsilon^2} - \frac{12C_F \mathcal{L}Q}{\epsilon^3} + \frac{4\pi^2 C_F \mathcal{L}Q}{3\epsilon} + \frac{166}{3} \zeta_3 C_F \mathcal{L}Q + 4C_F \mathcal{L}Q^4 - \frac{253}{12} C_F \mathcal{L}Q^3 \\
& - \frac{1201}{24} C_F \mathcal{L}Q^2 - \frac{4}{3} \pi^2 C_F \mathcal{L}Q^2 - \frac{20767}{72} C_F \mathcal{L}Q + \frac{81}{16} \pi^2 C_F \mathcal{L}Q - \frac{83\zeta_3 C_F}{3\epsilon} - \frac{166\zeta_3 C_F}{3} + \frac{32033C_F}{108} - \frac{59\pi^4 C_F}{120} - \frac{269\pi^2 C_F}{36} \tag{C.4}
\end{aligned}$$

$$\begin{aligned}
V_5^{(Q)} = & - \frac{4n_f C_A C_F T_f \mathcal{L}Q^2}{3\epsilon} + \frac{68n_f C_A C_F T_f \mathcal{L}Q}{9\epsilon} + \frac{4n_f C_A C_F T_f \mathcal{L}Q}{3\epsilon^2} + \frac{8}{9} n_f C_A C_F T_f \mathcal{L}Q^3 - \frac{68}{9} n_f C_A C_F T_f \mathcal{L}Q^2 + \frac{802}{27} n_f C_A C_F T_f \mathcal{L}Q \\
& + \frac{2}{9} \pi^2 n_f C_A C_F T_f \mathcal{L}Q - \frac{401n_f C_A C_F T_f}{27\epsilon} - \frac{34n_f C_A C_F T_f}{9\epsilon^2} - \frac{2n_f C_A C_F T_f}{3\epsilon^3} - \frac{\pi^2 n_f C_A C_F T_f}{9\epsilon} - \frac{1361}{162} n_f C_A C_F T_f \\
& - \frac{2}{27} \pi^2 n_f C_A C_F T_f + \frac{7649C_A C_F^2}{144\epsilon} - \frac{953C_A C_F}{432\epsilon} + \frac{301C_A C_F^2}{12\epsilon^2} - \frac{5C_A C_F}{18\epsilon^2} - \frac{69C_A C_F^2}{\epsilon^3} + \frac{477C_A^2 C_F}{16\epsilon^3} - \frac{C_A C_F}{24\epsilon^3} - \frac{20C_A C_F^2}{\epsilon^4} \\
& + \frac{6C_A^2 C_F}{\epsilon^4} + \frac{37\pi^2 C_A C_F^2}{48\epsilon} - \frac{\pi^2 C_A C_F}{144\epsilon} + \frac{\pi^2 C_A C_F^2}{2\epsilon^2} + \frac{32C_A C_F^2 \mathcal{L}Q^3}{3\epsilon} - \frac{75C_A C_F^2 \mathcal{L}Q^2}{4\epsilon} - \frac{C_A C_F \mathcal{L}Q^2}{12\epsilon} - \frac{16C_A C_F^2 \mathcal{L}Q^2}{\epsilon^2} \\
& - \frac{301C_A C_F^2 \mathcal{L}Q}{6\epsilon} + \frac{45C_A C_F \mathcal{L}Q}{9\epsilon} + \frac{75C_A C_F^2 \mathcal{L}Q}{4\epsilon^2} + \frac{C_A C_F \mathcal{L}Q}{12\epsilon^2} + \frac{40C_A C_F^2 \mathcal{L}Q}{\epsilon^3} - \frac{12C_A^2 C_F \mathcal{L}Q}{\epsilon^3} - \frac{\pi^2 C_A C_F^2 \mathcal{L}Q}{\epsilon} - \frac{16}{3} C_A C_F^2 \mathcal{L}Q^4 \\
& + \frac{25}{2} C_A C_F^2 \mathcal{L}Q^3 + \frac{1}{18} C_A C_F \mathcal{L}Q^3 + \frac{301}{6} C_A C_F^2 \mathcal{L}Q^2 - \frac{5}{9} C_A C_F \mathcal{L}Q^2 + \pi^2 C_A C_F^2 \mathcal{L}Q^2 - \frac{7649}{72} C_A C_F^2 \mathcal{L}Q + \frac{953}{216} C_A C_F \mathcal{L}Q \\
& - \frac{37}{24} \pi^2 C_A C_F^2 \mathcal{L}Q + \frac{1}{72} \pi^2 C_A C_F \mathcal{L}Q + \frac{62051}{864} C_A C_F^2 - \frac{20039C_A C_F}{2592} + \frac{41}{144} \pi^2 C_A C_F^2 - \frac{5}{432} \pi^2 C_A C_F + \frac{59201C_F}{288\epsilon} + \frac{116C_F}{3\epsilon^2} \\
& - \frac{355\pi^2 C_F}{96\epsilon} - \frac{2\pi^2 C_F}{3\epsilon^2} - \frac{8C_F \mathcal{L}Q^3}{\epsilon} + \frac{477C_F \mathcal{L}Q^2}{8\epsilon} + \frac{12C_F \mathcal{L}Q^2}{\epsilon^2} - \frac{232C_F \mathcal{L}Q}{3\epsilon} - \frac{477C_F \mathcal{L}Q}{8\epsilon^2} + \frac{4\pi^2 C_F \mathcal{L}Q}{3\epsilon} + \frac{166}{3} \zeta_3 C_F \mathcal{L}Q + 4C_F \mathcal{L}Q^4 \\
& - \frac{159}{4} C_F \mathcal{L}Q^3 + \frac{232}{3} C_F \mathcal{L}Q^2 - \frac{4}{3} \pi^2 C_F \mathcal{L}Q^2 - \frac{59201}{144} C_F \mathcal{L}Q + \frac{355}{48} \pi^2 C_F \mathcal{L}Q - \frac{83\zeta_3 C_F}{3\epsilon} - \frac{332\zeta_3 C_F}{3} + \frac{1265063C_F}{1728} - \frac{59\pi^4 C_F}{120} \\
& - \frac{4375\pi^2 C_F}{288} \tag{C.5}
\end{aligned}$$

$$\begin{aligned}
V_6^{(Q)} = & \frac{11}{8} C_A^3 \mathcal{L}Q^4 - \frac{1}{48} C_A^2 Y_f^2 \mathcal{L}Q^4 + \frac{1}{48} Y_f^2 \mathcal{L}Q^4 - \frac{11}{8} C_A \mathcal{L}Q^4 - \frac{35}{6} C_A^2 C_F \mathcal{L}Q^4 + \frac{35}{6} C_F \mathcal{L}Q^4 - \frac{427}{24} C_A^3 \mathcal{L}Q^3 + \frac{1}{36} C_A^2 \mathcal{L}Q^3 - \frac{7}{24} C_A^2 Y_f^2 \mathcal{L}Q^3 \\
& + \frac{C_A^2 Y_f^2 \mathcal{L}Q^3}{24\epsilon} - \frac{Y_f^2 \mathcal{L}Q^3}{24\epsilon} + \frac{7}{24} Y_f^2 \mathcal{L}Q^3 + \frac{427}{24} C_A \mathcal{L}Q^3 + \frac{58}{3} C_A^2 C_F \mathcal{L}Q^3 - \frac{58}{3} C_F \mathcal{L}Q^3 - \frac{11C_A^3 \mathcal{L}Q^3}{4\epsilon} + \frac{11C_A \mathcal{L}Q^3}{4\epsilon} + \frac{35C_A^2 C_F \mathcal{L}Q^3}{3\epsilon} \\
& - \frac{35C_F \mathcal{L}Q^3}{3\epsilon} - \frac{\mathcal{L}Q^3}{36} + \frac{4543}{96} C_A^3 \mathcal{L}Q^2 - \frac{7}{36} C_A^2 \mathcal{L}Q^2 + \frac{107}{64} C_A^2 Y_f^2 \mathcal{L}Q^2 + \frac{7C_A^2 Y_f^2 \mathcal{L}Q^2}{16\epsilon} - \frac{7Y_f^2 \mathcal{L}Q^2}{16\epsilon} - \frac{C_A^2 Y_f^2 \mathcal{L}Q^2}{16\epsilon^2} + \frac{Y_f^2 \mathcal{L}Q^2}{16\epsilon^2} - \frac{107}{64} Y_f^2 \mathcal{L}Q^2 \\
& - \frac{4543}{96} C_A \mathcal{L}Q^2 + \frac{103}{8} C_A^2 C_F \mathcal{L}Q^2 - \frac{103}{8} C_F \mathcal{L}Q^2 - \frac{1}{3} C_A^2 n_f T_f \mathcal{L}Q^2 + \frac{1}{3} n_f T_f \mathcal{L}Q^2 + \frac{427C_A^3 \mathcal{L}Q^2}{16\epsilon} - \frac{C_A^2 \mathcal{L}Q^2}{24\epsilon} - \frac{427C_A \mathcal{L}Q^2}{16\epsilon} \\
& - \frac{29C_A^2 C_F \mathcal{L}Q^2}{\epsilon} + \frac{29C_F \mathcal{L}Q^2}{\epsilon} + \frac{\mathcal{L}Q^2}{24\epsilon} + \frac{33C_A^3 \mathcal{L}Q^2}{8\epsilon^2} - \frac{33C_A \mathcal{L}Q^2}{8\epsilon^2} - \frac{35C_A^2 C_F \mathcal{L}Q^2}{2\epsilon^2} + \frac{35C_F \mathcal{L}Q^2}{2\epsilon^2} - \frac{3}{4} C_A^3 \pi^2 \mathcal{L}Q^2 \\
& + \frac{3}{4} C_A \pi^2 \mathcal{L}Q^2 + 2C_A^2 C_F \pi^2 \mathcal{L}Q^2 - 2C_F \pi^2 \mathcal{L}Q^2 + \frac{7\mathcal{L}Q^2}{36} - \frac{9865}{144} C_A^3 \mathcal{L}Q + \frac{77}{108} C_A^2 \mathcal{L}Q + \frac{623}{64} C_A^2 Y_f^2 \mathcal{L}Q - \frac{107C_A^2 Y_f^2 \mathcal{L}Q}{64\epsilon} + \frac{107Y_f^2 \mathcal{L}Q}{64\epsilon} \\
& - \frac{7C_A^2 Y_f^2 \mathcal{L}Q}{16\epsilon^2} + \frac{7Y_f^2 \mathcal{L}Q}{16\epsilon^2} + \frac{C_A^2 Y_f^2 \mathcal{L}Q}{16\epsilon^3} - \frac{Y_f^2 \mathcal{L}Q}{16\epsilon^3} - \frac{1}{32} C_A^2 \pi^2 Y_f^2 \mathcal{L}Q + \frac{1}{32} \pi^2 Y_f^2 \mathcal{L}Q - \frac{623}{64} Y_f^2 \mathcal{L}Q + \frac{9865}{144} C_A \mathcal{L}Q + \frac{1153}{4} C_A^2 C_F \mathcal{L}Q
\end{aligned}$$

$$\begin{aligned}
& -\frac{1153}{4}C_F\mathcal{L}Q + \frac{11}{9}C_A^2n_fT_f\mathcal{L}Q + \frac{C_A^2n_fT_f\mathcal{L}Q}{3\epsilon} - \frac{n_fT_f\mathcal{L}Q}{3\epsilon} - \frac{11}{9}n_fT_f\mathcal{L}Q - \frac{4543C_A^3\mathcal{L}Q}{96\epsilon} + \frac{7C_A^2\mathcal{L}Q}{36\epsilon} + \frac{4543C_A\mathcal{L}Q}{96\epsilon} \\
& - \frac{103C_A^2C_F\mathcal{L}Q}{8\epsilon} + \frac{103C_F\mathcal{L}Q}{8\epsilon} - \frac{7\mathcal{L}Q}{36\epsilon} - \frac{427C_A^3\mathcal{L}Q}{16\epsilon^2} + \frac{C_A^2\mathcal{L}Q}{24\epsilon^2} + \frac{427C_A\mathcal{L}Q}{16\epsilon^2} + \frac{29C_A^2C_F\mathcal{L}Q}{\epsilon^2} - \frac{29C_F\mathcal{L}Q}{\epsilon^2} - \frac{\mathcal{L}Q}{24\epsilon^2} - \frac{33C_A^3\mathcal{L}Q}{8\epsilon^3} \\
& + \frac{33C_A\mathcal{L}Q}{8\epsilon^3} + \frac{35C_A^2C_F\mathcal{L}Q}{2\epsilon^3} - \frac{35C_F\mathcal{L}Q}{2\epsilon^3} + \frac{83}{3}C_A^3\zeta_3\mathcal{L}Q - \frac{83}{3}C_A\zeta_3\mathcal{L}Q - \frac{166}{3}C_A^2C_F\zeta_3\mathcal{L}Q + \frac{166}{3}C_F\zeta_3\mathcal{L}Q + \frac{81}{32}C_A^3\pi^2\mathcal{L}Q \\
& + \frac{1}{144}C_A^2\pi^2\mathcal{L}Q - \frac{81}{32}C_A\pi^2\mathcal{L}Q - \frac{71}{12}C_A^2C_F\pi^2\mathcal{L}Q + \frac{71}{12}C_F\pi^2\mathcal{L}Q + \frac{3C_A^3\pi^2\mathcal{L}Q}{4\epsilon} - \frac{3C_A\pi^2\mathcal{L}Q}{4\epsilon} - \frac{2C_A^2C_F\pi^2\mathcal{L}Q}{\epsilon} + \frac{2C_F\pi^2\mathcal{L}Q}{\epsilon} \\
& - \frac{1}{144}\pi^2\mathcal{L}Q - \frac{77\mathcal{L}Q}{108} + \frac{684445C_A^3}{1728} + \frac{205C_A^2}{1296} + \frac{1147}{128}C_A^2Y_f^2 - \frac{623C_A^2Y_f^2}{128\epsilon} + \frac{623Y_f^2}{128\epsilon} + \frac{107C_A^2Y_f^2}{128\epsilon^2} - \frac{107Y_f^2}{128\epsilon^2} + \frac{7C_A^2Y_f^2}{32\epsilon^3} - \frac{7Y_f^2}{32\epsilon^3} \\
& - \frac{C_A^2Y_f^2}{32\epsilon^4} + \frac{Y_f^2}{32\epsilon^4} - \frac{7}{96}C_A^2\pi^2Y_f^2 + \frac{C_A^2\pi^2Y_f^2}{64\epsilon} - \frac{\pi^2Y_f^2}{64\epsilon} + \frac{7\pi^2Y_f^2}{96} - \frac{1147Y_f^2}{128} - \frac{684445C_A}{1728} - \frac{5179}{8}C_A^2C_F + \frac{5179C_F}{8} \\
& - \frac{56}{27}C_A^2n_fT_f - \frac{11C_A^2n_fT_f}{18\epsilon} + \frac{11n_fT_f}{18\epsilon} - \frac{C_A^2n_fT_f}{6\epsilon^2} + \frac{n_fT_f}{6\epsilon^2} + \frac{56n_fT_f}{27} - \frac{1}{36}C_A^2n_f\pi^2T_f + \frac{1}{36}n_f\pi^2T_f + \frac{9865C_A^3}{288\epsilon} - \frac{77C_A^2}{216\epsilon} \\
& - \frac{9865C_A}{288\epsilon} - \frac{1153C_A^2C_F}{8\epsilon} + \frac{1153C_F}{8\epsilon} + \frac{77}{216\epsilon} + \frac{4543C_A^3}{192\epsilon^2} - \frac{7C_A^2}{72\epsilon^2} - \frac{4543C_A}{192\epsilon^2} + \frac{103C_A^2C_F}{16\epsilon^2} - \frac{103C_F}{16\epsilon^2} + \frac{7}{72\epsilon^2} + \frac{427C_A^3}{32\epsilon^3} - \frac{C_A^2}{48\epsilon^3} \\
& - \frac{427C_A}{32\epsilon^3} - \frac{29C_A^2C_F}{2\epsilon^3} + \frac{29C_F}{48\epsilon^3} + \frac{1}{16\epsilon^4} - \frac{33C_A^3}{16\epsilon^4} - \frac{33C_A}{4\epsilon^4} - \frac{35C_A^2C_F}{4\epsilon^4} + \frac{35C_F}{2} - \frac{83C_A^3\zeta_3}{2} + \frac{83C_A\zeta_3}{2} + 83C_A^2C_F\zeta_3 - 83C_F\zeta_3 \\
& - \frac{83C_A^3\zeta_3}{6\epsilon} + \frac{83C_A\zeta_3}{6\epsilon} + \frac{83C_A^2C_F\zeta_3}{3\epsilon} - \frac{83C_F\zeta_3}{3\epsilon} - \frac{59}{240}C_A^3\pi^4 + \frac{59C_A\pi^4}{240} + \frac{59}{120}C_A^2C_F\pi^4 - \frac{59C_F\pi^4}{120} - \frac{3829}{576}C_A^3\pi^2 + \frac{1}{864}C_A^2\pi^2 \\
& + \frac{3829C_A\pi^2}{576} + \frac{12C_A^2C_F\pi^2}{12} - \frac{12C_F\pi^2}{12} - \frac{81C_A^3\pi^2}{64\epsilon} - \frac{C_A^2\pi^2}{288\epsilon} + \frac{81C_A\pi^2}{64\epsilon} + \frac{71C_A^2C_F\pi^2}{24\epsilon} - \frac{71C_F\pi^2}{24\epsilon} + \frac{\pi^2}{288\epsilon} - \frac{3C_A^3\pi^2}{8\epsilon^2} + \frac{3C_A\pi^2}{8\epsilon^2} \\
& + \frac{C_A^2C_F\pi^2}{\epsilon^2} - \frac{C_F\pi^2}{\epsilon^2} - \frac{\pi^2}{864} - \frac{205}{1296}
\end{aligned} \tag{C.6}$$

## C.2 Matching at $\mu \sim m_{1,2}$ :

$$\begin{aligned}
V_1^{(m_{1,2})} &= \frac{n_f C_A C_F T_f \mathcal{L} m_2}{2\epsilon} - \frac{1}{2} n_f C_A C_F T_f \mathcal{L} m_2^2 - \frac{n_f C_A C_F T_f}{4\epsilon^2} - \frac{5}{4} n_f \zeta_2 C_A C_F T_f - \frac{1}{2} n_f C_A C_F T_f - \frac{45 C_A C_F^2}{8\epsilon} + \frac{3 C_A C_F}{4\epsilon} + \\
& \frac{23 C_A C_F^2}{8\epsilon^2} + \frac{C_A C_F}{4\epsilon^2} - \frac{7 C_A C_F^2}{8\epsilon^3} - \frac{7 C_A C_F^2 \mathcal{L} m_2}{4\epsilon} - \frac{23 C_A C_F^2 \mathcal{L} m_2}{4\epsilon} - \frac{C_A C_F \mathcal{L} m_2}{2\epsilon} + \frac{7 C_A C_F^2 \mathcal{L} m_2}{4\epsilon^2} - \frac{6 Y_f^2 C_A C_F \mathcal{L} m_2}{\epsilon} - \\
& \frac{3 Y_f^2 C_A C_F^2 \mathcal{L} m_2}{2\epsilon} - \frac{Y_f^2 C_A C_F \mathcal{L} m_2}{\epsilon} + \frac{6 Y_f^2 C_A C_F \mathcal{L} m_2}{\epsilon^2} + \frac{3}{2} Y_f^2 C_A C_F^2 \mathcal{L} m_2 + 16 Y_f^2 C_A C_F \mathcal{L} m_2 - \frac{13}{2} Y_f^2 C_A C_F^2 \mathcal{L} m_2 - \\
& \frac{21}{2} Y_f^2 C_A C_F \mathcal{L} m_2 + \frac{81}{8} C_A C_F^2 \mathcal{L} m_2 + \frac{1}{2} C_A C_F \mathcal{L} m_2 + \frac{103}{16} C_A C_F^2 \mathcal{L} m_2 - \frac{3}{2} C_A C_F \mathcal{L} m_2 + \frac{13 Y_f^2 C_A C_F^2}{4\epsilon} - \frac{3 Y_f^2 C_A C_F}{\epsilon} + \\
& \frac{3 Y_f^2 C_A C_F^2}{4\epsilon^2} + \frac{Y_f^2 C_A C_F}{2\epsilon^2} - \frac{3 Y_f^2 C_A C_F}{\epsilon^3} - \frac{15 Y_f^2 \zeta_2 C_A C_F}{\epsilon} + \frac{15}{4} Y_f^2 \zeta_2 C_A C_F^2 + 40 Y_f^2 \zeta_2 C_A C_F + \frac{47}{4} Y_f^2 C_A C_F^2 - \frac{13}{8} Y_f^2 C_A C_F - \\
& \frac{35 \zeta_2 C_A C_F^2}{8\epsilon} + \frac{405}{16} \zeta_2 C_A C_F^2 + \frac{5}{4} \zeta_2 C_A C_F - \frac{3977}{64} C_A C_F^2 + \frac{9 C_A C_F}{4} + \frac{3 Y_f^4 C_A \mathcal{L} m_2}{4\epsilon} - \frac{3 Y_f^4 C_A \mathcal{L} m_2}{4\epsilon^2} - \frac{15}{8} Y_f^4 C_A \mathcal{L} m_2 + \\
& \frac{7}{4} Y_f^4 C_A \mathcal{L} m_2 + \frac{5 Y_f^4 C_A}{32\epsilon} + \frac{3 Y_f^4 C_A}{8\epsilon^3} + \frac{15 Y_f^4 \zeta_2 C_A}{8\epsilon} - \frac{75}{16} Y_f^4 \zeta_2 C_A - \frac{59 Y_f^4 C_A}{64} - \frac{237 C_F}{16\epsilon} - \frac{57 C_F}{16\epsilon^2} - \frac{55 C_F}{16\epsilon^3} - \frac{55 C_F \mathcal{L} m_2}{8\epsilon} + \\
& \frac{57 C_F \mathcal{L} m_2}{8\epsilon} + \frac{55 C_F \mathcal{L} m_2}{8\epsilon^2} - \frac{3 Y_f^2 C_F \mathcal{L} m_2}{4\epsilon} + \frac{3}{4} Y_f^2 C_F \mathcal{L} m_2 - \frac{13}{4} Y_f^2 C_F \mathcal{L} m_2 + \frac{161}{16} C_F \mathcal{L} m_2 + \frac{343}{32} C_F \mathcal{L} m_2 + \frac{13 Y_f^2 C_F}{8\epsilon} + \\
& \frac{3 Y_f^2 C_F}{8\epsilon^2} + \frac{15}{8} Y_f^2 \zeta_2 C_F + \frac{47 Y_f^2 C_F}{8} - \frac{275 \zeta_2 C_F}{16\epsilon} + \frac{805 \zeta_2 C_F}{32} - \frac{6489 C_F}{128}
\end{aligned} \tag{C.7}$$

$$\begin{aligned}
V_2^{(m_{1,2})} &= -\frac{9}{32} C_A \mathcal{L} m_2^2 Y_f^4 - \frac{C_A \mathcal{L} m_2^2 Y_f^4}{32\epsilon} - \frac{219 C_A Y_f^4}{512} - \frac{45}{64} C_A \zeta_2 Y_f^4 - \frac{5 C_A \zeta_2 Y_f^4}{64\epsilon} + \frac{41}{128} C_A \mathcal{L} m_2 Y_f^4 + \frac{23 C_A \mathcal{L} m_2 Y_f^4}{64\epsilon} + \frac{C_A \mathcal{L} m_2 Y_f^4}{32\epsilon^2} \\
& - \frac{13 C_A Y_f^4}{64\epsilon} - \frac{23 C_A Y_f^4}{128\epsilon^2} - \frac{C_A Y_f^4}{64\epsilon^3} + \frac{19}{8} C_A C_F^2 Y_f^2 + \frac{1}{4} C_A C_F^2 \mathcal{L} m_2 Y_f^2 + \frac{33}{8} C_A C_F \mathcal{L} m_2 Y_f^2 + \frac{1}{8} C_F \mathcal{L} m_2 Y_f^2 - \frac{C_A C_F \mathcal{L} m_2 Y_f^2}{4\epsilon} \\
& + \frac{1129}{64} C_A C_F Y_f^2 + \frac{19 C_F Y_f^2}{16} + \frac{5}{8} C_A C_F^2 \zeta_2 Y_f^2 + \frac{165}{16} C_A C_F \zeta_2 Y_f^2 + \frac{5}{16} C_F \zeta_2 Y_f^2 - \frac{5 C_A C_F \zeta_2 Y_f^2}{8\epsilon} - \frac{5}{4} C_A C_F^2 \mathcal{L} m_2 Y_f^2 \\
& - \frac{137}{16} C_A C_F \mathcal{L} m_2 Y_f^2 - \frac{5}{8} C_F \mathcal{L} m_2 Y_f^2 - \frac{C_A C_F^2 \mathcal{L} m_2 Y_f^2}{4\epsilon} - \frac{7 C_A C_F \mathcal{L} m_2 Y_f^2}{2\epsilon} - \frac{C_F \mathcal{L} m_2 Y_f^2}{8\epsilon} + \frac{C_A C_F \mathcal{L} m_2 Y_f^2}{4\epsilon^2} + \frac{5 C_A C_F^2 Y_f^2}{8\epsilon} \\
& + \frac{63 C_A C_F Y_f^2}{16\epsilon} + \frac{5 C_F Y_f^2}{16\epsilon} + \frac{C_A C_F^2 Y_f^2}{8\epsilon^2} + \frac{7 C_A C_F Y_f^2}{4\epsilon^2} + \frac{C_F Y_f^2}{16\epsilon^2} - \frac{C_A C_F Y_f^2}{8\epsilon^3} + \frac{C_A C_F M_W \mathcal{L} m_2 Y_f}{4m_-} + \frac{C_A C_F M_W \mathcal{L} m_2 Y_f}{2\epsilon m_-} \\
& + \frac{175 C_A C_F M_W Y_f}{32m_-} + \frac{7 C_A C_F M_W Y_f}{4\epsilon m_-} + \frac{3 C_A C_F M_W Y_f}{4\epsilon^2 m_-} + \frac{C_A C_F M_W Y_f}{4\epsilon^3 m_-} + \frac{5 C_A C_F M_W \zeta_2 Y_f}{8m_-} + \frac{5 C_A C_F M_W \zeta_2 Y_f}{4\epsilon m_-} \\
& - \frac{17 C_A C_F M_W \mathcal{L} m_2 Y_f}{8\epsilon m_-} - \frac{3 C_A C_F M_W \mathcal{L} m_2 Y_f}{2\epsilon m_-} - \frac{C_A C_F M_W \mathcal{L} m_2 Y_f}{2\epsilon^2 m_-} - \frac{495}{4} C_A C_F^2 - \frac{141}{16} C_A C_F^2 \mathcal{L} m_2 - \frac{7}{16} C_A C_F \mathcal{L} m_2^2 \\
& - \frac{53}{32} C_F \mathcal{L} m_2^2 + \frac{C_A C_F^2 m_2 \mathcal{L} m_2^2}{m_1 - m_2} + \frac{C_F m_2 \mathcal{L} m_2^2}{2m_-} + \frac{2 C_A C_F^2 m_2 \mathcal{L} m_2^2}{\epsilon m_-} + \frac{C_F m_2 \mathcal{L} m_2^2}{\epsilon m_-} - \frac{1}{4} C_A C_F n_f T_f \mathcal{L} m_2 - \frac{3 C_A C_F^2 \mathcal{L} m_2^2}{\epsilon}
\end{aligned}$$



$$\begin{aligned}
& -\frac{4C_F\mathcal{L}_2^2}{\epsilon} - \frac{41C_A C_F}{16} - \frac{1503C_F}{32} + \frac{175C_A C_F^2 m_2}{8m_-} + \frac{175C_F m_2}{16m_-} + \frac{7C_A C_F^2 m_2}{\epsilon m_-} + \frac{7C_F m_2}{2\epsilon m_-} + \frac{3C_A C_F^2 m_2}{\epsilon^2 m_-} + \frac{3C_F m_2}{2\epsilon^2 m_-} \\
& + \frac{C_A C_F^2 m_2}{\epsilon^3 m_-} + \frac{C_F m_2}{2\epsilon^3 m_-} - \frac{85}{32} C_A C_F n_f T_f - \frac{11C_A C_F n_f T_f}{32\epsilon} - \frac{C_A C_F n_f T_f}{8\epsilon^2} - \frac{705}{32} C_A C_F^2 \zeta_2 - \frac{35}{32} C_A C_F \zeta_2 - \frac{265C_F \zeta_2}{64} \\
& + \frac{5C_A C_F^2 m_2 \zeta_2}{2m_-} + \frac{5C_F m_2 \zeta_2}{4m_-} + \frac{5C_A C_F^2 m_2 \zeta_2}{\epsilon m_-} + \frac{5C_F m_2 \zeta_2}{2\epsilon m_-} - \frac{5}{8} C_A C_F n_f T_f \zeta_2 - \frac{15C_A C_F^2 \zeta_2}{2\epsilon} - \frac{10C_F \zeta_2}{\epsilon} + \frac{989}{16} C_A C_F^2 \mathcal{L}_{m_2} \\
& + \frac{19}{16} C_A C_F \mathcal{L}_{m_2} + \frac{833}{32} C_F \mathcal{L}_{m_2} - \frac{17C_A C_F^2 m_2 \mathcal{L}_{m_2}}{2m_-} - \frac{17C_F m_2 \mathcal{L}_{m_2}}{4m_-} - \frac{6C_A C_F^2 m_2 \mathcal{L}_{m_2}}{\epsilon m_-} - \frac{3C_F m_2 \mathcal{L}_{m_2}}{\epsilon m_-} - \frac{2C_A C_F^2 m_2 \mathcal{L}_{m_2}}{\epsilon^2 m_-} \\
& - \frac{C_F m_2 \mathcal{L}_{m_2}}{\epsilon^2 m_-} + \frac{11}{16} C_A C_F n_f T_f \mathcal{L}_{m_2} + \frac{C_A C_F n_f T_f \mathcal{L}_{m_2}}{4\epsilon} + \frac{261C_A C_F^2 \mathcal{L}_{m_2}}{16\epsilon} + \frac{7C_A C_F \mathcal{L}_{m_2}}{16\epsilon} + \frac{373C_F \mathcal{L}_{m_2}}{32\epsilon} + \frac{3C_A C_F^2 \mathcal{L}_{m_2}}{\epsilon^2} \\
& + \frac{4C_F \mathcal{L}_{m_2}}{\epsilon^2} - \frac{1121C_A C_F^2}{32\epsilon} - \frac{19C_A C_F}{32\epsilon} - \frac{1185C_F}{64\epsilon} - \frac{261C_A C_F^2}{32\epsilon^2} - \frac{7C_A C_F}{32\epsilon^2} - \frac{373C_F}{64\epsilon^2} - \frac{3C_A C_F^2}{2\epsilon^3} - \frac{2C_F}{\epsilon^3} \quad (C.8)
\end{aligned}$$

$$\begin{aligned}
V_3^{(m_{1,2})} &= \frac{n_f C_A C_F T_f \mathcal{L}_{m_2}}{2\epsilon} - \frac{1}{2} n_f C_A C_F T_f \mathcal{L}_{m_2}^2 + n_f C_A C_F T_f \mathcal{L}_{m_2} - \frac{n_f C_A C_F T_f}{2\epsilon} - \frac{n_f C_A C_F T_f}{4\epsilon^2} \\
& - \frac{5}{4} n_f \zeta_2 C_A C_F T_f - n_f C_A C_F T_f - \frac{24C_A C_F^2}{\epsilon} + \frac{3C_A C_F}{4\epsilon} - \frac{47C_A C_F^2}{8\epsilon^2} + \frac{C_A C_F}{4\epsilon^2} - \frac{3C_A C_F^2}{4\epsilon^3} - \frac{3C_A C_F^2 \mathcal{L}_{m_2}^2}{2\epsilon} \\
& + \frac{47C_A C_F^2 \mathcal{L}_{m_2}}{4\epsilon} - \frac{C_A C_F \mathcal{L}_{m_2}}{2\epsilon} + \frac{3C_A C_F^2 \mathcal{L}_{m_2}}{2\epsilon^2} - \frac{6Y_f^2 C_A C_F \mathcal{L}_{m_2}^2}{\epsilon} - \frac{3Y_f^2 C_A C_F^2 \mathcal{L}_{m_2}}{2\epsilon} + \frac{2Y_f^2 C_A C_F \mathcal{L}_{m_2}}{\epsilon} \\
& + \frac{6Y_f^2 C_A C_F \mathcal{L}_{m_2}}{\epsilon^2} + \frac{3}{2} Y_f^2 C_A C_F^2 \mathcal{L}_{m_2}^2 + 13Y_f^2 C_A C_F \mathcal{L}_{m_2}^2 - \frac{13}{2} Y_f^2 C_A C_F^2 \mathcal{L}_{m_2} - \frac{17}{2} Y_f^2 C_A C_F \mathcal{L}_{m_2} - 8C_A C_F^2 \mathcal{L}_{m_2}^2 \\
& + \frac{1}{2} C_A C_F \mathcal{L}_{m_2}^2 + \frac{351}{8} C_A C_F^2 \mathcal{L}_{m_2} - \frac{3}{2} C_A C_F \mathcal{L}_{m_2} + \frac{13Y_f^2 C_A C_F^2}{4\epsilon} - \frac{4Y_f^2 C_A C_F}{\epsilon} + \frac{3Y_f^2 C_A C_F^2}{4\epsilon^2} - \frac{Y_f^2 C_A C_F}{\epsilon^2} - \frac{3Y_f^2 C_A C_F}{\epsilon^3} \\
& - \frac{15Y_f^2 \zeta_2 C_A C_F}{\epsilon} + \frac{15}{4} Y_f^2 \zeta_2 C_A C_F^2 + \frac{65}{2} Y_f^2 \zeta_2 C_A C_F + \frac{47}{4} Y_f^2 C_A C_F^2 + \frac{7}{8} Y_f^2 C_A C_F - \frac{15\zeta_2 C_A C_F^2}{4\epsilon} - 20\zeta_2 C_A C_F^2 \\
& + \frac{5}{4} \zeta_2 C_A C_F - \frac{3165}{32} C_A C_F^2 + \frac{9C_A C_F}{4} + \frac{3Y_f^4 C_A \mathcal{L}_{m_2}^2}{4\epsilon} - \frac{3Y_f^4 C_A \mathcal{L}_{m_2}}{4\epsilon^2} - \frac{15}{8} Y_f^4 C_A \mathcal{L}_{m_2}^2 + \frac{7}{4} Y_f^4 C_A \mathcal{L}_{m_2} + \frac{5Y_f^4 C_A}{32\epsilon} \\
& + \frac{3Y_f^4 C_A}{8\epsilon^3} + \frac{15Y_f^4 \zeta_2 C_A}{8\epsilon} - \frac{75}{16} Y_f^4 \zeta_2 C_A - \frac{59Y_f^4 C_A}{64} - \frac{26C_F}{\epsilon} - \frac{95C_F}{16\epsilon^2} - \frac{27C_F}{8\epsilon^3} - \frac{27C_F \mathcal{L}_{m_2}^2}{4\epsilon} + \frac{95C_F \mathcal{L}_{m_2}}{8\epsilon} + \frac{27C_F \mathcal{L}_{m_2}}{4\epsilon^2} \\
& - \frac{3Y_f^2 C_F \mathcal{L}_{m_2}}{4\epsilon} + \frac{3}{4} Y_f^2 C_F \mathcal{L}_{m_2}^2 - \frac{13}{4} Y_f^2 C_F \mathcal{L}_{m_2} + 5C_F \mathcal{L}_{m_2}^2 + \frac{535}{16} C_F \mathcal{L}_{m_2} + \frac{13Y_f^2 C_F}{8\epsilon} + \frac{3Y_f^2 C_F}{8\epsilon^2} + \frac{15}{8} Y_f^2 \zeta_2 C_F + \frac{47Y_f^2 C_F}{8} \\
& - \frac{135\zeta_2 C_F}{8\epsilon} + \frac{25\zeta_2 C_F}{2} - \frac{4293C_F}{64} \quad (C.9)
\end{aligned}$$

$$\begin{aligned}
V_4^{(m_{1,2})} &= -\frac{27C_A C_F^2 Y_s^2 M_H^4}{512m_- m_2^3 M_W^2} - \frac{243C_A Y_s^2 M_H^4}{512m_- m_2^3 M_W^2} - \frac{27C_F Y_s^2 M_H^4}{1024m_- m_2^3 M_W^2} - \frac{C_A C_F^2 Y_s^2 M_H^4}{256\epsilon m_- m_2^3 M_W^2} - \frac{9C_A Y_s^2 M_H^4}{256\epsilon m_- m_2^3 M_W^2} - \frac{C_F Y_s^2 M_H^4}{512\epsilon m_- m_2^3 M_W^2} \\
& + \frac{C_A C_F^2 Y_s^2 M_H^4}{512\epsilon^2 m_- m_2^3 M_W^2} + \frac{9C_A Y_s^2 M_H^4}{512\epsilon^2 m_- m_2^3 M_W^2} + \frac{C_F Y_s^2 M_H^4}{1024\epsilon^2 m_- m_2^3 M_W^2} + \frac{31C_A C_F^2 Y_s^2 M_H^4}{1024m_2^4 M_W^2} + \frac{279C_A Y_s^2 M_H^4}{1024m_2^4 M_W^2} + \frac{31C_F Y_s^2 M_H^4}{2048m_2^4 M_W^2} \\
& + \frac{C_A C_F^2 Y_s^2 M_H^4}{128\epsilon m_2^4 M_W^2} + \frac{9C_A Y_s^2 M_H^4}{128\epsilon m_2^4 M_W^2} + \frac{C_F Y_s^2 M_H^4}{256\epsilon m_2^4 M_W^2} - \frac{5C_A C_F^2 Y_s^2 M_H^4}{1024\epsilon^2 m_2^4 M_W^2} - \frac{45C_A Y_s^2 M_H^4}{1024\epsilon^2 m_2^4 M_W^2} - \frac{5C_F Y_s^2 M_H^4}{2048\epsilon^2 m_2^4 M_W^2} + \frac{C_A C_F^2 Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{256m_- m_2^3 M_W^2} \\
& + \frac{9C_A Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{256m_- m_2^3 M_W^2} + \frac{C_F Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{512m_- m_2^3 M_W^2} - \frac{5C_A C_F^2 Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{512m_2^4 M_W^2} - \frac{45C_A Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{512m_2^4 M_W^2} - \frac{5C_F Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{1024m_2^4 M_W^2} + \frac{5C_A C_F^2 Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{512m_- m_2^3 M_W^2} \\
& + \frac{45C_A Y_s^2 \zeta_2 M_H^4}{512m_- m_2^3 M_W^2} + \frac{5C_F Y_s^2 \zeta_2 M_H^4}{1024m_- m_2^3 M_W^2} - \frac{25C_A C_F^2 Y_s^2 \zeta_2 M_H^4}{1024m_2^4 M_W^2} - \frac{225C_A Y_s^2 \zeta_2 M_H^4}{1024m_2^4 M_W^2} - \frac{25C_F Y_s^2 \zeta_2 M_H^4}{2048m_2^4 M_W^2} + \frac{C_A C_F^2 Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{128m_- m_2^3 M_W^2} \\
& + \frac{9C_A Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{128m_- m_2^3 M_W^2} + \frac{C_F Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{256\epsilon m_- m_2^3 M_W^2} - \frac{C_A C_F^2 Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{256\epsilon m_- m_2^3 M_W^2} - \frac{9C_A Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{256\epsilon m_- m_2^3 M_W^2} - \frac{C_F Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{512\epsilon m_- m_2^3 M_W^2} - \frac{C_A C_F^2 Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{64m_2^4 M_W^2} \\
& - \frac{9C_A Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{64m_2^4 M_W^2} - \frac{C_F Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{128m_2^4 M_W^2} + \frac{5C_A C_F^2 Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{512\epsilon m_2^4 M_W^2} + \frac{45C_A Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{512\epsilon m_2^4 M_W^2} + \frac{5C_F Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{1024\epsilon m_2^4 M_W^2} + \frac{2193C_A Y_s^3 M_H^2}{1024m_2^4 M_W} \\
& + \frac{141C_A Y_s^3 M_H^2}{128\epsilon m_2^4 M_W} + \frac{123C_A Y_s^3 M_H^2}{256\epsilon^2 m_2^4 M_W} - \frac{9C_A Y_s^3 M_H^2}{128\epsilon^3 m_2^4 M_W} - \frac{2943C_A m_1 Y_s^3 M_H^2}{2048m_2^5 M_W} - \frac{195C_A m_1 Y_s^3 M_H^2}{256\epsilon m_2^5 M_W} - \frac{21C_A m_1 Y_s^3 M_H^2}{64\epsilon^2 m_2^5 M_W} \\
& + \frac{15C_A m_1 Y_s^3 M_H^2}{256\epsilon^3 m_2^5 M_W} + \frac{15C_A Y_s^3 M_H^2}{128m_- m_2^3 M_W} + \frac{3C_A Y_s^3 M_H^2}{16\epsilon m_- m_2^3 M_W} + \frac{9C_A Y_s^3 M_H^2}{128\epsilon^2 m_- m_2^3 M_W} + \frac{3C_A Y_s^3 M_H^2}{128\epsilon^3 m_- m_2^3 M_W} - \frac{165C_A Y_s^3 M_H^2}{256m_2^6 M_W} \\
& - \frac{33C_A Y_s^3 M_H^2}{32\epsilon m_2^6 M_W} - \frac{99C_A Y_s^3 M_H^2}{256\epsilon^2 m_2^6 M_W} - \frac{33C_A Y_s^3 M_H^2}{256\epsilon^3 m_2^6 M_W} + \frac{21C_A Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{16m_2^4 M_W} - \frac{9C_A Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{64\epsilon m_2^4 M_W} - \frac{243C_A m_1 Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{256m_2^5 M_W} \\
& + \frac{15C_A m_1 Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{128\epsilon m_2^5 M_W} + \frac{3C_A Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{64\epsilon m_- m_2^3 M_W} - \frac{33C_A Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{128\epsilon m_2^5 M_W} + \frac{105C_A Y_s^3 \zeta_2 M_H^2}{32m_2^4 M_W} - \frac{45C_A Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{128\epsilon m_2^4 M_W} \\
& - \frac{1215C_A m_1 Y_s^3 \zeta_2 M_H^2}{512m_2^5 M_W} + \frac{75C_A m_1 Y_s^3 \zeta_2 M_H^2}{256\epsilon m_2^5 M_W} + \frac{3C_A Y_s^3 \zeta_2 M_H^2}{128\epsilon m_- m_2^3 M_W} - \frac{33C_A Y_s^3 \zeta_2 M_H^2}{256\epsilon m_2^6 M_W} - \frac{663C_A Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{256m_2^4 M_W} \\
& - \frac{123C_A Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{128\epsilon m_2^4 M_W} + \frac{9C_A Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{64\epsilon^2 m_2^4 M_W} + \frac{945C_A m_1 Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{512m_2^5 M_W} + \frac{21C_A m_1 Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{32\epsilon m_2^5 M_W} - \frac{15C_A m_1 Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{128\epsilon^2 m_2^5 M_W}
\end{aligned}$$

$$\begin{aligned}
& - \frac{3C_A Y_s^3 \mathcal{L} m_2 M_H^2}{64 m_- m_2^3 M_W} - \frac{9C_A Y_s^3 \mathcal{L} m_2 M_H^2}{64 \epsilon m_- m_2^5 M_W} - \frac{3C_A Y_s^3 \mathcal{L} m_2 M_H^2}{64 \epsilon^2 m_- m_2^5 M_W} + \frac{33C_A Y_s^3 \mathcal{L} m_2 M_H^2}{128 m_2^3 M_W} + \frac{99C_A Y_s^3 \mathcal{L} m_2 M_H^2}{128 \epsilon m_2^3 M_W} + \frac{33C_A Y_s^3 \mathcal{L} m_2 M_H^2}{128 \epsilon^2 m_2^3 M_W} \\
& + \frac{339C_A Y_s^4}{256 m_2^4} + \frac{99C_A Y_s^4}{256 \epsilon m_2^4} - \frac{3C_A Y_s^4}{256 \epsilon^2 m_2^4} - \frac{267C_A m_1 Y_s^4}{256 m_2^5} - \frac{75C_A m_1 Y_s^4}{256 \epsilon m_2^5} - \frac{3C_A m_1 Y_s^4}{256 \epsilon^2 m_2^5} - \frac{1449C_A Y_s^4}{1024 m_2^6} - \frac{289C_A Y_s^4}{1024 \epsilon m_2^6} - \frac{921C_A Y_s^4}{4096 \epsilon^2 m_2^6} \\
& - \frac{27C_A Y_s^4}{1024 \epsilon^3 m_2^6} + \frac{1315C_A m_1 Y_s^4}{1024 m_2^7} + \frac{339C_A m_1 Y_s^4}{1024 \epsilon m_2^7} + \frac{915C_A m_1 Y_s^4}{4096 \epsilon^2 m_2^7} + \frac{9C_A m_1 Y_s^4}{1024 \epsilon^3 m_2^7} - \frac{119}{2} C_A C_F^2 - \frac{13C_A C_F^2 M_W^2 Y_s^2}{4 m_- m_2^3} \\
& - \frac{13C_F M_W^2 Y_s^2}{8 m_- m_2^3} - \frac{5C_A C_F^2 M_W^2 Y_s^2}{16 \epsilon m_- m_2^3} - \frac{5C_F M_W^2 Y_s^2}{32 \epsilon m_- m_2^3} + \frac{C_A C_F^2 M_W^2 Y_s^2}{8 \epsilon^2 m_- m_2^3} + \frac{C_F M_W^2 Y_s^2}{16 \epsilon^2 m_- m_2^3} + \frac{27C_A C_F^2 M_W^2 Y_s^2}{16 m_2^4} + \frac{27C_F M_W^2 Y_s^2}{32 m_2^4} \\
& + \frac{21C_A C_F^2 M_W^2 Y_s^2}{32 \epsilon m_2^4} + \frac{21C_F M_W^2 Y_s^2}{64 \epsilon m_2^4} - \frac{5C_A C_F^2 M_W^2 Y_s^2}{16 \epsilon^2 m_2^4} - \frac{5C_F M_W^2 Y_s^2}{32 \epsilon^2 m_2^4} + \frac{3C_A C_F^2 Y_s^2}{4 m_2^2} + \frac{231C_A C_F Y_s^2}{64 m_2^2} + \frac{3C_F Y_s^2}{8 m_2^2} \\
& + \frac{3C_A C_F^2 Y_s^2}{8 \epsilon m_2^2} - \frac{5C_A C_F Y_s^2}{16 \epsilon m_2^2} + \frac{3C_F Y_s^2}{16 \epsilon m_2^2} + \frac{3C_A C_F^2 Y_s^2}{16 \epsilon^2 m_2^2} - \frac{11C_A C_F Y_s^2}{16 \epsilon^2 m_2^2} + \frac{3C_F Y_s^2}{32 \epsilon^2 m_2^2} - \frac{3C_A C_F Y_s^2}{8 \epsilon^3 m_2^2} - \frac{C_A C_F^2 m_1 Y_s^2}{4 m_2^3} \\
& - \frac{19C_A C_F m_1 Y_s^2}{16 m_2^3} - \frac{C_F m_1 Y_s^2}{8 m_2^3} - \frac{C_A C_F^2 m_1 Y_s^2}{8 \epsilon m_2^3} + \frac{13C_A C_F m_1 Y_s^2}{16 \epsilon m_2^3} - \frac{C_F m_1 Y_s^2}{16 \epsilon m_2^3} - \frac{C_A C_F^2 m_1 Y_s^2}{16 \epsilon^2 m_2^3} + \frac{23C_A C_F m_1 Y_s^2}{16 \epsilon^2 m_2^3} \\
& - \frac{C_F m_1 Y_s^2}{32 \epsilon^2 m_2^3} - \frac{565C_A C_F Y_s^2}{64 m_2^4} - \frac{287C_A C_F Y_s^2}{64 \epsilon m_2^4} - \frac{709C_A C_F Y_s^2}{128 \epsilon^2 m_2^4} + \frac{41C_A C_F Y_s^2}{64 \epsilon^3 m_2^4} - \frac{3C_A C_F Y_s^2}{16 \epsilon^4 m_2^4} - \frac{3C_A Y_s^4 \mathcal{L} m_2}{128 m_2^4} \\
& - \frac{3C_A m_1 Y_s^4 \mathcal{L} m_2}{128 m_2^5} - \frac{705C_A Y_s^4 \mathcal{L} m_2}{2048 m_2^6} - \frac{27C_A Y_s^4 \mathcal{L} m_2}{512 \epsilon m_2^6} + \frac{843C_A m_1 Y_s^4 \mathcal{L} m_2}{2048 m_2^7} + \frac{9C_A m_1 Y_s^4 \mathcal{L} m_2}{512 \epsilon m_2^7} - \frac{11}{4} C_A C_F^2 \mathcal{L} m_2 \\
& + \frac{C_A C_F^2 M_W^2 Y_s^2 \mathcal{L} m_2}{4 m_- m_2^3} + \frac{C_F M_W^2 Y_s^2 \mathcal{L} m_2}{8 m_- m_2^3} - \frac{5C_A C_F^2 M_W^2 Y_s^2 \mathcal{L} m_2}{8 m_2^4} - \frac{5C_F M_W^2 Y_s^2 \mathcal{L} m_2}{16 m_2^4} + \frac{3C_A C_F^2 Y_s^2 \mathcal{L} m_2}{8 m_2^2} + \frac{C_A C_F Y_s^2 \mathcal{L} m_2}{2 m_2^2} \\
& + \frac{3C_F Y_s^2 \mathcal{L} m_2}{16 m_2^2} - \frac{3C_A C_F Y_s^2 \mathcal{L} m_2}{4 \epsilon m_2^2} - \frac{C_A C_F^2 m_1 Y_s^2 \mathcal{L} m_2}{8 m_2^3} + \frac{23C_A C_F m_1 Y_s^2 \mathcal{L} m_2}{8 m_2^3} - \frac{C_F m_1 Y_s^2 \mathcal{L} m_2}{16 m_2^3} - \frac{963C_A C_F Y_s^2 \mathcal{L} m_2}{64 m_2^4} \\
& + \frac{65C_A C_F Y_s^2 \mathcal{L} m_2}{32 \epsilon m_2^4} - \frac{3C_A C_F Y_s^2 \mathcal{L} m_2}{8 \epsilon^2 m_2^4} - \frac{15}{4} C_A C_F \mathcal{L} m_2 + \frac{69}{8} C_F \mathcal{L} m_2 + C_A C_F n_f T_f \mathcal{L} m_2 - \frac{191C_A C_F M_W Y_s \mathcal{L} m_2}{16 m_2^2} \\
& + \frac{11C_A C_F M_W Y_s \mathcal{L} m_2}{8 \epsilon m_2^2} - \frac{2C_A C_F^2 \mathcal{L} m_2}{\epsilon} - \frac{7C_F \mathcal{L} m_2}{\epsilon} + \frac{43C_A C_F}{8} - \frac{447C_F}{8} + 7C_A C_F n_f T_f + \frac{2C_A C_F n_f T_f}{\epsilon} + \frac{C_A C_F n_f T_f}{2 \epsilon^2} \\
& - \frac{2931C_A C_F M_W Y_s}{128 m_2^5} - \frac{85C_A C_F M_W Y_s}{8 \epsilon m_2^5} - \frac{17C_A C_F M_W Y_s}{4 \epsilon^2 m_2^5} + \frac{11C_A C_F M_W Y_s}{16 \epsilon^3 m_2^5} - \frac{15C_A Y_s^4 \zeta_2}{256 m_2^4} - \frac{15C_A m_1 Y_s^4 \zeta_2}{256 m_2^5} \\
& + \frac{1239C_A Y_s^4 \zeta_2}{4096 m_2^6} - \frac{135C_A Y_s^4 \zeta_2}{1024 \epsilon m_2^6} + \frac{195C_A m_1 Y_s^4 \zeta_2}{4096 m_2^7} + \frac{45C_A m_1 Y_s^4 \zeta_2}{1024 \epsilon m_2^7} - \frac{55}{8} C_A C_F^2 \zeta_2 + \frac{5C_A C_F^2 M_W^2 Y_s^2 \zeta_2}{8 m_- m_2^3} + \frac{5C_F M_W^2 Y_s^2 \zeta_2}{16 m_- m_2^3} \\
& - \frac{25C_A C_F^2 M_W^2 Y_s^2 \zeta_2}{16 m_2^4} - \frac{25C_F M_W^2 Y_s^2 \zeta_2}{32 m_2^4} + \frac{15C_A C_F^2 Y_s^2 \zeta_2}{16 m_2^2} + \frac{5C_A C_F Y_s^2 \zeta_2}{4 m_2^2} + \frac{15C_F Y_s^2 \zeta_2}{32 m_2^2} - \frac{15C_A C_F Y_s^2 \zeta_2}{8 \epsilon m_2^2} \\
& - \frac{5C_A C_F^2 m_1 Y_s^2 \zeta_2}{16 m_2^3} + \frac{115C_A C_F m_1 Y_s^2 \zeta_2}{16 m_2^3} - \frac{5C_F m_1 Y_s^2 \zeta_2}{32 m_2^3} - \frac{4379C_A C_F Y_s^2 \zeta_2}{128 m_2^4} + \frac{289C_A C_F Y_s^2 \zeta_2}{64 \epsilon m_2^4} - \frac{15C_A C_F Y_s^2 \zeta_2}{16 \epsilon^2 m_2^4} \\
& - \frac{75}{8} C_A C_F \zeta_2 + \frac{345C_F \zeta_2}{16} + \frac{5}{2} C_A C_F n_f T_f \zeta_2 - \frac{955C_A C_F M_W Y_s \zeta_2}{32 m_2^2} + \frac{55C_A C_F M_W Y_s \zeta_2}{16 \epsilon m_2^2} - \frac{5C_A C_F^2 \zeta_2}{\epsilon} - \frac{35C_F \zeta_2}{2 \epsilon} \\
& - \frac{99C_A Y_s^4 \mathcal{L} m_2}{128 m_2^4} + \frac{3C_A Y_s^4 \mathcal{L} m_2}{128 \epsilon m_2^4} + \frac{75C_A m_1 Y_s^4 \mathcal{L} m_2}{128 m_2^5} + \frac{3C_A m_1 Y_s^4 \mathcal{L} m_2}{128 \epsilon m_2^5} + \frac{131C_A Y_s^4 \mathcal{L} m_2}{256 m_2^6} + \frac{921C_A Y_s^4 \mathcal{L} m_2}{2048 \epsilon m_2^6} + \frac{27C_A Y_s^4 \mathcal{L} m_2}{512 \epsilon^2 m_2^6} \\
& - \frac{165C_A m_1 Y_s^4 \mathcal{L} m_2}{256 m_2^7} - \frac{915C_A m_1 Y_s^4 \mathcal{L} m_2}{2048 \epsilon m_2^7} - \frac{9C_A m_1 Y_s^4 \mathcal{L} m_2}{512 \epsilon^2 m_2^7} + \frac{93}{4} C_A C_F^2 \mathcal{L} m_2 + \frac{5C_A C_F^2 M_W^2 Y_s^2 \mathcal{L} m_2}{8 m_- m_2^3} + \frac{5C_F M_W^2 Y_s^2 \mathcal{L} m_2}{16 m_- m_2^3} \\
& - \frac{C_A C_F^2 M_W^2 Y_s^2 \mathcal{L} m_2}{4 \epsilon m_- m_2^3} - \frac{C_F M_W^2 Y_s^2 \mathcal{L} m_2}{8 \epsilon m_- m_2^3} - \frac{21C_A C_F^2 M_W^2 Y_s^2 \mathcal{L} m_2}{16 m_2^4} - \frac{21C_F M_W^2 Y_s^2 \mathcal{L} m_2}{32 m_2^4} + \frac{5C_A C_F^2 M_W^2 Y_s^2 \mathcal{L} m_2}{8 \epsilon m_2^4} \\
& + \frac{5C_F M_W^2 Y_s^2 \mathcal{L} m_2}{16 \epsilon m_2^4} - \frac{3C_A C_F^2 Y_s^2 \mathcal{L} m_2}{4 m_2^2} - \frac{23C_A C_F Y_s^2 \mathcal{L} m_2}{16 m_2^2} - \frac{3C_F Y_s^2 \mathcal{L} m_2}{8 m_2^2} - \frac{3C_A C_F^2 Y_s^2 \mathcal{L} m_2}{8 \epsilon m_2^2} + \frac{11C_A C_F Y_s^2 \mathcal{L} m_2}{8 \epsilon m_2^2} \\
& - \frac{3C_F Y_s^2 \mathcal{L} m_2}{16 \epsilon m_2^2} + \frac{3C_A C_F Y_s^2 \mathcal{L} m_2}{4 \epsilon^2 m_2^2} + \frac{C_A C_F^2 m_1 Y_s^2 \mathcal{L} m_2}{4 m_2^3} - \frac{13C_A C_F m_1 Y_s^2 \mathcal{L} m_2}{8 m_2^3} + \frac{C_F m_1 Y_s^2 \mathcal{L} m_2}{8 m_2^3} + \frac{C_A C_F^2 m_1 Y_s^2 \mathcal{L} m_2}{8 \epsilon m_2^3} \\
& - \frac{23C_A C_F m_1 Y_s^2 \mathcal{L} m_2}{8 \epsilon m_2^3} + \frac{C_F m_1 Y_s^2 \mathcal{L} m_2}{16 \epsilon m_2^3} + \frac{203C_A C_F Y_s^2 \mathcal{L} m_2}{16 m_2^4} + \frac{685C_A C_F Y_s^2 \mathcal{L} m_2}{64 \epsilon m_2^4} - \frac{41C_A C_F Y_s^2 \mathcal{L} m_2}{32 \epsilon^2 m_2^4} + \frac{3C_A C_F Y_s^2 \mathcal{L} m_2}{8 \epsilon^3 m_2^4} \\
& + \frac{3}{4} C_A C_F \mathcal{L} m_2 + \frac{129}{8} C_F \mathcal{L} m_2 - 4C_A C_F n_f T_f \mathcal{L} m_2 - \frac{C_A C_F n_f T_f \mathcal{L} m_2}{\epsilon} + \frac{801C_A C_F M_W Y_s \mathcal{L} m_2}{32 m_2^2} + \frac{17C_A C_F M_W Y_s \mathcal{L} m_2}{2 \epsilon m_2^2} \\
& - \frac{11C_A C_F M_W Y_s \mathcal{L} m_2}{8 \epsilon^2 m_2^2} + \frac{31C_A C_F^2 \mathcal{L} m_2}{4 \epsilon} + \frac{15C_A C_F \mathcal{L} m_2}{4 \epsilon} + \frac{71C_F \mathcal{L} m_2}{8 \epsilon} + \frac{2C_A C_F^2 \mathcal{L} m_2}{\epsilon^2} + \frac{7C_F \mathcal{L} m_2}{\epsilon^2} - \frac{115C_A C_F^2}{8 \epsilon} \\
& - \frac{3C_A C_F}{8 \epsilon} - \frac{283C_F}{16 \epsilon} - \frac{31C_A C_F^2}{8 \epsilon^2} - \frac{15C_A C_F}{8 \epsilon^2} - \frac{71C_F}{16 \epsilon^2} - \frac{C_A C_F^2}{\epsilon^3} - \frac{7C_F}{2 \epsilon^3}
\end{aligned} \tag{C.10}$$

$$\begin{aligned}
V_5(m_{1,2}) = & - \frac{27C_A C_F^2 Y_s^2 M_H^4}{512 m_- m_2^3 M_W^2} - \frac{243C_A Y_s^2 M_H^4}{512 m_- m_2^3 M_W^2} - \frac{27C_F Y_s^2 M_H^4}{1024 m_- m_2^3 M_W^2} - \frac{C_A C_F^2 Y_s^2 M_H^4}{256 \epsilon m_- m_2^3 M_W^2} - \frac{9C_A Y_s^2 M_H^4}{256 \epsilon m_- m_2^3 M_W^2} - \frac{C_F Y_s^2 M_H^4}{512 \epsilon m_- m_2^3 M_W^2} \\
& + \frac{C_A C_F^2 Y_s^2 M_H^4}{512 \epsilon^2 m_- m_2^3 M_W^2} + \frac{9C_A Y_s^2 M_H^4}{512 \epsilon^2 m_- m_2^3 M_W^2} + \frac{C_F Y_s^2 M_H^4}{1024 \epsilon^2 m_- m_2^3 M_W^2} + \frac{17C_A C_F^2 Y_s^2 M_H^4}{256 m_2^4 M_W^2} + \frac{153C_A Y_s^2 M_H^4}{256 m_2^4 M_W^2} + \frac{17C_F Y_s^2 M_H^4}{512 m_2^4 M_W^2}
\end{aligned}$$

$$\begin{aligned}
& + \frac{3C_A C_F^2 Y_s^2 M_H^4}{512\epsilon m_s^2 M_W^2} + \frac{27C_A Y_s^2 M_H^4}{512\epsilon m_s^2 M_W^2} + \frac{3C_F Y_s^2 M_H^4}{1024\epsilon m_s^2 M_W^2} + \frac{C_A C_F^2 Y_s^2 M_H^4}{256\epsilon^2 m_s^2 M_W^2} + \frac{9C_A Y_s^2 M_H^4}{256\epsilon^2 m_s^2 M_W^2} + \frac{C_F Y_s^2 M_H^4}{512\epsilon^2 m_s^2 M_W^2} + \frac{C_A C_F^2 Y_s^2 \zeta_2 m_2^4 M_H^4}{256m_- m_s^2 M_W^2} \\
& + \frac{9C_A Y_s^2 \zeta_2 m_2^4 M_H^4}{256m_- m_s^2 M_W^2} + \frac{C_F Y_s^2 \zeta_2 m_2^4 M_H^4}{512m_- m_s^2 M_W^2} + \frac{C_A C_F^2 Y_s^2 \zeta_2 m_2^4 M_H^4}{128m_s^2 M_W^2} + \frac{9C_A Y_s^2 \zeta_2 m_2^4 M_H^4}{128m_s^2 M_W^2} + \frac{C_F Y_s^2 \zeta_2 m_2^4 M_H^4}{256m_s^2 M_W^2} + \frac{5C_A C_F^2 Y_s^2 \zeta_2 m_2^4 M_H^4}{512m_- m_s^2 M_W^2} \\
& + \frac{45C_A Y_s^2 \zeta_2 m_2^4 M_H^4}{1024m_- m_s^2 M_W^2} + \frac{5C_F Y_s^2 \zeta_2 m_2^4 M_H^4}{256m_s^2 M_W^2} + \frac{5C_A C_F^2 Y_s^2 \zeta_2 m_2^4 M_H^4}{256\epsilon m_- m_s^2 M_W^2} + \frac{45C_A Y_s^2 \zeta_2 m_2^4 M_H^4}{256\epsilon m_- m_s^2 M_W^2} + \frac{5C_F Y_s^2 \zeta_2 m_2^4 M_H^4}{512\epsilon m_- m_s^2 M_W^2} + \frac{C_A C_F^2 Y_s^2 \zeta_2 m_2^4 M_H^4}{256\epsilon m_- m_s^2 M_W^2} \\
& + \frac{9C_A Y_s^2 \zeta_2 m_2^4 M_H^4}{128m_- m_s^2 M_W^2} + \frac{C_F Y_s^2 \zeta_2 m_2^4 M_H^4}{256m_- m_s^2 M_W^2} - \frac{C_A C_F^2 Y_s^2 \zeta_2 m_2^4 M_H^4}{256\epsilon m_- m_s^2 M_W^2} - \frac{9C_A Y_s^2 \zeta_2 m_2^4 M_H^4}{256\epsilon m_- m_s^2 M_W^2} - \frac{C_F Y_s^2 \zeta_2 m_2^4 M_H^4}{512\epsilon m_- m_s^2 M_W^2} - \frac{3C_A C_F^2 Y_s^2 \zeta_2 m_2^4 M_H^4}{256m_- m_s^2 M_W^2} \\
& - \frac{27C_A Y_s^2 \zeta_2 m_2^4 M_H^4}{256m_s^2 M_W^2} - \frac{3C_F Y_s^2 \zeta_2 m_2^4 M_H^4}{512m_s^2 M_W^2} - \frac{C_A C_F^2 Y_s^2 \zeta_2 m_2^4 M_H^4}{128\epsilon m_s^2 M_W^2} - \frac{9C_A Y_s^2 \zeta_2 m_2^4 M_H^4}{128\epsilon m_s^2 M_W^2} - \frac{C_F Y_s^2 \zeta_2 m_2^4 M_H^4}{256\epsilon m_s^2 M_W^2} + \frac{369C_A Y_s^3 M_H^2}{4096m_s^2 M_W} \\
& + \frac{105C_A Y_s^3 M_H^2}{512\epsilon m_s^2 M_W} - \frac{21C_A Y_s^3 M_H^2}{512\epsilon^2 m_s^2 M_W} - \frac{81C_A Y_s^3 M_H^2}{512\epsilon^3 m_s^2 M_W} - \frac{1875C_A m_1 Y_s^3 M_H^2}{4096m_s^2 M_W} - \frac{123C_A m_1 Y_s^3 M_H^2}{512\epsilon m_s^2 M_W} + \frac{3C_A m_1 Y_s^3 M_H^2}{512\epsilon^2 m_s^2 M_W} \\
& + \frac{51C_A m_1 Y_s^3 M_H^2}{512\epsilon^3 m_s^2 M_W} + \frac{15C_A Y_s^3 M_H^2}{128m_- m_s^2 M_W} + \frac{3C_A Y_s^3 M_H^2}{16\epsilon m_- m_s^2 M_W} + \frac{9C_A Y_s^3 M_H^2}{128\epsilon^2 m_- m_s^2 M_W} + \frac{3C_A Y_s^3 M_H^2}{128\epsilon^3 m_- m_s^2 M_W} + \frac{51C_A Y_s^3 M_H^2}{128m_2^2 M_W} \\
& - \frac{9C_A Y_s^3 M_H^2}{32\epsilon m_2^2 M_W} - \frac{15C_A Y_s^3 M_H^2}{128\epsilon^2 m_2^2 M_W} - \frac{9C_A Y_s^3 M_H^2}{128\epsilon^3 m_2^2 M_W} + \frac{363C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{512m_s^2 M_W} - \frac{81C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{256\epsilon m_s^2 M_W} - \frac{249C_A m_1 Y_s^3 \zeta_2 m_2^2 M_H^2}{512m_s^2 M_W} \\
& + \frac{51C_A m_1 Y_s^3 \zeta_2 m_2^2 M_H^2}{256\epsilon m_s^2 M_W} + \frac{3C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{64\epsilon m_- m_s^2 M_W} + \frac{3C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{16m_2^2 M_W} - \frac{9C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{64\epsilon m_2^2 M_W} + \frac{1815C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{1024m_2^2 M_W} - \frac{405C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{512\epsilon m_2^2 M_W} \\
& - \frac{1245C_A m_1 Y_s^3 \zeta_2 m_2^2 M_H^2}{1024m_2^2 M_W} + \frac{255C_A m_1 Y_s^3 \zeta_2 m_2^2 M_H^2}{512\epsilon m_2^2 M_W} + \frac{3C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{128\epsilon m_- m_s^2 M_W} + \frac{3C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{32m_2^2 M_W} - \frac{9C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{128\epsilon m_2^2 M_W} - \frac{1311C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{1024m_2^2 M_W} \\
& + \frac{21C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{256\epsilon m_s^2 M_W} + \frac{81C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{256\epsilon^2 m_s^2 M_W} + \frac{1053C_A m_1 Y_s^3 \zeta_2 m_2^2 M_H^2}{1024m_s^2 M_W} - \frac{3C_A m_1 Y_s^3 \zeta_2 m_2^2 M_H^2}{256\epsilon m_s^2 M_W} - \frac{51C_A m_1 Y_s^3 \zeta_2 m_2^2 M_H^2}{256\epsilon^2 m_s^2 M_W} \\
& - \frac{3C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{64m_- m_s^2 M_W} - \frac{9C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{64\epsilon m_- m_s^2 M_W} - \frac{3C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{64\epsilon^2 m_- m_s^2 M_W} - \frac{27C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{64m_2^2 M_W} + \frac{15C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{64\epsilon m_2^2 M_W} + \frac{9C_A Y_s^3 \zeta_2 m_2^2 M_H^2}{64\epsilon^2 m_2^2 M_W} \\
& - \frac{111C_A Y_s^4}{256m_2^2} - \frac{3C_A Y_s^4}{512\epsilon m_2^2} + \frac{13C_A Y_s^4}{128\epsilon^2 m_2^2} + \frac{C_A Y_s^4}{32\epsilon^3 m_2^2} + \frac{151C_A m_1 Y_s^4}{512m_2^2} + \frac{5C_A m_1 Y_s^4}{512\epsilon m_2^2} - \frac{11C_A m_1 Y_s^4}{128\epsilon^2 m_2^2} - \frac{C_A m_1 Y_s^4}{64\epsilon^3 m_2^2} - \frac{177C_A Y_s^4}{128m_2^2} \\
& - \frac{175C_A Y_s^4}{512\epsilon m_2^2} + \frac{127C_A Y_s^4}{1024\epsilon^2 m_2^2} - \frac{29C_A Y_s^4}{1024\epsilon^3 m_2^2} + \frac{5C_A Y_s^4}{256\epsilon^4 m_2^2} + \frac{297C_A m_1 Y_s^4}{256m_2^2} + \frac{153C_A m_1 Y_s^4}{512\epsilon m_2^2} - \frac{171C_A m_1 Y_s^4}{2048\epsilon^2 m_2^2} - \frac{11C_A m_1 Y_s^4}{1024\epsilon^3 m_2^2} \\
& - \frac{3C_A m_1 Y_s^4}{256\epsilon^4 m_2^2} + \frac{2361}{32} C_A C_F^2 - \frac{13C_A C_F^2 M_W^2 Y_s^2}{4m_- m_s^2} - \frac{13C_F M_W^2 Y_s^2}{8m_- m_s^2} - \frac{5C_A C_F^2 M_W^2 Y_s^2}{16\epsilon m_- m_s^2} - \frac{5C_F M_W^2 Y_s^2}{32\epsilon m_- m_s^2} + \frac{C_A C_F^2 M_W^2 Y_s^2}{8\epsilon^2 m_- m_s^2} \\
& + \frac{C_F M_W^2 Y_s^2}{16\epsilon^2 m_- m_s^2} + \frac{65C_A C_F^2 M_W^2 Y_s^2}{16m_s^2} + \frac{65C_F M_W^2 Y_s^2}{32m_s^2} + \frac{C_A C_F^2 M_W^2 Y_s^2}{4\epsilon m_s^2} + \frac{C_F M_W^2 Y_s^2}{8\epsilon m_s^2} + \frac{C_A C_F^2 M_W^2 Y_s^2}{4\epsilon^2 m_s^2} + \frac{C_F M_W^2 Y_s^2}{8\epsilon^2 m_s^2} \\
& + \frac{89C_A C_F^2 Y_s^2}{16m_2^2} + \frac{247C_A C_F Y_s^2}{32m_2^2} + \frac{89C_F Y_s^2}{32m_2^2} + \frac{19C_A C_F^2 Y_s^2}{16\epsilon m_2^2} - \frac{313C_A C_F Y_s^2}{128\epsilon m_2^2} + \frac{19C_F Y_s^2}{32\epsilon m_2^2} + \frac{C_A C_F^2 Y_s^2}{16\epsilon^2 m_2^2} - \frac{133C_A C_F Y_s^2}{64\epsilon^2 m_2^2} \\
& + \frac{C_F Y_s^2}{32\epsilon^2 m_2^2} - \frac{3C_A C_F Y_s^2}{4\epsilon^3 m_2^2} - \frac{41C_A C_F^2 m_1 Y_s^2}{16m_2^2} - \frac{77C_A C_F m_1 Y_s^2}{32m_2^2} - \frac{41C_F m_1 Y_s^2}{32m_2^2} - \frac{7C_A C_F^2 m_1 Y_s^2}{16\epsilon m_2^2} + \frac{267C_A C_F m_1 Y_s^2}{128\epsilon m_2^2} \\
& - \frac{7C_F m_1 Y_s^2}{32\epsilon m_2^2} + \frac{C_A C_F^2 m_1 Y_s^2}{16\epsilon^2 m_2^2} + \frac{111C_A C_F m_1 Y_s^2}{64\epsilon^2 m_2^2} + \frac{C_F m_1 Y_s^2}{32\epsilon^2 m_2^2} + \frac{3C_A C_F m_1 Y_s^2}{8\epsilon^3 m_2^2} - \frac{65C_A C_F Y_s^2}{32m_2^2} - \frac{135C_A C_F Y_s^2}{32\epsilon m_2^2} \\
& - \frac{261C_A C_F Y_s^2}{32\epsilon^2 m_2^2} - \frac{C_A C_F Y_s^2}{64\epsilon^3 m_2^2} - \frac{3C_A C_F Y_s^2}{16\epsilon^4 m_2^2} + \frac{3C_A Y_s^4 \zeta_2 m_2}{64m_2^2} + \frac{C_A Y_s^4 \zeta_2 m_2}{16\epsilon m_2^2} - \frac{3C_A m_1 Y_s^4 \zeta_2 m_2}{32m_2^2} - \frac{C_A m_1 Y_s^4 \zeta_2 m_2}{32\epsilon m_2^2} \\
& + \frac{57C_A Y_s^4 \zeta_2 m_2}{128m_2^2} - \frac{69C_A Y_s^4 \zeta_2 m_2}{512\epsilon m_2^2} + \frac{5C_A Y_s^4 \zeta_2 m_2}{128\epsilon^2 m_2^2} - \frac{153C_A m_1 Y_s^4 \zeta_2 m_2}{1024m_2^2} + \frac{13C_A m_1 Y_s^4 \zeta_2 m_2}{512\epsilon m_2^2} \\
& - \frac{3C_A m_1 Y_s^4 \zeta_2 m_2}{128\epsilon^2 m_2^2} + 16C_A C_F^2 + \frac{C_A C_F^2 M_W^2 Y_s^2 \zeta_2 m_2}{4m_- m_s^2} + \frac{C_F M_W^2 Y_s^2 \zeta_2 m_2}{8m_- m_s^2} + \frac{C_A C_F^2 M_W^2 Y_s^2 \zeta_2 m_2}{2m_2^2} + \frac{C_F M_W^2 Y_s^2 \zeta_2 m_2}{4m_2^2} \\
& + \frac{C_A C_F^2 Y_s^2 \zeta_2 m_2}{8m_2^2} - \frac{13C_A C_F Y_s^2 \zeta_2 m_2}{32m_2^2} + \frac{C_F Y_s^2 \zeta_2 m_2}{16m_2^2} - \frac{3C_A C_F Y_s^2 \zeta_2 m_2}{2\epsilon m_2^2} + \frac{C_A C_F^2 m_1 Y_s^2 \zeta_2 m_2}{8m_2^2} + \frac{51C_A C_F m_1 Y_s^2 \zeta_2 m_2}{32m_2^2} \\
& + \frac{C_F m_1 Y_s^2 \zeta_2 m_2}{16m_2^2} + \frac{3C_A C_F m_1 Y_s^2 \zeta_2 m_2}{4\epsilon m_2^2} - \frac{565C_A C_F Y_s^2 \zeta_2 m_2}{32m_2^2} + \frac{23C_A C_F Y_s^2 \zeta_2 m_2}{32\epsilon m_2^2} - \frac{3C_A C_F Y_s^2 \zeta_2 m_2}{8\epsilon^2 m_2^2} + \frac{13}{2} C_F \zeta_2 m_2 \\
& - \frac{87C_A C_F M_W Y_s \zeta_2 m_2}{16m_2^2} + \frac{3C_A C_F M_W Y_s \zeta_2 m_2}{8\epsilon m_2^2} - \frac{2C_A C_F^2 \zeta_2 m_2}{\epsilon} - \frac{3C_F \zeta_2 m_2}{\epsilon} - \frac{51C_A C_F}{32} - \frac{887C_F}{64} + \frac{11}{8} C_A C_F n_f T_f \\
& + \frac{C_A C_F n_f T_f}{4\epsilon} - \frac{1515C_A C_F M_W Y_s}{128m_2^2} - \frac{41C_A C_F M_W Y_s}{8\epsilon m_2^2} - \frac{9C_A C_F M_W Y_s}{4\epsilon^2 m_2^2} + \frac{3C_A C_F M_W Y_s}{16\epsilon^3 m_2^2} + \frac{15C_A Y_s^4 \zeta_2}{128m_2^2} + \frac{5C_A Y_s^4 \zeta_2}{32\epsilon m_2^2} \\
& - \frac{15C_A m_1 Y_s^4 \zeta_2}{64m_2^2} - \frac{5C_A m_1 Y_s^4 \zeta_2}{64\epsilon m_2^2} + \frac{409C_A Y_s^4 \zeta_2}{256m_2^2} - \frac{413C_A Y_s^4 \zeta_2}{1024\epsilon m_2^2} + \frac{25C_A Y_s^4 \zeta_2}{256\epsilon^2 m_2^2} - \frac{1345C_A m_1 Y_s^4 \zeta_2}{2048m_2^2} + \frac{157C_A m_1 Y_s^4 \zeta_2}{1024\epsilon m_2^2} \\
& - \frac{15C_A m_1 Y_s^4 \zeta_2}{256\epsilon^2 m_2^2} + 40C_A C_F^2 \zeta_2 + \frac{5C_A C_F^2 M_W^2 Y_s^2 \zeta_2}{8m_- m_s^2} + \frac{5C_F M_W^2 Y_s^2 \zeta_2}{16m_- m_s^2} + \frac{5C_A C_F^2 M_W^2 Y_s^2 \zeta_2}{4m_2^2} + \frac{5C_F M_W^2 Y_s^2 \zeta_2}{8m_2^2} \\
& + \frac{5C_A C_F^2 Y_s^2 \zeta_2}{16m_2^2} - \frac{65C_A C_F Y_s^2 \zeta_2}{64m_2^2} + \frac{5C_F Y_s^2 \zeta_2}{32m_2^2} - \frac{15C_A C_F Y_s^2 \zeta_2}{4\epsilon m_2^2} + \frac{5C_A C_F^2 m_1 Y_s^2 \zeta_2}{16m_2^2} + \frac{255C_A C_F m_1 Y_s^2 \zeta_2}{64m_2^2} + \frac{5C_F m_1 Y_s^2 \zeta_2}{32m_2^2}
\end{aligned}$$

$$\begin{aligned}
& + \frac{15C_A C_F m_1 Y_s^2 \zeta_2}{8\epsilon m_2^3} - \frac{2397C_A C_F Y_s^2 \zeta_2}{64m_2^4} + \frac{79C_A C_F Y_s^2 \zeta_2}{64\epsilon m_2^4} - \frac{15C_A C_F Y_s^2 \zeta_2}{16\epsilon^2 m_2^4} + \frac{65C_F \zeta_2}{4} - \frac{435C_A C_F M_W Y_s \zeta_2}{32m_2^2} \\
& + \frac{15C_A C_F M_W Y_s \zeta_2}{16\epsilon m_2^2} - \frac{5C_A C_F^2 \zeta_2}{\epsilon} - \frac{15C_F \zeta_2}{2\epsilon} + \frac{47C_A Y_s^4 \zeta_2}{256m_2^4} - \frac{13C_A Y_s^4 \zeta_2}{64\epsilon m_2^4} - \frac{C_A Y_s^4 \zeta_2}{16\epsilon^2 m_2^4} - \frac{27C_A m_1 Y_s^4 \zeta_2}{256m_2^5} \\
& + \frac{11C_A m_1 Y_s^4 \zeta_2}{64\epsilon m_2^5} + \frac{C_A m_1 Y_s^4 \zeta_2}{32\epsilon^2 m_2^5} + \frac{383C_A Y_s^4 \zeta_2}{512m_2^6} - \frac{107C_A Y_s^4 \zeta_2}{512\epsilon m_2^6} + \frac{29C_A Y_s^4 \zeta_2}{512\epsilon^2 m_2^6} - \frac{5C_A Y_s^4 \zeta_2}{128\epsilon^3 m_2^6} \\
& - \frac{431C_A m_1 Y_s^4 \zeta_2}{512m_2^7} + \frac{147C_A m_1 Y_s^4 \zeta_2}{1024\epsilon m_2^7} + \frac{11C_A m_1 Y_s^4 \zeta_2}{512\epsilon^2 m_2^7} + \frac{3C_A m_1 Y_s^4 \zeta_2}{128\epsilon^3 m_2^7} - \frac{377}{8} C_A C_F^2 \zeta_2 + \frac{5C_A C_F^2 M_W^2 Y_s^2 \zeta_2}{8m_2^3} \\
& + \frac{5C_F M_W^2 Y_s^2 \zeta_2}{16m_2^3} - \frac{C_A C_F^2 M_W^2 Y_s^2 \zeta_2}{4\epsilon m_2^3} - \frac{C_F M_W^2 Y_s^2 \zeta_2}{8\epsilon m_2^3} - \frac{C_A C_F^2 M_W^2 Y_s^2 \zeta_2}{2m_2^4} - \frac{C_F M_W^2 Y_s^2 \zeta_2}{4m_2^4} \\
& - \frac{C_A C_F^2 M_W^2 Y_s^2 \zeta_2}{2\epsilon m_2^4} - \frac{C_F M_W^2 Y_s^2 \zeta_2}{4\epsilon m_2^4} - \frac{19C_A C_F^2 Y_s^2 \zeta_2}{8m_2^2} + \frac{49C_A C_F Y_s^2 \zeta_2}{64m_2^2} - \frac{19C_F Y_s^2 \zeta_2}{16m_2^2} - \frac{C_A C_F Y_s^2 \zeta_2}{8\epsilon m_2^2} \\
& + \frac{133C_A C_F Y_s^2 \zeta_2}{32\epsilon m_2^2} - \frac{C_F Y_s^2 \zeta_2}{16\epsilon m_2^2} + \frac{3C_A C_F Y_s^2 \zeta_2}{2\epsilon^2 m_2^2} + \frac{7C_A C_F^2 m_1 Y_s^2 \zeta_2}{8m_2^3} - \frac{135C_A C_F m_1 Y_s^2 \zeta_2}{64m_2^3} + \frac{7C_F m_1 Y_s^2 \zeta_2}{16m_2^3} \\
& - \frac{C_A C_F^2 m_1 Y_s^2 \zeta_2}{8\epsilon m_2^3} - \frac{111C_A C_F m_1 Y_s^2 \zeta_2}{32\epsilon m_2^3} - \frac{C_F m_1 Y_s^2 \zeta_2}{16\epsilon m_2^3} - \frac{3C_A C_F m_1 Y_s^2 \zeta_2}{4\epsilon^2 m_2^3} + \frac{347C_A C_F Y_s^2 \zeta_2}{32m_2^4} \\
& + \frac{255C_A C_F Y_s^2 \zeta_2}{16\epsilon m_2^4} + \frac{C_A C_F Y_s^2 \zeta_2}{32\epsilon^2 m_2^4} + \frac{3C_A C_F Y_s^2 \zeta_2}{8\epsilon^3 m_2^4} + \frac{15}{8} C_A C_F \zeta_2 + \frac{107}{16} C_F \zeta_2 - \frac{1}{2} C_A C_F n_f T_f \zeta_2 \\
& + \frac{361C_A C_F M_W Y_s \zeta_2}{32m_2^2} + \frac{9C_A C_F M_W Y_s \zeta_2}{2\epsilon m_2^2} - \frac{3C_A C_F M_W Y_s \zeta_2}{8\epsilon^2 m_2^2} - \frac{11C_A C_F^2 \zeta_2}{\epsilon} + \frac{C_F \zeta_2}{\epsilon} + \frac{2C_A C_F^2 \zeta_2}{\epsilon^2} \\
& + \frac{3C_F \zeta_2}{\epsilon^2} + \frac{333C_A C_F^2}{16\epsilon} - \frac{15C_A C_F}{16\epsilon} - \frac{239C_F}{32\epsilon} + \frac{11C_A C_F^2}{2\epsilon^2} - \frac{C_F}{2\epsilon^2} - \frac{C_A C_F^2}{\epsilon^3} - \frac{3C_F}{2\epsilon^3} \tag{C.11}
\end{aligned}$$

$$\begin{aligned}
V_6^{(m_1,2)} = & \frac{C_A C_F^2 Y_f Y_s \zeta_2^2 M_H^4}{128m_2^2 M_W^2} + \frac{9C_A Y_f Y_s \zeta_2^2 M_H^4}{128m_2^2 M_W^2} + \frac{C_F Y_f Y_s \zeta_2^2 M_H^4}{256m_2^2 M_W^2} - \frac{27C_A C_F^2 Y_f Y_s M_H^4}{256m_2^2 M_W^2} - \frac{243C_A Y_f Y_s M_H^4}{256m_2^2 M_W^2} - \frac{27C_F Y_f Y_s M_H^4}{512m_2^2 M_W^2} \\
& - \frac{C_A C_F^2 Y_f Y_s M_H^4}{128\epsilon m_2^2 M_W^2} - \frac{9C_A Y_f Y_s M_H^4}{128\epsilon m_2^2 M_W^2} - \frac{C_F Y_f Y_s M_H^4}{256\epsilon m_2^2 M_W^2} + \frac{C_A C_F^2 Y_f Y_s M_H^4}{256\epsilon^2 m_2^2 M_W^2} + \frac{9C_A Y_f Y_s M_H^4}{256\epsilon^2 m_2^2 M_W^2} + \frac{C_F Y_f Y_s M_H^4}{512\epsilon^2 m_2^2 M_W^2} \\
& + \frac{5C_A C_F^2 Y_f Y_s \zeta_2 M_H^4}{256m_2^2 M_W^2} + \frac{45C_A Y_f Y_s \zeta_2 M_H^4}{256m_2^2 M_W^2} + \frac{5C_F Y_f Y_s \zeta_2 M_H^4}{512m_2^2 M_W^2} + \frac{C_A C_F^2 Y_f Y_s \zeta_2 M_H^4}{64m_2^2 M_W^2} + \frac{9C_A Y_f Y_s \zeta_2 M_H^4}{64m_2^2 M_W^2} \\
& + \frac{C_F Y_f Y_s \zeta_2 M_H^4}{128m_2^2 M_W^2} - \frac{C_A C_F^2 Y_f Y_s \zeta_2 M_H^4}{128\epsilon m_2^2 M_W^2} - \frac{9C_A Y_f Y_s \zeta_2 M_H^4}{128\epsilon m_2^2 M_W^2} - \frac{C_F Y_f Y_s \zeta_2 M_H^4}{256\epsilon m_2^2 M_W^2} + \frac{357C_A Y_f Y_s^2 M_H^4}{512m_2^2 M_W} \\
& + \frac{3C_A Y_f Y_s^2 M_H^4}{16\epsilon m_2^2 M_W} + \frac{3C_A Y_f Y_s^2 M_H^4}{32\epsilon^2 m_2^2 M_W} + \frac{3C_A Y_f Y_s^2 M_H^4}{64\epsilon^3 m_2^2 M_W} + \frac{17973C_A Y_f Y_s^2 M_H^4}{8192m_2^3 M_W} + \frac{1509C_A Y_f Y_s^2 M_H^4}{1024\epsilon m_2^3 M_W} + \frac{513C_A Y_f Y_s^2 M_H^4}{1024\epsilon^2 m_2^3 M_W} \\
& - \frac{309C_A Y_f Y_s^2 M_H^4}{1024\epsilon^3 m_2^3 M_W} - \frac{16797C_A m_1 Y_f Y_s^2 M_H^4}{8192m_2^4 M_W} - \frac{1185C_A m_1 Y_f Y_s^2 M_H^4}{1024\epsilon m_2^4 M_W} - \frac{405C_A m_1 Y_f Y_s^2 M_H^4}{1024\epsilon^2 m_2^4 M_W} + \frac{189C_A m_1 Y_f Y_s^2 M_H^4}{1024\epsilon^3 m_2^4 M_W} \\
& - \frac{3C_A Y_f Y_s^2 \zeta_2 M_H^4}{64m_2^2 M_W} + \frac{3C_A Y_f Y_s^2 \zeta_2 M_H^4}{32\epsilon m_2^2 M_W} + \frac{2571C_A Y_f Y_s^2 \zeta_2 M_H^4}{1024m_2^3 M_W} - \frac{309C_A Y_f Y_s^2 \zeta_2 M_H^4}{512\epsilon m_2^3 M_W} - \frac{1755C_A m_1 Y_f Y_s^2 \zeta_2 M_H^4}{1024m_2^4 M_W} \\
& + \frac{189C_A m_1 Y_f Y_s^2 \zeta_2 M_H^4}{512\epsilon m_2^4 M_W} + \frac{3C_A Y_f^2 Y_s \zeta_2 M_H^4}{32m_2^2 M_W} - \frac{3C_A Y_f^2 Y_s \zeta_2 M_H^4}{16\epsilon m_2^2 M_W} - \frac{75C_A Y_f^2 Y_s \zeta_2 M_H^4}{256m_2^2 M_W} + \frac{21C_A Y_f^2 Y_s \zeta_2 M_H^4}{128\epsilon m_2^2 M_W} \\
& - \frac{357C_A Y_f^2 Y_s M_H^4}{256m_2^2 M_W} - \frac{3C_A Y_f^2 Y_s M_H^4}{8\epsilon m_2^2 M_W} - \frac{3C_A Y_f^2 Y_s M_H^4}{16\epsilon^2 m_2^2 M_W} - \frac{3C_A Y_f^2 Y_s M_H^4}{32\epsilon^3 m_2^2 M_W} + \frac{1755C_A Y_f^2 Y_s M_H^4}{2048m_2^2 M_W} + \frac{39C_A Y_f^2 Y_s M_H^4}{256\epsilon m_2^2 M_W} \\
& + \frac{15C_A Y_f^2 Y_s M_H^4}{256\epsilon^2 m_2^2 M_W} + \frac{21C_A Y_f^2 Y_s M_H^4}{128\epsilon m_2^2 M_W} - \frac{15C_A Y_f^2 Y_s \zeta_2 M_H^4}{128m_2^2 M_W} + \frac{15C_A Y_f^2 Y_s \zeta_2 M_H^4}{64\epsilon m_2^2 M_W} + \frac{12855C_A Y_f^2 Y_s \zeta_2 M_H^4}{2048m_2^3 M_W} \\
& - \frac{1545C_A Y_f^2 Y_s \zeta_2 M_H^4}{1024\epsilon m_2^3 M_W} - \frac{8775C_A m_1 Y_f Y_s^2 \zeta_2 M_H^4}{2048m_2^4 M_W} + \frac{945C_A m_1 Y_f Y_s^2 \zeta_2 M_H^4}{1024\epsilon m_2^4 M_W} + \frac{15C_A Y_f^2 Y_s \zeta_2 M_H^4}{64m_2^2 M_W} - \frac{15C_A Y_f^2 Y_s \zeta_2 M_H^4}{32\epsilon m_2^2 M_W} \\
& - \frac{375C_A Y_f^2 Y_s \zeta_2 M_H^4}{512m_2^2 M_W} + \frac{105C_A Y_f^2 Y_s \zeta_2 M_H^4}{256\epsilon m_2^2 M_W} - \frac{15C_A Y_f^2 Y_s \zeta_2 M_H^4}{128m_2^2 M_W} - \frac{3C_A Y_f^2 Y_s \zeta_2 M_H^4}{16\epsilon m_2^2 M_W} - \frac{3C_A Y_f^2 Y_s \zeta_2 M_H^4}{32\epsilon^2 m_2^2 M_W} \\
& - \frac{9435C_A Y_f^2 Y_s \zeta_2 M_H^4}{2048m_2^3 M_W} - \frac{513C_A Y_f^2 Y_s \zeta_2 M_H^4}{512\epsilon m_2^3 M_W} + \frac{309C_A Y_f^2 Y_s \zeta_2 M_H^4}{512\epsilon^2 m_2^3 M_W} + \frac{6819C_A m_1 Y_f Y_s^2 \zeta_2 M_H^4}{2048m_2^4 M_W} \\
& + \frac{405C_A m_1 Y_f Y_s^2 \zeta_2 M_H^4}{512\epsilon m_2^4 M_W} - \frac{189C_A m_1 Y_f Y_s^2 \zeta_2 M_H^4}{512\epsilon^2 m_2^4 M_W} + \frac{15C_A Y_f^2 Y_s \zeta_2 M_H^4}{64m_2^2 M_W} + \frac{3C_A Y_f^2 Y_s \zeta_2 M_H^4}{8\epsilon m_2^2 M_W} + \frac{3C_A Y_f^2 Y_s \zeta_2 M_H^4}{16\epsilon^2 m_2^2 M_W} \\
& + \frac{75C_A Y_f^2 Y_s \zeta_2 M_H^4}{512m_2^2 M_W} - \frac{15C_A Y_f^2 Y_s \zeta_2 M_H^4}{128\epsilon m_2^2 M_W} - \frac{21C_A Y_f^2 Y_s \zeta_2 M_H^4}{128\epsilon^2 m_2^2 M_W} + \frac{3721C_A Y_f Y_s^3}{2048m_2^3} + \frac{49C_A Y_f Y_s^3}{64\epsilon m_2^3} + \frac{3C_A Y_f Y_s^3}{16\epsilon^2 m_2^3} \\
& - \frac{C_A Y_f Y_s^3}{256\epsilon^3 m_2^3} - \frac{3099C_A m_1 Y_f Y_s^3}{2048m_2^4} - \frac{39C_A m_1 Y_f Y_s^3}{64\epsilon m_2^4} - \frac{3C_A m_1 Y_f Y_s^3}{16\epsilon^2 m_2^4} + \frac{3C_A m_1 Y_f Y_s^3}{256\epsilon^3 m_2^4} - \frac{27355C_A Y_f Y_s^3}{4096m_2^5} - \frac{1555C_A Y_f Y_s^3}{1024\epsilon m_2^5} \\
& - \frac{1649C_A Y_f Y_s^3}{4096\epsilon^2 m_2^5} + \frac{7C_A Y_f Y_s^3}{4096\epsilon^3 m_2^5} + \frac{C_A Y_f Y_s^3}{1024\epsilon^4 m_2^5} + \frac{21391C_A m_1 Y_f Y_s^3}{4096m_2^6} + \frac{1291C_A m_1 Y_f Y_s^3}{1024\epsilon m_2^6} + \frac{1491C_A m_1 Y_f Y_s^3}{4096\epsilon^2 m_2^6} - \frac{175C_A m_1 Y_f Y_s^3}{4096\epsilon^3 m_2^6} \\
& - \frac{C_A m_1 Y_f Y_s^3}{1024\epsilon^4 m_2^6} - \frac{3827}{64} C_A C_F^2 + \frac{245}{16} C_A C_F Y_f^2 + \frac{13C_A C_F Y_f^2}{4\epsilon} + \frac{7C_A C_F Y_f^2}{4\epsilon^2} - \frac{C_A C_F Y_f^2}{2\epsilon^3} - \frac{1145C_A Y_f^2 Y_s^2}{2048m_2^2} - \frac{145C_A Y_f^2 Y_s^2}{256\epsilon m_2^2}
\end{aligned}$$

$$\begin{aligned}
& -\frac{15C_A Y_f^2 Y_s^2}{256\epsilon^2 m_2^2} + \frac{9C_A Y_f^2 Y_s^2}{256\epsilon^3 m_2^2} + \frac{1711C_A m_1 Y_f^2 Y_s^2}{2048m_2^3} + \frac{65C_A m_1 Y_f^2 Y_s^2}{128\epsilon m_2^3} + \frac{33C_A m_1 Y_f^2 Y_s^2}{256\epsilon^2 m_2^3} - \frac{7C_A m_1 Y_f^2 Y_s^2}{256\epsilon^3 m_2^3} + \frac{1513C_A Y_f^2 Y_s^2}{1024m_2^4} \\
& + \frac{209C_A Y_f^2 Y_s^2}{512\epsilon m_2^4} + \frac{27C_A Y_f^2 Y_s^2}{256\epsilon^2 m_2^4} - \frac{7C_A Y_f^2 Y_s^2}{1024\epsilon^3 m_2^4} + \frac{3C_A Y_f^2 Y_s^2}{256\epsilon^4 m_2^4} - \frac{913C_A C_F Y_s^2}{256m_2^2} - \frac{77C_A C_F Y_s^2}{32\epsilon m_2^2} - \frac{31C_A C_F Y_s^2}{32\epsilon^2 m_2^2} + \frac{C_A C_F Y_s^2}{32\epsilon^3 m_2^2} \\
& + \frac{1059C_A C_F m_1 Y_s^2}{256m_2^3} + \frac{69C_A C_F m_1 Y_s^2}{32\epsilon m_2^3} + \frac{39C_A C_F m_1 Y_s^2}{32\epsilon^2 m_2^3} - \frac{3C_A C_F m_1 Y_s^2}{32\epsilon^3 m_2^3} + \frac{539C_A C_F Y_s^2}{128m_2^4} - \frac{65C_A C_F Y_s^2}{64\epsilon m_2^4} - \frac{643C_A C_F Y_s^2}{256\epsilon^2 m_2^4} \\
& + \frac{15C_A C_F Y_s^2}{64\epsilon^3 m_2^4} + \frac{101C_A Y_f Y_s^3 C_{m_2}^2}{256m_2^3} - \frac{C_A Y_f Y_s^3 C_{m_2}^2}{128\epsilon m_2^3} - \frac{111C_A m_1 Y_f Y_s^3 C_{m_2}^2}{256m_2^4} + \frac{3C_A m_1 Y_f Y_s^3 C_{m_2}^2}{128\epsilon m_2^4} - \frac{63C_A Y_f Y_s^3 C_{m_2}^2}{64m_2^5} \\
& - \frac{C_A Y_f Y_s^3 C_{m_2}^2}{2048\epsilon m_2^5} + \frac{C_A Y_f Y_s^3 C_{m_2}^2}{512\epsilon^2 m_2^5} + \frac{1065C_A m_1 Y_f Y_s^3 C_{m_2}^2}{1024m_2^6} - \frac{167C_A m_1 Y_f Y_s^3 C_{m_2}^2}{2048\epsilon m_2^6} - \frac{C_A m_1 Y_f Y_s^3 C_{m_2}^2}{512\epsilon^2 m_2^6} \\
& - \frac{23}{8} C_A C_F^2 C_{m_2}^2 + 6C_A C_F Y_f^2 C_{m_2}^2 - \frac{C_A C_F Y_f^2 C_{m_2}^2}{\epsilon} - \frac{75C_A Y_f^2 Y_s^2 C_{m_2}^2}{256m_2^2} + \frac{9C_A Y_f^2 Y_s^2 C_{m_2}^2}{128\epsilon m_2^2} + \frac{101C_A m_1 Y_f^2 Y_s^2 C_{m_2}^2}{256m_2^3} \\
& - \frac{7C_A m_1 Y_f^2 Y_s^2 C_{m_2}^2}{128\epsilon m_2^3} + \frac{199C_A Y_f^2 Y_s^2 C_{m_2}^2}{512m_2^4} - \frac{31C_A Y_f^2 Y_s^2 C_{m_2}^2}{512\epsilon m_2^4} + \frac{3C_A Y_f^2 Y_s^2 C_{m_2}^2}{128\epsilon^2 m_2^4} - \frac{67C_A C_F Y_s^2 C_{m_2}^2}{32m_2^2} + \frac{C_A C_F Y_s^2 C_{m_2}^2}{16\epsilon m_2^2} \\
& + \frac{93C_A C_F m_1 Y_s^2 C_{m_2}^2}{32m_2^3} - \frac{3C_A C_F m_1 Y_s^2 C_{m_2}^2}{16\epsilon m_2^3} - \frac{731C_A C_F Y_s^2 C_{m_2}^2}{128m_2^4} + \frac{15C_A C_F Y_s^2 C_{m_2}^2}{32\epsilon m_2^4} + \frac{1}{4} C_A C_F C_{m_2}^2 + \frac{117}{16} C_F C_{m_2}^2 \\
& + \frac{1}{4} C_A C_F^n T_f C_{m_2}^2 - \frac{C_A C_F M_W Y_f C_{m_2}^2}{2m_-} + \frac{C_A C_F M_W Y_f C_{m_2}^2}{\epsilon m_-} - \frac{7C_A Y_f^3 Y_s C_{m_2}^2}{8m_2} + \frac{C_A Y_f^3 Y_s C_{m_2}^2}{32\epsilon m_2} + \frac{C_A m_1 Y_f^3 Y_s C_{m_2}^2}{4m_2^2} \\
& + \frac{C_A m_1 Y_f^3 Y_s C_{m_2}^2}{32\epsilon m_2^2} + \frac{C_A C_F M_W Y_s C_{m_2}^2}{4m_- m_2} - \frac{C_A C_F M_W Y_s C_{m_2}^2}{2\epsilon m_- m_2} - \frac{13C_A C_F M_W Y_s C_{m_2}^2}{2m_2^2} + \frac{3C_A C_F M_W Y_s C_{m_2}^2}{4\epsilon m_2^2} \\
& + \frac{C_A C_F^2 M_W^2 Y_f Y_s C_{m_2}^2}{2m_- m_2^2} + \frac{C_F M_W^2 Y_f Y_s C_{m_2}^2}{4m_- m_2^2} + \frac{C_A C_F^2 Y_f Y_s C_{m_2}^2}{2m_2} + \frac{55C_A C_F Y_f Y_s C_{m_2}^2}{32m_2} + \frac{C_F Y_f Y_s C_{m_2}^2}{4m_2} \\
& - \frac{33C_A C_F Y_f Y_s C_{m_2}^2}{16\epsilon m_2} - \frac{27C_A C_F m_1 Y_f Y_s C_{m_2}^2}{32m_2^2} + \frac{13C_A C_F m_1 Y_f Y_s C_{m_2}^2}{16\epsilon m_2^2} - \frac{5C_A C_F^2 C_{m_2}^2}{4\epsilon} - \frac{49C_F C_{m_2}^2}{8\epsilon} + \frac{19C_A C_F}{8} \\
& - \frac{6167C_F}{128} + \frac{7}{4} C_A C_F^n T_f + \frac{C_A C_F^n T_f}{2\epsilon} + \frac{C_A C_F^n T_f}{8\epsilon^2} + \frac{119C_A C_F M_W Y_f}{16m_-} + \frac{2C_A C_F M_W Y_f}{\epsilon m_-} + \frac{C_A C_F M_W Y_f}{\epsilon^2 m_-} \\
& + \frac{C_A C_F M_W Y_f}{2\epsilon^3 m_-} - \frac{1245C_A Y_f^3 Y_s}{512m_2} - \frac{15C_A Y_f^3 Y_s}{32\epsilon m_2} - \frac{51C_A Y_f^3 Y_s}{128\epsilon^2 m_2} + \frac{C_A Y_f^3 Y_s}{64\epsilon^3 m_2} + \frac{403C_A m_1 Y_f^3 Y_s}{512m_2^2} + \frac{5C_A m_1 Y_f^3 Y_s}{32\epsilon m_2^2} \\
& + \frac{21C_A m_1 Y_f^3 Y_s}{128\epsilon^2 m_2^2} + \frac{C_A m_1 Y_f^3 Y_s}{64\epsilon^3 m_2^2} - \frac{119C_A C_F M_W Y_s}{32m_- m_2} - \frac{C_A C_F M_W Y_s}{\epsilon m_- m_2} - \frac{C_A C_F M_W Y_s}{2\epsilon^2 m_- m_2} - \frac{C_A C_F M_W Y_s}{4\epsilon^3 m_- m_2} \\
& - \frac{791C_A C_F M_W Y_s}{64m_2^2} - \frac{93C_A C_F M_W Y_s}{16\epsilon m_2^2} - \frac{37C_A C_F M_W Y_s}{16\epsilon^2 m_2^2} + \frac{3C_A C_F M_W Y_s}{8\epsilon^3 m_2^2} - \frac{13C_A C_F^2 M_W^2 Y_f Y_s}{2m_- m_2^2} - \frac{13C_F M_W^2 Y_f Y_s}{4m_- m_2^2} \\
& - \frac{5C_A C_F^2 M_W^2 Y_f Y_s}{8\epsilon m_- m_2^2} - \frac{5C_F M_W^2 Y_f Y_s}{16\epsilon m_- m_2^2} + \frac{C_A C_F^2 M_W^2 Y_f Y_s}{4\epsilon^2 m_- m_2^2} + \frac{C_F M_W^2 Y_f Y_s}{8\epsilon^2 m_- m_2^2} + \frac{81C_A C_F^2 Y_f Y_s}{16m_2} - \frac{1347C_A C_F Y_f Y_s}{256m_2} \\
& + \frac{81C_F Y_f Y_s}{32m_2} + \frac{21C_A C_F^2 Y_f Y_s}{16\epsilon m_2} - \frac{295C_A C_F Y_f Y_s}{64\epsilon m_2} + \frac{21C_F Y_f Y_s}{32\epsilon m_2} + \frac{C_A C_F^2 Y_f Y_s}{4\epsilon^2 m_2} - \frac{55C_A C_F Y_f Y_s}{32\epsilon^2 m_2} + \frac{C_F Y_f Y_s}{8\epsilon^2 m_2} \\
& - \frac{33C_A C_F Y_f Y_s}{32\epsilon^3 m_2} - \frac{25C_A C_F^2 m_1 Y_f Y_s}{16m_2^2} + \frac{135C_A C_F m_1 Y_f Y_s}{256m_2^2} - \frac{25C_F m_1 Y_f Y_s}{32m_2^2} - \frac{5C_A C_F^2 m_1 Y_f Y_s}{16\epsilon m_2^2} + \frac{187C_A C_F m_1 Y_f Y_s}{64\epsilon m_2^2} \\
& - \frac{5C_F m_1 Y_f Y_s}{32\epsilon m_2^2} + \frac{19C_A C_F m_1 Y_f Y_s}{32\epsilon^2 m_2^2} + \frac{13C_A C_F m_1 Y_f Y_s}{32\epsilon^3 m_2^2} + \frac{505C_A Y_f Y_s^3 C_2}{512m_2^2} - \frac{5C_A Y_f Y_s^3 C_2}{256\epsilon m_2^2} - \frac{555C_A m_1 Y_f Y_s^3 C_2}{512m_2^4} \\
& + \frac{15C_A m_1 Y_f Y_s^3 C_2}{256\epsilon m_2^4} + \frac{193C_A Y_f Y_s^3 C_2}{128m_2^5} - \frac{1465C_A Y_f Y_s^3 C_2}{4096\epsilon m_2^5} + \frac{5C_A Y_f Y_s^3 C_2}{1024\epsilon^2 m_2^5} - \frac{495C_A m_1 Y_f Y_s^3 C_2}{2048m_2^6} + \frac{369C_A m_1 Y_f Y_s^3 C_2}{4096\epsilon m_2^6} \\
& - \frac{5C_A m_1 Y_f Y_s^3 C_2}{1024\epsilon^2 m_2^6} - \frac{115}{16} C_A C_F^2 C_2 + 15C_A C_F Y_f^2 C_2 - \frac{5C_A C_F Y_f^2 C_2}{2\epsilon} - \frac{375C_A Y_f^2 Y_s^2 C_2}{512m_2^2} + \frac{45C_A Y_f^2 Y_s^2 C_2}{256\epsilon m_2^2} \\
& + \frac{505C_A m_1 Y_f^2 Y_s^2 C_2}{512m_2^3} - \frac{35C_A m_1 Y_f^2 Y_s^2 C_2}{256\epsilon m_2^3} - \frac{449C_A Y_f^2 Y_s^2 C_2}{1024m_2^4} + \frac{9C_A Y_f^2 Y_s^2 C_2}{1024\epsilon m_2^4} + \frac{15C_A Y_f^2 Y_s^2 C_2}{256\epsilon^2 m_2^4} - \frac{335C_A C_F Y_s^2 C_2}{64m_2^2} \\
& + \frac{5C_A C_F Y_s^2 C_2}{32\epsilon m_2^2} + \frac{465C_A C_F m_1 Y_s^2 C_2}{64m_2^3} - \frac{15C_A C_F m_1 Y_s^2 C_2}{32\epsilon m_2^3} - \frac{4947C_A C_F Y_s^2 C_2}{256m_2^4} + \frac{107C_A C_F Y_s^2 C_2}{64\epsilon m_2^4} + \frac{5}{8} C_A C_F C_2 + \frac{585C_F C_2}{32} \\
& + \frac{5}{8} C_A C_F^n T_f C_2 - \frac{5C_A C_F M_W Y_f C_2}{4m_-} + \frac{5C_A C_F M_W Y_f C_2}{2\epsilon m_-} - \frac{35C_A Y_f^3 Y_s C_2}{16m_2} + \frac{5C_A Y_f^3 Y_s C_2}{64\epsilon m_2} + \frac{5C_A m_1 Y_f^3 Y_s C_2}{8m_2^2} \\
& + \frac{5C_A m_1 Y_f^3 Y_s C_2}{64\epsilon m_2^2} + \frac{5C_A C_F M_W Y_s C_2}{8m_- m_2} - \frac{5C_A C_F M_W Y_s C_2}{4\epsilon m_- m_2} - \frac{65C_A C_F M_W Y_s C_2}{4m_2^2} + \frac{15C_A C_F M_W Y_s C_2}{8\epsilon m_2^2} \\
& + \frac{5C_A C_F^2 M_W^2 Y_f Y_s C_2}{4m_- m_2^2} + \frac{5C_F M_W^2 Y_f Y_s C_2}{8m_- m_2^2} + \frac{5C_A C_F^2 Y_f Y_s C_2}{4m_2} + \frac{275C_A C_F Y_f Y_s C_2}{64m_2} + \frac{5C_F Y_f Y_s C_2}{8m_2} - \frac{165C_A C_F Y_f Y_s C_2}{32\epsilon m_2} \\
& - \frac{135C_A C_F m_1 Y_f Y_s C_2}{64m_2^2} + \frac{65C_A C_F m_1 Y_f Y_s C_2}{32\epsilon m_2^2} - \frac{25C_A C_F^2 C_2}{8\epsilon} - \frac{245C_F C_2}{16\epsilon} - \frac{795C_A Y_f Y_s^3 C_{m_2}}{512m_2^3} - \frac{3C_A Y_f Y_s^3 C_{m_2}}{8\epsilon m_2^3} \\
& + \frac{C_A Y_f Y_s^3 C_{m_2}}{128\epsilon^2 m_2^3} + \frac{657C_A m_1 Y_f Y_s^3 C_{m_2}}{512m_2^4} + \frac{3C_A m_1 Y_f Y_s^3 C_{m_2}}{8\epsilon m_2^4} - \frac{3C_A m_1 Y_f Y_s^3 C_{m_2}}{128\epsilon^2 m_2^4} + \frac{8409C_A Y_f Y_s^3 C_{m_2}}{2048m_2^5}
\end{aligned}$$

$$\begin{aligned}
& + \frac{1653C_A Y_f Y_s^3 \mathcal{L} m_2}{2048\epsilon m_2^5} - \frac{7C_A Y_f Y_s^3 \mathcal{L} m_2}{2048\epsilon^2 m_2^5} - \frac{C_A Y_f Y_s^3 \mathcal{L} m_2}{512\epsilon^3 m_2^5} - \frac{7137C_A m_1 Y_f Y_s^3 \mathcal{L} m_2}{2048m_2^6} - \frac{1495C_A m_1 Y_f Y_s^3 \mathcal{L} m_2}{2048\epsilon m_2^6} \\
& + \frac{175C_A m_1 Y_f Y_s^3 \mathcal{L} m_2}{2048\epsilon^2 m_2^6} + \frac{C_A m_1 Y_f Y_s^3 \mathcal{L} m_2}{512\epsilon^3 m_2^6} + \frac{377}{16} C_A C_F^2 \mathcal{L} m_2 - \frac{37}{4} C_A C_F Y_f^2 \mathcal{L} m_2 - \frac{7C_A C_F Y_f^2 \mathcal{L} m_2}{2\epsilon} + \frac{C_A C_F Y_f^2 \mathcal{L} m_2}{\epsilon^2} \\
& + \frac{679C_A Y_f^2 Y_s^2 \mathcal{L} m_2}{512m_2^2} + \frac{15C_A Y_f^2 Y_s^2 \mathcal{L} m_2}{128\epsilon m_2^2} - \frac{9C_A Y_f^2 Y_s^2 \mathcal{L} m_2}{128\epsilon^2 m_2^2} - \frac{597C_A m_1 Y_f^2 Y_s^2 \mathcal{L} m_2}{512m_2^3} - \frac{33C_A m_1 Y_f^2 Y_s^2 \mathcal{L} m_2}{128\epsilon m_2^3} \\
& + \frac{7C_A m_1 Y_f^2 Y_s^2 \mathcal{L} m_2}{128\epsilon^2 m_2^3} - \frac{695C_A Y_f^2 Y_s^2 \mathcal{L} m_2}{512m_2^4} - \frac{3C_A Y_f^2 Y_s^2 \mathcal{L} m_2}{16\epsilon m_2^4} + \frac{7C_A Y_f^2 Y_s^2 \mathcal{L} m_2}{512\epsilon^2 m_2^4} - \frac{3C_A Y_f^2 Y_s^2 \mathcal{L} m_2}{128\epsilon^3 m_2^4} + \frac{319C_A C_F Y_s^2 \mathcal{L} m_2}{64m_2^2} \\
& + \frac{31C_A C_F Y_s^2 \mathcal{L} m_2}{16\epsilon m_2^2} - \frac{C_A C_F Y_s^2 \mathcal{L} m_2}{16\epsilon^2 m_2^2} - \frac{309C_A C_F m_1 Y_s^2 \mathcal{L} m_2}{64m_2^3} - \frac{39C_A C_F m_1 Y_s^2 \mathcal{L} m_2}{16\epsilon m_2^3} + \frac{3C_A C_F m_1 Y_s^2 \mathcal{L} m_2}{16\epsilon^2 m_2^3} \\
& + \frac{C_A C_F Y_s^2 \mathcal{L} m_2}{m_2^4} + \frac{643C_A C_F Y_s^2 \mathcal{L} m_2}{128\epsilon m_2^4} - \frac{15C_A C_F Y_s^2 \mathcal{L} m_2}{32\epsilon^2 m_2^4} - \frac{5}{4} C_A C_F \mathcal{L} m_2 + \frac{525}{32} C_F \mathcal{L} m_2 - C_A C_F n_f T_f \mathcal{L} m_2 \\
& - \frac{C_A C_F n_f T_f \mathcal{L} m_2}{4\epsilon} - \frac{5C_A C_F M_W Y_f \mathcal{L} m_2}{4m_-} - \frac{2C_A C_F M_W Y_f \mathcal{L} m_2}{\epsilon m_-} - \frac{C_A C_F M_W Y_f \mathcal{L} m_2}{\epsilon^2 m_-} + \frac{131C_A Y_f^3 Y_s \mathcal{L} m_2}{128m_2} \\
& + \frac{51C_A Y_f^3 Y_s \mathcal{L} m_2}{64\epsilon m_2} - \frac{C_A Y_f^3 Y_s \mathcal{L} m_2}{32\epsilon^2 m_2} - \frac{29C_A m_1 Y_f^3 Y_s \mathcal{L} m_2}{128m_2^2} - \frac{21C_A m_1 Y_f^3 Y_s \mathcal{L} m_2}{64\epsilon m_2^2} - \frac{C_A m_1 Y_f^3 Y_s \mathcal{L} m_2}{32\epsilon^2 m_2^2} \\
& + \frac{5C_A C_F M_W Y_s \mathcal{L} m_2}{8m_- m_2} + \frac{C_A C_F M_W Y_s \mathcal{L} m_2}{\epsilon m_- m_2} + \frac{C_A C_F M_W Y_s \mathcal{L} m_2}{2\epsilon^2 m_- m_2} + \frac{219C_A C_F M_W Y_s \mathcal{L} m_2}{16m_2^2} + \frac{37C_A C_F M_W Y_s \mathcal{L} m_2}{8\epsilon m_2^2} \\
& - \frac{3C_A C_F M_W Y_s \mathcal{L} m_2}{4\epsilon^2 m_2^2} + \frac{5C_A C_F^2 M_W^2 Y_f Y_s \mathcal{L} m_2}{4m_- m_2^2} + \frac{5C_F M_W^2 Y_f Y_s \mathcal{L} m_2}{8m_- m_2^2} - \frac{C_A C_F^2 M_W^2 Y_f Y_s \mathcal{L} m_2}{2\epsilon m_- m_2^2} - \frac{C_F M_W^2 Y_f Y_s \mathcal{L} m_2}{4\epsilon m_- m_2^2} \\
& - \frac{21C_A C_F^2 Y_f Y_s \mathcal{L} m_2}{8m_2} + \frac{227C_A C_F Y_f Y_s \mathcal{L} m_2}{64m_2} - \frac{21C_F Y_f Y_s \mathcal{L} m_2}{16m_2} - \frac{C_A C_F^2 Y_f Y_s \mathcal{L} m_2}{2m_2} + \frac{55C_A C_F Y_f Y_s \mathcal{L} m_2}{16\epsilon m_2} \\
& - \frac{C_F Y_f Y_s \mathcal{L} m_2}{4\epsilon m_2} + \frac{33C_A C_F Y_f Y_s \mathcal{L} m_2}{16\epsilon^2 m_2} + \frac{5C_A C_F^2 m_1 Y_f Y_s \mathcal{L} m_2}{8m_2^2} - \frac{231C_A C_F m_1 Y_f Y_s \mathcal{L} m_2}{64m_2^2} + \frac{5C_F m_1 Y_f Y_s \mathcal{L} m_2}{16m_2^2} \\
& - \frac{19C_A C_F m_1 Y_f Y_s \mathcal{L} m_2}{16\epsilon m_2^2} - \frac{13C_A C_F m_1 Y_f Y_s \mathcal{L} m_2}{16\epsilon^2 m_2^2} + \frac{6C_A C_F^2 \mathcal{L} m_2}{\epsilon} - \frac{C_A C_F \mathcal{L} m_2}{4\epsilon} + \frac{8C_F \mathcal{L} m_2}{\epsilon} + \frac{5C_A C_F^2 \mathcal{L} m_2}{4\epsilon^2} \\
& + \frac{49C_F \mathcal{L} m_2}{8\epsilon^2} - \frac{27C_A C_F^2}{2\epsilon} + \frac{5C_A C_F}{8\epsilon} - \frac{133C_F}{8\epsilon} - \frac{3C_A C_F^2}{\epsilon^2} + \frac{C_A C_F}{8\epsilon^2} - \frac{4C_F}{\epsilon^2} - \frac{5C_A C_F^2}{8\epsilon^3} - \frac{49C_F}{16\epsilon^3}
\end{aligned} \tag{C.12}$$

### C.3 Matching at $\mu \sim m_2$ :

$$\begin{aligned}
V_1^{(m_2)} = & - \frac{2n_f C_A C_F T_f \mathcal{L} m_2}{\epsilon} + 2n_f C_A C_F T_f \mathcal{L}^2 m_2 - 2n_f C_A C_F T_f \mathcal{L} m_2 + \frac{n_f C_A C_F T_f}{\epsilon} \\
& + \frac{n_f C_A C_F T_f}{\epsilon^2} + 5n_f \zeta_2 C_A C_F T_f + 6n_f C_A C_F T_f + \frac{107C_A C_F^2}{8\epsilon} + \frac{9C_A C_F}{16\epsilon} - \frac{429C_A C_F^2}{16\epsilon^2} + \frac{C_A C_F}{16\epsilon^2} + \frac{13C_A C_F^2 \mathcal{L}^2 m_2}{2\epsilon} \\
& + \frac{91C_A C_F^2 \mathcal{L} m_2}{8\epsilon} - \frac{C_A C_F \mathcal{L} m_2}{8\epsilon} - \frac{91Y_f^2 C_A C_F \mathcal{L}^2 m_2}{32\epsilon} - \frac{Y_f^2 C_A C_F^2 \mathcal{L} m_2}{2\epsilon} + \frac{369Y_f^2 C_A C_F \mathcal{L} m_2}{64\epsilon} + \frac{1}{2} Y_f^2 C_A C_F^2 \mathcal{L}^2 m_2 \\
& + \frac{561}{32} Y_f^2 C_A C_F \mathcal{L}^2 m_2 - \frac{7}{2} Y_f^2 C_A C_F^2 \mathcal{L} m_2 - \frac{5157}{64} Y_f^2 C_A C_F \mathcal{L} m_2 - \frac{723}{8} C_A C_F^2 \mathcal{L}^2 m_2 + \frac{3}{8} C_A C_F \mathcal{L}^2 m_2 + \frac{767}{2} C_A C_F^2 \mathcal{L} m_2 \\
& - \frac{11}{4} C_A C_F \mathcal{L} m_2 + \frac{135}{2} S_1 S_2 Y_f^2 C_A C_F - \frac{15}{2} S_1 Y_f^2 C_A C_F + \frac{7Y_f^2 C_A C_F^2}{4\epsilon} - \frac{919Y_f^2 C_A C_F}{256\epsilon} + \frac{Y_f^2 C_A C_F^2}{4\epsilon^2} + \frac{11Y_f^2 C_A C_F}{4\epsilon^2} \\
& - \frac{251Y_f^2 \zeta_2 C_A C_F}{64\epsilon} + \frac{5}{4} Y_f^2 \zeta_2 C_A C_F^2 - \frac{23}{64} Y_f^2 \zeta_2 C_A C_F + 5i\pi Y_f^2 \zeta_2 C_A C_F - 15Y_f^2 \zeta_3 C_A C_F + 8Y_f^2 C_A C_F^2 + \frac{36443}{256} Y_f^2 C_A C_F \\
& - \frac{3\zeta_2 C_A C_F^2}{4\epsilon} + \frac{849}{16} \zeta_2 C_A C_F^2 + \frac{15}{16} \zeta_2 C_A C_F - \frac{23147}{32} C_A C_F^2 + \frac{177C_A C_F}{32} + \frac{7Y_f^4 C_A \mathcal{L}^2 m_2}{32\epsilon} - \frac{33Y_f^4 C_A \mathcal{L} m_2}{64\epsilon} - \frac{133}{128} Y_f^4 C_A \mathcal{L}^2 m_2 \\
& + \frac{1515}{256} Y_f^4 C_A \mathcal{L} m_2 + \frac{27}{8} S_1 S_2 Y_f^4 C_A - \frac{9}{16} S_1 Y_f^4 C_A - \frac{77Y_f^4 C_A}{256\epsilon} - \frac{Y_f^4 C_A}{8\epsilon^2} + \frac{35Y_f^4 \zeta_2 C_A}{64\epsilon} + \frac{339}{256} Y_f^4 \zeta_2 C_A + \frac{1}{4} i\pi Y_f^4 \zeta_2 C_A \\
& - \frac{3}{4} Y_f^4 \zeta_3 C_A - \frac{11261Y_f^4 C_A}{1024} + \frac{3589C_F}{32\epsilon} - \frac{237C_F}{32\epsilon^2} + \frac{11C_F \mathcal{L}^2 m_2}{\epsilon} - \frac{823C_F \mathcal{L} m_2}{16\epsilon} - \frac{Y_f^2 C_F \mathcal{L} m_2}{4\epsilon} + \frac{1}{4} Y_f^2 C_F \mathcal{L}^2 m_2 - \frac{7}{4} Y_f^2 C_F \mathcal{L} m_2 \\
& - \frac{845}{16} C_F \mathcal{L}^2 m_2 + \frac{4295}{16} C_F \mathcal{L} m_2 - 81S_1 S_2 C_F + 45S_1 C_F + \frac{7Y_f^2 C_F}{8\epsilon} + \frac{Y_f^2 C_F}{8\epsilon^2} + \frac{5}{8} Y_f^2 \zeta_2 C_F + 4Y_f^2 C_F - \frac{13\zeta_2 C_F}{2\epsilon} \\
& - \frac{745\zeta_2 C_F}{32} - 6i\pi \zeta_2 C_F + 18\zeta_3 C_F - \frac{4085C_F}{8}
\end{aligned} \tag{C.13}$$

$$\begin{aligned}
V_2^{(m_2)} = & \frac{2n_f C_A C_F T_f \mathcal{L} m_2}{\epsilon} - 5n_f C_A C_F T_f \mathcal{L}^2 m_2 + \frac{47}{2} n_f C_A C_F T_f \mathcal{L} m_2 - \frac{2n_f C_A C_F T_f}{\epsilon} - \frac{n_f C_A C_F T_f}{\epsilon^2} - \frac{25}{2} n_f \zeta_2 C_A C_F T_f \\
& - \frac{429}{8} n_f C_A C_F T_f + \frac{53C_A C_F^2}{2\epsilon} - \frac{11C_A C_F}{16\epsilon} + \frac{115C_A C_F^2}{16\epsilon^2} - \frac{3C_A C_F}{16\epsilon^2} - \frac{C_A C_F^2 \mathcal{L}^2 m_2}{2\epsilon} - \frac{89C_A C_F^2 \mathcal{L} m_2}{8\epsilon} + \frac{3C_A C_F \mathcal{L} m_2}{8\epsilon} \\
& - \frac{5Y_f^2 C_A C_F \mathcal{L}^2 m_2}{8\epsilon} - \frac{43Y_f^2 C_A C_F \mathcal{L} m_2}{16\epsilon} + \frac{713}{96} Y_f^2 C_A C_F \mathcal{L}^2 m_2 - \frac{3}{2} Y_f^2 C_A C_F^2 \mathcal{L} m_2 - \frac{5195}{192} Y_f^2 C_A C_F \mathcal{L} m_2 + \frac{137}{8} C_A C_F^2 \mathcal{L}^2 m_2
\end{aligned}$$

$$\begin{aligned}
& -\frac{5}{8}C_A C_F \mathcal{L}_{m_2}^2 - \frac{305}{4}C_A C_F^2 \mathcal{L}_{m_2} + 3C_A C_F \mathcal{L}_{m_2} - \frac{27}{4}S_1 S_2 Y_f^2 C_A C_F + \frac{39}{8}S_1 Y_f^2 C_A C_F + \frac{3Y_f^2 C_A C_F^2}{4\epsilon} + \frac{1029Y_f^2 C_A C_F}{64\epsilon} \\
& + \frac{3Y_f^2 C_A C_F}{4\epsilon^2} - \frac{109Y_f^2 \zeta_2 C_A C_F}{16\epsilon} + \frac{313}{192}Y_f^2 \zeta_2 C_A C_F - \frac{1}{2}i\pi Y_f^2 \zeta_2 C_A C_F + \frac{3}{2}Y_f^2 \zeta_3 C_A C_F + \frac{19}{4}Y_f^2 C_A C_F^2 + \frac{13805}{768}Y_f^2 C_A C_F \\
& - \frac{17\zeta_2 C_A C_F^2}{4\epsilon} - \frac{131}{16}\zeta_2 C_A C_F^2 - \frac{25}{16}\zeta_2 C_A C_F + \frac{5191}{32}C_A C_F^2 - \frac{289C_A C_F}{32} - \frac{23Y_f^4 C_A \mathcal{L}_{m_2}}{64\epsilon} + \frac{25}{384}Y_f^4 C_A \mathcal{L}_{m_2}^2 \\
& + \frac{1205}{768}Y_f^4 C_A \mathcal{L}_{m_2} - \frac{27}{32}S_1 S_2 Y_f^4 C_A - \frac{9}{64}S_1 Y_f^4 C_A + \frac{235Y_f^4 C_A}{256\epsilon} + \frac{Y_f^4 C_A}{64\epsilon^2} - \frac{3Y_f^4 \zeta_2 C_A}{8\epsilon} + \frac{1373}{768}Y_f^4 \zeta_2 C_A - \frac{1}{16}i\pi Y_f^4 \zeta_2 C_A \\
& + \frac{3}{16}Y_f^4 \zeta_3 C_A - \frac{17099Y_f^4 C_A}{3072} - \frac{651C_F}{16\epsilon} + \frac{131C_F}{32\epsilon^2} - \frac{7C_F \mathcal{L}_{m_2}^2}{4\epsilon} + \frac{219C_F \mathcal{L}_{m_2}}{16\epsilon} - \frac{3}{4}Y_f^2 C_F \mathcal{L}_{m_2} + \frac{399}{16}C_F \mathcal{L}_{m_2}^2 - \frac{2051}{16}C_F \mathcal{L}_{m_2} \\
& - \frac{81}{2}S_1 S_2 C_F + \frac{69S_1 C_F}{2} + \frac{3Y_f^2 C_F}{8\epsilon} + \frac{19Y_f^2 C_F}{8} + \frac{121\zeta_2 C_F}{8\epsilon} + \frac{451\zeta_2 C_F}{32} - 3i\pi \zeta_2 C_F + 9\zeta_3 C_F + \frac{8231C_F}{32} \tag{C.14}
\end{aligned}$$

$$\begin{aligned}
V_3(m_2) = & -\frac{2n_f C_A C_F T_f \mathcal{L}_{m_2}}{\epsilon} + 10n_f C_A C_F T_f \mathcal{L}_{m_2} + \frac{3n_f C_A C_F T_f}{2\epsilon} + \frac{n_f C_A C_F T_f}{\epsilon^2} - \frac{89}{4}n_f C_A C_F T_f - \frac{7853C_A C_F^2}{32\epsilon} + \frac{9C_A C_F}{16\epsilon} \\
& - \frac{45C_A C_F^2}{16\epsilon^2} + \frac{C_A C_F}{16\epsilon^2} - \frac{65C_A C_F^2 \mathcal{L}_{m_2}^2}{4\epsilon} + \frac{445C_A C_F^2 \mathcal{L}_{m_2}}{4\epsilon} - \frac{C_A C_F \mathcal{L}_{m_2}}{8\epsilon} - \frac{29Y_f^2 C_A C_F \mathcal{L}_{m_2}^2}{16\epsilon} - \frac{Y_f^2 C_A C_F^2 \mathcal{L}_{m_2}}{2\epsilon} \\
& - \frac{229Y_f^2 C_A C_F \mathcal{L}_{m_2}}{32\epsilon} + \frac{1}{2}Y_f^2 C_A C_F^2 \mathcal{L}_{m_2}^2 + \frac{373}{32}Y_f^2 C_A C_F \mathcal{L}_{m_2}^2 - \frac{7}{2}Y_f^2 C_A C_F^2 \mathcal{L}_{m_2} - \frac{2575}{64}Y_f^2 C_A C_F \mathcal{L}_{m_2} + \frac{81}{8}C_A C_F^2 \mathcal{L}_{m_2}^2 \\
& + \frac{3}{8}C_A C_F \mathcal{L}_{m_2}^2 - \frac{377}{4}C_A C_F^2 \mathcal{L}_{m_2} - \frac{11}{4}C_A C_F \mathcal{L}_{m_2} + \frac{135}{4}S_1 S_2 Y_f^2 C_A C_F - \frac{9}{4}S_1 Y_f^2 C_A C_F + \frac{7Y_f^2 C_A C_F^2}{4\epsilon} + \frac{4611Y_f^2 C_A C_F}{128\epsilon} \\
& + \frac{Y_f^2 C_A C_F^2}{4\epsilon^2} + \frac{9Y_f^2 C_A C_F}{4\epsilon^2} - \frac{541Y_f^2 \zeta_2 C_A C_F}{32\epsilon} + \frac{5}{4}Y_f^2 \zeta_2 C_A C_F^2 - \frac{127}{64}Y_f^2 \zeta_2 C_A C_F + \frac{5}{2}i\pi Y_f^2 \zeta_2 C_A C_F \\
& - \frac{15}{2}Y_f^2 \zeta_3 C_A C_F + 8Y_f^2 C_A C_F^2 + \frac{9617}{256}Y_f^2 C_A C_F + \frac{223\zeta_2 C_A C_F^2}{8\epsilon} + \frac{53}{16}\zeta_2 C_A C_F^2 + \frac{15}{16}\zeta_2 C_A C_F + \frac{6811}{32}C_A C_F^2 + \frac{177C_A C_F}{32} \\
& + \frac{33Y_f^4 C_A \mathcal{L}_{m_2}^2}{128\epsilon} - \frac{155Y_f^4 C_A \mathcal{L}_{m_2}}{256\epsilon} - \frac{69}{64}Y_f^4 C_A \mathcal{L}_{m_2}^2 + \frac{175}{32}Y_f^4 C_A \mathcal{L}_{m_2} + \frac{27}{8}S_1 S_2 Y_f^4 C_A - \frac{9}{16}S_1 Y_f^4 C_A - \frac{339Y_f^4 C_A}{1024\epsilon} - \frac{Y_f^4 C_A}{8\epsilon^2} \\
& + \frac{177Y_f^4 \zeta_2 C_A}{256\epsilon} + \frac{117}{128}Y_f^4 \zeta_2 C_A + \frac{1}{4}i\pi Y_f^4 \zeta_2 C_A - \frac{3}{4}Y_f^4 \zeta_3 C_A - \frac{1153Y_f^4 C_A}{128} + \frac{1951C_F}{32\epsilon} - \frac{109C_F}{32\epsilon^2} + \frac{C_F \mathcal{L}_{m_2}^2}{4\epsilon} - \frac{211C_F \mathcal{L}_{m_2}}{16\epsilon} \\
& - \frac{Y_f^2 C_F \mathcal{L}_{m_2}}{4\epsilon} + \frac{1}{4}Y_f^2 C_F \mathcal{L}_{m_2}^2 - \frac{7}{4}Y_f^2 C_F \mathcal{L}_{m_2} - \frac{185}{16}C_F \mathcal{L}_{m_2}^2 + \frac{819}{16}C_F \mathcal{L}_{m_2} - 27S_1 S_2 C_F - \frac{75S_1 C_F}{2} + \frac{7Y_f^2 C_F}{8\epsilon} + \frac{Y_f^2 C_F}{8\epsilon^2} \\
& + \frac{5}{8}Y_f^2 \zeta_2 C_F + 4Y_f^2 C_F - \frac{183\zeta_2 C_F}{8\epsilon} - \frac{333\zeta_2 C_F}{32} - 2i\pi \zeta_2 C_F + 6\zeta_3 C_F - \frac{875C_F}{8} \tag{C.15}
\end{aligned}$$

$$\begin{aligned}
V_4(m_2) = & \frac{127C_A C_F^2 Y_s^2 M_H^4}{256m_2^4 M_W^2} + \frac{1143C_A Y_s^2 M_H^4}{256m_2^4 M_W^2} + \frac{127C_F Y_s^2 M_H^4}{512m_2^4 M_W^2} - \frac{C_A C_F^2 Y_s^2 M_H^4}{32\epsilon m_2^4 M_W^2} - \frac{9C_A Y_s^2 M_H^4}{32\epsilon m_2^4 M_W^2} - \frac{C_F Y_s^2 M_H^4}{64\epsilon m_2^4 M_W^2} - \frac{C_A C_F^2 Y_s^2 M_H^4}{128\epsilon^2 m_2^4 M_W^2} \\
& - \frac{9C_A Y_s^2 M_H^4}{128\epsilon^2 m_2^4 M_W^2} - \frac{C_F Y_s^2 M_H^4}{256\epsilon^2 m_2^4 M_W^2} + \frac{C_A C_F^2 Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{64m_2^4 M_W^2} + \frac{9C_A Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{64m_2^4 M_W^2} + \frac{C_F Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{128m_2^4 M_W^2} + \frac{5C_A C_F^2 Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{128m_2^4 M_W^2} \\
& + \frac{45C_A Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{128m_2^4 M_W^2} + \frac{5C_F Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{256m_2^4 M_W^2} - \frac{9C_A C_F^2 Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{64m_2^4 M_W^2} - \frac{81C_A Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{64m_2^4 M_W^2} - \frac{9C_F Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{128m_2^4 M_W^2} + \frac{C_A C_F^2 Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{64\epsilon m_2^4 M_W^2} \\
& + \frac{9C_A Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{64\epsilon m_2^4 M_W^2} + \frac{C_F Y_s^2 \mathcal{L}_{m_2}^2 M_H^4}{128\epsilon m_2^4 M_W^2} - \frac{5373C_A Y_s^3 M_H^2}{512m_2^4 M_W} + \frac{345C_A Y_s^3 M_H^2}{64\epsilon m_2^4 M_W} - \frac{45C_A Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{64m_2^4 M_W} + \frac{3C_A Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{8\epsilon m_2^4 M_W} \\
& + \frac{63C_A Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{128m_2^4 M_W} - \frac{3C_A Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{8\epsilon m_2^4 M_W} + \frac{603C_A Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{128m_2^4 M_W} - \frac{39C_A Y_s^3 \mathcal{L}_{m_2}^2 M_H^2}{16\epsilon m_2^4 M_W} - \frac{5695C_A Y_s^4}{2048\epsilon m_2^4} + \frac{3457C_A Y_s^4}{2048\epsilon m_2^4} - \frac{C_A Y_s^4}{16\epsilon^2 m_2^4} \\
& + \frac{3289}{32}C_A C_F^2 + \frac{131C_A C_F^2 M_W^2 Y_s^2}{4m_2^4} + \frac{131C_F M_W^2 Y_s^2}{8m_2^4} - \frac{7C_A C_F^2 M_W^2 Y_s^2}{4\epsilon m_2^4} - \frac{7C_F M_W^2 Y_s^2}{8\epsilon m_2^4} - \frac{C_A C_F^2 M_W^2 Y_s^2}{2\epsilon^2 m_2^4} - \frac{C_F M_W^2 Y_s^2}{4\epsilon^2 m_2^4} \\
& + \frac{81C_A C_F S_1 S_2 Y_s^2}{8m_2^2} + \frac{115C_A C_F^2 Y_s^2}{16m_2^2} + \frac{11063C_A C_F Y_s^2}{256m_2^2} + \frac{115C_F Y_s^2}{32m_2^2} - \frac{3857C_A C_F Y_s^2}{256\epsilon m_2^2} + \frac{C_A C_F Y_s^2}{4\epsilon^2 m_2^2} - \frac{105C_A Y_s^4 \mathcal{L}_{m_2}^2}{256m_2^4} \\
& + \frac{37C_A Y_s^4 \mathcal{L}_{m_2}^2}{256\epsilon m_2^4} + \frac{17}{4}C_A C_F^2 \mathcal{L}_{m_2}^2 + \frac{C_A C_F^2 M_W^2 Y_s^2 \mathcal{L}_{m_2}^2}{m_2^4} + \frac{C_F M_W^2 Y_s^2 \mathcal{L}_{m_2}^2}{2m_2^4} + \frac{C_A C_F^2 Y_s^2 \mathcal{L}_{m_2}^2}{2m_2^2} + \frac{161C_A C_F Y_s^2 \mathcal{L}_{m_2}^2}{32m_2^2} \\
& + \frac{C_F Y_s^2 \mathcal{L}_{m_2}^2}{4m_2^2} - \frac{61C_A C_F Y_s^2 \mathcal{L}_{m_2}^2}{32\epsilon m_2^2} + \frac{1}{4}C_A C_F \mathcal{L}_{m_2}^2 + \frac{57}{8}C_F \mathcal{L}_{m_2}^2 + 2C_A C_F n_f T_f \mathcal{L}_{m_2}^2 - \frac{125C_A C_F M_W Y_s \mathcal{L}_{m_2}^2}{16m_2^2} \\
& + \frac{19C_A C_F M_W Y_s \mathcal{L}_{m_2}^2}{8\epsilon m_2^2} - \frac{7C_A C_F^2 \mathcal{L}_{m_2}^2}{\epsilon} - \frac{35C_F \mathcal{L}_{m_2}^2}{4\epsilon} + \frac{319C_A C_F}{32} + \frac{5913C_F}{64} - \frac{27C_F S_1}{2} + \frac{243}{4}C_F S_1 S_2 + 21C_A C_F n_f T_f \\
& + \frac{5C_A C_F n_f T_f}{\epsilon} + \frac{C_A C_F n_f T_f}{\epsilon^2} - \frac{11033C_A C_F M_W Y_s}{128m_2^2} + \frac{2185C_A C_F M_W Y_s}{64\epsilon m_2^2} - \frac{57C_A Y_s^4 \zeta_2}{512m_2^4} + \frac{5C_A Y_s^4 \zeta_2}{512\epsilon m_2^4} - \frac{35}{8}C_A C_F^2 \zeta_2 \\
& + \frac{5C_A C_F^2 M_W^2 Y_s^2 \zeta_2}{2m_2^4} + \frac{5C_F M_W^2 Y_s^2 \zeta_2}{4m_2^4} + \frac{5C_A C_F^2 Y_s^2 \zeta_2}{4m_2^2} - \frac{15C_A C_F Y_s^2 \zeta_2}{64m_2^2} + \frac{5C_F Y_s^2 \zeta_2}{8m_2^2} - \frac{125C_A C_F Y_s^2 \zeta_2}{64\epsilon m_2^2} + \frac{3iC_A C_F \pi Y_s^2 \zeta_2}{4m_2^2} \\
& + \frac{5}{8}C_A C_F \zeta_2 - \frac{59C_F \zeta_2}{16} + 5C_A C_F n_f T_f \zeta_2 - \frac{25C_A C_F M_W Y_s \zeta_2}{32m_2^2} + \frac{15C_A C_F M_W Y_s \zeta_2}{16\epsilon m_2^2} + \frac{29C_A C_F^2 \zeta_2}{2\epsilon} - \frac{35C_F \zeta_2}{8\epsilon} + \frac{9}{2}iC_F \pi \zeta_2
\end{aligned}$$

$$\begin{aligned}
& -\frac{9C_A C_F Y_s^2 \zeta_3}{4m_2^2} - \frac{27C_F \zeta_3}{2} + \frac{833C_A Y_s^4 L m_2}{512m_2^4} - \frac{375C_A Y_s^4 L m_2}{512\epsilon m_2^4} - \frac{315}{8} C_A C_F^2 L m_2 - \frac{19C_A C_F^2 M_W^2 Y_s^2 L m_2}{2m_2^4} \\
& - \frac{19C_F M_W^2 Y_s^2 L m_2}{4m_2^4} + \frac{C_A C_F^2 M_W^2 Y_s^2 L m_2}{\epsilon m_2^4} + \frac{C_F M_W^2 Y_s^2 L m_2}{2\epsilon m_2^4} - \frac{13C_A C_F^2 Y_s^2 L m_2}{4m_2^2} - \frac{1545C_A C_F Y_s^2 L m_2}{64m_2^2} - \frac{13C_F Y_s^2 L m_2}{8m_2^2} \\
& + \frac{551C_A C_F Y_s^2 L m_2}{64\epsilon m_2^2} + \frac{223}{8} C_A C_F L m_2 - \frac{707}{16} C_F L m_2 - 10C_A C_F n_f T_f L m_2 - \frac{2C_A C_F n_f T_f L m_2}{\epsilon} + \frac{1391C_A C_F M_W Y_s L m_2}{32m_2^2} \\
& - \frac{247C_A C_F M_W Y_s L m_2}{16\epsilon m_2^2} + \frac{49C_A C_F^2 L m_2}{\epsilon} + \frac{363C_F L m_2}{8\epsilon} - \frac{905C_A C_F^2}{8\epsilon} - \frac{59C_A C_F}{4\epsilon} - \frac{2655C_F}{32\epsilon} - \frac{7C_A C_F^2}{4\epsilon^2} - \frac{17C_F}{8\epsilon^2} \quad (C.16)
\end{aligned}$$

$$\begin{aligned}
V_5^{(m_2)} = & \frac{405C_A C_F^2 Y_s^2 M_H^4}{512m_2^4 M_W^2} + \frac{3645C_A Y_s^2 M_H^4}{512m_2^4 M_W^2} + \frac{405C_F Y_s^2 M_H^4}{1024m_2^4 M_W^2} - \frac{C_A C_F^2 Y_s^2 M_H^4}{32\epsilon m_2^4 M_W^2} - \frac{9C_A Y_s^2 M_H^4}{32\epsilon m_2^4 M_W^2} - \frac{C_F Y_s^2 M_H^4}{64\epsilon m_2^4 M_W^2} + \frac{C_A C_F^2 Y_s^2 M_H^4}{256\epsilon^2 m_2^4 M_W^2} \\
& + \frac{9C_A Y_s^2 M_H^4}{256\epsilon^2 m_2^4 M_W^2} + \frac{C_F Y_s^2 M_H^4}{512\epsilon^2 m_2^4 M_W^2} + \frac{7C_A C_F^2 Y_s^2 L m_2 M_H^4}{128m_2^4 M_W^2} + \frac{63C_A Y_s^2 L m_2 M_H^4}{128m_2^4 M_W^2} + \frac{7C_F Y_s^2 L m_2 M_H^4}{256m_2^4 M_W^2} + \frac{35C_A C_F^2 Y_s^2 L m_2 M_H^4}{256m_2^4 M_W^2} \\
& + \frac{315C_A Y_s^2 L m_2 M_H^4}{256m_2^4 M_W^2} + \frac{35C_F Y_s^2 L m_2 M_H^4}{512m_2^4 M_W^2} - \frac{31C_A C_F^2 Y_s^2 L m_2 M_H^4}{128m_2^4 M_W^2} - \frac{279C_A Y_s^2 L m_2 M_H^4}{128m_2^4 M_W^2} - \frac{31C_F Y_s^2 L m_2 M_H^4}{256m_2^4 M_W^2} \\
& - \frac{C_A C_F^2 Y_s^2 L m_2 M_H^4}{128\epsilon m_2^4 M_W^2} - \frac{9C_A Y_s^2 L m_2 M_H^4}{128\epsilon m_2^4 M_W^2} - \frac{C_F Y_s^2 L m_2 M_H^4}{256\epsilon m_2^4 M_W^2} - \frac{3657C_A Y_s^3 M_H^2}{1024m_2^4 M_W} + \frac{345C_A Y_s^3 M_H^2}{128\epsilon m_2^4 M_W} - \frac{9C_A Y_s^3 L m_2 M_H^2}{128m_2^4 M_W} \\
& + \frac{3C_A Y_s^3 L m_2 M_H^2}{16\epsilon m_2^4 M_W} - \frac{189C_A Y_s^3 L m_2 M_H^2}{256m_2^4 M_W} - \frac{9C_A Y_s^3 L m_2 M_H^2}{16\epsilon m_2^4 M_W} + \frac{303C_A Y_s^3 L m_2 M_H^2}{256m_2^4 M_W} - \frac{39C_A Y_s^3 L m_2 M_H^2}{32\epsilon m_2^4 M_W} - \frac{3C_A S_1 Y_s^4}{16m_2^4} \\
& - \frac{12611C_A Y_s^4}{4096m_2^4} + \frac{3333C_A Y_s^4}{4096\epsilon m_2^4} - \frac{C_A Y_s^4}{64\epsilon^2 m_2^4} + \frac{78671}{192} C_A C_F^2 + \frac{413C_A C_F^2 M_W^2 Y_s^2}{8m_2^4} + \frac{413C_F M_W^2 Y_s^2}{16m_2^4} - \frac{17C_A C_F^2 M_W^2 Y_s^2}{8\epsilon m_2^4} \\
& - \frac{17C_F M_W^2 Y_s^2}{16\epsilon m_2^4} + \frac{C_A C_F^2 M_W^2 Y_s^2}{4\epsilon^2 m_2^4} + \frac{C_F M_W^2 Y_s^2}{8\epsilon^2 m_2^4} - \frac{69C_A C_F S_1 Y_s^2}{16m_2^2} + \frac{135C_A C_F S_1 S_2 Y_s^2}{4m_2^2} + \frac{219C_A C_F^2 Y_s^2}{32m_2^2} + \frac{46149C_A C_F Y_s^2}{1024m_2^2} \\
& + \frac{219C_F Y_s^2}{64m_2^2} + \frac{C_A C_F^2 Y_s^2}{2\epsilon m_2^2} - \frac{2523C_A C_F Y_s^2}{1024\epsilon m_2^2} + \frac{C_F Y_s^2}{4\epsilon m_2^2} + \frac{C_A C_F Y_s^2}{4\epsilon^2 m_2^2} - \frac{197C_A Y_s^4 L m_2}{512m_2^4} + \frac{57C_A Y_s^4 L m_2}{512\epsilon m_2^4} + \frac{731}{24} C_A C_F^2 L m_2 \\
& + \frac{7C_A C_F^2 M_W^2 Y_s^2 L m_2}{2m_2^4} + \frac{7C_F M_W^2 Y_s^2 L m_2}{4m_2^4} + \frac{C_A C_F^2 Y_s^2 L m_2}{4m_2^2} + \frac{691C_A C_F Y_s^2 L m_2}{128m_2^2} + \frac{C_F Y_s^2 L m_2}{8m_2^2} - \frac{263C_A C_F Y_s^2 L m_2}{128\epsilon m_2^2} \\
& + \frac{1487}{24} C_A C_F L m_2 + \frac{503}{48} C_F L m_2 + \frac{1}{3} C_A C_F n_f T_f L m_2 - \frac{323C_A C_F M_W Y_s L m_2}{32m_2^2} + \frac{45C_A C_F M_W Y_s L m_2}{16\epsilon m_2^2} - \frac{19C_A C_F^2 L m_2}{2\epsilon} \\
& - \frac{101C_F L m_2}{8\epsilon} + \frac{43073C_A C_F}{192} + \frac{68407C_F}{384} - \frac{108}{5} C_A C_F S_1 - \frac{117C_F S_1}{2} + \frac{405}{2} C_F S_1 S_2 + \frac{55}{6} C_A C_F n_f T_f + \frac{23C_A C_F n_f T_f}{2\epsilon} \\
& + \frac{3C_A C_F n_f T_f}{2\epsilon^2} - \frac{29295C_A C_F M_W Y_s}{256m_2^2} + \frac{4791C_A C_F M_W Y_s}{128\epsilon m_2^2} - \frac{53C_A Y_s^4 \zeta_2}{1024m_2^4} + \frac{153C_A Y_s^4 \zeta_2}{1024\epsilon m_2^4} + \frac{1975}{48} C_A C_F^2 \zeta_2 \\
& + \frac{35C_A C_F^2 M_W^2 Y_s^2 \zeta_2}{4m_2^4} + \frac{35C_F M_W^2 Y_s^2 \zeta_2}{8m_2^4} + \frac{5C_A C_F^2 Y_s^2 \zeta_2}{8m_2^2} + \frac{99C_A C_F Y_s^2 \zeta_2}{256m_2^2} + \frac{5C_F Y_s^2 \zeta_2}{16m_2^2} - \frac{1447C_A C_F Y_s^2 \zeta_2}{256\epsilon m_2^2} \\
& + \frac{5iC_A C_F \pi Y_s^2 \zeta_2}{2m_2^2} + \frac{1579}{48} C_A C_F \zeta_2 + \frac{3083C_F \zeta_2}{96} + \frac{5}{6} C_A C_F n_f T_f \zeta_2 + \frac{737C_A C_F M_W Y_s \zeta_2}{64m_2^2} + \frac{81C_A C_F M_W Y_s \zeta_2}{32\epsilon m_2^2} \\
& + \frac{85C_A C_F^2 \zeta_2}{4\epsilon} - \frac{205C_F \zeta_2}{16\epsilon} + 15iC_F \pi \zeta_2 - \frac{15C_A C_F Y_s^2 \zeta_3}{2m_2^2} - 45C_F \zeta_3 + \frac{1853C_A Y_s^4 L m_2}{1024m_2^4} - \frac{499C_A Y_s^4 L m_2}{1024\epsilon m_2^4} - \frac{8789}{48} C_A C_F^2 L m_2 \\
& - \frac{61C_A C_F^2 M_W^2 Y_s^2 L m_2}{4m_2^4} - \frac{61C_F M_W^2 Y_s^2 L m_2}{8m_2^4} - \frac{C_A C_F^2 M_W^2 Y_s^2 L m_2}{2\epsilon m_2^4} - \frac{C_F M_W^2 Y_s^2 L m_2}{4\epsilon m_2^4} - \frac{21C_A C_F^2 Y_s^2 L m_2}{8m_2^2} \\
& - \frac{6683C_A C_F Y_s^2 L m_2}{256m_2^2} - \frac{21C_F Y_s^2 L m_2}{16m_2^2} + \frac{1485C_A C_F Y_s^2 L m_2}{256\epsilon m_2^2} - \frac{4799}{48} C_A C_F L m_2 - \frac{7373}{96} C_F L m_2 - \frac{17}{3} C_A C_F n_f T_f L m_2 \\
& - \frac{3C_A C_F n_f T_f L m_2}{\epsilon} + \frac{3737C_A C_F M_W Y_s L m_2}{64m_2^2} - \frac{585C_A C_F M_W Y_s L m_2}{32\epsilon m_2^2} + \frac{103C_A C_F^2 L m_2}{2\epsilon} - \frac{107C_A C_F L m_2}{2\epsilon} + \frac{969C_F L m_2}{16\epsilon} \\
& - \frac{1971C_A C_F^2}{16\epsilon} + \frac{87C_A C_F}{2\epsilon} - \frac{6521C_F}{64\epsilon} + \frac{41C_A C_F^2}{8\epsilon^2} + \frac{107C_A C_F}{4\epsilon^2} - \frac{59C_F}{16\epsilon^2} \quad (C.17)
\end{aligned}$$

$$\begin{aligned}
V_6^{(m_2)} = & -\frac{89C_A Y_s L m_2 Y_f^3}{64m_2} + \frac{5C_A Y_s L m_2 Y_f^3}{16\epsilon m_2} + \frac{39C_A S_1 Y_s Y_f^3}{32m_2} - \frac{81C_A S_1 S_2 Y_s Y_f^3}{8m_2} - \frac{2413C_A Y_s Y_f^3}{512m_2} - \frac{123C_A Y_s Y_f^3}{128\epsilon m_2} - \frac{3C_A Y_s Y_f^3}{16\epsilon^2 m_2} \\
& - \frac{257C_A Y_s \zeta_2 Y_f^3}{128m_2} + \frac{37C_A Y_s \zeta_2 Y_f^3}{32\epsilon m_2} - \frac{3iC_A \pi Y_s \zeta_2 Y_f^3}{4m_2} + \frac{9C_A Y_s \zeta_3 Y_f^3}{4m_2} + \frac{603C_A Y_s L m_2 Y_f^3}{128m_2} - \frac{11C_A Y_s L m_2 Y_f^3}{32\epsilon m_2} - \frac{27C_A S_1 Y_s^2 Y_f^2}{32m_2^2} \\
& + \frac{351C_A S_1 S_2 Y_s^2 Y_f^2}{64m_2^2} + \frac{3277C_A Y_s^2 Y_f^2}{512m_2^2} - \frac{345C_A Y_s^2 Y_f^2}{128\epsilon m_2^2} + \frac{19C_A Y_s^2 L m_2 Y_f^2}{64m_2^2} - \frac{3C_A Y_s^2 L m_2 Y_f^2}{16\epsilon m_2^2} + \frac{89}{16} C_A C_F L m_2 Y_f^2 \\
& + \frac{69C_A M_H^2 Y_s L m_2 Y_f^2}{64m_2^2 M_W} - \frac{33C_A M_H^2 Y_s L m_2 Y_f^2}{64\epsilon m_2^2 M_W} - \frac{5C_A C_F L m_2 Y_f^2}{2\epsilon} + \frac{4201}{128} C_A C_F Y_f^2 - \frac{309}{16} C_A C_F S_1 Y_f^2 + \frac{513}{8} C_A C_F S_1 S_2 Y_f^2 \\
& + \frac{6663C_A M_H^2 Y_s Y_f^2}{512m_2^2 M_W} - \frac{3795C_A M_H^2 Y_s Y_f^2}{512\epsilon m_2^2 M_W} + \frac{129C_A Y_s^2 \zeta_2 Y_f^2}{128m_2^2} + \frac{21C_A Y_s^2 \zeta_2 Y_f^2}{32\epsilon m_2^2} + \frac{13iC_A \pi Y_s^2 \zeta_2 Y_f^2}{32m_2^2} + \frac{323}{32} C_A C_F \zeta_2 Y_f^2
\end{aligned}$$



$$\begin{aligned}
& - \frac{63C_A M_H^2 Y_s \zeta_2 Y_f^2}{128m_2^3 M_W} + \frac{63C_A M_H^2 Y_s \zeta_2 Y_f^2}{128\epsilon m_2^3 M_W} - \frac{37C_A C_F \zeta_2 Y_f^2}{4\epsilon} + \frac{19}{4} i C_A C_F \pi \zeta_2 Y_f^2 - \frac{39C_A Y_s^2 \zeta_3 Y_f^2}{32m_2^2} - \frac{57}{4} C_A C_F \zeta_3 Y_f^2 \\
& - \frac{319C_A Y_s^2 L_{m_2} Y_f^2}{128m_2^3} + \frac{39C_A Y_s^2 L_{m_2} Y_f^2}{32\epsilon m_2^3} - \frac{763}{32} C_A C_F L_{m_2} Y_f^2 - \frac{801C_A M_H^2 Y_s L_{m_2} Y_f^2}{128m_2^3 M_W} + \frac{429C_A M_H^2 Y_s L_{m_2} Y_f^2}{128\epsilon m_2^3 M_W} \\
& + \frac{17C_A C_F L_{m_2} Y_f^2}{4\epsilon} + \frac{87C_A C_F Y_f^2}{16\epsilon} + \frac{3C_A C_F Y_f^2}{4\epsilon^2} - \frac{83C_A Y_s^3 Y_f}{8m_2^3} + \frac{1273C_A Y_s^3 Y_f}{256\epsilon m_2^3} - \frac{C_A Y_s^3 Y_f}{32\epsilon^2 m_2^3} - \frac{3651C_A M_H^2 Y_s^2 Y_f}{128m_2^3 M_W} \\
& + \frac{3795C_A M_H^2 Y_s^2 Y_f}{256\epsilon m_2^3 M_W} - \frac{3C_A Y_s^3 L_{m_2}^2 Y_f}{4m_2^3} + \frac{11C_A Y_s^3 L_{m_2}^2 Y_f}{32\epsilon m_2^3} - \frac{15C_A M_H^2 Y_s^2 L_{m_2}^2 Y_f}{8m_2^3 M_W} + \frac{33C_A M_H^2 Y_s^2 L_{m_2}^2 Y_f}{32\epsilon m_2^3 M_W} \\
& + \frac{C_A C_F^2 Y_s L_{m_2}^2 Y_f}{m_2} + \frac{21C_A C_F Y_s L_{m_2}^2 Y_f}{2m_2} + \frac{C_F Y_s L_{m_2}^2 Y_f}{2m_2} - \frac{5C_A C_F Y_s L_{m_2}^2 Y_f}{2\epsilon m_2} + \frac{99C_A C_F S_1 Y_s Y_f}{8m_2} + \frac{189C_A C_F S_1 S_2 Y_s Y_f}{8m_2} \\
& + \frac{141C_A C_F^2 Y_s Y_f}{8m_2} + \frac{1595C_A C_F Y_s Y_f}{32m_2} + \frac{141C_F Y_s Y_f}{16m_2} + \frac{C_A C_F^2 Y_s Y_f}{2m_2} - \frac{193C_A C_F Y_s Y_f}{32\epsilon m_2} + \frac{C_F Y_s Y_f}{4\epsilon m_2} - \frac{C_A Y_s^3 \zeta_2 Y_f}{16m_2^3} \\
& - \frac{29C_A Y_s^3 \zeta_2 Y_f}{64\epsilon m_2^3} + \frac{3C_A M_H^2 Y_s^2 \zeta_2 Y_f}{8m_2^3 M_W} - \frac{87C_A M_H^2 Y_s^2 \zeta_2 Y_f}{64\epsilon m_2^3 M_W} + \frac{5C_A C_F^2 Y_s \zeta_2 Y_f}{2m_2} - \frac{11C_A C_F Y_s \zeta_2 Y_f}{4m_2} + \frac{5C_F Y_s \zeta_2 Y_f}{4m_2} \\
& - \frac{31C_A C_F Y_s \zeta_2 Y_f}{4\epsilon m_2} + \frac{7iC_A C_F \pi Y_s \zeta_2 Y_f}{4m_2} - \frac{21C_A C_F Y_s \zeta_3 Y_f}{4m_2} + \frac{73C_A Y_s^3 L_{m_2} Y_f}{16m_2^3} - \frac{139C_A Y_s^3 L_{m_2} Y_f}{64\epsilon m_2^3} + \frac{405C_A M_H^2 Y_s^2 L_{m_2} Y_f}{32m_2^3 M_W} \\
& - \frac{429C_A M_H^2 Y_s^2 L_{m_2} Y_f}{64\epsilon m_2^3 M_W} - \frac{15C_A C_F^2 Y_s L_{m_2} Y_f}{2m_2} - \frac{329C_A C_F Y_s L_{m_2} Y_f}{8m_2} - \frac{15C_F Y_s L_{m_2} Y_f}{4m_2} + \frac{67C_A C_F Y_s L_{m_2} Y_f}{8\epsilon m_2} \\
& + \frac{1413}{64} C_A C_F^2 + \frac{1403C_A C_F Y_s^2}{32m_2^2} - \frac{1305C_A C_F Y_s^2}{64\epsilon m_2^2} - \frac{C_A C_F Y_s^2}{8\epsilon^2 m_2^2} + C_A C_F^2 L_{m_2}^2 + \frac{11C_A C_F Y_s^2 L_{m_2}^2}{4m_2^2} - \frac{11C_A C_F Y_s^2 L_{m_2}^2}{8\epsilon m_2^2} \\
& + \frac{1}{4} C_A C_F L_{m_2}^2 + \frac{159}{8} C_F L_{m_2}^2 + 5C_A C_F n_f T_f L_{m_2}^2 - \frac{25C_A C_F M_W Y_s L_{m_2}^2}{2m_2^2} + \frac{5C_A C_F M_W Y_s L_{m_2}^2}{\epsilon m_2^2} - \frac{4C_A C_F^2 L_{m_2}^2}{\epsilon} \\
& - \frac{91C_F L_{m_2}^2}{8\epsilon} + \frac{465C_A C_F}{64} + \frac{20327C_F}{128} - 36C_F S_1 + 108C_F S_1 S_2 - \frac{7}{4} C_A C_F n_f T_f + \frac{6C_A C_F n_f T_f}{\epsilon} + \frac{7C_A C_F n_f T_f}{2\epsilon^2} \\
& - \frac{657C_A C_F M_W Y_s}{4m_2^2} + \frac{575C_A C_F M_W Y_s}{8\epsilon m_2^2} - \frac{143}{4} C_A C_F^2 \zeta_2 - \frac{5C_A C_F Y_s^2 \zeta_2}{8m_2^2} + \frac{29C_A C_F Y_s^2 \zeta_2}{16\epsilon m_2^2} + \frac{5}{8} C_A C_F \zeta_2 - \frac{107C_F \zeta_2}{16} \\
& + \frac{25}{2} C_A C_F n_f T_f \zeta_2 + \frac{33C_A C_F M_W Y_s \zeta_2}{4m_2^2} - \frac{7C_A C_F M_W Y_s \zeta_2}{2\epsilon m_2^2} + \frac{6C_A C_F^2 \zeta_2}{\epsilon} - \frac{315C_F \zeta_2}{16\epsilon} + 8iC_F \pi \zeta_2 - 24C_F \zeta_3 \\
& - \frac{65}{16} C_A C_F^2 L_{m_2} - \frac{151C_A C_F Y_s^2 L_{m_2}}{8m_2^2} + \frac{147C_A C_F Y_s^2 L_{m_2}}{16\epsilon m_2^2} - \frac{49}{16} C_A C_F L_{m_2} - \frac{2991}{32} C_F L_{m_2} + C_A C_F n_f T_f L_{m_2} \\
& - \frac{7C_A C_F n_f T_f L_{m_2}}{\epsilon} + \frac{77C_A C_F M_W Y_s L_{m_2}}{m_2^2} - \frac{65C_A C_F M_W Y_s L_{m_2}}{2\epsilon m_2^2} + \frac{263C_A C_F^2 L_{m_2}}{8\epsilon} + \frac{C_A C_F L_{m_2}}{8\epsilon} + \frac{97C_F L_{m_2}}{2\epsilon} \\
& - \frac{1309C_A C_F^2}{16\epsilon} + \frac{5C_A C_F}{16\epsilon} - \frac{4593C_F}{64\epsilon} - \frac{55C_A C_F^2}{16\epsilon^2} - \frac{C_A C_F}{16\epsilon^2} - \frac{55C_F}{32\epsilon^2} \tag{C.18}
\end{aligned}$$

$$\begin{aligned}
V_7^{(m_2)} = & \frac{21C_A Y_s L_{m_2}^2 Y_f^3}{64m_2} + \frac{5C_A Y_s L_{m_2}^2 Y_f^3}{64\epsilon m_2} - \frac{2741C_A Y_s Y_f^3}{512m_2} + \frac{527C_A Y_s Y_f^3}{512\epsilon m_2} + \frac{79C_A Y_s \zeta_2 Y_f^3}{128m_2} - \frac{11C_A Y_s \zeta_2 Y_f^3}{128\epsilon m_2} + \frac{299C_A Y_s L_{m_2} Y_f^3}{128m_2} \\
& - \frac{65C_A Y_s L_{m_2} Y_f^3}{128\epsilon m_2} - \frac{3C_A S_1 Y_s^2 Y_f^2}{32m_2^2} + \frac{27C_A S_1 S_2 Y_s^2 Y_f^2}{64m_2^2} - \frac{119C_A Y_s^2 Y_f^2}{512m_2^2} - \frac{C_A Y_s^2 L_{m_2}^2 Y_f^2}{64m_2^2} + \frac{1}{2} C_A C_F L_{m_2}^2 Y_f^2 \\
& + \frac{9C_A M_H^2 Y_s L_{m_2}^2 Y_f^2}{32m_2^3 M_W} - \frac{15C_A M_H^2 Y_s L_{m_2}^2 Y_f^2}{64\epsilon m_2^3 M_W} - \frac{C_A C_F L_{m_2}^2 Y_f^2}{4\epsilon} + \frac{41}{4} C_A C_F Y_f^2 + \frac{267C_A M_H^2 Y_s Y_f^2}{64m_2^3 M_W} - \frac{1725C_A M_H^2 Y_s Y_f^2}{512\epsilon m_2^3 M_W} \\
& + \frac{29C_A Y_s^2 \zeta_2 Y_f^2}{128m_2^2} + \frac{iC_A \pi Y_s^2 \zeta_2 Y_f^2}{32m_2^2} - \frac{9}{16} C_A C_F \zeta_2 Y_f^2 - \frac{3C_A M_H^2 Y_s \zeta_2 Y_f^2}{64m_2^3 M_W} + \frac{33C_A M_H^2 Y_s \zeta_2 Y_f^2}{128\epsilon m_2^3 M_W} + \frac{7C_A C_F \zeta_2 Y_f^2}{8\epsilon} - \frac{3C_A Y_s^2 \zeta_3 Y_f^2}{32m_2^2} \\
& + \frac{13C_A Y_s^2 L_{m_2} Y_f^2}{128m_2^2} - \frac{33}{8} C_A C_F L_{m_2} Y_f^2 - \frac{15C_A M_H^2 Y_s L_{m_2} Y_f^2}{8m_2^3 M_W} + \frac{195C_A M_H^2 Y_s L_{m_2} Y_f^2}{128\epsilon m_2^3 M_W} + \frac{13C_A C_F L_{m_2} Y_f^2}{8\epsilon} - \frac{113C_A C_F Y_f^2}{32\epsilon} \\
& + \frac{3C_A S_1 Y_s^3 Y_f}{16m_2^3} - \frac{27C_A S_1 S_2 Y_s^3 Y_f}{32m_2^3} - \frac{525C_A Y_s^3 Y_f}{1024m_2^3} - \frac{95C_A Y_s^3 Y_f}{1024\epsilon m_2^3} + \frac{699C_A M_H^2 Y_s^2 Y_f}{256m_2^3 M_W} - \frac{345C_A M_H^2 Y_s^2 Y_f}{256\epsilon m_2^3 M_W} - \frac{15C_A Y_s^3 L_{m_2}^2 Y_f}{128m_2^3} \\
& + \frac{5C_A Y_s^3 L_{m_2}^2 Y_f}{128\epsilon m_2^3} + \frac{9C_A M_H^2 Y_s^2 L_{m_2}^2 Y_f}{32m_2^3 M_W} - \frac{3C_A M_H^2 Y_s^2 L_{m_2}^2 Y_f}{32\epsilon m_2^3 M_W} + \frac{21C_A C_F Y_s L_{m_2}^2 Y_f}{8m_2} - \frac{5C_A C_F Y_s L_{m_2}^2 Y_f}{16\epsilon m_2} \\
& - \frac{21C_A C_F S_1 Y_s Y_f}{8m_2} + \frac{27C_A C_F S_1 S_2 Y_s Y_f}{16m_2} + \frac{13C_A C_F^2 Y_s Y_f}{4m_2} + \frac{1021C_A C_F Y_s Y_f}{64m_2} + \frac{13C_F Y_s Y_f}{8m_2} + \frac{C_A C_F^2 Y_s Y_f}{2\epsilon m_2} \\
& + \frac{1435C_A C_F Y_s Y_f}{128\epsilon m_2} + \frac{C_F Y_s Y_f}{4\epsilon m_2} - \frac{63C_A Y_s^3 \zeta_2 Y_f}{256m_2^3} + \frac{37C_A Y_s^3 \zeta_2 Y_f}{256\epsilon m_2^3} - \frac{iC_A \pi Y_s^3 \zeta_2 Y_f}{16m_2^3} - \frac{27C_A M_H^2 Y_s^2 \zeta_2 Y_f}{64m_2^3 M_W} - \frac{3C_A M_H^2 Y_s^2 \zeta_2 Y_f}{64\epsilon m_2^3 M_W} \\
& - \frac{29C_A C_F Y_s \zeta_2 Y_f}{8m_2} - \frac{181C_A C_F Y_s \zeta_2 Y_f}{32\epsilon m_2} + \frac{iC_A C_F \pi Y_s \zeta_2 Y_f}{8m_2} + \frac{3C_A Y_s^3 \zeta_3 Y_f}{16m_2^3} - \frac{3C_A C_F Y_s \zeta_3 Y_f}{8m_2} + \frac{131C_A Y_s^3 L_{m_2} Y_f}{256m_2^3} \\
& - \frac{23C_A Y_s^3 L_{m_2} Y_f}{256\epsilon m_2^3} - \frac{93C_A M_H^2 Y_s^2 L_{m_2} Y_f}{64m_2^3 M_W} + \frac{39C_A M_H^2 Y_s^2 L_{m_2} Y_f}{64\epsilon m_2^3 M_W} - \frac{C_A C_F^2 Y_s L_{m_2} Y_f}{m_2} - \frac{189C_A C_F Y_s L_{m_2} Y_f}{16m_2} \\
& - \frac{C_F Y_s L_{m_2} Y_f}{2m_2} - \frac{61C_A C_F Y_s L_{m_2} Y_f}{32\epsilon m_2} - \frac{3177}{32} C_A C_F^2 - \frac{3C_A C_F S_1 Y_s^2}{4m_2^2} + \frac{27C_A C_F S_1 S_2 Y_s^2}{8m_2^2} + \frac{907C_A C_F Y_s^2}{512m_2^2} + \frac{287C_A C_F Y_s^2}{512\epsilon m_2^2}
\end{aligned}$$

$$\begin{aligned}
& + \frac{C_A C_F Y_s^2}{4\epsilon^2 m_2^2} - 12 C_A C_F^2 \mathcal{L}_{m_2}^2 + \frac{37 C_A C_F Y_s^2 \mathcal{L}_{m_2}^2}{64 m_2^2} - \frac{5 C_A C_F Y_s^2 \mathcal{L}_{m_2}^2}{64 \epsilon m_2^2} + \frac{9}{8} C_A C_F \mathcal{L}_{m_2}^2 - \frac{157}{16} C_F \mathcal{L}_{m_2}^2 - \frac{1}{2} C_A C_F n_f T_f \mathcal{L}_{m_2}^2 \\
& + \frac{19 C_A C_F M_W Y_s \mathcal{L}_{m_2}^2}{8 m_2^2} - \frac{C_A C_F M_W Y_s \mathcal{L}_{m_2}^2}{\epsilon m_2^2} + \frac{3 C_A C_F^2 \mathcal{L}_{m_2}^2}{4 \epsilon} + \frac{49 C_F \mathcal{L}_{m_2}^2}{16 \epsilon} + \frac{255 C_A C_F}{64} - \frac{14029 C_F}{128} - \frac{3 C_F S_1}{8} + 27 C_F S_1 S_2 \\
& - \frac{309}{16} C_A C_F n_f T_f + \frac{C_A C_F n_f T_f}{\epsilon} + \frac{C_A C_F n_f T_f}{2\epsilon^2} + \frac{2655 C_A C_F M_W Y_s}{64 m_2^2} - \frac{115 C_A C_F M_W Y_s}{8 \epsilon m_2^2} + \frac{47}{4} C_A C_F^2 \zeta_2 + \frac{85 C_A C_F Y_s^2 \zeta_2}{128 m_2^2} \\
& - \frac{37 C_A C_F Y_s^2 \zeta_2}{128 \epsilon m_2^2} + \frac{i C_A C_F \pi Y_s^2 \zeta_2}{4 m_2^2} + \frac{13}{16} C_A C_F \zeta_2 + \frac{7 C_F \zeta_2}{32} - \frac{5}{4} C_A C_F n_f T_f \zeta_2 - \frac{25 C_A C_F M_W Y_s \zeta_2}{16 m_2^2} + \frac{2 C_A C_F M_W Y_s \zeta_2}{\epsilon m_2^2} \\
& - \frac{9 C_A C_F^2 \zeta_2}{8 \epsilon} - \frac{343 C_F \zeta_2}{32 \epsilon} + 2 i C_F \pi \zeta_2 - \frac{3 C_A C_F Y_s^2 \zeta_3}{4 m_2^2} - 6 C_F \zeta_3 + \frac{447}{8} C_A C_F^2 \mathcal{L}_{m_2} - \frac{181 C_A C_F Y_s^2 \mathcal{L}_{m_2}}{128 m_2^2} - \frac{41 C_A C_F Y_s^2 \mathcal{L}_{m_2}}{128 \epsilon m_2^2} \\
& - \frac{1}{16} C_A C_F \mathcal{L}_{m_2} + \frac{1747}{32} C_F \mathcal{L}_{m_2} + \frac{31}{4} C_A C_F n_f T_f \mathcal{L}_{m_2} - \frac{C_A C_F n_f T_f \mathcal{L}_{m_2}}{\epsilon} - \frac{281 C_A C_F M_W Y_s \mathcal{L}_{m_2}}{16 m_2^2} + \frac{13 C_A C_F M_W Y_s \mathcal{L}_{m_2}}{2 \epsilon m_2^2} \\
& + \frac{33 C_A C_F^2 \mathcal{L}_{m_2}}{8 \epsilon} - \frac{C_A C_F \mathcal{L}_{m_2}}{\epsilon} - \frac{619 C_F \mathcal{L}_{m_2}}{32 \epsilon} - \frac{299 C_A C_F^2}{32 \epsilon} - \frac{3 C_A C_F}{8 \epsilon} + \frac{6197 C_F}{128 \epsilon} - \frac{9 C_A C_F^2}{2 \epsilon^2} + \frac{C_A C_F}{2 \epsilon^2} - \frac{9 C_F}{4 \epsilon^2} \quad (C.19)
\end{aligned}$$

#### C.4 Matching at $\mu \sim m_1$ :

The contributions from two-loop vertex corrections,  $V_i^{(m_1)}$ , with unevaluated MIs are too large to present here. We thus include the full expressions with description in an ancillary file.

#### C.5 Matching at $\mu \sim M$ :

The contributions from two-loop vertex corrections,  $V_i^{(M)}$ , with unevaluated MIs are too large to present here. We thus include the full expressions with description in an ancillary file.

## D Wave-function Corrections

### D.1 Full theory field at $m = 0$ and $M = 0$ :

$$\begin{aligned}
F_\psi^{(0,0)} &= \frac{2\Delta_M Y_f^2 C_A C_F}{M_W \epsilon_{UV}} - \frac{\Delta_M Y_f^4 C_A}{4M_W \epsilon_{UV}} + \frac{n_f C_A C_F T_f}{\epsilon_{IR} \epsilon_{UV}} - \frac{n_f C_A C_F T_f}{2\epsilon_{UV}} - \frac{n_f C_A C_F T_f}{2\epsilon_{UV}^2} - \frac{8C_A C_F^2}{\epsilon_{IR} \epsilon_{UV}} + \frac{Y_f^2 C_A C_F}{2\epsilon_{IR} \epsilon_{UV}} + \frac{4C_A C_F^2}{\epsilon_{UV}^2} \\
& - \frac{Y_f^2 C_A C_F}{4\epsilon_{UV}} - \frac{Y_f^2 C_A C_F}{4\epsilon_{UV}^2} - \frac{Y_f^4 C_A}{8\epsilon_{IR} \epsilon_{UV}} + \frac{5Y_f^4 C_A}{32\epsilon_{UV}} + \frac{Y_f^4 C_A}{16\epsilon_{UV}^2} - \frac{4C_F}{\epsilon_{IR} \epsilon_{UV}} + \frac{2C_F}{\epsilon_{UV}^2} \quad (D.1)
\end{aligned}$$

$$\begin{aligned}
F_X^{(0,0)} &= -\frac{10\Delta_M Y_s^2 C_A C_F}{M_W^3 \epsilon_{IR} \epsilon_{UV}} - \frac{14\Delta_M Y_s C_A C_F}{M_W^2 \epsilon_{IR} \epsilon_{UV}} + \frac{153\Delta_M Y_s^2 C_A C_F^2}{16M_W^3 \epsilon_{UV}} + \frac{37\Delta_M Y_s^2 C_A C_F}{4M_W^3 \epsilon_{UV}} + \frac{5\Delta_M Y_s^2 C_A C_F}{M_W^3 \epsilon_{UV}^2} + \frac{11\Delta_M Y_s C_A C_F}{M_W^2 \epsilon_{UV}} + \\
& \frac{7\Delta_M Y_s C_A C_F}{M_W^2 \epsilon_{UV}^2} + \frac{\Delta_M C_A C_F}{M_W \epsilon_{UV}} - \frac{5\Delta_M Y_s^4 C_A}{4M_W^5 \epsilon_{IR} \epsilon_{UV}} + \frac{7\Delta_M Y_s^4 C_A}{4M_W^5 \epsilon_{UV}} + \frac{5\Delta_M Y_s^4 C_A}{8M_W^5 \epsilon_{UV}^2} - \frac{3\Delta_M Y_s^3 C_A}{4M_W^4 \epsilon_{UV}} + \frac{9\Delta_M Y_s^2 C_A}{16M_W^3 \epsilon_{UV}} + \frac{2n_f C_A C_F T_f}{\epsilon_{IR} \epsilon_{UV}} - \\
& \frac{n_f C_A C_F T_f}{\epsilon_{UV}} - \frac{n_f C_A C_F T_f}{\epsilon_{UV}^2} + \frac{3Y_s^2 C_A C_F}{M_W^2 \epsilon_{IR} \epsilon_{UV}} + \frac{7Y_s C_A C_F}{M_W \epsilon_{IR} \epsilon_{UV}} - \frac{19C_A C_F^2}{2\epsilon_{IR} \epsilon_{UV}} - \frac{7C_A C_F}{2\epsilon_{IR} \epsilon_{UV}} - \frac{65Y_s^2 C_A C_F^2}{32M_W^2 \epsilon_{UV}} - \frac{2Y_s^2 C_A C_F}{M_W^2 \epsilon_{UV}} - \\
& \frac{3Y_s^2 C_A C_F}{2M_W^2 \epsilon_{UV}^2} - \frac{2Y_s C_A C_F}{M_W \epsilon_{UV}} - \frac{7Y_s C_A C_F}{2M_W \epsilon_{UV}^2} + \frac{49C_A C_F^2}{2\epsilon_{UV}} - \frac{5C_A C_F}{2\epsilon_{UV}} + \frac{19C_A C_F^2}{4\epsilon_{UV}^2} + \frac{7C_A C_F}{4\epsilon_{UV}^2} + \frac{5Y_s^4 C_A}{16M_W^4 \epsilon_{IR} \epsilon_{UV}} - \frac{9Y_s^4 C_A}{32M_W^4 \epsilon_{UV}} - \\
& \frac{5Y_s^4 C_A}{32M_W^4 \epsilon_{UV}^2} + \frac{3Y_s^3 C_A}{8M_W^3 \epsilon_{UV}} - \frac{9Y_s^2 C_A}{32M_W^2 \epsilon_{UV}} + \frac{153\Delta_M Y_s^2 C_F}{32M_W^3 \epsilon_{UV}} - \frac{47C_F}{4\epsilon_{IR} \epsilon_{UV}} - \frac{65Y_s^2 C_F}{64M_W^2 \epsilon_{UV}} + \frac{73C_F}{4\epsilon_{UV}} + \frac{47C_F}{8\epsilon_{UV}^2} \quad (D.2)
\end{aligned}$$

## D.2 Full theory field at $m = 0$ and $M \neq 0$ ( $\Delta_M \equiv M_H - M_W$ ):

$$\begin{aligned}
F_\psi^{(0,M)} = & -\frac{4\Delta_M Y_f^2 C_A C_F \mathcal{L} M_W}{M_W} - \frac{9\Delta_M S_2 C_A C_F}{M_W} + \frac{2\Delta_M Y_f^2 C_A C_F}{M_W \epsilon_{UV}} - \frac{3\Delta_M Y_f^2 \zeta_2 C_A C_F}{M_W} + \frac{4\Delta_M Y_f^2 C_A C_F}{M_W} + \frac{2\Delta_M C_A C_F}{M_W} + \\
& \frac{\Delta_M Y_f^4 C_A \mathcal{L} M_W}{2M_W} - \frac{\Delta_M Y_f^4 C_A}{4M_W \epsilon_{UV}} - \frac{5\Delta_M Y_f^4 C_A}{8M_W} + \frac{n_f C_A C_F T_f \mathcal{L} M_W}{\epsilon_{UV}} - n_f C_A C_F T_f \mathcal{L}_M^2 M_W + n_f C_A C_F T_f \mathcal{L} M_W - \\
& \frac{n_f C_A C_F T_f}{2\epsilon_{UV}} - \frac{n_f C_A C_F T_f}{2\epsilon_{UV}^2} - \frac{3}{2} n_f \zeta_2 C_A C_F T_f - \frac{1}{2} n_f C_A C_F T_f - \frac{8C_A C_F^2 \mathcal{L} M_W}{\epsilon_{UV}} + \frac{Y_f^2 C_A C_F \mathcal{L} M_W}{2\epsilon_{UV}} - \frac{1}{2} Y_f^2 C_A C_F \mathcal{L}_M^2 M_W + \\
& \frac{1}{2} Y_f^2 C_A C_F \mathcal{L} M_W + 8C_A C_F^2 \mathcal{L}_M^2 M_W + 18S_2 C_A C_F^2 + \frac{4C_A C_F^2}{\epsilon_{UV}^2} - \frac{Y_f^2 C_A C_F}{4\epsilon_{UV}} - \frac{Y_f^2 C_A C_F}{4\epsilon_{UV}^2} - \\
& \frac{3}{4} Y_f^2 \zeta_2 C_A C_F - Y_f^2 C_A C_F + 8\zeta_2 C_A C_F^2 - 12C_A C_F^2 - \frac{Y_f^4 C_A \mathcal{L} M_W}{8\epsilon_{UV}} + \frac{1}{8} Y_f^4 C_A \mathcal{L}_M^2 M_W - \frac{5}{16} Y_f^4 C_A \mathcal{L} M_W + \frac{5Y_f^4 C_A}{32\epsilon_{UV}} + \frac{Y_f^4 C_A}{16\epsilon_{UV}^2} + \\
& \frac{Y_f^4 C_A}{4} - \frac{4C_F \mathcal{L} M_W}{\epsilon_{UV}} + 4C_F \mathcal{L}_M^2 M_W + 9S_2 C_F + \frac{2C_F}{\epsilon_{UV}^2} + 4\zeta_2 C_F - 4C_F
\end{aligned} \tag{D.3}$$

$$\begin{aligned}
F_X^{(0,M)} = & -\frac{5C_A \mathcal{L}_M^2 M_W Y_s^4}{16M_W^4} + \frac{5C_A \Delta_M \mathcal{L}_M^2 M_W Y_s^4}{4M_W^5} - \frac{3C_A \zeta_2 Y_s^4}{32M_W^4} + \frac{3C_A \Delta_M \zeta_2 Y_s^4}{8M_W^5} + \frac{9C_A \mathcal{L} M_W Y_s^4}{16M_W^4} + \frac{5C_A \mathcal{L} M_W Y_s^4}{16\epsilon_{UV} M_W^4} - \frac{7C_A \Delta_M \mathcal{L} M_W Y_s^4}{2M_W^5} - \\
& \frac{5C_A \Delta_M \mathcal{L} M_W Y_s^4}{4\epsilon_{UV} M_W^5} - \frac{17C_A Y_s^4}{32M_W^4} - \frac{9C_A Y_s^4}{32\epsilon_{UV} M_W^4} - \frac{5C_A Y_s^4}{32\epsilon_{UV}^2 M_W^4} + \frac{13C_A \Delta_M Y_s^4}{4M_W^5} + \frac{7C_A \Delta_M Y_s^4}{4\epsilon_{UV} M_W^5} + \frac{5C_A \Delta_M Y_s^4}{8\epsilon_{UV}^2 M_W^5} + \frac{81C_A S_2 Y_s^3}{16M_W^3} - \\
& \frac{81C_A \Delta_M S_2 Y_s^3}{8M_W^4} - \frac{3C_A \mathcal{L} M_W Y_s^3}{4M_W^3} + \frac{3C_A \Delta_M \mathcal{L} M_W Y_s^3}{2M_W^4} - \frac{3C_A Y_s^3}{8M_W^3} + \frac{3C_A Y_s^3}{8\epsilon_{UV} M_W^3} - \frac{3C_A \Delta_M Y_s^3}{4M_W^4} - \frac{3C_A \Delta_M Y_s^3}{4\epsilon_{UV} M_W^4} - \\
& \frac{3C_A C_F \mathcal{L}_M^2 M_W Y_s^2}{M_W^2} + \frac{10C_A C_F \Delta_M \mathcal{L}_M^2 M_W Y_s^2}{M_W^3} + \frac{1005C_A C_F^2 S_2 Y_s^2}{64M_W^2} + \frac{189C_A S_2 Y_s^2}{64M_W^2} + \frac{1005C_F S_2 Y_s^2}{128M_W^2} - \frac{417C_A C_F^2 \Delta_M S_2 Y_s^2}{4M_W^3} - \\
& \frac{27C_A \Delta_M S_2 Y_s^2}{16M_W^3} - \frac{417C_F \Delta_M S_2 Y_s^2}{8M_W^3} - \frac{5C_A C_F \zeta_2 Y_s^2}{4M_W^2} + \frac{11C_A C_F \Delta_M \zeta_2 Y_s^2}{2M_W^3} + \frac{65C_A C_F^2 \mathcal{L} M_W Y_s^2}{16M_W^2} + \frac{9C_A \mathcal{L} M_W Y_s^2}{16M_W^2} + \\
& \frac{4C_A C_F \mathcal{L} M_W Y_s^2}{M_W^2} + \frac{65C_F \mathcal{L} M_W Y_s^2}{32M_W^2} + \frac{3C_A C_F \mathcal{L} M_W Y_s^2}{\epsilon_{UV} M_W^2} - \frac{153C_A C_F^2 \Delta_M \mathcal{L} M_W Y_s^2}{8M_W^3} - \frac{9C_A \Delta_M \mathcal{L} M_W Y_s^2}{8M_W^3} - \\
& \frac{37C_A C_F \Delta_M \mathcal{L} M_W Y_s^2}{2M_W^3} - \frac{153C_F \Delta_M \mathcal{L} M_W Y_s^2}{16M_W^3} - \frac{10C_A C_F \Delta_M \mathcal{L} M_W Y_s^2}{\epsilon_{UV} M_W^3} - \frac{213C_A C_F^2 Y_s^2}{32M_W^2} - \frac{45C_A Y_s^2}{32M_W^2} - \frac{3C_A C_F Y_s^2}{M_W^2} - \\
& \frac{213C_F Y_s^2}{64M_W^2} - \frac{65C_A C_F^2 Y_s^2}{32\epsilon_{UV} M_W^2} - \frac{9C_A Y_s^2}{32\epsilon_{UV} M_W^2} - \frac{2C_A C_F Y_s^2}{\epsilon_{UV} M_W^2} - \frac{65C_F Y_s^2}{64\epsilon_{UV} M_W^2} - \frac{3C_A C_F Y_s^2}{2\epsilon_{UV}^2 M_W^2} + \frac{645C_A C_F^2 \Delta_M Y_s^2}{16M_W^3} + \frac{45C_A \Delta_M Y_s^2}{16M_W^3} + \\
& \frac{45C_A C_F \Delta_M Y_s^2}{4M_W^3} + \frac{645C_F \Delta_M Y_s^2}{32M_W^3} + \frac{153C_A C_F^2 \Delta_M Y_s^2}{16\epsilon_{UV} M_W^3} + \frac{9C_A \Delta_M Y_s^2}{16\epsilon_{UV} M_W^3} + \frac{37C_A C_F \Delta_M Y_s^2}{4\epsilon_{UV} M_W^3} + \frac{153C_F \Delta_M Y_s^2}{32\epsilon_{UV} M_W^3} + \\
& \frac{5C_A C_F \Delta_M Y_s^2}{\epsilon_{UV}^2 M_W^3} - \frac{7C_A C_F \mathcal{L}_M^2 M_W Y_s}{M_W} + \frac{14C_A C_F \Delta_M \mathcal{L}_M^2 M_W Y_s}{M_W^2} + \frac{261C_A C_F S_2 Y_s}{4M_W} - \frac{192C_A C_F \Delta_M S_2 Y_s}{M_W^2} - \frac{4C_A C_F \zeta_2 Y_s}{M_W} + \\
& \frac{9C_A C_F \Delta_M \zeta_2 Y_s}{M_W^2} + \frac{4C_A C_F \mathcal{L} M_W Y_s}{M_W} + \frac{7C_A C_F \mathcal{L} M_W Y_s}{\epsilon_{UV} M_W} - \frac{22C_A C_F \Delta_M \mathcal{L} M_W Y_s}{M_W^2} - \frac{14C_A C_F \Delta_M \mathcal{L} M_W Y_s}{\epsilon_{UV} M_W^2} - \frac{35C_A C_F Y_s}{2M_W} - \\
& \frac{2C_A C_F Y_s}{\epsilon_{UV} M_W} - \frac{7C_A C_F Y_s}{2\epsilon_{UV}^2 M_W} + \frac{54C_A C_F \Delta_M Y_s}{M_W^2} + \frac{11C_A C_F \Delta_M Y_s}{\epsilon_{UV} M_W^2} + \frac{7C_A C_F \Delta_M Y_s}{\epsilon_{UV}^2 M_W^2} + \frac{81}{2} C_A C_F^2 + \frac{19}{2} C_A C_F^2 \mathcal{L}_M^2 M_W + \\
& \frac{7}{2} C_A C_F \mathcal{L}_M^2 M_W + \frac{47}{4} C_F \mathcal{L}_M^2 M_W - 2C_A C_F n_f T_f \mathcal{L}_M^2 M_W + \frac{19C_A C_F}{2} + \frac{97C_F}{4} - 141C_A C_F^2 S_2 - 6C_A C_F S_2 - \frac{141C_F S_2}{2} - \\
& \frac{21C_A C_F \Delta_M S_2}{M_W} - C_A C_F n_f T_f - \frac{C_A C_F n_f T_f}{\epsilon_{UV}} - \frac{C_A C_F n_f T_f}{\epsilon_{UV}^2} - \frac{17}{4} C_A C_F^2 \zeta_2 + \frac{7}{4} C_A C_F \zeta_2 + \\
& \frac{39C_F \zeta_2}{8} - 3C_A C_F n_f T_f \zeta_2 - 49C_A C_F^2 \mathcal{L} M_W + 5C_A C_F \mathcal{L} M_W - \frac{73}{2} C_F \mathcal{L} M_W + 2C_A C_F n_f T_f \mathcal{L} M_W + \frac{2C_A C_F n_f T_f \mathcal{L} M_W}{\epsilon_{UV}} - \\
& \frac{19C_A C_F^2 \mathcal{L} M_W}{2\epsilon_{UV}} - \frac{7C_A C_F \mathcal{L} M_W}{2\epsilon_{UV}} - \frac{47C_F \mathcal{L} M_W}{4\epsilon_{UV}} - \frac{2C_A C_F \Delta_M \mathcal{L} M_W}{M_W} + \frac{49C_A C_F^2}{2\epsilon_{UV}} - \frac{5C_A C_F}{2\epsilon_{UV}} + \frac{73C_F}{4\epsilon_{UV}} + \frac{4C_A C_F \Delta_M}{M_W} + \\
& \frac{C_A C_F \Delta_M}{\epsilon_{UV} M_W} + \frac{19C_A C_F^2}{4\epsilon_{UV}^2} + \frac{7C_A C_F}{4\epsilon_{UV}^2} + \frac{47C_F}{8\epsilon_{UV}^2}
\end{aligned} \tag{D.4}$$

## D.3 Full theory field at $m \neq 0$ and $M = 0$ :

$$\begin{aligned}
F_\psi^{(m,0)} = & -\frac{47}{16} C_A \mathcal{L}_m^2 Y_f^4 + \frac{5C_A \mathcal{L}_m^2 Y_f^4}{16\epsilon_{IR}} - \frac{2889C_A Y_f^4}{256} - \frac{47}{16} C_A \zeta_2 Y_f^4 + \frac{37C_A \zeta_2 Y_f^4}{32\epsilon_{IR}} + \frac{543}{64} C_A \mathcal{L}_m Y_f^4 + \frac{77C_A \mathcal{L}_m Y_f^4}{64\epsilon_{IR}} - \frac{769C_A Y_f^4}{256\epsilon_{IR}} \\
& + \frac{5C_A Y_f^4}{32\epsilon_{UV}} - \frac{C_A Y_f^4}{8\epsilon_{IR} \epsilon_{UV}} - \frac{115C_A Y_f^4}{128\epsilon_{IR}^2} + \frac{C_A Y_f^4}{16\epsilon_{UV}^2} - \frac{9131}{256} C_A^2 Y_f^2 - \frac{363}{32} C_A C_F^2 Y_f^2 - \frac{363}{32} C_A^2 \mathcal{L}_m^2 Y_f^2 - \frac{3}{4} C_A C_F^2 \mathcal{L}_m^2 Y_f^2
\end{aligned}$$

$$\begin{aligned}
& + \frac{3}{4} C_A^2 C_F L_m^2 Y_f^2 + \frac{363}{8} C_A C_F L_m^2 Y_f^2 - \frac{3}{8} C_F L_m^2 Y_f^2 - \frac{5 C_A C_F L_m^2 Y_f^2}{4 \epsilon_{\text{IR}}} + \frac{363}{32} L_m^2 Y_f^2 + \frac{363}{32} C_A^2 C_F Y_f^2 + \frac{9131}{64} C_A C_F Y_f^2 \\
& - \frac{363 C_F Y_f^2}{64} - \frac{699}{64} C_A^2 \zeta_2 Y_f^2 - \frac{15}{8} C_A C_F \zeta_2 Y_f^2 + \frac{15}{8} C_A^2 C_F \zeta_2 Y_f^2 + \frac{699}{16} C_A C_F \zeta_2 Y_f^2 - \frac{15}{16} C_F \zeta_2 Y_f^2 - \frac{37 C_A C_F \zeta_2 Y_f^2}{8 \epsilon_{\text{IR}}} \\
& + \frac{699 \zeta_2 Y_f^2}{64} + \frac{1501}{64} C_A^2 L_m Y_f^2 + \frac{41}{8} C_A C_F^2 L_m Y_f^2 - \frac{41}{8} C_A^2 C_F L_m Y_f^2 - \frac{1501}{16} C_A C_F L_m Y_f^2 + \frac{41}{16} C_F L_m Y_f^2 - \frac{3 C_A C_F^2 L_m Y_f^2}{4 \epsilon_{\text{IR}}} \\
& - \frac{53 C_A C_F L_m Y_f^2}{4 \epsilon_{\text{IR}}} - \frac{3 C_F L_m Y_f^2}{8 \epsilon_{\text{IR}}} - \frac{1501}{64} L_m Y_f^2 + \frac{41 C_A C_F^2 Y_f^2}{16 \epsilon_{\text{IR}}} + \frac{61 C_A C_F Y_f^2}{4 \epsilon_{\text{IR}}} + \frac{41 C_F Y_f^2}{32 \epsilon_{\text{IR}}} - \frac{C_A C_F Y_f^2}{4 \epsilon_{\text{UV}}} + \frac{C_A C_F Y_f^2}{2 \epsilon_{\text{IR}} \epsilon_{\text{UV}}} \\
& + \frac{3 C_A C_F^2 Y_f^2}{8 \epsilon_{\text{IR}}} + \frac{125 C_A C_F Y_f^2}{16 \epsilon_{\text{IR}}} + \frac{3 C_F Y_f^2}{16 \epsilon_{\text{IR}}} - \frac{C_A C_F Y_f^2}{4 \epsilon_{\text{UV}}} + \frac{9131 Y_f^2}{256} - \frac{81 C_A^2}{64} + \frac{8835}{32} C_A C_F^2 - \frac{1}{8} C_A^2 L_m^2 + \frac{267}{4} C_A C_F^2 L_m^2 \\
& - \frac{267}{4} C_A^2 C_F L_m^2 + \frac{1}{2} C_A C_F L_m^2 + \frac{375}{8} C_F L_m^2 - \frac{1}{2} C_A^2 n_f T_f L_m^2 + 2 C_A C_F n_f T_f L_m^2 + \frac{1}{2} n_f T_f L_m^2 + \frac{L_m^2}{8} - \frac{8835}{32} C_A^2 C_F \\
& + \frac{81 C_A C_F}{16} + \frac{10615 C_F}{64} - \frac{59}{16} C_A^2 n_f T_f + \frac{59}{4} C_A C_F n_f T_f + \frac{11 C_A C_F n_f T_f}{4 \epsilon_{\text{IR}}} - \frac{C_A C_F n_f T_f}{2 \epsilon_{\text{UV}}} + \frac{C_A C_F n_f T_f}{\epsilon_{\text{IR}} \epsilon_{\text{UV}}} - \frac{C_A C_F n_f T_f}{2 \epsilon_{\text{UV}}^2} \\
& + \frac{59 n_f T_f}{16} - \frac{5 C_A^2 \zeta_2}{16} + \frac{407}{8} C_A C_F^2 \zeta_2 - \frac{407}{8} C_A^2 C_F \zeta_2 + \frac{5}{4} C_A C_F \zeta_2 + \frac{659 C_F \zeta_2}{16} - \frac{5}{4} C_A^2 n_f T_f \zeta_2 + 5 C_A C_F n_f T_f \zeta_2 + \frac{5}{4} n_f T_f \zeta_2 \\
& + \frac{5 \zeta_2}{16} + \frac{11}{16} C_A^2 L_m - \frac{1169}{8} C_A C_F^2 L_m + \frac{1169}{8} C_A^2 C_F L_m - \frac{11}{4} C_A C_F L_m - \frac{1229}{16} C_F L_m + \frac{9}{4} C_A^2 n_f T_f L_m - 9 C_A C_F n_f T_f L_m \\
& - \frac{C_A C_F n_f T_f L_m}{\epsilon_{\text{IR}}} - \frac{9}{4} n_f T_f L_m + \frac{267 C_A C_F^2 L_m}{4 \epsilon_{\text{IR}}} - \frac{C_A C_F L_m}{4 \epsilon_{\text{IR}}} + \frac{159 C_F L_m}{8 \epsilon_{\text{IR}}} - \frac{11 L_m}{16} - \frac{1169 C_A C_F^2}{16 \epsilon_{\text{IR}}} + \frac{11 C_A C_F}{16 \epsilon_{\text{IR}}} - \frac{1109 C_F}{32 \epsilon_{\text{IR}}} \\
& - \frac{8 C_A C_F^2}{\epsilon_{\text{IR}} \epsilon_{\text{UV}}} - \frac{4 C_F}{\epsilon_{\text{IR}} \epsilon_{\text{UV}}} - \frac{235 C_A C_F^2}{8 \epsilon_{\text{IR}}^2} + \frac{C_A C_F}{8 \epsilon_{\text{IR}}^2} - \frac{127 C_F}{16 \epsilon_{\text{IR}}^2} + \frac{4 C_A C_F^2}{\epsilon_{\text{UV}}^2} + \frac{2 C_F}{\epsilon_{\text{UV}}^2} + \frac{81}{64} \tag{D.5}
\end{aligned}$$

$$\begin{aligned}
F_X(m,0) = & \frac{15 Y_s^2 C_A^5}{8 m^2} + \frac{Y_s^2 L_m^2 C_A^5}{8 m^2} - \frac{27}{8} L_m^2 C_A^5 - \frac{27 \zeta_2 C_A^5}{16} - \frac{13 Y_s^2 L_m C_A^5}{16 m^2} + \frac{191}{16} L_m C_A^5 - \frac{883 C_A^5}{32} - \frac{15 C_F Y_s^2 C_A^4}{4 m^2} + \frac{8121 Y_s^2 C_A^4}{1024 m^2} - \\
& \frac{C_F Y_s^2 L_m^2 C_A^4}{4 m^2} + \frac{239 Y_s^2 L_m^2 C_A^4}{128 m^2} + \frac{27}{4} C_F L_m^2 C_A^4 + 6 L_m^2 C_A^4 + \frac{883}{16} C_F C_A^4 + \frac{3}{2} n_f T_f C_A^4 + \frac{151 Y_s^2 \zeta_2 C_A^4}{256 m^2} + \frac{27}{8} C_F \zeta_2 C_A^4 + 3 \zeta_2 C_A^4 + \\
& \frac{13 C_F Y_s^2 L_m C_A^4}{8 m^2} - \frac{1111 Y_s^2 L_m C_A^4}{256 m^2} - \frac{191}{8} C_F L_m C_A^4 - \frac{1}{2} n_f T_f L_m C_A^4 - \frac{1}{8} L_m C_A^4 + \frac{175 C_A^4}{16} - \frac{3341 Y_s^4 C_A^3}{8192 m^4} - \frac{8121 C_F Y_s^2 C_A^3}{512 m^2} - \\
& \frac{15 Y_s^2 C_A^3}{8 m^2} - \frac{249 Y_s^4 L_m^2 C_A^3}{1024 m^4} - \frac{239 C_F Y_s^2 L_m^2 C_A^3}{64 m^2} - \frac{Y_s^2 L_m^2 C_A^3}{8 m^2} - 12 C_F L_m^2 C_A^3 + \frac{21}{8} L_m^2 C_A^3 - \frac{175}{8} C_F C_A^3 - 3 C_F n_f T_f C_A^3 - \\
& \frac{121 Y_s^4 \zeta_2 C_A^3}{2048 m^4} - \frac{151 C_F Y_s^2 \zeta_2 C_A^3}{128 m^2} - 6 C_F \zeta_2 C_A^3 + \frac{33 \zeta_2 C_A^3}{16} + \frac{779 Y_s^4 L_m C_A^3}{2048 m^4} + \frac{1111 C_F Y_s^2 L_m C_A^3}{128 m^2} + \frac{13 Y_s^2 L_m C_A^3}{16 m^2} + \\
& \frac{1}{4} C_F L_m C_A^3 + C_F n_f T_f L_m C_A^3 - \frac{47}{16} L_m C_A^3 + \frac{463 C_A^3}{32} + \frac{3341 C_F Y_s^4 C_A^2}{4096 m^4} + \frac{15 C_F Y_s^2 C_A^2}{4 m^2} - \frac{8121 Y_s^2 C_A^2}{1024 m^2} + \frac{249 C_F Y_s^4 L_m^2 C_A^2}{512 m^4} + \\
& \frac{C_F Y_s^2 L_m^2 C_A^2}{4 m^2} - \frac{239 Y_s^2 L_m^2 C_A^2}{128 m^2} - \frac{21}{4} C_F L_m^2 C_A^2 - 6 L_m^2 C_A^2 - \frac{463}{16} C_F C_A^2 - \frac{3}{2} n_f T_f C_A^2 + \frac{121 C_F Y_s^4 \zeta_2 C_A^2}{1024 m^4} - \frac{151 Y_s^2 \zeta_2 C_A^2}{256 m^2} - \\
& \frac{33}{8} C_F \zeta_2 C_A^2 - 3 \zeta_2 C_A^2 - \frac{779 C_F Y_s^4 L_m C_A^2}{1024 m^4} - \frac{13 C_F Y_s^2 L_m C_A^2}{8 m^2} + \frac{1111 Y_s^2 L_m C_A^2}{256 m^2} + \frac{47}{8} C_F L_m C_A^2 + \frac{1}{2} n_f T_f L_m C_A^2 + \frac{1}{8} L_m C_A^2 - \\
& \frac{175 C_A^2}{16} + \frac{241 Y_s^4 C_A}{2048 \epsilon_{\text{IR}} m^4} - \frac{Y_s^4 C_A}{64 \epsilon_{\text{IR}}^2 m^4} + \frac{Y_s^4 C_A}{64 \epsilon_{\text{IR}}^3 m^4} + \frac{9 Y_s^4 C_A}{32 \epsilon_{\text{IR}} M_W^4} + \frac{5 Y_s^4 C_A}{16 \epsilon_{\text{IR}} \epsilon_{\text{UV}} M_W^4} - \frac{9 Y_s^4 C_A}{32 \epsilon_{\text{UV}} M_W^4} - \frac{5 Y_s^4 C_A}{32 \epsilon_{\text{IR}}^2 M_W^4} - \frac{5 Y_s^4 C_A}{32 \epsilon_{\text{UV}}^2 M_W^4} - \\
& \frac{3 Y_s^3 C_A}{8 \epsilon_{\text{IR}} M_W^3} + \frac{3 Y_s^3 C_A}{8 \epsilon_{\text{UV}} M_W^3} + \frac{8121 C_F Y_s^2 C_A}{512 m^2} + \frac{13 C_F^2 Y_s^2 C_A}{8 \epsilon_{\text{IR}} m^2} + \frac{733 C_F Y_s^2 C_A}{512 \epsilon_{\text{IR}} m^2} + \frac{C_F^2 Y_s^2 C_A}{4 \epsilon_{\text{IR}}^2 m^2} + \frac{15 C_F Y_s^2 C_A}{16 \epsilon_{\text{IR}}^2 m^2} - \frac{3 C_F Y_s^2 C_A}{16 \epsilon_{\text{IR}}^3 m^2} + \\
& \frac{65 C_F^2 Y_s^2 C_A}{32 \epsilon_{\text{IR}} M_W^2} + \frac{2 C_F Y_s^2 C_A}{\epsilon_{\text{IR}} M_W^2} + \frac{9 Y_s^2 C_A}{32 \epsilon_{\text{IR}} M_W^2} - \frac{65 C_F^2 Y_s^2 C_A}{32 \epsilon_{\text{UV}} M_W^2} - \frac{2 C_F Y_s^2 C_A}{\epsilon_{\text{UV}} M_W^2} + \frac{3 C_F Y_s^2 C_A}{\epsilon_{\text{IR}} \epsilon_{\text{UV}} M_W^2} - \frac{9 Y_s^2 C_A}{32 \epsilon_{\text{UV}} M_W^2} - \frac{3 C_F Y_s^2 C_A}{2 \epsilon_{\text{IR}}^2 M_W^2} - \\
& \frac{3 C_F Y_s^2 C_A}{2 \epsilon_{\text{UV}}^2 M_W^2} + \frac{13 Y_s^4 L_m^2 C_A}{256 \epsilon_{\text{IR}} m^4} + \frac{239 C_F Y_s^2 L_m^2 C_A}{64 m^2} - \frac{39 C_F Y_s^2 L_m^2 C_A}{64 \epsilon_{\text{IR}} m^2} + 12 C_F L_m^2 C_A + \frac{3}{4} L_m^2 C_A + \frac{175 C_F C_A}{8} + 3 C_F n_f T_f C_A + \\
& \frac{3 C_F n_f T_f C_A}{2 \epsilon_{\text{IR}}} - \frac{C_F n_f T_f C_A}{\epsilon_{\text{UV}}} + \frac{2 C_F n_f T_f C_A}{\epsilon_{\text{IR}} \epsilon_{\text{UV}}} - \frac{C_F n_f T_f C_A}{\epsilon_{\text{IR}}^2} - \frac{C_F n_f T_f C_A}{\epsilon_{\text{UV}}^2} + \frac{2 C_F Y_s C_A}{\epsilon_{\text{IR}} M_W} - \frac{2 C_F Y_s C_A}{\epsilon_{\text{UV}} M_W} + \frac{7 C_F Y_s C_A}{\epsilon_{\text{IR}} \epsilon_{\text{UV}} M_W} - \\
& \frac{7 C_F Y_s C_A}{2 \epsilon_{\text{IR}}^2 M_W} - \frac{7 C_F Y_s C_A}{2 \epsilon_{\text{UV}}^2 M_W} - \frac{3 Y_s^4 \zeta_2 C_A}{512 \epsilon_{\text{IR}} m^4} + \frac{151 C_F Y_s^2 \zeta_2 C_A}{128 m^2} + \frac{9 C_F Y_s^2 \zeta_2 C_A}{128 \epsilon_{\text{IR}} m^2} + 6 C_F \zeta_2 C_A - \frac{3 \zeta_2 C_A}{8} - \frac{7 Y_s^4 L_m C_A}{512 \epsilon_{\text{IR}} m^4} - \\
& \frac{Y_s^4 L_m C_A}{32 \epsilon_{\text{IR}} m^4} - \frac{1111 C_F Y_s^2 L_m C_A}{128 m^2} - \frac{C_F^2 Y_s^2 L_m C_A}{2 \epsilon_{\text{IR}} m^2} - \frac{171 C_F Y_s^2 L_m C_A}{128 \epsilon_{\text{IR}} m^2} + \frac{3 C_F Y_s^2 L_m C_A}{8 \epsilon_{\text{IR}}^2 m^2} - \frac{1}{4} C_F L_m C_A - C_F n_f T_f L_m C_A + \\
& \frac{27 C_F^2 L_m C_A}{2 \epsilon_{\text{IR}}} - \frac{12 C_F L_m C_A}{\epsilon_{\text{IR}}} - 9 L_m C_A - \frac{387 C_F^2 C_A}{8 \epsilon_{\text{IR}}} + \frac{21 C_F C_A}{8 \epsilon_{\text{IR}}} + \frac{49 C_F^2 C_A}{2 \epsilon_{\text{UV}}} - \frac{5 C_F C_A}{2 \epsilon_{\text{UV}}} - \frac{19 C_F^2 C_A}{2 \epsilon_{\text{IR}} \epsilon_{\text{UV}}} - \frac{7 C_F C_A}{2 \epsilon_{\text{IR}} \epsilon_{\text{UV}}} - \\
& \frac{2 C_F^2 C_A}{\epsilon_{\text{IR}}^2} + \frac{31 C_F C_A}{4 \epsilon_{\text{IR}}} + \frac{19 C_F^2 C_A}{4 \epsilon_{\text{UV}}} + \frac{7 C_F C_A}{4 \epsilon_{\text{UV}}} + \frac{105 C_A}{8} + \frac{13 C_F Y_s^2}{16 \epsilon_{\text{IR}} m^2} + \frac{C_F Y_s^2}{8 \epsilon_{\text{IR}}^2 m^2} + \frac{65 C_F Y_s^2}{64 \epsilon_{\text{IR}} M_W^2} - \frac{65 C_F Y_s^2}{64 \epsilon_{\text{UV}} M_W^2} - \frac{3}{2} C_F L_m^2 - \\
& \frac{105 C_F}{4} + \frac{3 C_F \zeta_2}{4} - \frac{C_F Y_s^2 L_m}{4 \epsilon_{\text{IR}} m^2} + 18 C_F L_m + \frac{33 C_F L_m}{4 \epsilon_{\text{IR}}} - \frac{627 C_F}{16 \epsilon_{\text{IR}}} + \frac{73 C_F}{4 \epsilon_{\text{UV}}} - \frac{47 C_F}{4 \epsilon_{\text{IR}} \epsilon_{\text{UV}}} + \frac{7 C_F}{4 \epsilon_{\text{IR}}^2} + \frac{47 C_F}{8 \epsilon_{\text{UV}}^2} \tag{D.6}
\end{aligned}$$

#### D.4 Full theory field at $m \neq 0$ and $M \neq 0$ ( $\Delta_M \equiv M_H - M_W$ , $\Delta_{m,M} \equiv M_W - m$ ):

$$\begin{aligned}
F_\psi^{(m,M)} = & \frac{5}{64} C_A C_m^2 Y_f^4 - \frac{2229 C_A Y_f^4}{512} + \frac{9}{4} C_A S_1 Y_f^4 - \frac{39}{64} C_A \log(3) S_1 Y_f^4 + \frac{39 C_A S_1 Y_f^4}{64 \epsilon_{UV}} + \frac{25 C_A \Delta_M S_1 Y_f^4}{12m} + \frac{25 C_A \Delta_{m,M} S_1 Y_f^4}{12m} \\
& - \frac{3 C_A \Delta_M \log(3) S_1 Y_f^4}{16m} - \frac{3 C_A \Delta_{m,M} \log(3) S_1 Y_f^4}{16m} + \frac{3 C_A \Delta_M S_1 Y_f^4}{16 \epsilon_{UV} m} + \frac{3 C_A \Delta_{m,M} S_1 Y_f^4}{16 \epsilon_{UV} m} + \frac{2295}{256} C_A S_2 Y_f^4 + \frac{81}{64} C_A S_1 S_2 Y_f^4 \\
& - \frac{117 C_A \Delta_M S_1 S_2 Y_f^4}{16m} - \frac{117 C_A \Delta_{m,M} S_1 S_2 Y_f^4}{16m} - \frac{351 C_A \Delta_M S_2 Y_f^4}{16m} - \frac{351 C_A \Delta_{m,M} S_2 Y_f^4}{16m} - \frac{25}{32} C_A \zeta_2 Y_f^4 + \frac{9 C_A \Delta_M \zeta_2 Y_f^4}{8m} \\
& + \frac{9 C_A \Delta_{m,M} \zeta_2 Y_f^4}{8m} - \frac{9}{32} C_A \zeta_3 Y_f^4 + \frac{13 C_A \Delta_M \zeta_3 Y_f^4}{8m} + \frac{13 C_A \Delta_{m,M} \zeta_3 Y_f^4}{8m} + \frac{271}{128} C_A C_m Y_f^4 - \frac{39}{32} C_A S_1 C_m Y_f^4 \\
& - \frac{3 C_A \Delta_M S_1 C_m Y_f^4}{8m} - \frac{3 C_A \Delta_{m,M} S_1 C_m Y_f^4}{8m} - \frac{5 C_A C_m Y_f^4}{64 \epsilon_{UV}} + \frac{7 C_A \Delta_M C_m Y_f^4}{8m} + \frac{7 C_A \Delta_{m,M} C_m Y_f^4}{8m} - \frac{271 C_A Y_f^4}{256 \epsilon_{UV}} \\
& + \frac{37 C_A \Delta_M Y_f^4}{16m} + \frac{37 C_A \Delta_{m,M} Y_f^4}{16m} - \frac{7 C_A \Delta_M Y_f^4}{16 \epsilon_{UV} m} - \frac{7 C_A \Delta_{m,M} Y_f^4}{16 \epsilon_{UV} m} + \frac{5 C_A Y_f^4}{128 \epsilon_{UV}^2} - \frac{3 C_A Y_f^3}{8} - \frac{3}{8} C_A S_1 Y_f^3 - \frac{13 C_A \Delta_M S_1 Y_f^3}{4m} \\
& - \frac{23 C_A \Delta_{m,M} S_1 Y_f^3}{8m} + \frac{27}{8} C_A S_1 S_2 Y_f^3 - \frac{27 C_A \Delta_{m,M} S_1 S_2 Y_f^3}{8m} + \frac{405 C_A \Delta_M S_2 Y_f^3}{32m} + \frac{405 C_A \Delta_{m,M} S_2 Y_f^3}{32m} + \frac{33 C_A \Delta_M \zeta_2 Y_f^3}{16m} \\
& + \frac{33 C_A \Delta_{m,M} \zeta_2 Y_f^3}{16m} - \frac{3}{4} C_A \zeta_3 Y_f^3 + \frac{3 C_A \Delta_{m,M} \zeta_3 Y_f^3}{4m} - \frac{3 C_A \Delta_M Y_f^3}{4m} - \frac{3 C_A \Delta_{m,M} Y_f^3}{8m} + \frac{861}{64} C_A C_F^2 Y_f^2 - \frac{1}{4} C_A C_F^2 C_m^2 Y_f^2 \\
& - \frac{3}{8} C_A C_F C_m^2 Y_f^2 - \frac{1}{8} C_F C_m^2 Y_f^2 - \frac{45 C_A C_F \Delta_M C_m^2 Y_f^2}{16m} - \frac{45 C_A C_F \Delta_{m,M} C_m^2 Y_f^2}{8m} + \frac{135 C_A Y_f^2}{64} + \frac{2327}{64} C_A C_F Y_f^2 \\
& + \frac{861 C_F Y_f^2}{128} - \frac{521}{64} C_A C_F^2 S_1 Y_f^2 - \frac{105}{64} C_A S_1 Y_f^2 - \frac{19}{2} C_A C_F S_1 Y_f^2 - \frac{521}{128} C_F S_1 Y_f^2 + \frac{57}{32} C_A C_F^2 \log(3) S_1 Y_f^2 \\
& + \frac{9}{32} C_A \log(3) S_1 Y_f^2 + \frac{41}{8} C_A C_F \log(3) S_1 Y_f^2 + \frac{57}{64} C_F \log(3) S_1 Y_f^2 - \frac{57 C_A C_F^2 S_1 Y_f^2}{32 \epsilon_{UV}} - \frac{9 C_A S_1 Y_f^2}{32 \epsilon_{UV}} - \frac{41 C_A C_F S_1 Y_f^2}{8 \epsilon_{UV}} \\
& - \frac{57 C_F S_1 Y_f^2}{64 \epsilon_{UV}} + \frac{1891 C_A C_F^2 \Delta_M S_1 Y_f^2}{96m} - \frac{193 C_A \Delta_M S_1 Y_f^2}{32m} + \frac{5 C_A C_F \Delta_M S_1 Y_f^2}{6m} + \frac{1891 C_F \Delta_M S_1 Y_f^2}{192m} + \frac{C_A C_F^2 \Delta_{m,M} S_1 Y_f^2}{4m} \\
& - \frac{11 C_A \Delta_{m,M} S_1 Y_f^2}{4m} - \frac{245 C_A C_F \Delta_{m,M} S_1 Y_f^2}{36m} + \frac{C_F \Delta_{m,M} S_1 Y_f^2}{8m} + \frac{25 C_A C_F^2 \Delta_M \log(3) S_1 Y_f^2}{16m} + \frac{33 C_A \Delta_M \log(3) S_1 Y_f^2}{16m} \\
& + \frac{5 C_A C_F \Delta_M \log(3) S_1 Y_f^2}{2m} + \frac{25 C_F \Delta_M \log(3) S_1 Y_f^2}{32m} + \frac{9 C_A C_F^2 \Delta_{m,M} \log(3) S_1 Y_f^2}{2m} + \frac{3 C_A \Delta_{m,M} \log(3) S_1 Y_f^2}{2m} \\
& + \frac{15 C_A C_F \Delta_{m,M} \log(3) S_1 Y_f^2}{2m} + \frac{9 C_F \Delta_{m,M} \log(3) S_1 Y_f^2}{4m} - \frac{25 C_A C_F^2 \Delta_M S_1 Y_f^2}{16 \epsilon_{UV} m} - \frac{33 C_A \Delta_M S_1 Y_f^2}{16 \epsilon_{UV} m} - \frac{5 C_A C_F \Delta_M S_1 Y_f^2}{2 \epsilon_{UV} m} \\
& - \frac{25 C_F \Delta_M S_1 Y_f^2}{32 \epsilon_{UV} m} - \frac{9 C_A C_F^2 \Delta_{m,M} S_1 Y_f^2}{2 \epsilon_{UV} m} - \frac{3 C_A \Delta_{m,M} S_1 Y_f^2}{2 \epsilon_{UV} m} - \frac{15 C_A C_F \Delta_{m,M} S_1 Y_f^2}{2 \epsilon_{UV} m} - \frac{9 C_F \Delta_{m,M} S_1 Y_f^2}{4 \epsilon_{UV} m} \\
& + \frac{171}{32} C_A C_F^2 S_2 Y_f^2 + \frac{81}{32} C_A S_2 Y_f^2 - \frac{3501}{32} C_A C_F S_2 Y_f^2 + \frac{171}{64} C_F S_2 Y_f^2 - \frac{63}{4} C_A C_F S_1 S_2 Y_f^2 + \frac{81 C_A C_F \Delta_M S_1 S_2 Y_f^2}{4m} \\
& - \frac{9 C_A C_F \Delta_{m,M} S_1 S_2 Y_f^2}{4m} - \frac{10317 C_A C_F^2 \Delta_M S_2 Y_f^2}{64m} - \frac{513 C_A \Delta_M S_2 Y_f^2}{32m} - \frac{99 C_A C_F \Delta_M S_2 Y_f^2}{2m} - \frac{10317 C_F \Delta_M S_2 Y_f^2}{128m} \\
& - \frac{3051 C_A C_F^2 \Delta_{m,M} S_2 Y_f^2}{32m} - \frac{675 C_A \Delta_{m,M} S_2 Y_f^2}{32m} - \frac{639 C_A C_F \Delta_{m,M} S_2 Y_f^2}{4m} - \frac{3051 C_F \Delta_{m,M} S_2 Y_f^2}{64m} - \frac{57}{32} C_A C_F^2 \zeta_2 Y_f^2 \\
& - \frac{9}{32} C_A \zeta_2 Y_f^2 + \frac{5}{2} C_A C_F \zeta_2 Y_f^2 - \frac{57}{64} C_F \zeta_2 Y_f^2 + \frac{101 C_A C_F^2 \Delta_M \zeta_2 Y_f^2}{32m} - \frac{9 C_A \Delta_M \zeta_2 Y_f^2}{8m} - \frac{301 C_A C_F \Delta_M \zeta_2 Y_f^2}{32m} \\
& + \frac{101 C_F \Delta_M \zeta_2 Y_f^2}{64m} - \frac{33 C_A C_F^2 \Delta_{m,M} \zeta_2 Y_f^2}{16m} - \frac{9 C_A \Delta_{m,M} \zeta_2 Y_f^2}{16m} - \frac{439 C_A C_F \Delta_{m,M} \zeta_2 Y_f^2}{48m} - \frac{33 C_F \Delta_{m,M} \zeta_2 Y_f^2}{32m} \\
& + \frac{7}{2} C_A C_F \zeta_3 Y_f^2 - \frac{9 C_A C_F \Delta_M \zeta_3 Y_f^2}{2m} + \frac{C_A C_F \Delta_{m,M} \zeta_3 Y_f^2}{2m} - \frac{59}{16} C_A C_F^2 C_m Y_f^2 - \frac{9}{16} C_A C_m Y_f^2 - \frac{293}{16} C_A C_F C_m Y_f^2 \\
& - \frac{59}{32} C_F C_m Y_f^2 + \frac{57}{16} C_A C_F^2 S_1 C_m Y_f^2 + \frac{9}{16} C_A S_1 C_m Y_f^2 + \frac{41}{4} C_A C_F S_1 C_m Y_f^2 + \frac{57}{32} C_F S_1 C_m Y_f^2 + \frac{25 C_A C_F^2 \Delta_M S_1 C_m Y_f^2}{8m} \\
& + \frac{33 C_A \Delta_M S_1 C_m Y_f^2}{8m} + \frac{5 C_A C_F \Delta_M S_1 C_m Y_f^2}{m} + \frac{25 C_F \Delta_M S_1 C_m Y_f^2}{16m} + \frac{9 C_A C_F^2 \Delta_{m,M} S_1 C_m Y_f^2}{m} + \frac{3 C_A \Delta_{m,M} S_1 C_m Y_f^2}{m} \\
& + \frac{15 C_A C_F \Delta_{m,M} S_1 C_m Y_f^2}{m} + \frac{9 C_F \Delta_{m,M} S_1 C_m Y_f^2}{2m} + \frac{C_A C_F^2 C_m Y_f^2}{4 \epsilon_{UV}} + \frac{3 C_A C_F C_m Y_f^2}{8 \epsilon_{UV}} + \frac{C_F C_m Y_f^2}{8 \epsilon_{UV}} - \frac{51 C_A C_F^2 \Delta_M C_m Y_f^2}{4m} \\
& - \frac{27 C_A \Delta_M C_m Y_f^2}{4m} - \frac{81 C_A C_F \Delta_M C_m Y_f^2}{32m} - \frac{51 C_F \Delta_M C_m Y_f^2}{8m} - \frac{149 C_A C_F^2 \Delta_{m,M} C_m Y_f^2}{8m} - \frac{45 C_A \Delta_{m,M} C_m Y_f^2}{8m} \\
& - \frac{225 C_A C_F \Delta_{m,M} C_m Y_f^2}{16m} - \frac{149 C_F \Delta_{m,M} C_m Y_f^2}{16m} + \frac{59 C_A C_F^2 Y_f^2}{32 \epsilon_{UV}} + \frac{9 C_A Y_f^2}{32 \epsilon_{UV}} + \frac{293 C_A C_F Y_f^2}{32 \epsilon_{UV}} + \frac{59 C_F Y_f^2}{64 \epsilon_{UV}} - \frac{29 C_A C_F^2 \Delta_M Y_f^2}{32m} \\
& + \frac{435 C_A \Delta_M Y_f^2}{32m} + \frac{23 C_A C_F \Delta_M Y_f^2}{128m} - \frac{29 C_F \Delta_M Y_f^2}{64m} + \frac{191 C_A C_F^2 \Delta_{m,M} Y_f^2}{8m} + \frac{75 C_A \Delta_{m,M} Y_f^2}{8m} + \frac{3989 C_A C_F \Delta_{m,M} Y_f^2}{192m} \\
& + \frac{191 C_F \Delta_{m,M} Y_f^2}{16m} + \frac{51 C_A C_F^2 \Delta_M Y_f^2}{8 \epsilon_{UV} m} + \frac{27 C_A \Delta_M Y_f^2}{8 \epsilon_{UV} m} + \frac{9 C_A C_F \Delta_M Y_f^2}{2 \epsilon_{UV} m} + \frac{51 C_F \Delta_M Y_f^2}{16 \epsilon_{UV} m} + \frac{149 C_A C_F^2 \Delta_{m,M} Y_f^2}{16 \epsilon_{UV} m} \\
& + \frac{45 C_A \Delta_{m,M} Y_f^2}{16 \epsilon_{UV} m} + \frac{27 C_A C_F \Delta_{m,M} Y_f^2}{2 \epsilon_{UV} m} + \frac{149 C_F \Delta_{m,M} Y_f^2}{32 \epsilon_{UV} m} - \frac{C_A C_F^2 Y_f^2}{8 \epsilon_{UV}^2} - \frac{3 C_A C_F Y_f^2}{16 \epsilon_{UV}^2} - \frac{C_F Y_f^2}{16 \epsilon_{UV}^2} + \frac{5 C_A C_F \Delta_M C_m^2 Y_f}{2m}
\end{aligned}$$

$$\begin{aligned}
& + \frac{5C_A C_F \Delta_{m,M} \mathcal{L}_m^2 Y_f}{m} - \frac{22C_A C_F \Delta_M S_1 Y_f}{3m} + \frac{8C_A C_F \Delta_{m,M} S_1 Y_f}{m} - 54C_A C_F S_1 S_2 Y_f + \frac{72C_A C_F \Delta_M S_1 S_2 Y_f}{m} \\
& - \frac{54C_A C_F \Delta_{m,M} S_1 S_2 Y_f}{m} + \frac{297C_A C_F \Delta_M S_2 Y_f}{2m} + \frac{108C_A C_F \Delta_{m,M} S_2 Y_f}{m} + 8C_A C_F \zeta_2 Y_f - \frac{197C_A C_F \Delta_M \zeta_2 Y_f}{12m} \\
& - \frac{17C_A C_F \Delta_{m,M} \zeta_2 Y_f}{6m} + 12C_A C_F \zeta_3 Y_f - \frac{16C_A C_F \Delta_M \zeta_3 Y_f}{m} + \frac{12C_A C_F \Delta_{m,M} \zeta_3 Y_f}{m} - \frac{23C_A C_F \Delta_M \mathcal{L}_m Y_f}{4m} \\
& - \frac{23C_A C_F \Delta_{m,M} \mathcal{L}_m Y_f}{2m} - \frac{31C_A C_F \Delta_M Y_f}{16m} - \frac{31C_A C_F \Delta_{m,M} Y_f}{8m} - \frac{45485}{96} C_A C_F^2 - 5C_A C_F^2 \mathcal{L}_m^2 - 3C_F \mathcal{L}_m^2 \\
& - \frac{5C_A C_F \Delta_M \mathcal{L}_m^2}{12m} - \frac{5C_A C_F \Delta_{m,M} \mathcal{L}_m^2}{6m} - \frac{15C_F \Delta_{m,M} \mathcal{L}_m^2}{2m} + \frac{1231C_A C_F}{96} - \frac{45233C_F}{192} + \frac{2945}{18} C_A C_F^2 S_1 - \frac{29}{18} C_A C_F S_1 \\
& + \frac{3461C_F S_1}{36} - \frac{1061}{12} C_A C_F^2 \log(3) S_1 + \frac{85}{36} C_A C_F \log(3) S_1 - \frac{1025}{24} C_F \log(3) S_1 + \frac{1061C_A C_F^2 S_1}{12\epsilon_{UV}} - \frac{85C_A C_F S_1}{36\epsilon_{UV}} + \frac{1025C_F S_1}{24\epsilon_{UV}} \\
& + \frac{79C_A C_F \Delta_M S_1}{18m} + \frac{569C_A C_F^2 \Delta_{m,M} S_1}{6m} + \frac{233C_A C_F \Delta_{m,M} S_1}{54m} + \frac{2779C_F \Delta_{m,M} S_1}{36m} - \frac{5C_A C_F \Delta_M \log(3) S_1}{3m} \\
& - \frac{379C_A C_F^2 \Delta_{m,M} \log(3) S_1}{3m} + \frac{19C_A C_F \Delta_{m,M} \log(3) S_1}{9m} - \frac{93C_F \Delta_{m,M} \log(3) S_1}{2m} + \frac{5C_A C_F \Delta_M S_1}{3\epsilon_{UV} m} + \frac{379C_A C_F^2 \Delta_{m,M} S_1}{3\epsilon_{UV} m} \\
& - \frac{19C_A C_F \Delta_{m,M} S_1}{9\epsilon_{UV} m} + \frac{93C_F \Delta_{m,M} S_1}{2\epsilon_{UV} m} + \frac{23163}{16} C_A C_F^2 S_2 - \frac{817}{16} C_A C_F S_2 + \frac{19167C_F S_2}{32} + 144C_A C_F^2 S_1 S_2 + 72C_F S_1 S_2 \\
& + \frac{450C_A C_F^2 \Delta_{m,M} S_1 S_2}{m} + \frac{225C_F \Delta_{m,M} S_1 S_2}{m} + \frac{27C_A C_F \Delta_M S_2}{2m} + \frac{8769C_A C_F^2 \Delta_{m,M} S_2}{4m} - \frac{427C_A C_F \Delta_{m,M} S_2}{4m} \\
& + \frac{4113C_F \Delta_{m,M} S_2}{8m} + \frac{377}{24} C_A C_F \text{nf} T_f + \frac{55C_A C_F \text{nf} T_f}{12\epsilon_{UV}} + \frac{134C_A C_F \Delta_{m,M} \text{nf} T_f}{9m} + \frac{6C_A C_F \Delta_{m,M} \text{nf} T_f}{\epsilon_{UV} m} - 5C_A C_F \text{nf} S_1 T_f \\
& + \frac{25}{9} C_A C_F \log(3) \text{nf} S_1 T_f - \frac{25C_A C_F \text{nf} S_1 T_f}{9\epsilon_{UV}} + \frac{46C_A C_F \Delta_{m,M} \text{nf} S_1 T_f}{27m} + \frac{28C_A C_F \Delta_{m,M} \log(3) \text{nf} S_1 T_f}{9m} \\
& - \frac{28C_A C_F \Delta_{m,M} \text{nf} S_2 T_f}{9\epsilon_{UV} m} - \frac{175}{4} C_A C_F \text{nf} S_2 T_f - \frac{49C_A C_F \Delta_{m,M} \text{nf} S_2 T_f}{m} - \frac{351}{8} C_A C_F^2 \zeta_2 - \frac{5}{24} C_A C_F \zeta_2 - \frac{287C_F \zeta_2}{16} \\
& - \frac{1}{2} C_A C_F \text{nf} T_f \zeta_2 - \frac{6C_A C_F \Delta_{m,M} \text{nf} T_f \zeta_2}{m} - \frac{55C_A C_F \Delta_M \zeta_2}{72m} - \frac{559C_A C_F^2 \Delta_{m,M} \zeta_2}{6m} - \frac{C_A C_F \Delta_{m,M} \zeta_2}{36m} - \frac{139C_F \Delta_{m,M} \zeta_2}{3m} \\
& - 32C_A C_F^2 \zeta_3 - 16C_F \zeta_3 - \frac{100C_A C_F^2 \Delta_{m,M} \zeta_3}{m} - \frac{50C_F \Delta_{m,M} \zeta_3}{m} + \frac{2367}{8} C_A C_F^2 \mathcal{L}_m - \frac{199}{24} C_A C_F \mathcal{L}_m + \frac{2291}{16} C_F \mathcal{L}_m \\
& - \frac{1061}{6} C_A C_F^2 S_1 \mathcal{L}_m + \frac{85}{18} C_A C_F S_1 \mathcal{L}_m - \frac{1025}{12} C_F S_1 \mathcal{L}_m - \frac{10C_A C_F \Delta_M S_1 \mathcal{L}_m}{3m} - \frac{758C_A C_F^2 \Delta_{m,M} S_1 \mathcal{L}_m}{3m} \\
& + \frac{38C_A C_F \Delta_{m,M} S_1 \mathcal{L}_m}{9m} - \frac{93C_F \Delta_{m,M} S_1 \mathcal{L}_m}{m} - \frac{55}{6} C_A C_F \text{nf} T_f \mathcal{L}_m - \frac{12C_A C_F \Delta_{m,M} \text{nf} T_f \mathcal{L}_m}{m} + \frac{50}{9} C_A C_F \text{nf} S_1 T_f \mathcal{L}_m \\
& + \frac{56C_A C_F \Delta_{m,M} \text{nf} S_1 T_f \mathcal{L}_m}{9m} + \frac{5C_A C_F^2 \mathcal{L}_m}{\epsilon_{UV}} + \frac{3C_F \mathcal{L}_m}{\epsilon_{UV}} + \frac{167C_A C_F \Delta_M \mathcal{L}_m}{24m} + \frac{449C_A C_F^2 \Delta_{m,M} \mathcal{L}_m}{m} - \frac{61C_A C_F \Delta_{m,M} \mathcal{L}_m}{12m} \\
& + \frac{727C_F \Delta_{m,M} \mathcal{L}_m}{4m} - \frac{2367C_A C_F^2}{16\epsilon_{UV}} + \frac{199C_A C_F}{48\epsilon_{UV}} - \frac{2291C_F}{32\epsilon_{UV}} - \frac{315C_A C_F \Delta_M}{32m} - \frac{1328C_A C_F^2 \Delta_{m,M}}{3m} + \frac{1757C_A C_F \Delta_{m,M}}{144m} \\
& - \frac{9737C_F \Delta_{m,M}}{48m} - \frac{3C_A C_F \Delta_M}{\epsilon_{UV} m} - \frac{449C_A C_F^2 \Delta_{m,M}}{2\epsilon_{UV} m} + \frac{7C_A C_F \Delta_{m,M}}{2\epsilon_{UV} m} - \frac{329C_F \Delta_{m,M}}{4\epsilon_{UV} m} - \frac{5C_A C_F^2}{2\epsilon_{UV}^2} - \frac{3C_F}{2\epsilon_{UV}^2} \quad (D.7)
\end{aligned}$$

$$\begin{aligned}
F_X^{(m,M)} = & \frac{5C_A \Delta_M \mathcal{L}_m^2 Y_s^4}{64m^5} + \frac{5C_A \Delta_{m,M} \mathcal{L}_m^2 Y_s^4}{64m^5} - \frac{5C_A S_1 Y_s^4}{144m^4} - \frac{C_A \log(3) S_1 Y_s^4}{144m^4} + \frac{C_A S_1 Y_s^4}{144\epsilon_{UV} m^4} + \frac{31C_A \Delta_M S_1 Y_s^4}{216m^5} + \frac{31C_A \Delta_{m,M} S_1 Y_s^4}{216m^5} \\
& - \frac{C_A \Delta_M \log(3) S_1 Y_s^4}{18m^5} - \frac{C_A \Delta_{m,M} \log(3) S_1 Y_s^4}{18m^5} + \frac{C_A \Delta_M S_1 Y_s^4}{18\epsilon_{UV} m^5} + \frac{C_A \Delta_{m,M} S_1 Y_s^4}{18\epsilon_{UV} m^5} + \frac{9C_A S_1 S_2 Y_s^4}{32m^4} + \frac{31C_A S_2 Y_s^4}{64m^4} \\
& + \frac{61C_A \Delta_M S_2 Y_s^4}{32m^5} + \frac{61C_A \Delta_{m,M} S_2 Y_s^4}{32m^5} - \frac{5C_A \zeta_2 Y_s^4}{96m^4} + \frac{7C_A \Delta_M \zeta_2 Y_s^4}{384m^5} + \frac{7C_A \Delta_{m,M} \zeta_2 Y_s^4}{384m^5} - \frac{C_A \zeta_3 Y_s^4}{16m^4} - \frac{C_A S_1 \mathcal{L}_m Y_s^4}{72m^4} \\
& - \frac{C_A \Delta_M S_1 \mathcal{L}_m Y_s^4}{9m^5} - \frac{C_A \Delta_{m,M} S_1 \mathcal{L}_m Y_s^4}{9m^5} + \frac{C_A \mathcal{L}_m Y_s^4}{24m^4} - \frac{5C_A \Delta_M \mathcal{L}_m Y_s^4}{384m^5} - \frac{5C_A \Delta_{m,M} \mathcal{L}_m Y_s^4}{384m^5} - \frac{C_A Y_s^4}{24m^4} - \frac{C_A Y_s^4}{48\epsilon_{UV} m^4} \\
& - \frac{919C_A \Delta_M Y_s^4}{4608m^5} - \frac{919C_A \Delta_{m,M} Y_s^4}{4608m^5} - \frac{C_A \Delta_M Y_s^4}{12\epsilon_{UV} m^5} - \frac{C_A \Delta_{m,M} Y_s^4}{12\epsilon_{UV} m^5} - \frac{5C_A \Delta_M \mathcal{L}_m^2 Y_s^3}{64m^4} - \frac{5C_A \Delta_{m,M} \mathcal{L}_m^2 Y_s^3}{64m^4} - \frac{C_A \Delta_M S_1 Y_s^3}{8m^4} \\
& - \frac{C_A \Delta_{m,M} S_1 Y_s^3}{8m^4} + \frac{27C_A S_1 S_2 Y_s^3}{32m^3} + \frac{27C_A \Delta_M S_1 S_2 Y_s^3}{16m^4} + \frac{27C_A \Delta_{m,M} S_1 S_2 Y_s^3}{32m^4} - \frac{27C_A \Delta_M S_2 Y_s^3}{16m^4} - \frac{27C_A \Delta_{m,M} S_2 Y_s^3}{16m^4} \\
& - \frac{C_A \zeta_2 Y_s^3}{8m^3} - \frac{31C_A \Delta_M \zeta_2 Y_s^3}{384m^4} + \frac{17C_A \Delta_{m,M} \zeta_2 Y_s^3}{384m^4} - \frac{3C_A \zeta_3 Y_s^3}{16m^3} - \frac{3C_A \Delta_M \zeta_3 Y_s^3}{8m^4} - \frac{3C_A \Delta_{m,M} \zeta_3 Y_s^3}{16m^4} + \frac{23C_A \Delta_M \mathcal{L}_m Y_s^3}{128m^4} \\
& + \frac{23C_A \Delta_{m,M} \mathcal{L}_m Y_s^3}{128m^4} + \frac{31C_A \Delta_M Y_s^3}{512m^4} + \frac{31C_A \Delta_{m,M} Y_s^3}{512m^4} + \frac{5C_A \Delta_M \mathcal{L}_m^2 Y_s^2}{256m^3} - \frac{5C_A C_F \Delta_M \mathcal{L}_m^2 Y_s^2}{64m^3} - \frac{35C_A C_F^2 \Delta_{m,M} \mathcal{L}_m^2 Y_s^2}{2304m^3} \\
& + \frac{5C_A \Delta_{m,M} \mathcal{L}_m^2 Y_s^2}{256m^3} - \frac{5C_A C_F \Delta_{m,M} \mathcal{L}_m^2 Y_s^2}{32m^3} - \frac{35C_F \Delta_{m,M} \mathcal{L}_m^2 Y_s^2}{4608m^3} - \frac{191C_A C_F^2 S_1 Y_s^2}{96m^2} - \frac{5C_A S_1 Y_s^2}{32m^2} - \frac{101C_A C_F S_1 Y_s^2}{36m^2} \\
& - \frac{191C_F S_1 Y_s^2}{192m^2} + \frac{7C_A C_F^2 \log(3) S_1 Y_s^2}{32m^2} - \frac{C_A \log(3) S_1 Y_s^2}{32m^2} + \frac{3C_A C_F \log(3) S_1 Y_s^2}{4m^2} + \frac{7C_F \log(3) S_1 Y_s^2}{64m^2} - \frac{7C_A C_F^2 S_1 Y_s^2}{32\epsilon_{UV} m^2} \\
& + \frac{C_A S_1 Y_s^2}{32\epsilon_{UV} m^2} - \frac{3C_A C_F S_1 Y_s^2}{4\epsilon_{UV} m^2} - \frac{7C_F S_1 Y_s^2}{64\epsilon_{UV} m^2} + \frac{59C_A C_F^2 \Delta_M S_1 Y_s^2}{24m^3} - \frac{19C_A \Delta_M S_1 Y_s^2}{24m^3} - \frac{35C_A C_F \Delta_M S_1 Y_s^2}{18m^3} \\
& + \frac{59C_F \Delta_M S_1 Y_s^2}{48m^3} - \frac{109C_A C_F^2 \Delta_{m,M} S_1 Y_s^2}{144m^3} - \frac{23C_A \Delta_{m,M} S_1 Y_s^2}{48m^3} - \frac{29C_A C_F \Delta_{m,M} S_1 Y_s^2}{6m^3} - \frac{109C_F \Delta_{m,M} S_1 Y_s^2}{288m^3}
\end{aligned}$$

$$\begin{aligned}
& + \frac{9C_A C_F^2 \Delta_M \log(3) S_1 Y_s^2}{8m^3} + \frac{C_A \Delta_M \log(3) S_1 Y_s^2}{8m^3} + \frac{C_A C_F \Delta_M \log(3) S_1 Y_s^2}{6m^3} + \frac{9C_F \Delta_M \log(3) S_1 Y_s^2}{16m^3} \\
& + \frac{115C_A C_F^2 \Delta_{m,M} \log(3) S_1 Y_s^2}{144m^3} + \frac{3C_A \Delta_{m,M} \log(3) S_1 Y_s^2}{16m^3} + \frac{C_A C_F \Delta_{m,M} \log(3) S_1 Y_s^2}{2m^3} + \frac{115C_F \Delta_{m,M} \log(3) S_1 Y_s^2}{288m^3} \\
& - \frac{9C_A C_F^2 \Delta_M S_1 Y_s^2}{8\epsilon_{UV} m^3} - \frac{C_A \Delta_M S_1 Y_s^2}{8\epsilon_{UV} m^3} - \frac{C_A C_F \Delta_M S_1 Y_s^2}{6\epsilon_{UV} m^3} - \frac{9C_F \Delta_M S_1 Y_s^2}{16\epsilon_{UV} m^3} - \frac{115C_A C_F^2 \Delta_{m,M} S_1 Y_s^2}{144\epsilon_{UV} m^3} - \frac{3C_A \Delta_{m,M} S_1 Y_s^2}{16\epsilon_{UV} m^3} \\
& - \frac{C_A C_F \Delta_{m,M} S_1 Y_s^2}{2\epsilon_{UV} m^3} - \frac{115C_F \Delta_{m,M} S_1 Y_s^2}{288\epsilon_{UV} m^3} + \frac{9C_A C_F \Delta_M S_1 S_2 Y_s^2}{2m^3} + \frac{27C_A C_F \Delta_{m,M} S_1 S_2 Y_s^2}{4m^3} + \frac{855C_A C_F^2 S_2 Y_s^2}{128m^2} \\
& + \frac{279C_A S_2 Y_s^2}{128m^2} - \frac{63C_A C_F S_2 Y_s^2}{4m^2} + \frac{855C_F S_2 Y_s^2}{256m^2} - \frac{1467C_A C_F^2 \Delta_M S_2 Y_s^2}{32m^3} + \frac{63C_A \Delta_M S_2 Y_s^2}{16m^3} + \frac{27C_A C_F \Delta_M S_2 Y_s^2}{8m^3} \\
& - \frac{1467C_F \Delta_M S_2 Y_s^2}{64m^3} - \frac{523C_A C_F^2 \Delta_{m,M} S_2 Y_s^2}{64m^3} - \frac{27C_A \Delta_{m,M} S_2 Y_s^2}{64m^3} + \frac{147C_A C_F \Delta_{m,M} S_2 Y_s^2}{8m^3} - \frac{523C_F \Delta_{m,M} S_2 Y_s^2}{128m^3} \\
& - \frac{73C_A C_F^2 \zeta_2 Y_s^2}{192m^2} - \frac{3C_A \zeta_2 Y_s^2}{64m^2} + \frac{3C_A C_F \zeta_2 Y_s^2}{4m^2} - \frac{73C_F \zeta_2 Y_s^2}{384m^2} + \frac{5C_A C_F^2 \Delta_M \zeta_2 Y_s^2}{16m^3} - \frac{305C_A \Delta_M \zeta_2 Y_s^2}{1536m^3} - \frac{431C_A C_F \Delta_M \zeta_2 Y_s^2}{384m^3} \\
& + \frac{5C_F \Delta_M \zeta_2 Y_s^2}{32m^3} - \frac{8089C_A C_F^2 \Delta_{m,M} \zeta_2 Y_s^2}{13824m^3} - \frac{161C_A \Delta_{m,M} \zeta_2 Y_s^2}{1536m^3} - \frac{415C_A C_F \Delta_{m,M} \zeta_2 Y_s^2}{192m^3} - \frac{8089C_F \Delta_{m,M} \zeta_2 Y_s^2}{27648m^3} \\
& - \frac{C_A C_F \Delta_M \zeta_3 Y_s^2}{m^3} - \frac{3C_A C_F \Delta_{m,M} \zeta_3 Y_s^2}{2m^3} + \frac{7C_A C_F^2 S_1 L_m Y_s^2}{16m^2} - \frac{C_A S_1 L_m Y_s^2}{16m^2} + \frac{3C_A C_F S_1 L_m Y_s^2}{2m^2} + \frac{7C_F S_1 L_m Y_s^2}{32m^2} \\
& + \frac{9C_A C_F^2 \Delta_M S_1 L_m Y_s^2}{4m^3} + \frac{C_A \Delta_M S_1 L_m Y_s^2}{4m^3} + \frac{C_A C_F \Delta_M S_1 L_m Y_s^2}{3m^3} + \frac{9C_F \Delta_M S_1 L_m Y_s^2}{8m^3} + \frac{115C_A C_F^2 \Delta_{m,M} S_1 L_m Y_s^2}{72m^3} \\
& + \frac{3C_A \Delta_{m,M} S_1 L_m Y_s^2}{8m^3} + \frac{C_A C_F \Delta_{m,M} S_1 L_m Y_s^2}{m^3} + \frac{115C_F \Delta_{m,M} S_1 L_m Y_s^2}{144m^3} - \frac{5C_A C_F^2 L_m Y_s^2}{16m^2} + \frac{3C_A C_F L_m Y_s^2}{16m^2} \\
& - \frac{3C_A C_F L_m Y_s^2}{m^2} - \frac{5C_F L_m Y_s^2}{32m^2} - \frac{43C_A C_F^2 \Delta_M L_m Y_s^2}{8m^3} - \frac{215C_A \Delta_M L_m Y_s^2}{512m^3} - \frac{105C_A C_F \Delta_M L_m Y_s^2}{128m^3} - \frac{43C_F \Delta_M L_m Y_s^2}{16m^3} \\
& - \frac{15583C_A C_F^2 \Delta_{m,M} L_m Y_s^2}{4608m^3} - \frac{407C_A \Delta_{m,M} L_m Y_s^2}{512m^3} - \frac{41C_A C_F \Delta_{m,M} L_m Y_s^2}{64m^3} - \frac{15583C_F \Delta_{m,M} L_m Y_s^2}{9216m^3} + \frac{33C_A C_F^2 Y_s^2}{16m^2} \\
& - \frac{3C_A Y_s^2}{16m^2} + \frac{20C_A C_F Y_s^2}{3m^2} + \frac{33C_F Y_s^2}{32m^2} + \frac{5C_A C_F^2 Y_s^2}{32\epsilon_{UV} m^2} - \frac{3C_A Y_s^2}{32\epsilon_{UV} m^2} + \frac{3C_A C_F Y_s^2}{2\epsilon_{UV} m^2} + \frac{5C_F Y_s^2}{64\epsilon_{UV} m^2} + \frac{83C_A C_F^2 \Delta_M Y_s^2}{16m^3} \\
& + \frac{1377C_A \Delta_M Y_s^2}{2048m^3} + \frac{4189C_A C_F \Delta_M Y_s^2}{1536m^3} + \frac{83C_F \Delta_M Y_s^2}{32m^3} + \frac{65881C_A C_F^2 \Delta_{m,M} Y_s^2}{18432m^3} + \frac{2145C_A \Delta_{m,M} Y_s^2}{2048m^3} \\
& + \frac{2653C_A C_F \Delta_{m,M} Y_s^2}{768m^3} + \frac{65881C_F \Delta_{m,M} Y_s^2}{36864m^3} + \frac{43C_A C_F^2 \Delta_M Y_s^2}{16\epsilon_{UV} m^3} + \frac{3C_A \Delta_M Y_s^2}{16\epsilon_{UV} m^3} + \frac{C_A C_F \Delta_M Y_s^2}{2\epsilon_{UV} m^3} + \frac{43C_F \Delta_M Y_s^2}{32\epsilon_{UV} m^3} \\
& + \frac{41C_A C_F^2 \Delta_{m,M} Y_s^2}{24\epsilon_{UV} m^3} + \frac{3C_A \Delta_{m,M} Y_s^2}{8\epsilon_{UV} m^3} + \frac{C_A C_F \Delta_{m,M} Y_s^2}{2\epsilon_{UV} m^3} + \frac{41C_F \Delta_{m,M} Y_s^2}{48\epsilon_{UV} m^3} + \frac{5C_A C_F \Delta_{m,M} L_m^2 Y_s^2}{8m^2} + \frac{43C_A C_F S_1 Y_s}{6m} \\
& - \frac{2C_A C_F \log(3) S_1 Y_s}{m} + \frac{2C_A C_F S_1 Y_s}{\epsilon_{UV} m} - \frac{17C_A C_F \Delta_M S_1 Y_s}{18m^2} + \frac{107C_A C_F \Delta_{m,M} S_1 Y_s}{9m^2} - \frac{2C_A C_F \Delta_M \log(3) S_1 Y_s}{3m^2} \\
& - \frac{8C_A C_F \Delta_{m,M} \log(3) S_1 Y_s}{3m^2} + \frac{2C_A C_F \Delta_M S_1 Y_s}{3\epsilon_{UV} m^2} + \frac{8C_A C_F \Delta_{m,M} S_1 Y_s}{3\epsilon_{UV} m^2} - \frac{63C_A C_F S_1 S_2 Y_s}{8m} + \frac{27C_A C_F \Delta_M S_1 S_2 Y_s}{4m^2} \\
& + \frac{99C_A C_F \Delta_{m,M} S_1 S_2 Y_s}{8m^2} + \frac{18C_A C_F S_2 Y_s}{m} + \frac{24C_A C_F \Delta_M S_2 Y_s}{m^2} + \frac{39C_A C_F \Delta_{m,M} S_2 Y_s}{2m^2} + \frac{C_A C_F \zeta_2 Y_s}{m} \\
& - \frac{7C_A C_F \Delta_M \zeta_2 Y_s}{3m^2} - \frac{289C_A C_F \Delta_{m,M} \zeta_2 Y_s}{48m^2} + \frac{7C_A C_F \zeta_3 Y_s}{4m} - \frac{3C_A C_F \Delta_M \zeta_3 Y_s}{2m^2} - \frac{11C_A C_F \Delta_{m,M} \zeta_3 Y_s}{4m^2} \\
& - \frac{4C_A C_F S_1 L_m Y_s}{m} - \frac{4C_A C_F \Delta_M S_1 L_m Y_s}{3m^2} - \frac{16C_A C_F \Delta_{m,M} S_1 L_m Y_s}{3m^2} + \frac{6C_A C_F L_m Y_s}{m} + \frac{4C_A C_F \Delta_M L_m Y_s}{m^2} \\
& + \frac{137C_A C_F \Delta_{m,M} L_m Y_s}{16m^2} - \frac{23C_A C_F Y_s}{2m} - \frac{3C_A C_F Y_s}{\epsilon_{UV} m} - \frac{8C_A C_F \Delta_M Y_s}{3m^2} - \frac{2813C_A C_F \Delta_{m,M} Y_s}{192m^2} - \frac{2C_A C_F \Delta_M Y_s}{\epsilon_{UV} m^2} \\
& - \frac{5C_A C_F \Delta_{m,M} Y_s}{\epsilon_{UV} m^2} + \frac{1521}{32} C_A C_F^2 + \frac{73}{4} C_A C_F^2 L_m^2 - \frac{1}{4} C_A C_F L_m^2 + \frac{105}{8} C_F L_m^2 - C_A C_F n_f T_f L_m^2 - \frac{5C_A C_F^2 \Delta_{m,M} L_m^2}{8m} \\
& + \frac{5C_A C_F \Delta_{m,M} L_m^2}{2m} + \frac{5C_F \Delta_{m,M} L_m^2}{8m} + \frac{659C_A C_F}{96} + \frac{2045C_F}{64} - \frac{493}{12} C_A C_F^2 S_1 - \frac{89}{12} C_A C_F S_1 - \frac{493C_F S_1}{24} \\
& + \frac{53}{4} C_A C_F^2 \log(3) S_1 + \frac{41}{12} C_A C_F \log(3) S_1 + \frac{53}{8} C_F \log(3) S_1 - \frac{53C_A C_F^2 S_1}{4\epsilon_{UV}} - \frac{41C_A C_F S_1}{12\epsilon_{UV}} - \frac{53C_F S_1}{8\epsilon_{UV}} - \frac{5C_A C_F \Delta_M S_1}{3m} \\
& - \frac{841C_A C_F^2 \Delta_{m,M} S_1}{9m} + \frac{16C_A C_F \Delta_{m,M} S_1}{3m} - \frac{1243C_F \Delta_{m,M} S_1}{18m} + \frac{C_A C_F \Delta_M \log(3) S_1}{2m} + \frac{179C_A C_F^2 \Delta_{m,M} \log(3) S_1}{3m} \\
& + \frac{29C_A C_F \Delta_{m,M} \log(3) S_1}{9m} + \frac{185C_F \Delta_{m,M} \log(3) S_1}{6m} - \frac{C_A C_F \Delta_M S_1}{2\epsilon_{UV} m} - \frac{179C_A C_F^2 \Delta_{m,M} S_1}{3\epsilon_{UV} m} - \frac{29C_A C_F \Delta_{m,M} S_1}{9\epsilon_{UV} m} \\
& - \frac{185C_F \Delta_{m,M} S_1}{6\epsilon_{UV} m} - \frac{441}{16} C_A C_F^2 S_2 - \frac{609}{16} C_A C_F S_2 - \frac{1089C_F S_2}{32} + \frac{63}{4} C_A C_F^2 S_1 S_2 + \frac{63}{8} C_F S_1 S_2 - \frac{9C_A C_F^2 \Delta_{m,M} S_1 S_2}{m} \\
& - \frac{9C_F \Delta_{m,M} S_1 S_2}{2m} + \frac{3C_A C_F \Delta_M S_2}{8m} - \frac{2985C_A C_F^2 \Delta_{m,M} S_2}{4m} - \frac{53C_A C_F \Delta_{m,M} S_2}{4m} - \frac{2787C_F \Delta_{m,M} S_2}{8m} - \frac{55}{24} C_A C_F n_f T_f \\
& - \frac{3C_A C_F n_f T_f}{4\epsilon_{UV}} - \frac{8C_A C_F \Delta_{m,M} n_f T_f}{m} - \frac{4C_A C_F \Delta_{m,M} n_f T_f}{3\epsilon_{UV} m} - \frac{C_A C_F n_f T_f}{2\epsilon_{UV}^2} + \frac{4}{3} C_A C_F n_f S_1 T_f \\
& - \frac{1}{3} C_A C_F \log(3) n_f S_1 T_f + \frac{C_A C_F n_f S_1 T_f}{3\epsilon_{UV}} + \frac{8C_A C_F \Delta_{m,M} n_f S_1 T_f}{3m} - \frac{16C_A C_F \Delta_{m,M} \log(3) n_f S_1 T_f}{9m} \\
& + \frac{16C_A C_F \Delta_{m,M} n_f S_1 T_f}{9\epsilon_{UV} m} + \frac{21}{4} C_A C_F n_f S_2 T_f + \frac{28C_A C_F \Delta_{m,M} n_f S_2 T_f}{m} - C_A C_F^2 \zeta_2 + \frac{1}{4} C_A C_F \zeta_2 - 2C_A C_F n_f T_f \zeta_2
\end{aligned}$$

$$\begin{aligned}
& + \frac{C_A C_F \Delta_M \zeta_2}{12m} + \frac{115 C_A C_F^2 \Delta_{m,M} \zeta_2}{16m} + \frac{25 C_A C_F \Delta_{m,M} \zeta_2}{12m} + \frac{153 C_F \Delta_{m,M} \zeta_2}{16m} - \frac{7}{2} C_A C_F^2 \zeta_3 - \frac{7 C_F \zeta_3}{4} + \frac{2 C_A C_F^2 \Delta_{m,M} \zeta_3}{m} \\
& + \frac{C_F \Delta_{m,M} \zeta_3}{m} - \frac{307}{8} C_A C_F^2 \mathcal{L}_m + \frac{21}{8} C_A C_F \mathcal{L}_m - \frac{343}{16} C_F \mathcal{L}_m + \frac{53}{2} C_A C_F^2 S_1 \mathcal{L}_m + \frac{41}{6} C_A C_F S_1 \mathcal{L}_m + \frac{53}{4} C_F S_1 \mathcal{L}_m \\
& + \frac{C_A C_F \Delta_M S_1 \mathcal{L}_m}{m} + \frac{358 C_A C_F^2 \Delta_{m,M} S_1 \mathcal{L}_m}{3m} + \frac{58 C_A C_F \Delta_{m,M} S_1 \mathcal{L}_m}{9m} + \frac{185 C_F \Delta_{m,M} S_1 \mathcal{L}_m}{3m} + \frac{3}{2} C_A C_F n_f T_f \mathcal{L}_m \\
& + \frac{C_A C_F n_f T_f \mathcal{L}_m}{\varepsilon_{UV}} + \frac{8 C_A C_F \Delta_{m,M} n_f T_f \mathcal{L}_m}{3m} - \frac{2}{3} C_A C_F n_f S_1 T_f \mathcal{L}_m - \frac{32 C_A C_F \Delta_{m,M} n_f S_1 T_f \mathcal{L}_m}{9m} - \frac{73 C_A C_F^2 \mathcal{L}_m}{4\varepsilon_{UV}} \\
& + \frac{C_A C_F \mathcal{L}_m}{4\varepsilon_{UV}} - \frac{105 C_F \mathcal{L}_m}{8\varepsilon_{UV}} - \frac{2601 C_A C_F^2 \Delta_{m,M} \mathcal{L}_m}{16m} - \frac{565 C_A C_F \Delta_{m,M} \mathcal{L}_m}{12m} - \frac{1047 C_F \Delta_{m,M} \mathcal{L}_m}{16m} + \frac{307 C_A C_F^2}{16\varepsilon_{UV}} - \frac{21 C_A C_F}{16\varepsilon_{UV}} \\
& + \frac{343 C_F}{32\varepsilon_{UV}} + \frac{C_A C_F \Delta_M}{2m} + \frac{48925 C_A C_F^2 \Delta_{m,M}}{192m} + \frac{497 C_A C_F \Delta_{m,M}}{16m} + \frac{29123 C_F \Delta_{m,M}}{192m} + \frac{82 C_A C_F^2 \Delta_{m,M}}{\varepsilon_{UV} m} + \frac{62 C_A C_F \Delta_{m,M}}{3\varepsilon_{UV} m} \\
& + \frac{32 C_F \Delta_{m,M}}{\varepsilon_{UV} m} + \frac{73 C_A C_F^2}{8\varepsilon_{UV}^2} - \frac{C_A C_F}{8\varepsilon_{UV}^2} + \frac{105 C_F}{16\varepsilon_{UV}^2} \tag{D.8}
\end{aligned}$$

## D.5 HPET field at $M \neq 0$ :

The contributions from two-loop wavefunction corrections,  $F_h^{(M)}$ , with unevaluated MIs are too large to present here. We thus include the full expressions with description in an ancillary file.

## D.6 Parametric Integrals

$$\begin{aligned}
P(z) &= - \int_0^1 dx \left\{ 2(1-x) \log \left( \frac{1-x+z^2x^2}{1-x} \right) + \frac{4z^2x(1-x^2)}{1-x+z^2x^2} \right\} \\
&= \frac{3}{z^2} + \left\{ \frac{3}{2z^4} - 3 \right\} \log z^2 + \frac{(3-6z^2-12z^4)}{z^4\sqrt{1-4z^2}} \tanh^{-1}(\sqrt{1-4z^2}) \tag{D.9}
\end{aligned}$$

$$\begin{aligned}
P'(z) &= - \int_0^1 dx \left\{ (1-x) \log \left( \frac{1-x+z^2x^2}{1-x} \right) - \frac{2z^2x(1-x)(2-x)}{1-x+z^2x^2} \right\} \\
&= \frac{3}{2z^2} - \left\{ \frac{3}{z^2} - \frac{3}{4z^4} - \frac{3}{2} \right\} \log z^2 + \frac{(3-6z^2)\sqrt{1-4z^2}}{2z^4} \tanh^{-1}(\sqrt{1-4z^2}) \tag{D.10}
\end{aligned}$$

$$\begin{aligned}
S(z) &= \int_0^1 dx \left\{ (3x^2-6x+4) \log \left( \frac{1-x+z^2x^2}{1-x} \right) - \frac{z^2x(1-x^2)}{(1-x)(2-x)^2} \right\} \\
&= -\frac{1}{z^2} + \left\{ \frac{3}{2z^2} - \frac{1}{2z^4} \right\} \log z^2 + \frac{(z^2-1)\sqrt{1-4z^2}}{z^4} \tanh^{-1}(\sqrt{1-4z^2}) \tag{D.11}
\end{aligned}$$

$$\begin{aligned}
S'(z) &= - \int_0^1 dx \left\{ \frac{z^2x^3}{1-x+z^2x^2} \right\} \\
&= -\frac{1}{z^2} + \left\{ \frac{1}{2z^2} - \frac{1}{2z^4} \right\} \log z^2 + \frac{3z^2-1}{z^4\sqrt{1-4z^2}} \tanh^{-1}(\sqrt{1-4z^2}) \tag{D.12}
\end{aligned}$$

The integrals can be analytically continued in the regime,  $4z^2 \geq 1$ , using  $\sqrt{1-4z^2} \mapsto i\sqrt{4z^2-1}$  and therefore  $\tanh^{-1}(\sqrt{1-4z^2}) \mapsto i \tanh^{-1}(\sqrt{4z^2-1})$ , and plugging this back into the integrals one can verify that the integral remains real.

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