

DESY SR 87-04
July 1987

STRUCTURE DETERMINATION OF THE CoSi_2 :Si(111) INTERFACE
BY X-RAY STANDING WAVE ANALYSIS

by

A. E. M. J. Fischer, E. Vlieg, J. F. van der Veen
FOM-Institute for Atomic and Molecular Physics, Amsterdam

M. Clausnitzer
II. Inst. f. Experimentalphysik, Universität Hamburg

G. Materlik
Hamburger Synchrotronstrahlungslabor HASYLAB at DESY

Eigentum der Property of	DESY	Bibliothek library
Zugang Accession:	10. AUG. 1987	
Leihfrist: Loan period:	7	Tage days

ISSN 0723-7979

DESY behält sich alle Rechte für den Fall der Schutzrechtserteilung und für die wirtschaftliche Verwertung der in diesem Bericht enthaltenen Informationen vor.

DESY reserves all rights for commercial use of information included in this report, especially in case of filing application for or grant of patents.

To be sure that your preprints are promptly included in the
HIGH ENERGY PHYSICS INDEX,
send them to the following address (if possible by air mail):

**DESY
Bibliothek
Notkestrasse 85
2 Hamburg 52
Germany**

**STRUCTURE DETERMINATION OF THE CoSi_2 :Si(111) INTERFACE
BY X-RAY STANDING WAVE ANALYSIS**

A.E.M.J. Fischer, E. Vlieg and J.F. van der Veen
FOM-Institute for Atomic and Molecular Physics,
Kruislaan 407, 1098 SJ Amsterdam, The Netherlands

and

M. Cläusnitzer

II. Institut für Experimentalphysik, Universität Hamburg, Luruper Chaussee 149,
D-2000 Hamburg 52, F.R.G.

and

G. Materlik

Hamburger Synchrotronstrahlungslabor HASYLAB at DESY,
Notkestrasse 85, D-2000 Hamburg 52, F.R.G.

ISSN 0723-7979

submitted to Phys. Rev. B

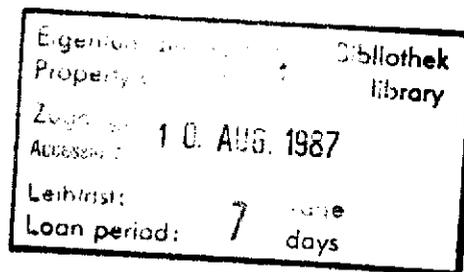
Abstract

The atomic structure at the interface of CoSi_2 films on Si(111) has been investigated with X-ray standing waves. The Co atoms at the interface are 5-fold coordinated. The bonds across the interface are dilated by $0.05 \pm 0.03 \text{ \AA}$. For the range of film thicknesses studied (9-28 \AA), the CoSi_2 lattice was measured to be almost free of strain in the direction perpendicular to the interface plane.

1. Introduction

Epitaxial silicide layers of the near-noble metals Co and Ni can be grown on Si(111) with a high degree of perfection and with abrupt interfaces [1]. These heterostructures are extensively studied because of their fundamental interest and possible applications in diversified fields [2]. They serve as model systems for Schottky barrier formation [3]; the Schottky barrier height may directly be correlated to the electronic bonding arrangement at the interface, provided the geometric interface structure is accurately known. In this paper we present a structural investigation of the CoSi_2 :Si(111) interface with X-ray standing waves (XSW) using synchrotron radiation. In studies of epitaxial layers, this method has previously been applied to determine the atomic bonding arrangement at the NiSi_2 :Si(111) interface [4].

The silicides CoSi_2 and NiSi_2 have the cubic CaF_2 structure with bulk lattice parameters smaller than that of Si by 1.2% and 0.46%, respectively. For CoSi_2 the commonly observed epitaxial relationship with Si(111) is one in which the film is (111) oriented but 180° rotated about the [111] substrate normal, labeled as type-B epitaxy [5]. The silicide may also have the same in-plane orientation as the Si(111) substrate (type-A epitaxy), but it appears that such a heterostructure can only be obtained free of B-type grains for NiSi_2 [1]. In modeling the corresponding interface structures one is left with only two distinct possibilities, if the tetrahedral coordination of the Si atoms is to be preserved in passing over from the silicon to the silicide lattice. Figure 1 shows these two possibilities for a B-type film. The two models are distinguished by the silicide being terminated with either 5-fold or 7-fold coordinated metal atoms, as opposed to 8-fold coordination for metal atoms in bulk disilicide. For the NiSi_2 :Si(111) system the results of various techniques agree on both A- and B-type interfaces being 7-fold coordinated [4,6,7]. The chemical



and structural similarity of nickel and cobalt silicides has led several authors to believe that also for the $\text{CoSi}_2:\text{Si}(111)$ interface a 7-fold coordination is favoured [8,9]. But by comparing measured cross-section TEM lattice images with calculated ones, Gibson et al. [10] identified 5-fold coordination as the most likely model (a definite structure assignment was hampered by uncertainties in the TEM image parameters). Recently we reported preliminary ion scattering data which gave evidence for a 5-fold coordinated $\text{CoSi}_2:\text{Si}(111)$ interface [11]. The XSW measurements presented here rule out the 7-fold model and show that the model based on 5-fold coordination is indeed correct. Furthermore, the (111) interplanar distance has been determined at the interface and in the CoSi_2 overlayer. Across the interface this distance is slightly dilated. In the overlayer lattice the interplanar distance is equal to that in bulk CoSi_2 , i.e. essentially free of strain for the range of thicknesses studied (9-28 Å).

2. X-ray standing wave analysis

With regard to the XSW method, the relevant difference between 5- and 7-fold coordination is the distance d_{IF} of the first Co layer to the last (111) diffraction plane in the substrate (fig. 1). The (111) interplanar spacing is $d_{\text{H}}=3.135$ Å for Si and 3.097 Å for bulk CoSi_2 . The bond across the interface can be assumed to be part of either the silicon substrate or the silicide; in what follows the former is assumed. Then, using bulk bond lengths, the interface distance is calculated to be 2.74 Å and 3.52 Å for unrelaxed 5- and 7-fold coordinated interfaces, respectively.

XSW is by now well-established as a structure probe of surface adsorbates [12] and epitaxial overlayer systems [4]. When a monochromatic X-ray beam is Bragg-reflected from the Si substrate lattice, the interference between the incident and reflected waves gives rise to a standing wave field which extends across the interface. According to the dynamical theory of X-ray diffraction the phase of the standing wave field changes by π radians as the crystal is rocked through the Bragg reflection [13]. On the low-angle side of the reflectivity curve the anti-nodal planes of the standing wave field lie halfway between the diffraction planes and move with increasing angle towards a position coinciding with the diffraction planes. The positions of the overlayer lattice planes are determined by measuring as a function of rocking angle the fluorescence radiation yield emitted by the Co atoms, which is proportional to the local E field intensity of the standing wave field.

Henceforth the overlayer is assumed to be composed of N lattice planes of the fluorescing element Co with a spacing d_{OL} and starting at a distance d_{IF} from the topmost diffraction plane of the substrate (fig. 1). For convenience we introduce the structural parameters $\delta_{\text{IF}} \equiv d_{\text{IF}}/d_{\text{H}}$ and $\gamma \equiv d_{\text{OL}}/d_{\text{H}}-1$, as in ref.[4].

In principle, any periodic overlayer influences the standing wave pattern generated in the substrate. Here this effect can be neglected, since the layers are very thin compared to the extinction depth in the layer material. Under this condition the total fluorescence yield $Y(\theta)$ emitted by the N overlayer planes as a function of reflection angle θ is given by [4]

$$Y(\theta)/NC = 1 + R(\theta) + 2\sqrt{R(\theta)} \cdot F \cos[v(\theta) - 2\pi D] , \quad (1)$$

where the reflectivity $R(\theta)$ and the phase $v(\theta)$ can be calculated from dynamic diffraction theory [13]. D and F are defined by

$$D \equiv \delta_{\text{IF}} + (N-1)\gamma/2 ; \quad (1a)$$

$$F \equiv f_c \sin(\pi N\gamma)/N \sin(\pi\gamma), \quad (1b)$$

and C includes constant geometry parameters. If the fluorescence selected overlayer atoms occupy just one site coherently, the parameter f_c in eq.(1b) gives the coherent fraction, i.e., the fraction of overlayer atoms which are positioned exactly in the N planes. The parameters D and F describe all possible yield curves and are determined experimentally by fitting the measured yield curves to eq.(1). The quantities δ_{IF} and γ then follow from determining D for (at least) two different coverages N and solving eq.(1a).

Once γ is known, the coherent fraction f_c for each sample can be calculated from eq.(1b). The coherent fraction is a measure of the crystalline quality of the overlayer.

The above XSW analysis depends on precise knowledge of the number of overlayer planes N. Since N cannot be uniquely determined from the shape of a fluorescence yield curve, it should be acquired from another independent method. Here Rutherford backscattering spectrometry (RBS) is employed.

3. Experimental

The silicide overlayers were prepared in ultrahigh vacuum at the FOM-Institute, where also the RBS analyses were performed. The XSW measurements were performed at HASYLAB. The Si(111) substrates ($6 \times 10 \times 0.38 \text{ mm}^3$ in dimension) were cleaned and annealed following the procedure outlined in ref. 4. Cobalt was then deposited at room temperature by sublimation from resistively heated Co wires. The deposition rate was $\sim 5 \times 10^{13} \text{ atoms cm}^{-2} \text{ s}^{-1}$. The pressure rose to $7 \times 10^{-8} \text{ Pa}$ during deposition, but recovered quickly to base pressure ($7 \times 10^{-9} \text{ Pa}$) once the deposition was stopped.

It is well-known that upon heating to $\sim 550^\circ\text{C}$ for a few minutes a solid-state reaction takes place at the interface between the deposited Co layer and the Si substrate under UHV conditions. An epitaxial CoSi_2 layer is then formed [14]. However, unlike the Ni:Si(111) system, for thin films such a treatment generally yields discontinuous CoSi_2 layers with sizeable holes in them (pinholes). This introduces a large uncertainty in the number of Co planes N contributing to the fluorescence yield. In the present study a substantial improvement in layer morphology was obtained by first depositing at room temperature a thin layer of Si on top of the unreacted Co film and then reacting the resulting Si/Co/Si(111) sandwich structure at elevated temperature. The Si layers were deposited *in situ* with use of an e-beam evaporator and were typically chosen to be of the same thickness as the Co films. The sandwich structures were heated for 3 minutes isochronically at stepwise increased temperatures starting at 300°C . In-situ reflection high-energy electron diffraction (RHEED) was used after each anneal step to verify the resulting surface periodicity. Once the surface exhibited the sharp 1×1 RHEED pattern which is typical for the formation of a well-ordered epitaxial disilicide [15], the annealing procedure was stopped. The temperature at which the 1×1 pattern appeared was $\sim 400^\circ\text{C}$.

In the above manner we have prepared overlayer systems of different thickness, containing, on average, $N \cong 4.7, 7.9,$ and 9.2 Co planes. One sample, with $N = 3.0$, was prepared by annealing a sandwich structure having a top Si layer 10 times thicker than the Co layer at a temperature of 550°C . The N value for each sample was determined by measuring with RBS the areal density of deposited Co atoms and noting that a single Co plane in bulk silicide ($N = 1$) contains $8.02 \times 10^{14} \text{ atoms cm}^{-2}$. Clearly, the conversion of areal densities into N values is accurate only if the CoSi_2

layer is reasonably uniform and covers most of the substrate surface area. In order to check upon the latter point, the overlayer morphology has been studied using high-resolution ion backscattering [16] and high-resolution transmission electron microscopy (TEM) in planar and cross-sectional view. All silicide layers are found to be continuous and free of pinholes, justifying the direct conversion of areal density to number of Co planes. For the layers under investigation the epitaxial relationship observed is entirely type B, in good agreement with other results [17]. Details on morphology and structure of the CoSi_2 layers will be published elsewhere. The TEM micrographs show that for $N > 4.7$ the interface contains steps of height $3d_H$. On either side of these distinct steps the interface is structurally perfect and flat. The inclusion of these steps in structure models of the interface hardly affects the XSW analysis and can therefore be ignored.

In order to protect the ultrathin silicide overlayers from contamination during transport to HASYLAB and during XSW analysis (all in air), the samples were capped by an amorphous layer of Si with a thickness of 50 to 80 Å. (The amorphous caps were "densified" by annealing at 200°C for 2 min.) The presence of this layer does not influence the XSW analyses.

The experimental arrangement at HASYLAB for XSW measurements has been described elsewhere [18,19,20]. With use of an asymmetrically cut double-crystal monochromator the synchrotron radiation from DORIS was collimated at an energy of 14.2 keV and directed onto the sample. Scans through the rocking curve were made by varying the energy of the monochromatized beam, while the sample remained fixed in position. The Co K_α fluorescence yield (at 6.9 keV photon energy) was measured with a Si(Li) solid-state detector at a distance of 10-20 mm from the sample. The (111) Bragg-reflected X-ray intensity was detected by a NaI scintillation counter. The quality of each sample as generator of standing waves was examined through high-resolution double-crystal topographs [21]. Perfect regions on the sample were then selected for XSW measurements. The reflectivity and fluorescence yield were measured simultaneously as a function of $\lambda - \lambda_B$, where λ_B is the wavelength determined by Bragg's equation $2d_H \sin \theta_B = \lambda_B$.

4. Results and discussion

A typical reflectivity curve is shown in fig. 2a. The angular profiles of the measured fluorescence yields are shown in fig. 2(b-e) in order of increasing overlayer thickness. For each of

the samples the parameters D and F (eqs.(1a) and (1b)) were determined by least-squares fitting the theoretical expression for the fluorescence yield given by eq.(1) to the experimental data. Prior to the fits the theoretical yield curves were convoluted with the experimental resolution function, as determined from the shape of the reflectivity curve [22]. The results of the fits are shown in fig. 2(b-e) as solid curves. The corresponding optimal values for D and F are listed in table 1. The D values are plotted in fig. 3 against N-1. The data are consistent with a linear dependence on N-1, as in eq.(1a). The straight solid line in fig. 3 represents the best fit to the data points, which is obtained for

$$\delta_{IF} = 0.891 \pm 0.011 ,$$

and

$$\gamma = (-11 \pm 5) \times 10^{-3} .$$

Note, that this δ_{IF} value corresponds to the intercept of the solid line with the ordinate axis and that $\gamma/2$ is given by the slope of the line. For comparison, fig. 3 shows also the expected dependencies of D on layer thickness for unrelaxed 5- and 7-fold coordinated interfaces, assuming for both interface models a bulklike interplanar spacing in the overlayer (dashed lines). The data are very close to the prediction for 5-fold coordination. Clearly, a 7-fold coordinated interface is ruled out. The ease with which the two interface models can be distinguished from one another is illustrated for the sample with N = 7.9 in fig. 2d, where the angle-dependent fluorescence yields for the unrelaxed 5- and 7-fold coordinated interfaces are shown as dashed curves.

The sensitivity of the XSW method to the structural parameters δ_{IF} and γ depends on the perfection of the CoSi_2 overlayers. The perfection in turn is described by the coherent fraction f_c . The measured coherent fractions are listed in table 1 in order of increasing overlayer thickness. They have been derived from eq.(1b). Since both γ and N are very small for the samples under investigation, f_c is equal to F. The coherent fractions ranged from 0.8 to 0.4, with the exception of the thickest overlayer (N = 9.2), for which f_c is only 0.19. The values suggest a correlation with layer thickness and possibly with annealing temperature. Presumably, for the thickest overlayer the annealing procedure described in section 3 does not adequately order the lattice. The much higher coherent fractions for the other samples are still lower than those for NiSi_2 overlayers [4]. The

lower quality is possibly related to local disorder in the CoSi_2 lattice at atomic steps at the interface. This corroborates the findings of other studies [17] that CoSi_2 layers of good crystal quality are more difficult to grow.

The overlayers were found to be completely relieved of strain along the direction perpendicular to the interface plane. If the CoSi_2 layers were to grow pseudomorphic on Si(111), as is the case for NiSi_2 layers of similar thickness [23], then a substantial perpendicular strain would be present. Assuming the same ratio of perpendicular to parallel strain as for NiSi_2 [4], a pseudomorphic growth of CoSi_2 would correspond to $\gamma = -23 \times 10^{-3}$. On the other hand, complete absence of strain implies $\gamma = -12 \times 10^{-3}$. These values are to be compared with the experimental one, $\gamma = (-11 \pm 5) \times 10^{-3}$. It is concluded that the 1.2% mismatch between CoSi_2 and Si is largely accommodated by defects (e.g. dislocations) in the layer rather than by lattice strain. It was noted in section 3 that steps are present at the interface. Their existence are known to be one of the mechanisms for strain relief [10].

The lowered crystal symmetry at the 5-fold coordinated interface may give rise to a local lattice distortion in the form of bond-length or bond-angle changes. In our XSW analysis such lattice distortions would be evident as a change in d_{IF} . A distance parameter of $\delta_{IF} = d_{IF}/d_H = 0.891 \pm 0.011$ was measured, yielding $d_{IF}(\text{exp}) = 2.79 \pm 0.03 \text{ \AA}$. Comparing this value with $d_{IF} = 2.74 \text{ \AA}$ expected for an unrelaxed 5-fold interface, we conclude that the interface bonding arrangement, if distorted at all, is such that d_{IF} is dilated by $0.05 \pm 0.03 \text{ \AA}$. This may be the result of an expansion of the Si-Co bond across the interface and/or an angular distortion of the Si-Si bonds at the substrate side of it (fig. 1a).

In summary, the geometric structure of the epitaxial CoSi_2 :Si(111) interface has been determined by X-ray standing wave analysis. The Si atoms in the top layer of the Si(111) substrate are bonded to the Co atoms in the silicide, corresponding to 5-fold coordination of the Co atoms at the interface. This is in striking contrast with XSW [4] and ion scattering [7] results for NiSi_2 :Si(111), which both gave conclusive evidence for 7-fold coordination at that interface. The reason for the difference in coordination number between Co and Ni is not yet known. Quantum chemical calculations on cluster models of the different interfaces are now in progress in an attempt to resolve this issue [24].

Acknowledgements

We thank C.W. Bulle-Lieuwma from the Philips Research Laboratories for performing the TEM-analyses. Three of us (Fischer, Vlieg and van der Veen) are grateful for the hospitality extended to them while at HASYLAB. This work is part of the research program of the Stichting voor Fundamenteel Onderzoek der Materie (Foundation for Fundamental Research on Matter) and was made possible by financial support from the Nederlandse Organisatie voor Zuiver-Wetenschappelijk Onderzoek (Netherlands Organization for the Advancement of Pure Research).

References

- [1] R.T. Tung, J.M. Gibson, and J.M. Poate, *Appl. Phys. Lett.* **42**, 888 (1983).
- [2] E. Rosencher, S. Delage, Y. Campidelli, and F. Arnaud d'Avitaya, *Electr. Lett.* **20**, 762 (1984).
- [3] R.T. Tung, *Phys. Rev. Lett.* **52**, 461 (1984).
- [4] E. Vlieg, A.E.M.J. Fischer, J.F. van der Veen, B.N. Dev, and G. Materlik, *Surf. Sci.* **178**, 36 (1986).
- [5] R.T. Tung, J.M. Gibson, and J.M. Poate, *Phys. Rev. Lett.* **50**, 429 (1983).
- [6] D. Chems, G.R. Anstis, J.L. Hutchison, and J.C.H. Spence, *Phil. Mag.* **464A**, 849 (1982).
- [7] E.J. van Loenen, J.W.M. Frenken, J.F. van der Veen, and S. Valeri, *Phys. Rev. Lett.* **54**, 827 (1985).
- [8] C. Pirri, J.C. Peruchetti, D. Bolmont, and G. Gewinner, *Phys. Rev. B* **33**, 4108 (1986).
- [9] S.A. Chambers, S.B. Anderson, H.W. Chen, and J.H. Weaver, *Phys. Rev. B* **34**, 913 (1986).
- [10] J.M. Gibson, J.C. Bean, J.M. Poate, and R.T. Tung, *Appl. Phys. Lett.* **41**, 818 (1982).
- [11] J.F. van der Veen, and A.E.M.J. Fischer, in *Proceedings of the 18th International Conference on the Physics of Semiconductors, August 1986, Stockholm*, edited by O. Engström (World Scientific Publ. Co, Singapore, 1987).
- [12] G. Materlik, J. Zegenhagen, and W. Uelhoff, *Phys. Rev. B* **32**, 5502 (1985). For a

- summary of recent work, see: G. Materlik, *Z. Phys. B-Condensed Matter* **61**, 405 (1986).
- [13] Z.G. Pinsker, *Dynamical Scattering of X-rays in Crystals*, (Springer, Berlin-Heidelberg-New York, 1978).
- [14] F. Arnaud d'Avitaya, S. Delage, E. Rosencher, and J. Derrien, *J. Vac. Sci. Technol. B* **3**, 770 (1985).
- [15] S. Saitoh, H. Ishiwara, and S. Furukawa, *Appl. Phys. Lett.* **37**, 203 (1980).
- [16] J.F. van der Veen, *Surf. Sci. Rept.* **5**, 199 (1986).
- [17] R.T. Tung, A.F.J. Levi, and J.M. Gibson, *Appl. Phys. Lett.* **48**, 635 (1986).
- [18] P. Funke, and G. Materlik, *Solid State Comm.* **54**, 921 (1985).
- [19] A. Krolzig, G. Materlik, and J. Zegenhagen, *Nucl. Instr. Meth.* **208**, 618 (1983).
- [20] A. Krolzig, G. Materlik, M. Swars, and J. Zegenhagen, *Nucl. Instr. Meth.* **219**, 430 (1984).
- [21] U. Bonse, and E. Kappler, *Z. Naturforsch.* **13a**, 792 (1958).
- [22] N. Hertel, G. Materlik, and J. Zegenhagen, *Z. Physik B* **58**, 199 (1985).
- [23] E.J. van Loenen, A.E.M.J. Fischer, J.F. van der Veen, and F. Legoues, *Surf. Sci.* **154**, 52 (1985).
- [24] P. J. van den Hoek, private communication.

Table 1

The parameters D and F measured on $\text{CoSi}_2:\text{Si}(111)$ samples with different number of Co planes N in the silicide. The error in the number of Co planes is determined by the accuracy of the RBS measurements. The coherent fraction f_c of Co atoms is approximately equal to F. Sample nr. 1 was annealed at 550°C , the other samples were annealed at 400°C .

Sample	nr. of Co planes N	D	F = f_c
1	3.0 ± 0.2	0.883 ± 0.002	0.79 ± 0.01
2	4.7 ± 0.3	0.859 ± 0.006	0.40 ± 0.02
3	7.9 ± 0.5	0.863 ± 0.006	0.46 ± 0.01
4	9.2 ± 0.6	0.847 ± 0.016	0.19 ± 0.03

Figure captions

- Fig.1 Ball-and-stick models of the B-type $\text{CoSi}_2:\text{Si}(111)$ interface in a projected view down the $[\bar{1}\bar{1}0]$ direction. Open and filled circles represent Si and Co atoms, respectively. The Co atoms are (a) 5-fold or (b) 7-fold coordinated at the interface.
- Fig.2 Angular profile of the reflectivity and the Co K_α fluorescence yield for $\text{CoSi}_2:\text{Si}(111)$ samples of different thickness. The reflection angle scale refers to the crystal angle of the monochromator. The circles are the data points, the solid lines are the best fits. The data were obtained using the (111) Bragg reflection at an X-ray energy of 14.2 keV. Curves *b* to *e* have an accumulative vertical offset by 1 unit. Curve *a*: typical reflectivity curve of the $\text{Si}(111)$ substrate, taken on the sample with $N = 7.9$. Curves *b* to *e*: fluorescence curves for samples with overlayer thicknesses of (b) $N = 3.0$, (c) $N = 4.7$, (d) $N = 7.9$, and (e) $N = 9.2$. The dash-dotted and dashed curves represent the yields calculated for unrelaxed 5- and 7-fold coordinated interfaces, assuming for both interface models a bulklike interplanar spacing in an overlayer, containing $N = 7.9$ planes of Co atoms.
- Fig.3 The fitted D-values as a function of overlayer thickness $N-1$. Closed circles denote the data points. The solid line represents a least squares fit to the data points. The dashed lines show the expected dependencies of D on $N-1$ for unrelaxed 7- and 5-fold coordinated interfaces, assuming for both interface models a bulklike interplanar spacing in the overlayer.

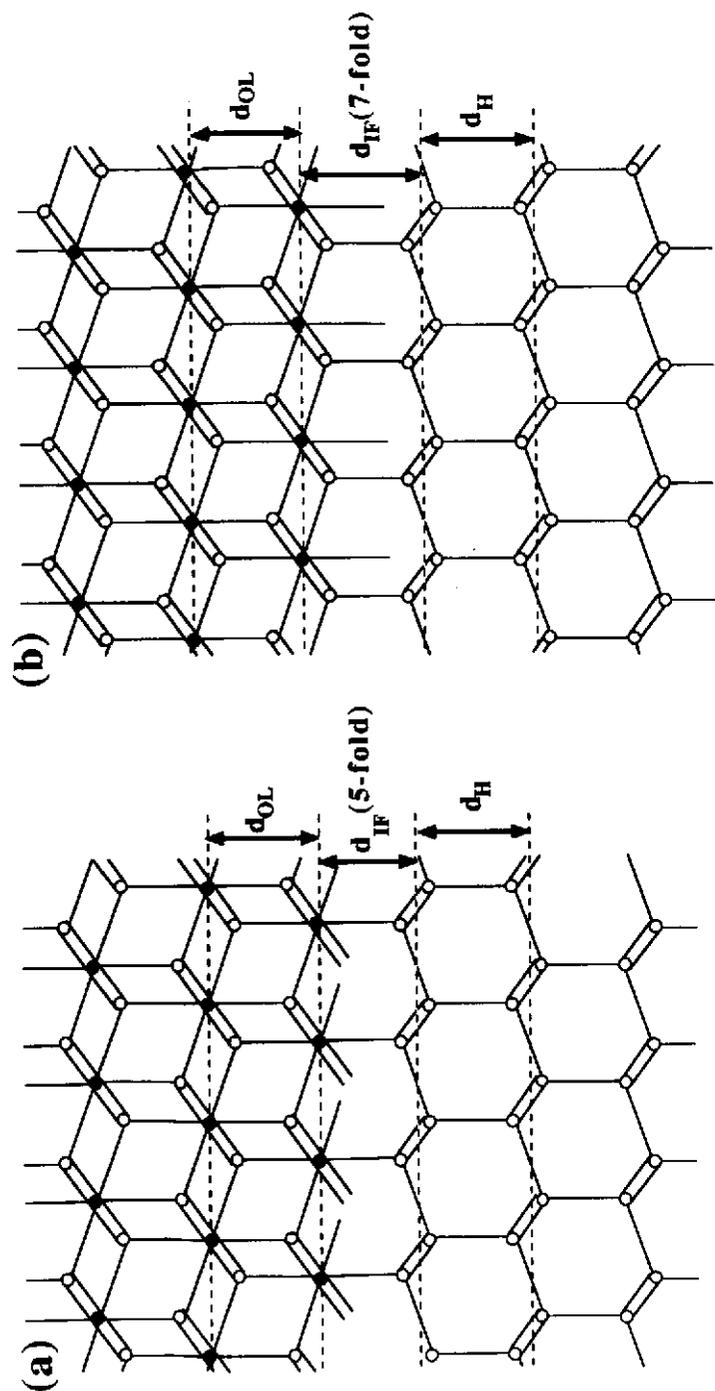


FIG. 1

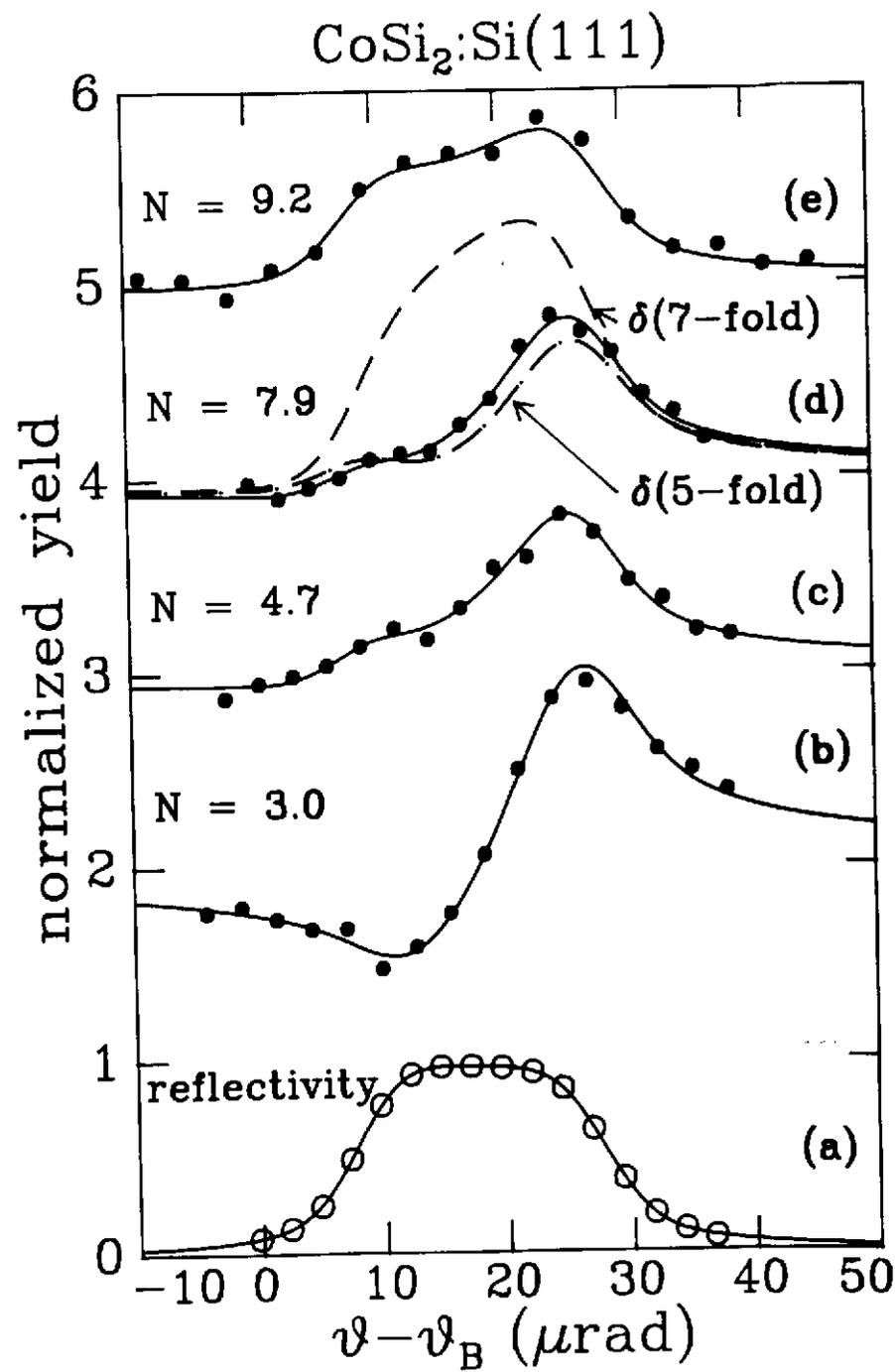


FIG. 2

CoSi₂:Si(111)

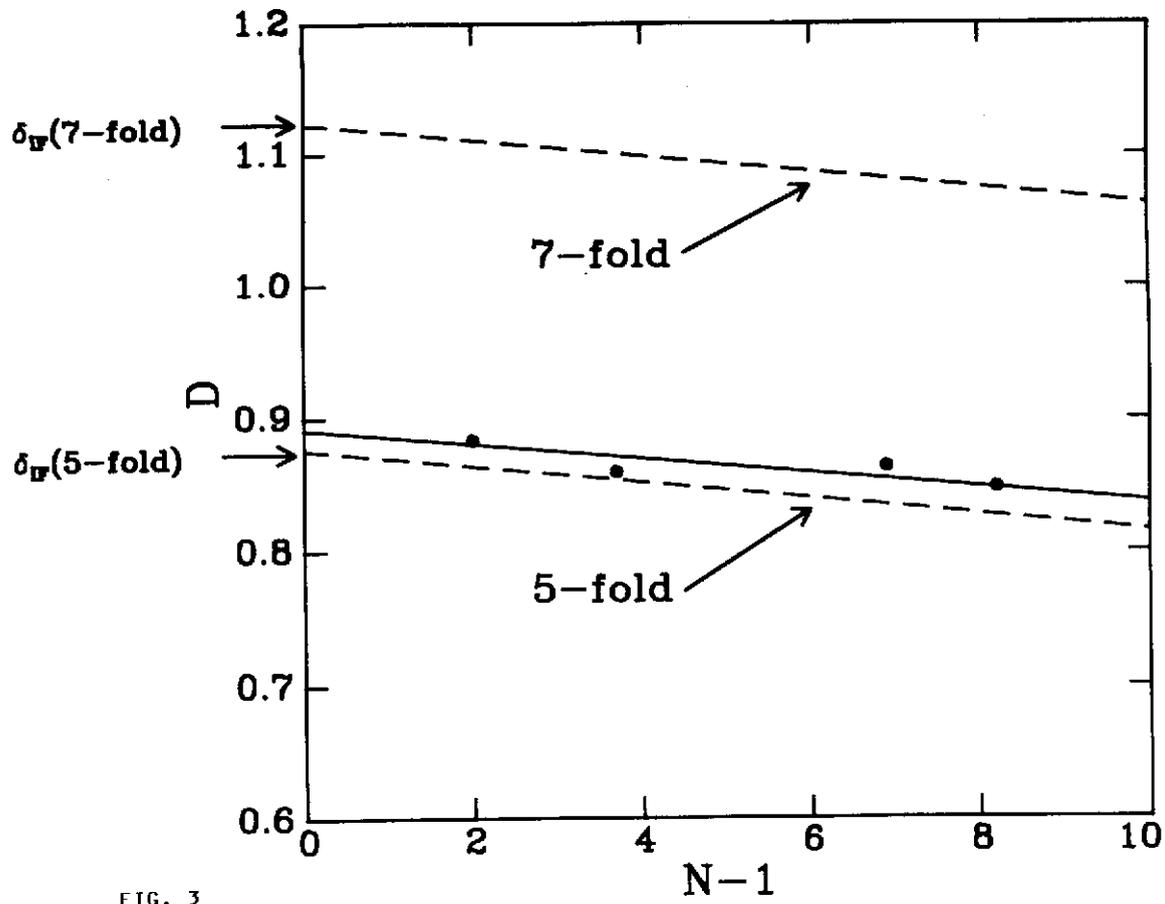


FIG. 3