PHYSICS LETTERS

## AN UPPER LIMIT ON THE MASS OF THE TAU NEUTRINO

The ARGUS collaboration

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Using the ARGUS detector at the  $e^+e^-$  storage ring DORIS II, we have measured the  $\nu_{\tau}$  energy spectrum in the decay  $\pi^+\pi^-\pi^\pm\nu_{\tau}$  of  $\tau$  leptons produced near  $\sqrt{s} = 10$  GeV From this energy spectrum, we derive an upper limit of  $m(\nu_{\tau}) < 70$  MeV/ $c^2$  at the 95% confidence level

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0370-2693/85/\$ 03.30 © Elsevier Science Publishers B.V. (North-Holland Physics Publishing Division) Upper limits on the mass of the tau neutrino have been determined previously from the energy spectrum of the  $\tau$  decay modes with one charged particle [1] and from the invariant mass spectrum of the decay modes  $\pi^{+}\pi^{-}\pi^{\pm}\pi^{0}\nu_{\tau}$  [2,3] and  $K^{+}K^{-}\pi^{\pm}\nu_{\tau}$  [4]. The best limit, reported in ref. [3], is 143 MeV/ $c^{2}$  at the 95% confidence level.

In this letter we report a new limit from the study of about 1500  $\pi^+\pi^-\pi^\pm\nu_{\tau}$  decays of  $\tau$  leptons produced in  $e^+e^-$  annihilation near  $\sqrt{s} = 10$  GeV. High statistics and good momentum resolution allow us to improve the limit on  $m(\nu_{\tau})$  by a factor of two using an analysis of the energy spectrum of the three-pion system. The data have been collected with the ARGUS detector at the electron—positron storage ring DORIS II at DESY. The centre-of-mass energy varied from 9.4 to 10.6 GeV.

A short description of the detector and the trigger conditions is given in ref. [5]. The event sample used in this analysis corresponds to an integrated luminosity of  $61.4 \text{ pb}^{-1}$ . Tau pair events of the 1-3 topology were selected by requiring:

- Exactly four charged particles from the main vertex with a total charge zero, and no more than two additional charged particles;

- the momentum sum  $\sum_{i=1}^{4} |p_i| \ge 2.7 \text{ GeV/}c$ , to suppress beam-gas and photon-photon reactions;

- the momentum sum  $\sum_{i=1}^{4} |p_i| \leq 0.92 \sqrt{s}$ , to suppress exclusive events like  $\Upsilon' \rightarrow \pi^+ \pi^- \ell^+ \ell^-$ ;

- a hemisphere cut  $\cos \theta_{1i} \leq 0$ , where  $\theta_{1i}$  is the angle between particle 1 and particle *i*, *i* = 2, 3, 4;

- an opening angle of less than  $90^{\circ}$  between each pair of particles on the three-prong side;

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- a polar angle cut of  $|\cos \theta_1| \le 0.75$  on the oneprong  $\tau$  decay, to ensure good momentum resolution and trigger conditions;

- no photons with  $E_{\gamma} \ge 50$  MeV in the shower counters, for an efficient suppression of the  $\pi^{+}\pi^{-}\pi^{\pm}\pi^{0}\nu_{\tau}$  mode, or exactly one  $\pi^{0}$ , with an opening angle with respect to the charged pion on the oneprong side of less than 90°, and which, when combined with this track, yields a  $\rho$ -meson candidate with mass between 0.57 and 1.07 GeV/ $c^{2}$ , and momentum larger than 0.9 GeV/c;

- sufficiently large opening angles  $\cos \theta_{ij} \le 0.992$ to reject radiative Bhabha and  $\mu\mu\gamma$  events with a converted photon (*i*, *j* are the opposite sign particles on the three-prong side);

- the energy deposited by particle 1, and the sum of that deposited by particles 2 to 4 in the electromagnetic calorimeter should be less than 4 GeV, to further suppress radiative Bhabha events;

- agreement with the pion hypothesis from timeof-flight and dE/dx measurements for all three particles i = 2, 3, 4 and with either the electron, muon, pion or kaon hypothesis for particle 1.

One exclusive decay  $\Upsilon' \rightarrow \pi^+\pi^-\Upsilon \rightarrow \pi^+\pi^-\mu^+\mu^$ passed these cuts but was reconstructed unambiguously and rejected.

The 1536 events which satisfy these selection criteria are predominantly three-prong  $\tau$  decays into  $\pi^+\pi^-\pi^\pm\nu_{\tau}$ . The background from KK $\pi\nu$  and K $\pi\pi\nu$  decays [4] is estimated to be less than 4%. Assuming all hadrons to be pions, we obtain the invariant mass spectrum shown in fig. 1. This spectrum is not cor-



Fig. 1. Invariant mass of the three charged pions from the decay  $\tau^{\pm} \rightarrow \pi^{+}\pi^{-}\pi^{\pm}\nu_{\tau}$ , uncorrected for acceptance.



Fig. 2. The tau-neutrino energy spectrum for the decay  $\tau^{\pm} \rightarrow \pi^{+}\pi^{-}\pi^{\pm}\nu_{\tau}$ , uncorrected for acceptance.

rected for acceptance. The broad structure visible is found to decay mainly into  $\rho^0 \pi^{\pm}$ .

For further studies we rejected the few events with the three-pion invariant mass above  $m_{\tau}$ . Fig. 2 shows the  $\nu_{\tau}$  energy spectrum derived from

$$E_{\nu} = E_{\text{beam}} - \sum_{i=2}^{4} E_i$$
,

where the energies of the three decay hadrons are calculated from the measured momenta using the pion mass assignment. The shape of the spectrum near  $E_v$ 



Fig. 3. Experimental data and Monte Carlo expectations for  $m(\nu_{\tau}) = 0$  and 140 MeV/ $c^2$  in the mass-sensitive region normalized to the total number of events in the entire  $E_{\nu}$  spectrum.

= 0 depends sensitively on the mass of the tau neutrino. It depends also on the invariant mass distribution of the three-pion system. However, this distribution is measured in the experiment and there is no need for model-dependent assumptions as required in other approaches [2,3]. In fig. 3, we compare the experimental  $E_{\mu}$  spectrum in the mass sensitive region with the predictions of a Monte Carlo calculation assuming various neutrino masses and a constant level of background. The background in the spectrum is estimated from the number of events with  $m_{3\pi} > m_{\tau}$ to be 5 events/GeV. The calculation uses the observed mass spectrum of fig. 1 with masses less than  $m_{\pi}$ , assumes that  $v_{\tau}$  is produced isotropically in the  $\tau$  rest frame, and includes the effects of the momentum resolution of the detector, the beam energy spread, radiative corrections [6], and detector acceptance. The momentum resolution used in the Monte Carlo calculation is confirmed at high momenta by muon pair events [7], where  $\sigma(p_t)/p_t = 0.012p_t/(\text{GeV}/c)$ , and at low momenta by the width of the  $\Upsilon$  peak in the missing mass spectrum of  $\Upsilon' \rightarrow \pi^+\pi^- X$  events [5], where the average  $\sigma(p_t)/p_t$  is 0.009.

Fig. 3 shows that the observed  $E_{\nu}$  spectrum is in agreement with  $m(\nu_{\tau}) = 0$ . To obtain a confidence interval, we have calculated the likelihood function in the mass sensitive interval from  $E_{\nu} = -100$  MeV to  $E_{\nu} = 300$  MeV as a function of  $m(\nu_{\tau})$ . There are 102 events in the fit region. The result is  $m(\nu_{\tau}) < 56$  MeV/ $c^2$  at the 95% confidence level.

Systematic effects were studied by varying the upper and lower limits of the fit interval by  $\pm 100$  MeV, by degrading the momentum resolution by 10%, by increasing the background level in the  $E_v$  spectrum by a factor of 4, and by considering the uncertainty in the absolute momentum scale. The momentum scale is known to  $\pm 0.15\%$  from reconstructed  $K_s^0$  in various momentum and angular intervals. Taking into account these sources of systematic uncertainty, we obtain an upper limit of 70 MeV/ $c^2$  on the tau-neutrino mass at the 95% confidence level. Misidentified tau-lepton decays into KK $\pi$  and K $\pi\pi$  in the sample can only lead to an overestimation of the upper limit.

To conclude, we have measured the energy spectrum of the  $\nu_{\tau}$  in about 1500 decays  $\tau^{\pm} \rightarrow \pi^{+}\pi^{-}\pi^{\pm}\nu_{\tau}$ . The low-energy part of the spectrum depends sensitively on the mass of the  $\nu_{\tau}$ . We have used this dependence to determine an upper limit on the tau-neutrino mass of 70  $MeV/c^2$  at the 95% confidence level. This result improves the best previous limit [3] by a factor of two.

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