

# Bilarge Leptonic Mixing and Lepton Masses from Flavor Symmetries

Gerhart Seidl



In collaboration with  
T. Ohlsson

Phys. Lett. **B** 537, 95 (2002), hep-ph/0203117,  
hep-ph/0206087

# Outline

- Mixing Parameters
- Models for  $\theta_{12} \simeq \theta_{23} \simeq 1$
- Extension of the SM
- Yukawa interactions of the Leptons
- Charged Lepton and Neutrino Mass Matrices
- Bilarge Mixing Pattern
- Summary & Conclusions

# Mixing Parameters

SM: 13 of 18 parameters in the fermion sector  
→ understanding necessary

Small quark  
mixings  $\lesssim 0.2$

MSW LMA

U(1) charges,  
Irreps for 1st & 2nd  
generations

solar:  $\theta_{12} \approx \pi/6$   
atmospheric:  $\theta_{23} \approx \pi/4$ ,  
reactor:  $\theta_{13} \leq 0.2$

Nearly maximal  
 $\nu_\mu$ - $\nu_\tau$  mixing:  
 $|\theta_{23} - \pi/4| \ll 1$

Non-Abelian flavor  
symmetry acting on  
2nd & 3rd generations

↪ New symmetries and breaking patterns needed

R. N. Mohapatra and S. Nussinov, PRD **60**, 013002 (1999)

C. Wetterich, PLB **451**, 397 (1999)

## Models for $\theta_{12} \simeq \theta_{23} \simeq 1$

In the basis where  $M_\ell$  is diagonal:

$$\text{inverted hierarchy} \quad M_\nu = \begin{pmatrix} m_{11} & cM & sM \\ cM & m_{22} & m_{23} \\ sM & m_{23} & m_{33} \end{pmatrix} \longrightarrow \begin{array}{l} \text{VAC} \\ \text{or} \\ \text{LOW} \end{array}$$

where  $c = \cos(\theta_{23})$ ,  $s = \sin(\theta_{23})$ , and  $m_{ij} \ll M$ .

R.N.Mohapatra, A.Pérez-Lorenzana, and C.A. de S.Pires, PLB **474**, 355 (2000); L. Lavoura and W. Grimus, JHEP **0009**, 007 (2000); . . .

$$\text{normal hierarchy} \quad M_\nu = \begin{pmatrix} m_{11} & m_{12} & m_{13} \\ m_{12} & s^2M & scM \\ m_{13} & scM & c^2M \end{pmatrix} \longrightarrow \theta_{23} \simeq 1$$

S. Davidson and S.L. King, PLB **445**, 191 (1998);

Q. Shafi and Z. Tavartkiladze, PLB **451**, 129 (1999); . . .

Mixing entirely from  $M_\ell$ :

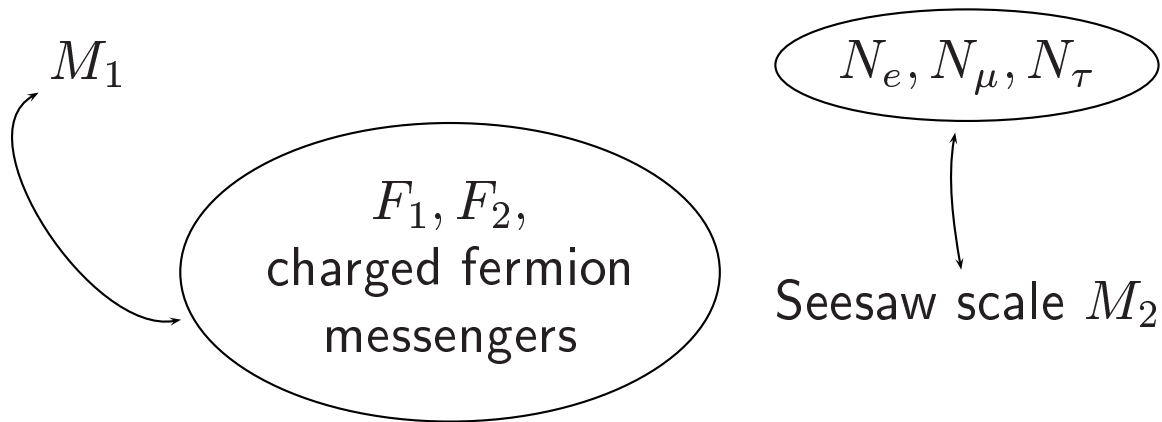
$$\text{flavor democracy} \quad M_\ell \simeq \begin{pmatrix} 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{pmatrix} \longrightarrow \begin{array}{l} \theta_{12} = \pi/4 \\ \sin^2(2\theta_{23}) = 8/9 \end{array}$$

H. Fritzsch and Z. Z. Xing, PLB **372**, 265 (1996);

M. Tanimoto, T. Watari, and T. Yanagida, PLB **461**, 345 (1999); . . .

# Extension of the SM

Fermions:



SM singlet scalars:

$$\begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix}, \begin{pmatrix} \phi_3 \\ \phi_4 \end{pmatrix}, \begin{pmatrix} \phi_5 \\ \phi_6 \end{pmatrix}, \begin{pmatrix} \phi_7 \\ \phi_8 \end{pmatrix}, \phi_9, \phi_{10}, \theta$$

$$\begin{pmatrix} \phi'_1 \\ \phi'_2 \end{pmatrix}, \begin{pmatrix} \phi'_3 \\ \phi'_4 \end{pmatrix}, \begin{pmatrix} \phi'_5 \\ \phi'_6 \end{pmatrix}$$

Leptonic mass and mixing parameters from **approximately conserved** flavor symmetries:

Symmetry breaking parameter  $\epsilon \simeq \frac{\langle \theta \rangle}{M_1} \simeq \frac{\langle \phi_i \rangle}{M_1} \simeq \frac{\langle \phi'_j \rangle}{M_1} \simeq 10^{-1}$

C. D. Froggatt and H.B. Nielsen, NPB **147**, 277 (1979);

S. Dimopoulos, PLB **129**, 417 (1983); J. Bagger, S. Dimopoulos, H.

Georgi and S. Raby, *5th workshop on grand unification*, Providence,

95 (1984); E. Witten, PLB **105**, 267 (1981)

# U(1) Charges

Assignment of **anomaly-free** U(1) charges:

| Fermions                       | $(Q_1, Q_2, Q_3)$ |
|--------------------------------|-------------------|
| $L_e, E_e$                     | $(1, 0, 0)$       |
| $L_\mu, L_\tau, E_\mu, E_\tau$ | $(0, 1, 0)$       |
| $N_e$                          | $(1, 0, 0)$       |
| $N_\mu, N_\tau$                | $(0, 1, 0)$       |
| $F_1$                          | $(1, 0, 0)$       |
| $F_2$                          | $(-1, 0, 1)$      |

| Scalars  | $(Q_1, Q_2, Q_3)$ |
|--|-------------------|
| $H_1, H_2$   | $(0, 0, 0)$       |
| $\phi_1, \phi_2$                                       | $(1, -1, 2)$      |
| $\phi_3, \phi_4$                                       | $(0, 0, 0)$       |
| $\phi_5, \phi_6$                                       | $(0, 0, 1)$       |
| $\phi'_1, \phi'_2, \phi'_3, \phi'_4, \phi'_5, \phi'_6$ | $(0, 0, 0)$       |
| $\phi_7, \phi_8$                                       | $(-1, -1, 0)$     |
| $\phi_9$   | $(-2, 0, 1)$      |
| $\phi_{10}$  | $(0, 0, 0)$       |
| $\theta$   | $(0, 0, -1)$      |

Note: SSB close to TeV-scale possible

M. Leurer, Y. Nir and N. Seiberg, NPB **398**, 319 (1993); NPB **420**, 468 (1994)

# Discrete Symmetries

Distinguishing the **first** generation with  $\mathbb{Z}_n$  symmetry:

$$\mathcal{D}_1 : \begin{cases} E_e \rightarrow P^{-4}E_e, & E_\mu \rightarrow P^{-1}E_\mu, & E_\tau \rightarrow P^{-1}E_\tau, \\ \phi_1 \rightarrow P\phi_1, & \phi_2 \rightarrow P\phi_2, & \phi_3 \rightarrow P\phi_3, \\ \phi_4 \rightarrow P\phi_4, & \phi_5 \rightarrow P\phi_5, & \phi_6 \rightarrow P\phi_6, \end{cases}$$

where  $P \equiv \exp(2\pi i/n)$ ,  $n \geq 5$  & RH  $\nu$ s  $\mathcal{D}_1$ -**singlets**.

$|\theta_{23} - \pi/4| \ll 1$  from  $\mathbb{Z}_2$  permutation symmetries:

$$\mathcal{D}_2 : \begin{cases} L_\mu \rightarrow -L_\mu, & E_\mu \rightarrow -E_\mu, & N_\mu \rightarrow -N_\mu, \\ \phi'_1 \leftrightarrow \phi'_2, & \phi_1 \leftrightarrow \phi_2, & \phi_7 \rightarrow -\phi_7, \end{cases}$$

$$\mathcal{D}_3 : \begin{cases} L_\mu \rightarrow -L_\mu, & N_\mu \rightarrow -N_\mu, \\ \phi'_3 \leftrightarrow \phi'_4, & \phi_3 \leftrightarrow \phi_4, \\ \phi'_5 \leftrightarrow \phi'_6, & \phi_5 \leftrightarrow \phi_6, \\ \phi_7 \rightarrow -\phi_7, \end{cases}$$

$$\mathcal{D}_4 : \begin{cases} L_\mu \leftrightarrow L_\tau, & E_\mu \leftrightarrow E_\tau, & N_\mu \leftrightarrow N_\tau, \\ \phi_2 \rightarrow -\phi_2, & \phi_4 \rightarrow -\phi_4, & \phi_6 \rightarrow -\phi_6, \\ \phi_7 \leftrightarrow \phi_8. \end{cases}$$

- Degenerate Yukawa couplings in  $E_\alpha$  and  $\nu_\alpha$  sector
- Degenerate scalar-scalar couplings

## Forbidding Terms

Terms **spoiling** the vacuum alignment are forbidden by

$$\mathcal{D}_5 : \begin{cases} \phi'_1 \rightarrow -\phi'_1, & \phi'_3 \rightarrow -\phi'_3, & \phi'_5 \rightarrow -\phi'_5 \\ \phi_1 \rightarrow -\phi_1, & \phi_3 \rightarrow -\phi_3, & \phi_5 \rightarrow -\phi_5, \end{cases}$$

$$\mathcal{D}_6 : \begin{cases} E_e \rightarrow P^{-(4l+1)} E_e, & N_e \rightarrow P N_e, \\ \phi'_1 \rightarrow P^{-k} \phi'_1, & \phi'_2 \rightarrow P^{-k} \phi'_2, \\ \phi'_3 \rightarrow P^{-l} \phi'_3, & \phi'_4 \rightarrow P^{-l} \phi'_4, \\ \phi'_5 \rightarrow P^{-m} \phi'_5, & \phi'_6 \rightarrow P^{-m} \phi'_6 \\ \phi_1 \rightarrow P^k \phi_1, & \phi_2 \rightarrow P^k \phi_2, \\ \phi_3 \rightarrow P^l \phi_3, & \phi_4 \rightarrow P^l \phi_4, \\ \phi_5 \rightarrow P^m \phi_5, & \phi_6 \rightarrow P^m \phi_6, \\ \phi_9 \rightarrow P^{-1} \phi_9, & \phi_{10} \rightarrow P \phi_{10}. \end{cases}$$

- $k, l, m$ : integers with  $|k - l|, |k - m|, |l - m| \gg 1$
- $\mathcal{D}_6$ -charges of charged fermion messengers:  
 $\mathbb{N} \times \{2\pi k/n, 2\pi l/n, 2\pi m/n\}$



# Charged Lepton Yukawa Interactions

Charged lepton mass terms:

$$\mathcal{L}_Y^\ell = \overline{L}_\alpha H_2 [(Y_{\text{eff}}^1)_{\alpha\beta} + (Y_{\text{eff}}^2)_{\alpha\beta}] E_\beta + \text{h.c.}$$

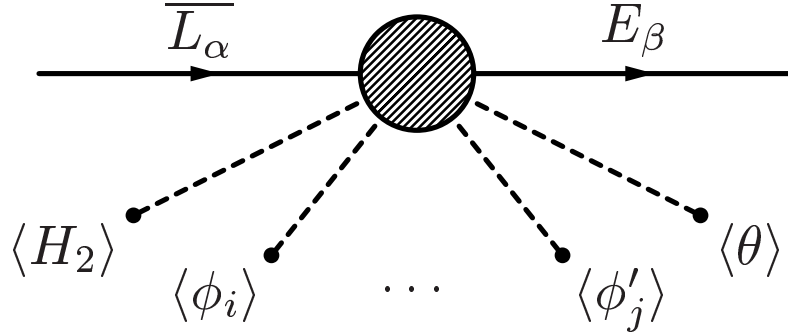


Fig. 1: Non-renormalizable terms generating  $Y_{\text{eff}}^1$  and  $Y_{\text{eff}}^2$ .

$$Y_{\text{eff}}^1 = \begin{pmatrix} A_1 & B_1 - B_2 & B_1 + B_2 \\ 0 & C_1 - C_2 & 0 \\ 0 & 0 & C_1 + C_2 \end{pmatrix},$$

$$Y_{\text{eff}}^2 = \text{diag}(0, D_1 - D_2, D_2 + D_2)$$

$$A_1 = Y_a^\ell \frac{\phi_{10}}{(M_1)^5} [(\phi_3)^4 + (\phi_4)^4],$$

$$B_1 = Y_b^\ell \frac{\theta^2}{(M_1)^4} \phi_1 \phi'_1, \quad C_1 = Y_c^\ell \frac{\phi'_3 \phi_3}{(M_1)^2}, \quad D_1 = Y_d^\ell \frac{\phi'_5 \phi_5}{(M_1)^3},$$

$$B_2 = Y_b^\ell \frac{\theta^2}{(M_1)^4} \phi_2 \phi'_2, \quad C_2 = Y_c^\ell \frac{\phi'_4 \phi_4}{(M_1)^2}, \quad D_2 = Y_d^\ell \frac{\phi'_6 \phi_6}{(M_1)^3}$$

# Neutrino Yukawa Interactions

$$\mathcal{L}_Y^\nu = \overline{L}_\alpha^c \frac{(H_1)^2}{M_2} (Y_{\text{eff}}^3)_{\alpha\beta} L_\beta + \text{h.c.}, \quad Y_{\text{eff}}^3 = \begin{pmatrix} A_2 & B_3 & B_4 \\ B_3 & 0 & 0 \\ B_4 & 0 & 0 \end{pmatrix},$$

Dimensionful Yukawa couplings:

$$A_2 = Y_a^\nu \frac{\phi_9 \phi_{10} \theta}{(M_1)^3}, \quad B_3 = Y_b^\nu \frac{\phi_7}{M_1}, \quad B_4 = Y_b^\nu \frac{\phi_8}{M_1},$$

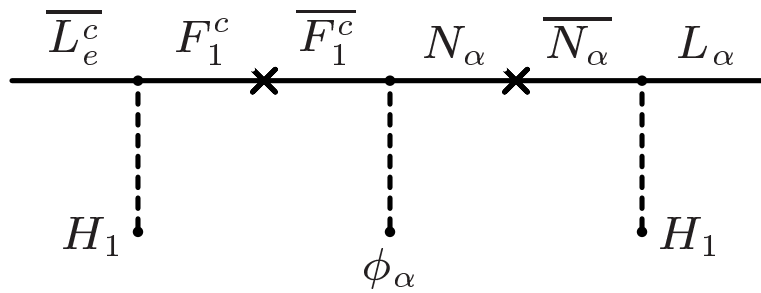


Fig. 2: The dimension six operators for  $\alpha = e, \mu$  and  $\phi_\alpha = \phi_7, \phi_8$  generating  $M_{e\mu}^\nu$  and  $M_{e\tau}^\nu$ .

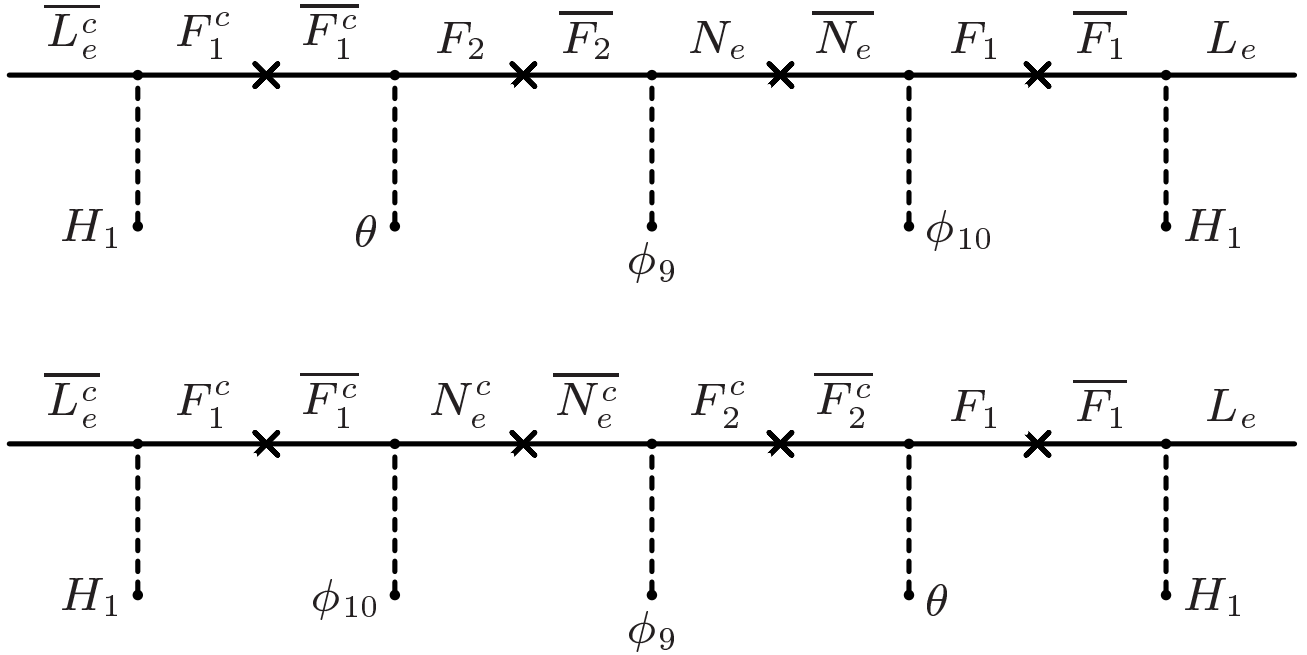


Fig. 3: The dimension eight operators generating  $M_{ee}^\nu$ .

## Charged Lepton Masses

For a range of parameters **vacuum alignment mechanism** :

$$\begin{pmatrix} \langle \phi_1 \rangle & \langle \phi_3 \rangle & \langle \phi_5 \rangle \\ \langle \phi_2 \rangle & \langle \phi_4 \rangle & \langle \phi_6 \rangle \end{pmatrix} = \pm \begin{pmatrix} +1 & -1 & +1 \end{pmatrix}$$

$$\begin{pmatrix} \langle \phi'_1 \rangle & \langle \phi'_3 \rangle & \langle \phi'_5 \rangle \\ \langle \phi'_2 \rangle & \langle \phi'_4 \rangle & \langle \phi'_6 \rangle \end{pmatrix} = \pm \begin{pmatrix} +1 & +1 & +1 \end{pmatrix}$$

↓

$$\begin{pmatrix} B_1 & C_1 & D_1 \\ B_2 & C_2 & D_2 \end{pmatrix} = \pm \begin{pmatrix} +1 & -1 & +1 \end{pmatrix}$$

↪ Leading order effective Yukawa couplings:  
**Exactly degenerate** with **relative sign**

Strictly **hierarchical** charged lepton mass matrix patterns:

$$M_\ell \simeq m_\tau \begin{pmatrix} \epsilon^3 & \epsilon^2 & \epsilon^4 \\ \epsilon^3 & \epsilon & \epsilon^2 \\ \epsilon^3 & \epsilon^2 & 1 \end{pmatrix} \leftrightarrow m_\tau \begin{pmatrix} \epsilon^3 & \epsilon^4 & \epsilon^2 \\ \epsilon^3 & 1 & \epsilon^2 \\ \epsilon^3 & \epsilon^2 & \epsilon \end{pmatrix}$$

Mass ratios:

$$m_e/m_\mu \simeq \epsilon^2 \simeq 10^{-2} \quad \text{and} \quad m_\mu/m_\tau \simeq \epsilon \simeq 10^{-1}$$

Mixing parameters:

$$\theta_{12}^\ell \simeq 6^\circ \quad \text{and} \quad \theta_{13}^\ell \approx \theta_{23}^\ell \approx 0^\circ$$

# Neutrino Masses

For a range of parameters: vacuum alignment  $\curvearrowright$

$$\frac{\langle \phi_7 \rangle}{\langle \phi_8 \rangle} = \frac{B_3}{B_4} = \pm 1$$

$\curvearrowright$  Effective neutrino mass matrix:

$$M_\nu = \frac{\langle H_1 \rangle^2}{M_2} \begin{pmatrix} Y_a^\nu \epsilon^3 & Y_b^\nu \epsilon & -Y_b^\nu \epsilon \\ Y_b^\nu \epsilon & \epsilon^5 & \epsilon^5 \\ -Y_b^\nu \epsilon & \epsilon^5 & \epsilon^5 \end{pmatrix} = m_\nu \begin{pmatrix} \epsilon^2 & \mathbf{1} & \mathbf{-1} \\ \mathbf{1} & \epsilon^4 & \epsilon^4 \\ \mathbf{-1} & \epsilon^4 & \epsilon^4 \end{pmatrix}$$

$\curvearrowright$   $M_\nu$  on **bimaximal mixing** form

V. Barger *et al.*, PLB **437**, 107 (1998);

R. N. Mohapatra and S. Nussinov, PRD **60**, 013002 (1999); ...

$$\theta_{12}^\nu = \pi/4 + \mathcal{O}(\epsilon^2), \quad \theta_{13}^\nu \approx 0, \quad \text{and} \quad \theta_{23}^\nu = \pi/4,$$

$$m_1 \simeq \sqrt{2}m_\nu, \quad m_2 \simeq \sqrt{2}m_\nu, \quad \text{and} \quad m_3 \simeq 2\epsilon^4 m_\nu \approx 0$$

**Inverse** hierarchical spectrum

Mass squared differences:

$$\begin{aligned} \Delta m_{\odot}^2 &\equiv |\Delta m_{21}^2| \simeq 2\sqrt{2}\epsilon^2 m_\nu^2 \\ \Delta m_{\text{atm}}^2 &\equiv |\Delta m_{32}^2| \simeq 2m_\nu^2 \end{aligned}$$

MSW LMA values  $\curvearrowright$   $m_\nu \simeq 0.04$  eV and  $M_2 \simeq 10^{14}$  GeV

## Bilarge Mixing

Charged lepton & neutrino mixing  $\curvearrowright$

$$\begin{aligned}\theta_{12} &= \pi/4 - \epsilon/\sqrt{2} + \mathcal{O}(\epsilon^2), \\ \theta_{13} &= \epsilon/\sqrt{2} + \mathcal{O}(\epsilon^3), \\ \theta_{23} &= \pi/4 + \mathcal{O}(\epsilon^2)\end{aligned}$$

For  $\epsilon \simeq 10^{-1}$ :

$$\theta_{12} \simeq 41^\circ, \quad \theta_{13} \simeq 4^\circ, \quad \text{and} \quad \theta_{23} \simeq 44^\circ$$

Mild fine-tuning  $Y_b^\ell/Y_c^\ell \simeq 2$ :

$$\theta_{12} \simeq 37^\circ, \quad \theta_{13} \simeq 8^\circ, \quad \text{and} \quad \theta_{23} \simeq 44^\circ$$

$\curvearrowright$  Compatible with **MSW LMA** and CHOOZ upper bound

# Summary and Conclusions

- Strict **hierarchical** charged lepton mass pattern
- **Large**, but not necessarily close to maximal  $\theta_{12}$
- **Relation**  $\theta_{12} \simeq \pi/4 - \theta_{13}$
- $|\theta_{23} - \pi/4| \ll 1$  due to symmetries
- Inverted neutrino mass hierarchy
- Can give the **MSW LMA** solution

## References:

T. Ohlsson and G. Seidl, Phys. Lett. **B** 537, 95 (2002), hep-ph/0203117;  
T. Ohlsson and G. Seidl, *A Flavor Symmetry Model for Bilarge Leptonic Mixing and the Lepton Masses*, hep-ph/0206087