

GRAND UNIFICATION WITH ANOMALOUS U(1) SYMMETRY.

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<u>SO(10)</u>	<u>hep-ph/0104200 (PTP)</u>	N.M.
μ	hep-ph/0107313 (PLB)	N.M.
E_6 (matter)	hep-ph/0109018 (PTP)	M.Bando & N.M.
<u>Coupling Unification</u>	<u>hep-ph/0111205 (PTP)</u>	N.M.
E_6 (Higgs)	hep-ph/0202050 (PTP)	N.M. & T.Yamashita
$SU(3)^3$	hep-ph/0204030	N.M. & Q.Shafi
2loop RGE	hep-ph/0205185	N.M. & T.Yamashita

I OVERVIEW OF OUR SCENARIO

II. ANOMALOUS U(1) GAUGE SYMMETRY

III. DOUBLET-TRIPLET SPLITTING

IV. GAUGE COUPLING UNIFICATION

V. QUARK & LEPTON MASS MATRICES

VI. PREDICTION & SUMMARY

I Overview of our scenario

* "Generic" interactions.

Symmetry $\downarrow \begin{pmatrix} SO(10) \\ E_6 \end{pmatrix} \times \underline{U(1)_A}$ Anomalous U(1) Gauge Symmetry

Definition of the theory (except O(1) coefficients).

Input (# Anomalous U(1) charges (integer)).

	matter	Higgs
SO(10)	$4[3 \times 16 + 1 \times 10]$	$8[2 \times 45 + 2 \times (16 + \bar{16}) + 2 \times 1]$
E_6	$3[3 \times 27]$	$8[2 \times 78 + 3 \times (27 + \bar{27})]$

Output

- GUT scales, other symmetry breaking scales
- Doublet-triplet splitting.
- Proton decay via dim. 5 operators is suppressed.
- MSSM is obtained at the low energy scale.
- Mass spectrum of superheavy fields. $m_{\text{SH}} < \Delta \epsilon$
- Gauge coupling unification in MSSM is naturally explained.
- Realistic quark & lepton masses & mixings (no GUT relation)

Neutrino Bi-large mixing (LMA)

- μ problem
- Suppression of FCNC (in E_6)

$\downarrow + \text{SUSY}$

Predictions.

- * Neutrino \rightarrow bi-large neutrino mixing
 - $\lambda m_{\nu_e} \sim m_{\nu_\mu}$ (LMA)
 - $\lambda m_{\nu_\mu} \sim m_{\nu_\tau}$ $\lambda \sim \sin \theta_c \sim 0.22$
 - $U_{e3} \sim \lambda \sim 0(0.1)$, $\delta_{\text{CP}} \sim 0(1)$

* Small $\tan \beta \equiv \frac{v_u}{v_d}$

- ④ Unification scale $\Delta_u < \Delta_G = 2 \times 10^{16} \text{ GeV}$
 \Rightarrow Proton decay via dimension 6 operators

$$\tau(P \rightarrow e\pi) \sim \begin{cases} 5 \times 10^{33} \text{ years } (a=-1) \\ 8 \times 10^{34} \text{ years } (a=-\frac{1}{2}) \end{cases}$$

$$\tau_{\text{exp}}(P \rightarrow e\pi) \gtrsim 5 \times 10^{33} \text{ years.}$$

- ④ Cutoff scale $\Lambda \sim \Delta_G = 2 \times 10^{16} \text{ GeV} < \Delta_P$

\Rightarrow Extra-dimensions?

Horava - Witten



- * SUSY

II ANOMALOUS U(1) GAUGE SYMMETRY

- The gauge anomaly is cancelled by Green-Schwarz mechanism.

- D flatness.

charge $\theta = -1$

$$D_A = \tilde{\Xi}_{\text{FI}}^2 - |\Theta|^2 = 0 \rightarrow \langle \Theta \rangle = \tilde{\Xi}_{\text{FI}} \equiv \lambda \Lambda \quad \lambda < 1.$$

Anomalous U(1) is broken at just below the cutoff.

- Yukawa hierarchy in SSM. Ibanez - Ross.

$\lambda \sim 0.2$

$$W = \left(\frac{\Theta}{\Lambda}\right)^{x+y+z} XYZ \rightarrow \lambda^{x+y+z} XYZ$$

(Small letter charges.)

- SUSY zero

$x+y+z < 0 \Rightarrow$ No gauge invariant interaction with positive powered Θ .

- Vacuum expectation value (VEV)

Z : GUT gauge singlet operator

$$\langle Z \rangle = \begin{cases} 0 & z > 0 \\ \lambda^z & z < 0 \end{cases}$$

negative ↑

"The GUT scale is determined by the Higgs charge"

III DOUBLET-TRIPLET SPLITTING.

One interesting point

- 2 adjoint fields. 45

$$A'(\alpha'=3) \quad A(\alpha=-1)$$

$$W_{A'} = \lambda^{0+\alpha} A'A + \lambda^{3\alpha+\alpha'} A'A^3$$

$$\Rightarrow \langle A' \rangle = 0$$

$$\langle A \rangle = \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix} \otimes \begin{pmatrix} x_1 & & & & \\ & x_2 & & & \\ & & x_3 & & \\ & & & x_4 & \\ & & & & x_5 \end{pmatrix}$$

$$\Delta = 1 \text{ unit}$$

$$\frac{\partial W_{A'}}{\partial A'} = 0 \Rightarrow \lambda^{0+\alpha} x_i (1 + x_i^2 \lambda^{2\alpha}) = 0 \text{ for all } i$$

$$\Rightarrow x_i = 0 \text{ or } \lambda^{-\alpha} \quad \text{only two solutions.}$$

The vacuum is classified by # of zero solution.

$$N_0 = 2$$

$$\langle A \rangle = \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix} \otimes \begin{pmatrix} v & & & & \\ & v & & & \\ & & v & & \\ & & & 0 & \\ & & & & 0 \end{pmatrix}, \quad v \sim \lambda^{-\alpha}$$

Dimopoulos-Wilczek type VEV

- $A'A^{2L+1}$ $(L \geq 2)$ $\rightarrow \frac{\partial W}{\partial A'} = 0$ have more solutions
 \Rightarrow less natural to obtain DW VEV

SUSY zero forbids these terms.

$$v \sim \lambda^{-\alpha}$$

GUT scale is determined by the charge of A .

- Doublet Higgs is included in $\textbf{10}$ $H(h=3)$ +
 $\langle H' \rangle = 0$ $\textbf{10}$ $H'(h=4)$ -

$$W = \lambda^{h+h'+a} H' A H + \lambda^{2h'} H'^2 \quad A(a=-1) -$$

$$\Rightarrow M_H = \begin{pmatrix} H & H' \\ H' & \lambda^{h+h'+a} \langle A \rangle \end{pmatrix} \quad \text{SUSY zero.}$$

$$\langle A \rangle = \begin{cases} 0 & \text{for doublet} \\ \lambda^a & \text{for triplet} \end{cases}$$

- Only 1 pair of doublet Higgs is massless.

- Proton decay (dim. 5) is naturally suppressed.

$$m_{H^0}^{\text{eff}} \sim \lambda^{2h} > \Delta \quad (\oplus h < 0)$$

$$(2h' > h+h' \rightarrow \lambda^{2h'} < \lambda^{h+h'} \sim \lambda^{h+h'+a} \langle A \rangle)$$

IV GAUGE COUPLING UNIFICATION

Anomalous U(1) charges \Rightarrow mass spectrum (superheavy)

$$SO(10) \supset SU(5)$$

$$\mathbf{45} := \mathbf{24} + \mathbf{10} + \overline{\mathbf{10}} + \mathbf{1}$$

$$\mathbf{24}, \quad I_A \quad I_{A'}$$

$$M_I = \frac{I_A}{I_{A'}} \begin{pmatrix} 0 & \lambda^{a+a'} \\ \lambda^{a+a'} & \lambda^{2a'} \end{pmatrix} \quad I = G(8,1)_0, W(1,3)_0, X(3,2)_E$$

$\alpha_X = 0$

c.f. $\langle A \rangle \sim \lambda^a \sim \lambda$

$$I = G, W \quad (\lambda^{a+a'}, \lambda^{a+a'}) = (\lambda^2, \lambda^2)$$

$$I = X \quad (\underbrace{0,}_{\text{L}} \lambda^{2a'}) = (0, \lambda^6)$$

absorbed by Higgs mechanism.

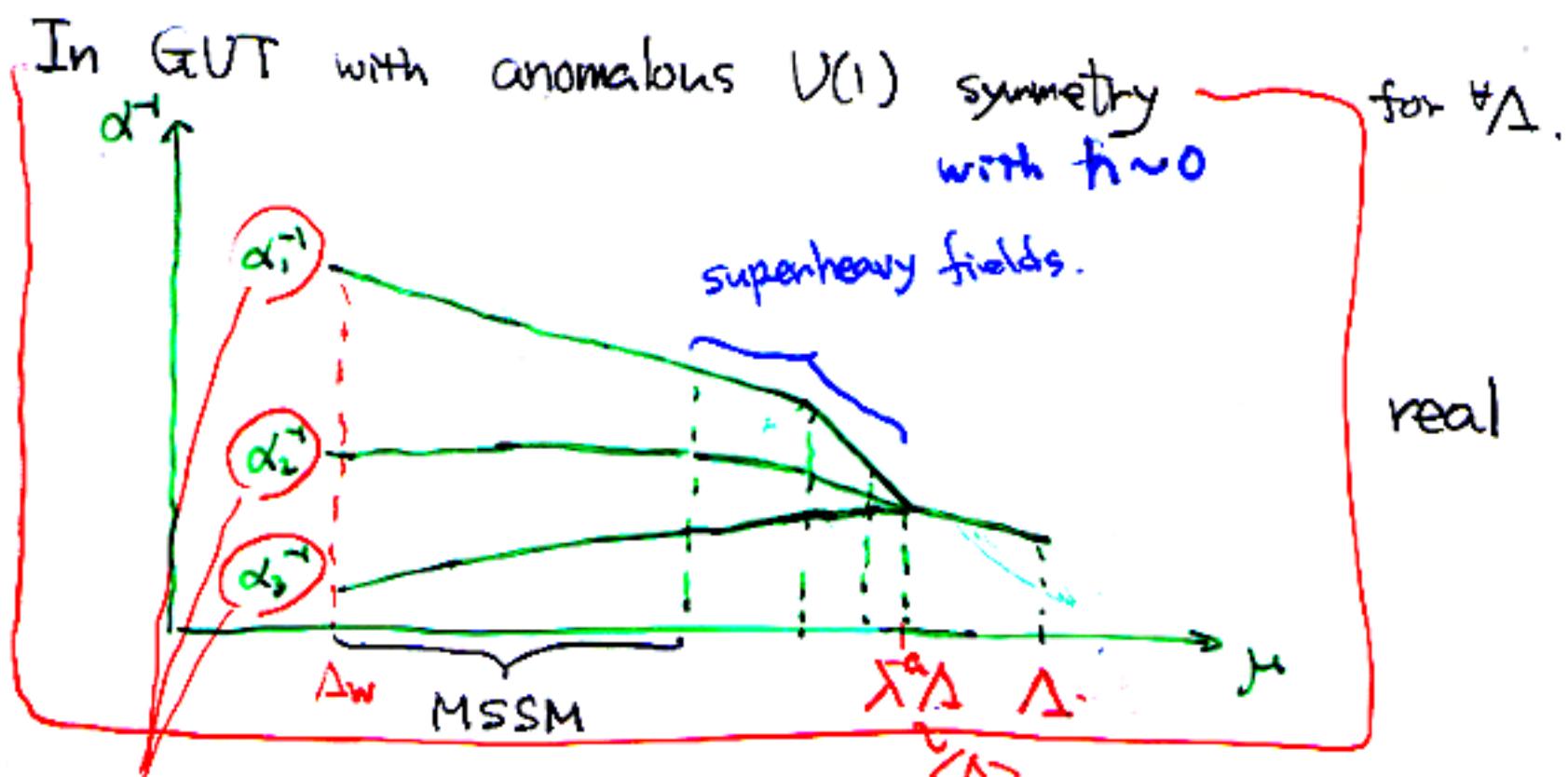
\Rightarrow The spectrum does not respect $SU(5)$.

\Rightarrow Is the success of gauge coupling unification in MSSM spoiled in our scenario?

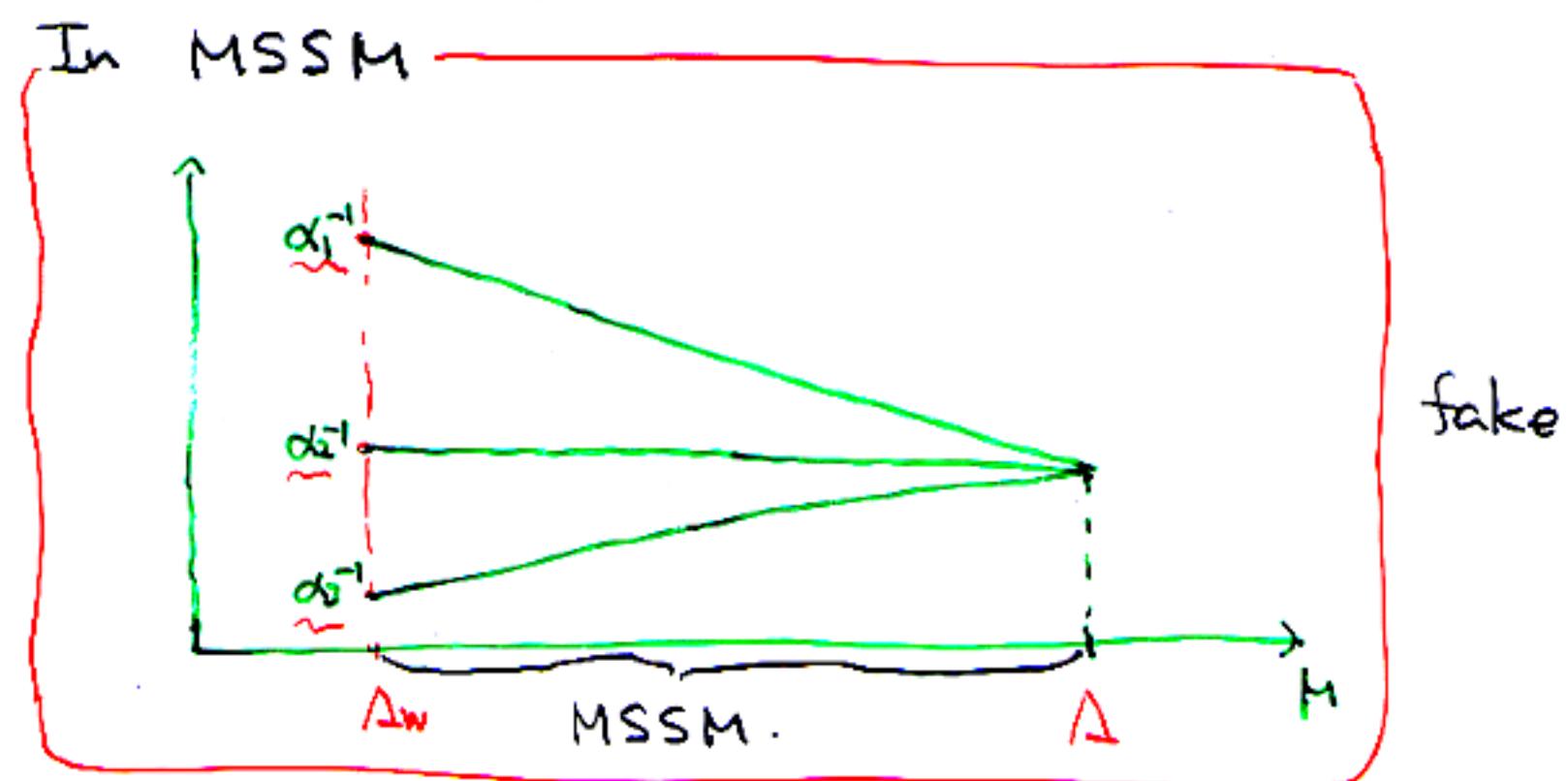
No!!

mass spectrum } are determined by anomalous U(1)
 GUT breaking scales } charges.

Straightforward calculation of running gauge couplings.



By using these values



The real unification scale $\Delta_W \sim \lambda^a \Delta_G < \Delta_G \sim 2 \times 10^{16} \text{ GeV}$
 \Rightarrow Proton decay $\tau(P \rightarrow e\pi) \sim \begin{cases} 3 \times 10^{33} \text{ yrs} & a=1 \\ 5 \times 10^{34} \text{ yrs} & a=-\frac{1}{2} \end{cases}$

V QUARK & LEPTON MASS MATRICES

$$M_u = \begin{pmatrix} \lambda^6 & \lambda^5 & \lambda^3 \\ \lambda^5 & \lambda^4 & \lambda^2 \\ \lambda^3 & \lambda^2 & 1 \end{pmatrix} \langle H_u \rangle, \quad M_d = \begin{pmatrix} \lambda^6 & \lambda^{6.5} & \lambda^5 \\ \lambda^5 & \lambda^{4.5} & \lambda^4 \\ \lambda^3 & \lambda^{2.5} & \lambda^2 \end{pmatrix} \langle H_d \rangle$$

$$M_\nu = \lambda^{-3/4} \begin{pmatrix} \lambda^2 & \lambda^{1.5} & \lambda \\ \lambda^{1.5} & \lambda & \lambda^{0.5} \\ \lambda & \lambda^{0.5} & 1 \end{pmatrix} \frac{\eta \langle H_b \rangle}{\Delta}, \quad M_e^\top = \begin{pmatrix} \lambda^6 & \lambda^{6.5} & \lambda^5 \\ \lambda^5 & \lambda^{4.5} & \lambda^4 \\ \lambda^3 & \lambda^{2.5} & \lambda^2 \end{pmatrix} \langle H_d \rangle$$

$$U_{CKM} = \begin{pmatrix} 1 & \lambda & \lambda^3 \\ \lambda & 1 & \lambda^2 \\ \lambda^3 & \lambda^2 & 1 \end{pmatrix} \quad \lambda \sim \sin \theta_C \sim 0.22.$$

$$U_{MNS} = \begin{pmatrix} 1 & \bar{\lambda} & \lambda \\ \bar{\lambda} & 1 & \bar{\lambda} \\ \lambda & \bar{\lambda} & 1 \end{pmatrix} \begin{pmatrix} 1 & \bar{\lambda} & \lambda \\ \bar{\lambda} & 1 & \bar{\lambda} \\ \lambda & \bar{\lambda} & 1 \end{pmatrix} \sim \begin{pmatrix} 1 & \bar{\lambda} & \lambda \\ \bar{\lambda} & 1 & \bar{\lambda} \\ \lambda & \bar{\lambda} & 1 \end{pmatrix} \quad \bar{\lambda} + \bar{\lambda} \sim \bar{\lambda}$$

* bi-large neutrino mixing

$$\bar{\lambda} \sim 0.5, \quad \bar{\lambda} + \bar{\lambda} \sim \bar{\lambda}$$

* LMA solution for solar neutrino problem.

$$\lambda m_{\nu_e} \sim m_{\nu_\mu}$$

* leptogenesis? $\lambda m_{\nu_\mu} \sim m_{\nu_e}$ $[(M_{\nu_1}, M_{\nu_2}, M_{\nu_3}) = (\lambda^2, \lambda^0, \lambda^6)]$

* $U_{e3} \sim \lambda \sim O(0.1)$

* $S_{\text{CP}} \sim 0.0$

* small $\tan \beta$

* unrealistic GUT relation $Y_u = Y_d = Y_e$ is avoided.

How to avoid the unrealistic GUT relation?

$$\bar{\Psi}_i : \mathbf{16} = \mathbf{10} + \bar{\mathbf{5}} + \mathbf{1} \quad H : \mathbf{10} = \mathbf{5} + \bar{\mathbf{5}}$$

$i=1,2,3$

$$W_Y = \lambda^{\Phi_i + \Phi_j + h} \bar{\Psi}_i \bar{\Psi}_j H.$$

$$\rightarrow \underbrace{Y_u}_{SO(10)} = \underbrace{Y_d}_{SU(5)} = Y_e$$

1. An additional $\mathbf{10}$ matter T

$$\mathbf{16} = \mathbf{10} + \bar{\mathbf{5}} + \mathbf{1} \qquad \text{Nomura-Yanagida.}$$

$$\mathbf{10} = \underline{\mathbf{5}} + \bar{\underline{\mathbf{5}}}$$

We can pick up $SO(10)$ breaking VEV in
the mass matrix of $(\bar{5}_T, (\bar{5}_T, \bar{5}_1, \bar{5}_2, \bar{5}_3))$.

2. Non-renormalizable operations.

$$\Delta W_Y = \lambda^{\Phi_i + \Phi_j + h + 2a} \bar{\Psi}_i \langle A^2 \rangle^{\overbrace{T}} \bar{\Psi}_j H. \quad \frac{\langle A \rangle}{\Lambda} \sim \tilde{\lambda}^a$$

Usually the effect from non-renormalizable terms is suppressed. But in our scenario, this suppression is cancelled by enhancement factor $\tilde{\lambda}^{2a}$

VI PREDICTION & SUMMARY

* Completely consistent GUT with anomalous U(1).

"Generic" \rightarrow input: 3~4 (matter) + 8 (Higgs)

It is surprising that such a simple set up can solve most of the problems of GUT. (DT splitting, unrealistic GUT relation...)

* Predictions.

→ Neutrino. \rightarrow bi-large neutrino mixing.

$$\cdot \lambda m_{\nu_e} \sim m_{\nu_\mu} \text{ (LMA)} \quad \cdot \lambda m_{\nu_\mu} \sim m_{\nu_\tau}$$

$$\cdot U_{e3} \sim \lambda$$

→ Small $\tan\beta$

→ Proton decay (dim 6 though SUSY).

rough estimation

$$\tau(P \rightarrow e\bar{\nu}) \sim \begin{cases} 3 \times 10^{33} \text{ years} & (\alpha = -1) \\ 5 \times 10^{34} \text{ years} & (\alpha = -\frac{1}{2}) \end{cases}$$

* How is the GUT induced from superstring theory?

{ • 2 adjoint Higgs.

{ • "Generic" superpotential.

{ x negative powered interactions. H^{-n}

{ x non-perturbative interaction of dilaton. e^{-D}

\Rightarrow These terms spoil SUSY zero mechanism.