

The MSSM without μ term

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Contents

- The μ SSM and its low energy spectrum
- $U(1)_R$ Symmetric Gauge Mediation
- Contribution to precision electroweak parameters
 - T parameter
 - Muon anomalous magnetic moment
- Unification of couplings

μ SSM

The μ SSM

- No supersymmetric μ term in the superpotential
- Approximate $U(1)_R$ symmetry forbids Majorana masses
- Additional chiral superfields: **S**, **T** (adjoint of $SU(2)$) and **O** (adjoint of $SU(3)$)
- Superpotential couplings

$$\int d^2\theta \quad h_S S H_1 H_2 + h_T H_1 T H_2$$

- Usual **MSSM** superpotential for quarks and leptons allowed

μ SSM

$U(1)_R$ charges

Ψ_{H_1}	Ψ_{H_2}	$\Psi_T^\pm$	H_1	H_2	$\lambda^\pm$
1	-1	-1	2	0	1

Trilinear terms prohibited by $U(1)_R$ symmetry $\implies$ suppresses tadpole for the T scalar

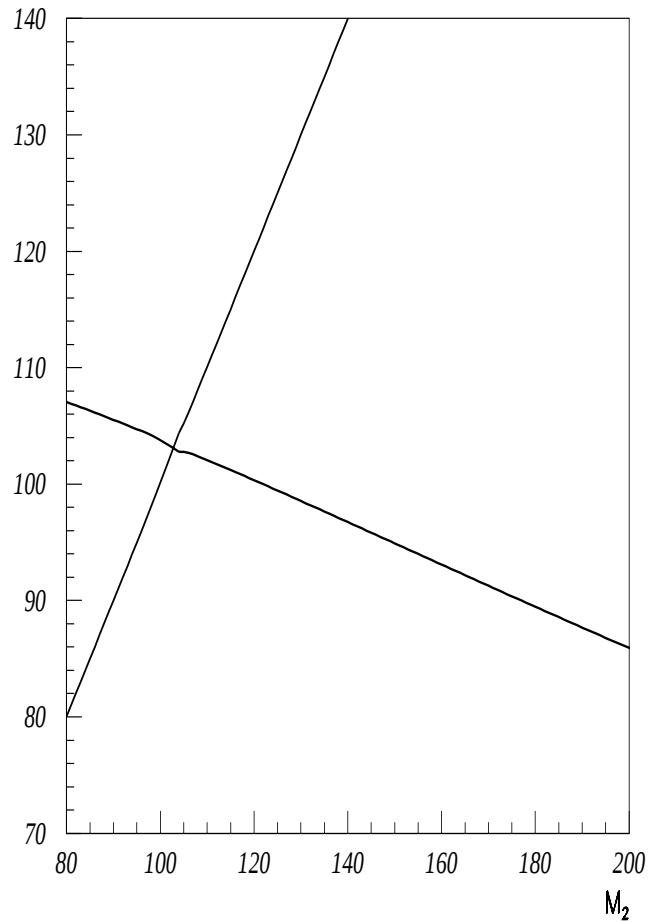
Small $U(1)_R$ breaking effects $\mu B H_1 H_2$
 set $\tan \beta \sim 60 \implies$ explains top/bottom mass hierachy

Chargino mass matrix

	Ψ_T^+	$-i\lambda^+$	$\Psi_{H_2}^+$
Ψ_T^-	0	$\tilde{M}_2$	$-h_T v_1$
$-i\lambda^-$	$\tilde{M}_2$	$\tilde{m}_2$	$\sqrt{2} m_W s_\beta$
$\Psi_{H_1}^-$	$h_T v_2$	$\sqrt{2} m_W c_\beta$	0

μ SSM

Lighter chargino masses ($\tilde{m}_2 = 5 \text{ GeV}$)



LEP II limit $m_{\chi^\pm} > 104 \text{ GeV} \implies$

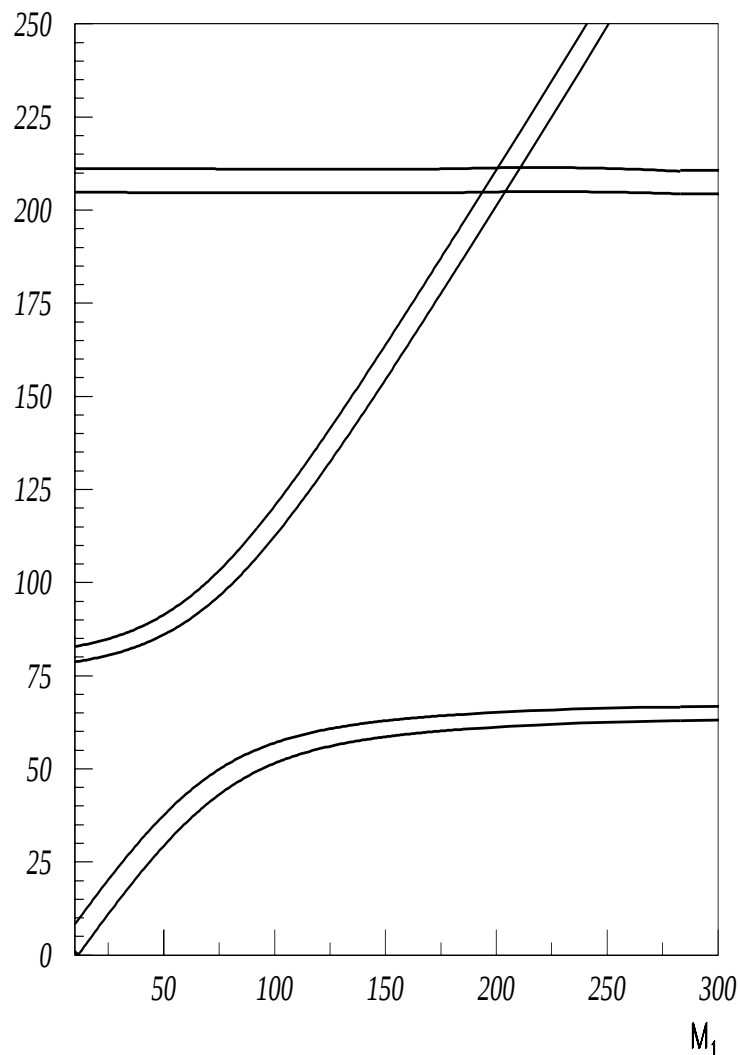
$104 \text{ GeV} < \tilde{M}_2 < 120 \text{ GeV}$ and $h_T \sim 1$

μ SSM

μ SSM

	Ψ_T^3	Ψ_S	$-i\lambda'$	$-i\lambda^3$	$\Psi_{H_1}^1$	$\Psi_{H_2}^2$
Ψ_T^3	0	0	0	$\tilde{M}_2$	$h_T \frac{v_2}{\sqrt{2}}$	$h_T \frac{v_1}{\sqrt{2}}$
Ψ_S	0	0	$\tilde{M}_1$	0	$h_S \frac{v_2}{\sqrt{2}}$	$h_S \frac{v_1}{\sqrt{2}}$
$-i\lambda'$	0	$\tilde{M}_1$	$\tilde{m}_1$	0	$-m_Z s_W c_\beta$	$m_Z s_W s_\beta$
$-i\lambda^3$	$\tilde{M}_2$	0	0	$\tilde{m}_2$	$m_Z c_W c_\beta$	$-m_Z c_W s_\beta$
$\Psi_{H_1}^1$	$h_T \frac{v_2}{\sqrt{2}}$	$h_S \frac{v_2}{\sqrt{2}}$	$-m_Z s_W c_\beta$	$m_Z c_W c_\beta$	0	0
$\Psi_{H_2}^2$	$h_T \frac{v_1}{\sqrt{2}}$	$h_S \frac{v_1}{\sqrt{2}}$	$m_Z s_W s_\beta$	$-m_Z c_W s_\beta$	0	0

Neutralino masses



$$h_T = 1, h_S = 0.1, \tilde{M}_2 = 104 \text{ GeV}$$

$$\tilde{m}_1 = \tilde{m}_2 = 5 \text{ GeV}, \tan \beta = 60$$

μ SSM

$U(1)_R$ Gauge Mediation

- No singlet messengers
- SB mediated by some gauge $U(1)$ group also carried by the messengers $\implies$ nonholomorphic scalar SUSY breaking masses for the messengers
- $\text{Tr } T_Y T_{new} = 0 \implies$ no D term for hypercharge

Messengers: $L, \bar{L}$ ($D, \bar{D}$) transform as the fundamental of $SU(2)$ ($SU(3)$)

$$\lambda_S S \bar{L} L + \lambda'_S S \bar{D} D + \lambda_T T \bar{L} L + \lambda_O O \bar{D} D + M_L \bar{L} L + M_D \bar{D} D$$

$\not\mu$ SSM

$L, \bar{L}$ mass matrix

$$\begin{pmatrix} M_L^2 + \tilde{m}_L^2 & 0 \\ 0 & M_{\bar{L}}^2 + \tilde{m}_{\bar{L}}^2 \end{pmatrix}$$

soft SUSY breaking masses $\tilde{m}_L^2 + \tilde{m}_{\bar{L}}^2 < 0$

$\Rightarrow$ scalar superpartner masses

Large soft masses for the scalar S, T, O

$$\tilde{m}_a \sim \frac{g_a^2}{16\pi^2} M_L$$

$$\tilde{m}_a \lesssim 10 \text{ TeV} \rightarrow M_L \sim 10^6 \text{ GeV}$$

Dirac SUSY breaking masses couple
gauginos and fermionic components of

S, T, O :

$$\tilde{M}_2 \sim S_L \frac{g_2 \lambda_T}{4\pi^2} \frac{\tilde{m}_{\bar{L}}^2 - \tilde{m}_L^2}{M_L}$$

- small μ and μB terms induced, without
the usual problem of gauge mediated
models

μ SSM

The μ SSM model may also arise in
hidden sector models with gravity
mediated SUSY breaking

- hidden SUSY breaking sector does
NOT contain any SINGLET
- gauged U(1) with a non-vanishing
 D -term

μ SSM

T parameter

h_T and h_S superpotential couplings break
custodial SU(2) symmetry $\Rightarrow$ large
one-loop contributions to T parameter

$$T \sim \frac{h_T^2 v^2}{16\pi m_Z^2 s_W^2 c_W^2} \log \left(\frac{h_T^2 v^2}{m_Z^2} \right)$$

$$m_{\chi^\pm} > 104 \text{ GeV} \rightarrow h_T \sim 1 \rightarrow T \sim 2.5$$

$$m_{\chi^\pm} > 90 \text{ GeV} \rightarrow h_T \sim 0.6 \rightarrow T \sim 0.6$$

But ...

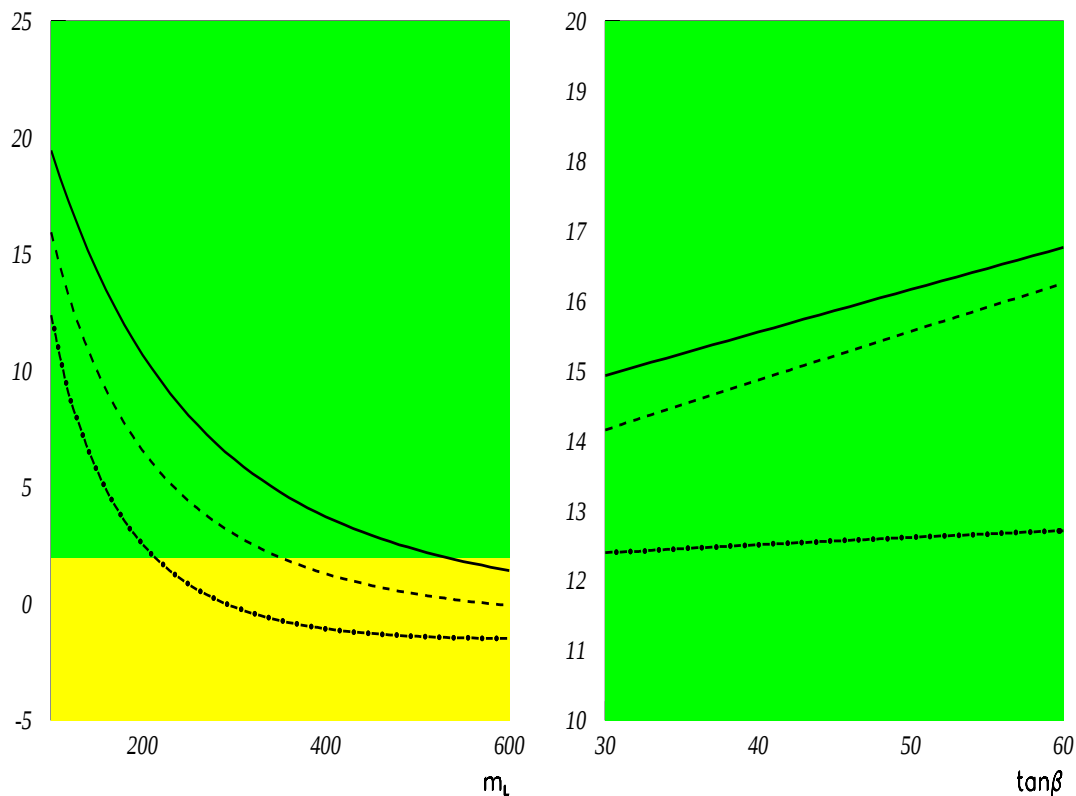
- Non standard contributions to Z width
from lightest nearly Dirac neutralinos
- Oblique approximation not appropriate
for light superpartners
- More careful study needed

μ SSM

Maximum $\delta a_\mu \times 10^{10}$

$$\tan \beta = 60$$

$$m_L = 100 \text{ GeV}$$



$$\tilde{m}_2 = 0 \text{ GeV}$$

$$\tilde{m}_2 = 5 \text{ GeV}$$

$$\tilde{m}_2 = 10 \text{ GeV}$$

Strong suppression due to small Majorana
masses

μ SSM

Unification of couplings

One can embed the T, S, O fields into a complete adjoint multiplet of $SU(3)^3$ or $SU(5)$

$$(1, 2, \pm 1/2) + (1, 1, \pm 1) + (1, 1, \pm 1)$$

$$(3, 2, -5/6) + (\bar{3}, 2, 5/6)$$

$$\text{GUT scale} \sim 10^{18}$$

$$\text{messenger scale} \sim 10^7$$

μ_{SSM}