

SUSY 02  
DESY

GRAND UNIFICATION

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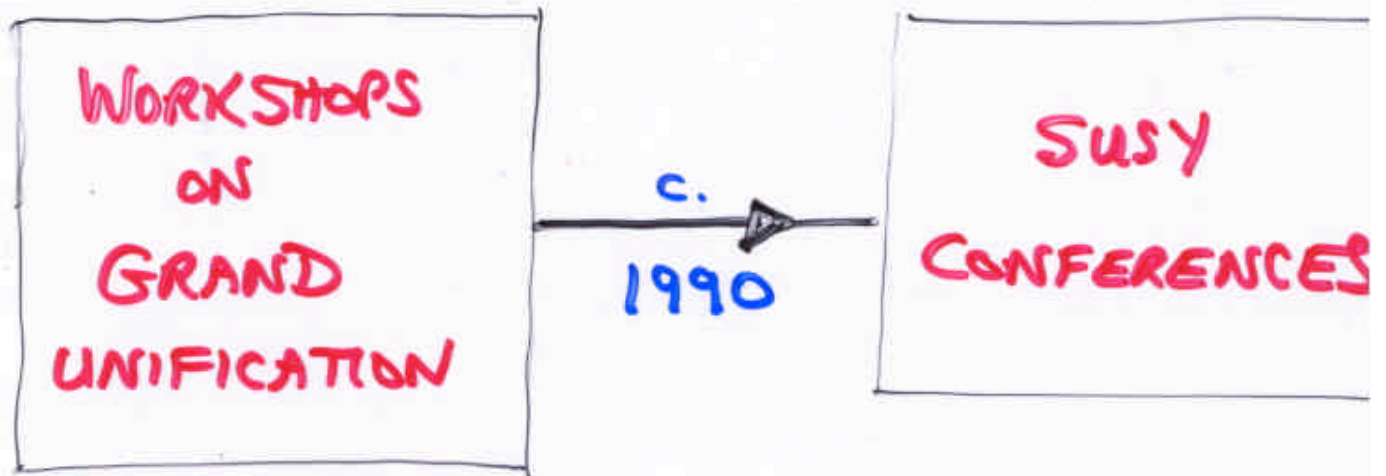
IN

HIGHER DIMENSIONS.

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Lawrence Hall

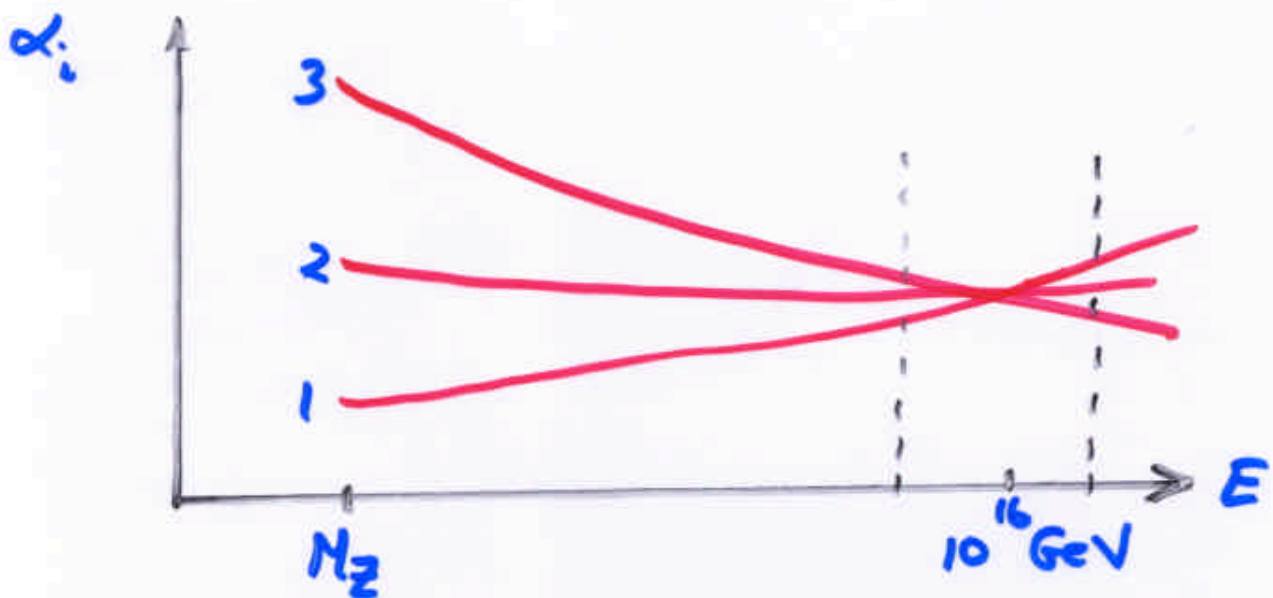
UC Berkeley



Evidence for  
weak scale susy



Evidence for  
grand unification



1. WHAT IS THE PHYSICS AT UNIFICATION?
2. HOW CAN IT BE TESTED?

(I) FEATURES OF 4D GUTS

(II) NEW VIEWPOINT FOR UNIFICATION

hep-ph/0103125

(III) A MINIMAL MODEL

hep-ph/0111068

(IV) EXPERIMENTAL SIGNALS

hep-ph/0205067

with Yasunori  
Nomura.

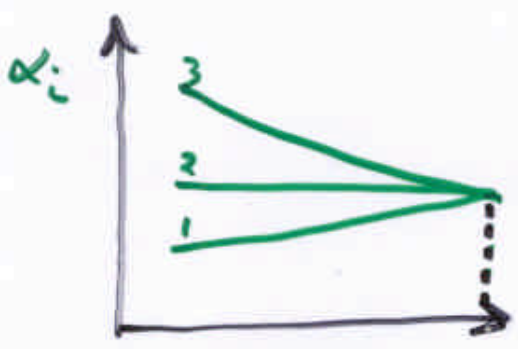
# THE DICHO TOMY

## UNIFICATION FITS

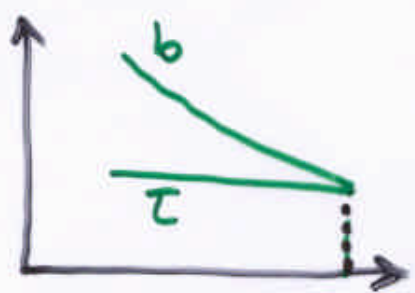
Light matter



## Numerical Prediction

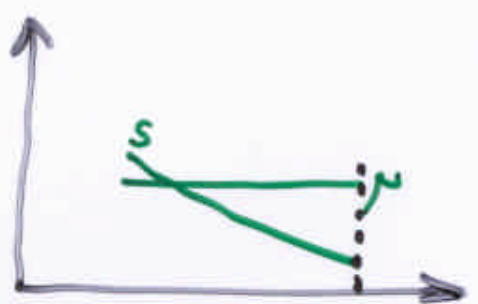
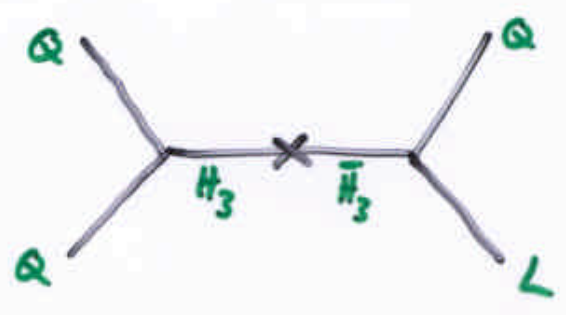


## Quark-lepton masses



## ABHORS UNIFICATION

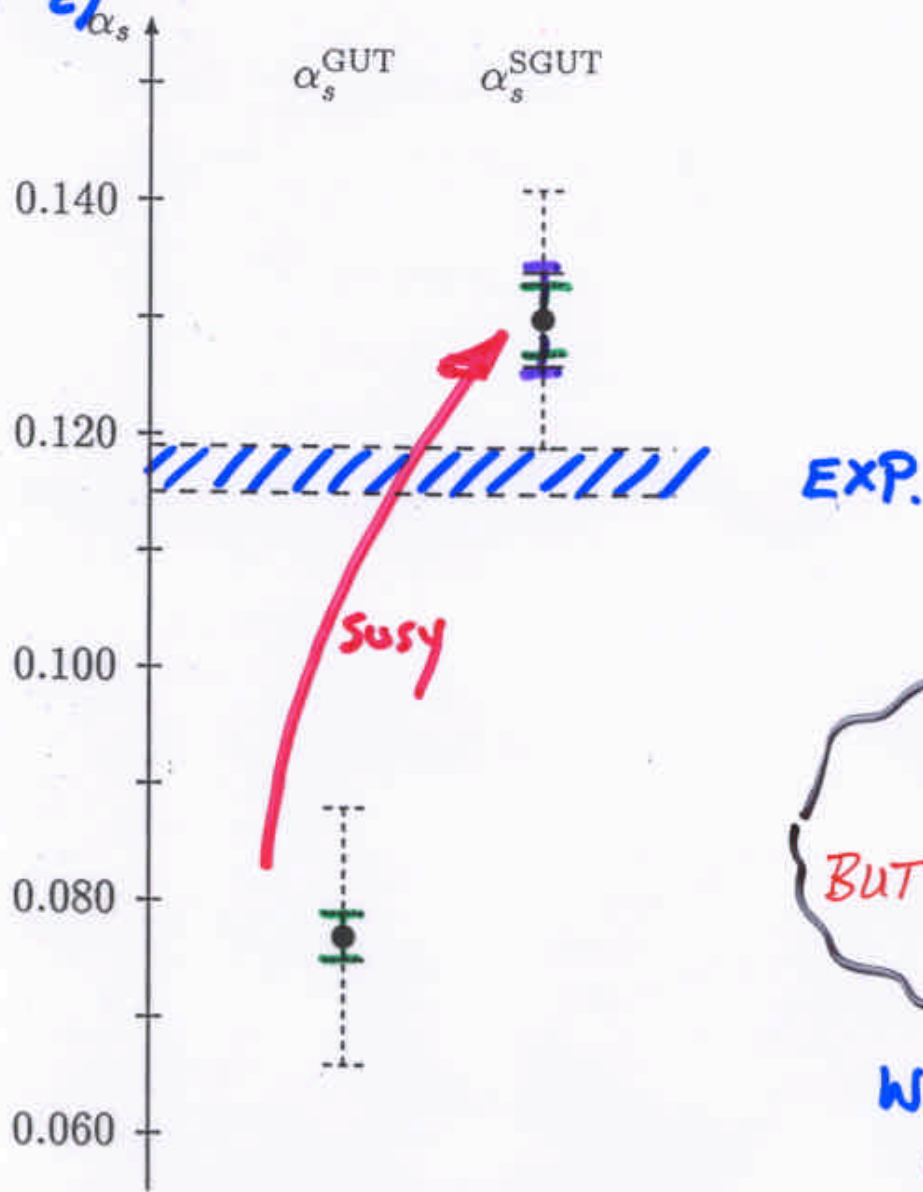
Light Higgs



RECONCILIATION ?

# PREDICTION OF $\alpha_s$

$\alpha_s(M_Z)$



GOOD  
BUT NOT PERFECT

WHAT IS THIS  
TELLING US?

I GUT threshold corr:  $(5+\bar{5})$ , unit log spl.

I Estimate of supersym. threshold corr.



MINIMAL 4D SUPERSYM. SU(5):

- No understanding of why  $h_2$  light.
- Excluded by experimental limit on  $\tau_p$

One is forced to invent mechanisms . . . . .

... a better way forward?

## II

# A NEW PHYSICS OF UNIFICATION

OLD  
TOOLS

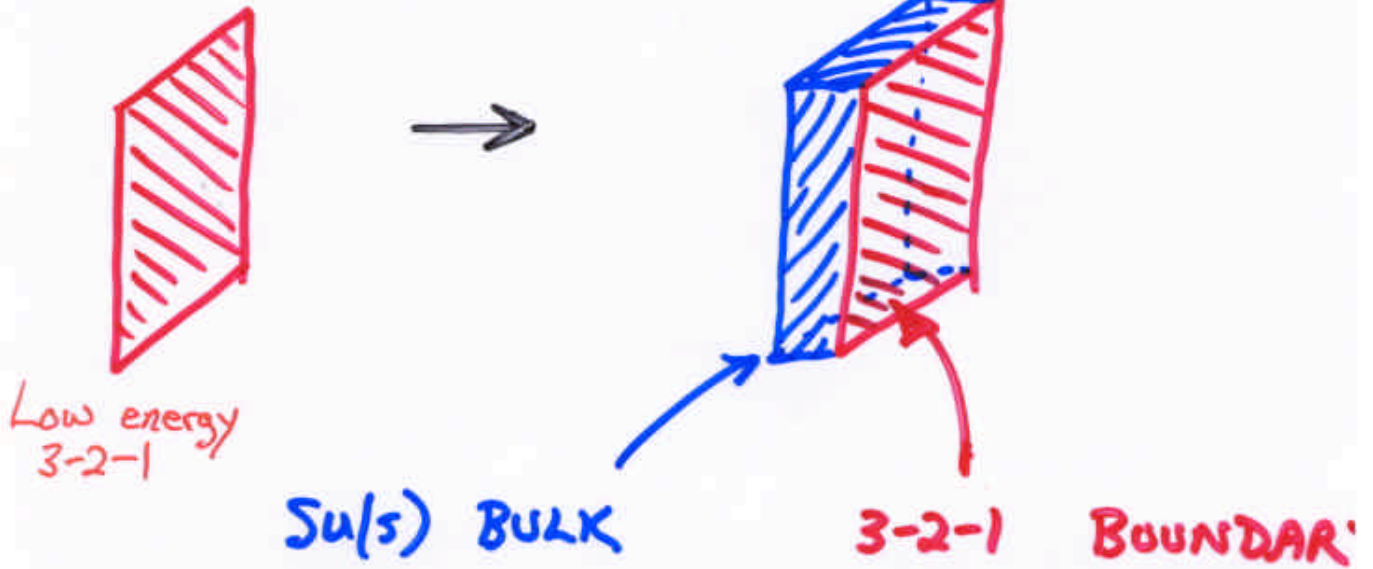


A NEW THEORY  
FROM  $10^{15}$  —  $10^{17}$  GeV

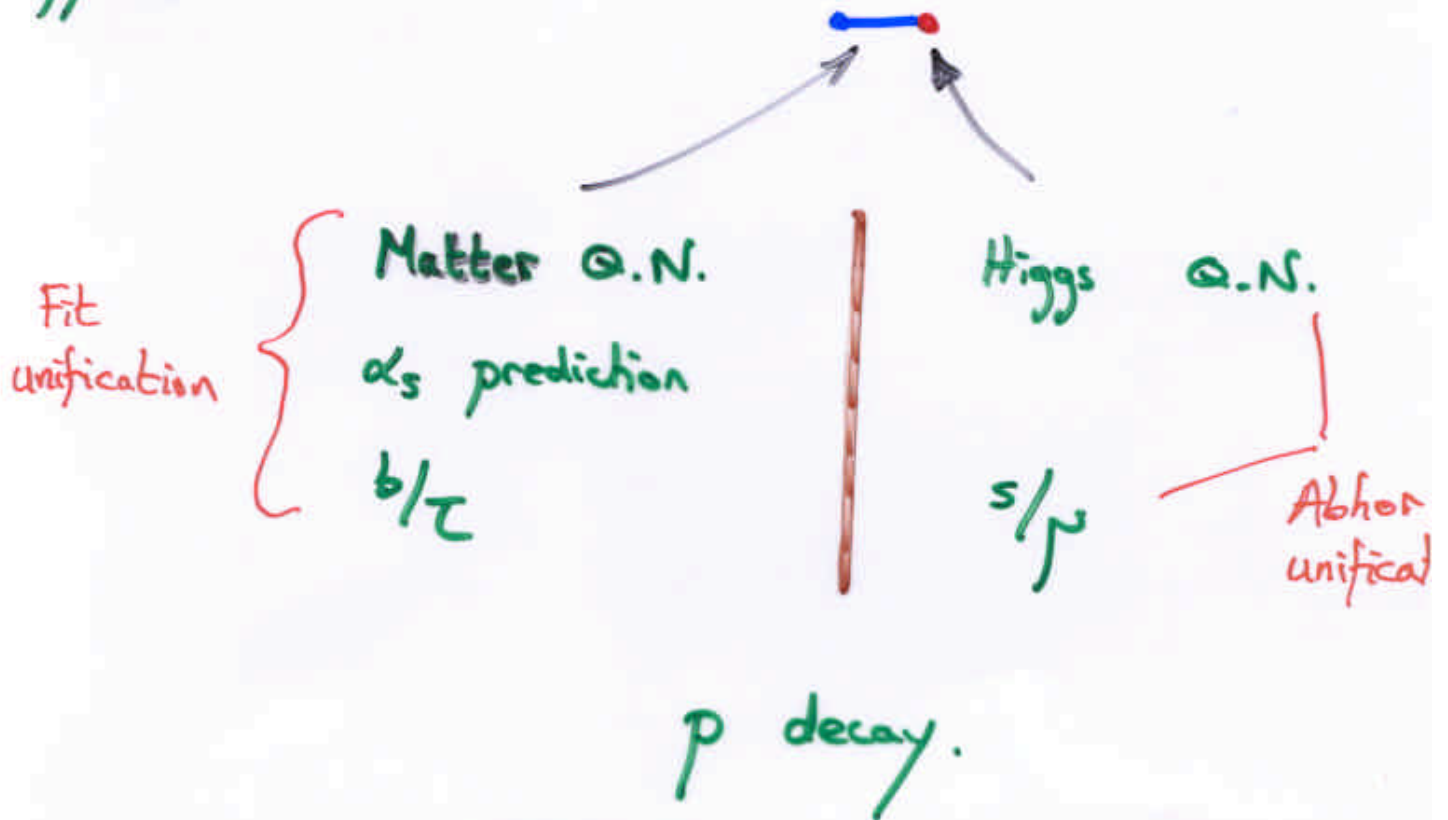
# HIGHER DIMENSIONS

I

$$x \rightarrow (x, y)$$



Suppress  $x$ :

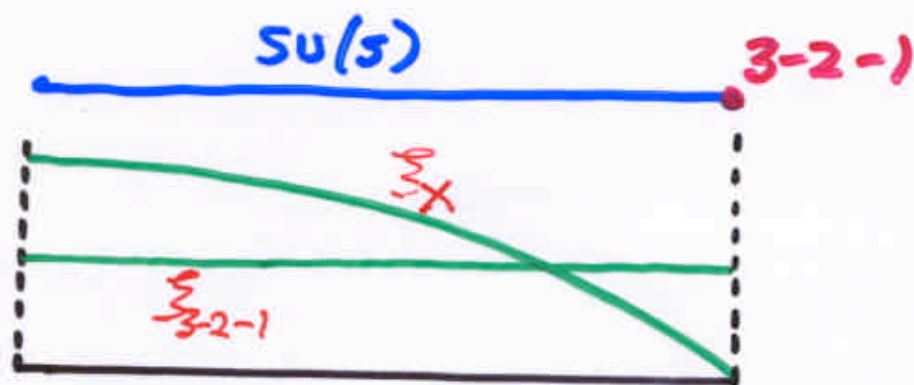




II.2

# ORIGIN OF 3-2-1 POINT DEFECT

RESTRICTED GAUGE SYMMETRY FROM BOUNDARY CONDITIONS : FIELD THEORY IN A BOX :



↑ LOCAL EXPLICIT  
 $SU(5)$  BREAKING

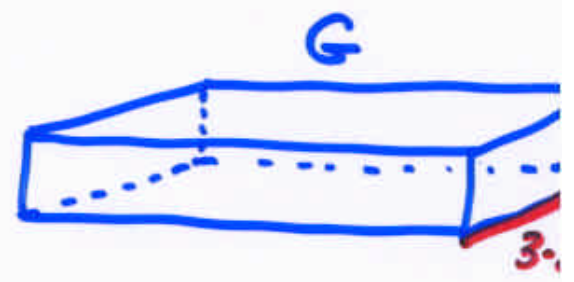
## A NEW CLASS OF PREDICTIVE GAUGE THEORIES.

- SAME UNITARITY BEHAVIOUR AS  $SU(5)$
- \* NO NEED TO BREAK THE UNIFIED GAUGE SYM

# THE SETUP

- Assume

Fixed geometry  
Boundary conditions

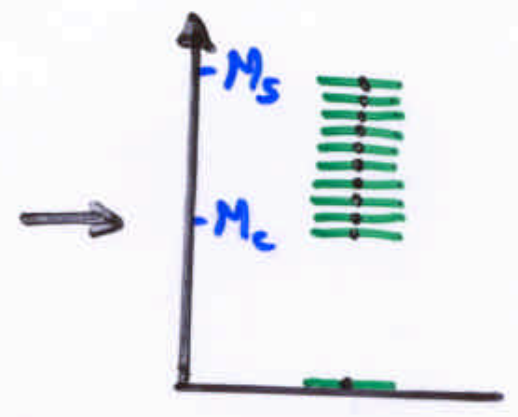


- Introduce

$\phi(y)$  & general action

- "Machine"

$\phi(y)$



GEOMETRY / B.C.



KK STRUCTURE



- No need for mass terms or spontaneous symmetry break

## CONSEQUENCES FOR :

- Gauge coupling unification
- Why  $h_2$  light
- $p$  decay
- quark-lepton mass ratios.

In each case the situation is quite unlike 4D:

HIGHER DIM  $\neq$  4D  
GUTs GUTs

# CONSEQUENCES FOR GAUGE COUPLING UNIF<sup>2</sup>

Cutoff for 4+d GUT  $\rightarrow$

$M_s$

4+d  
GUT

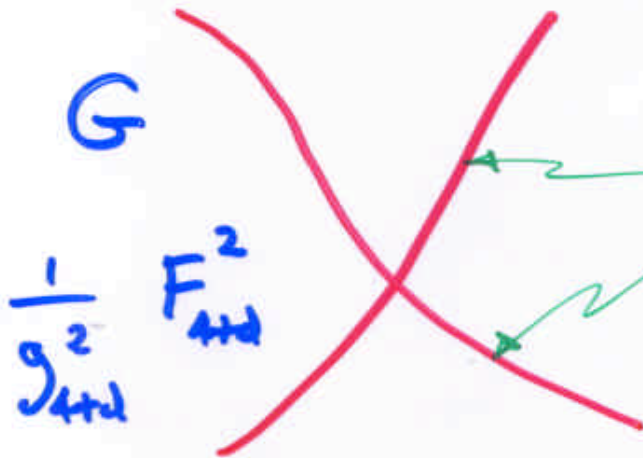
$\frac{1}{R}$  (could be several)  $\rightarrow$

$M_c$

susy  
desert

G

3-2-

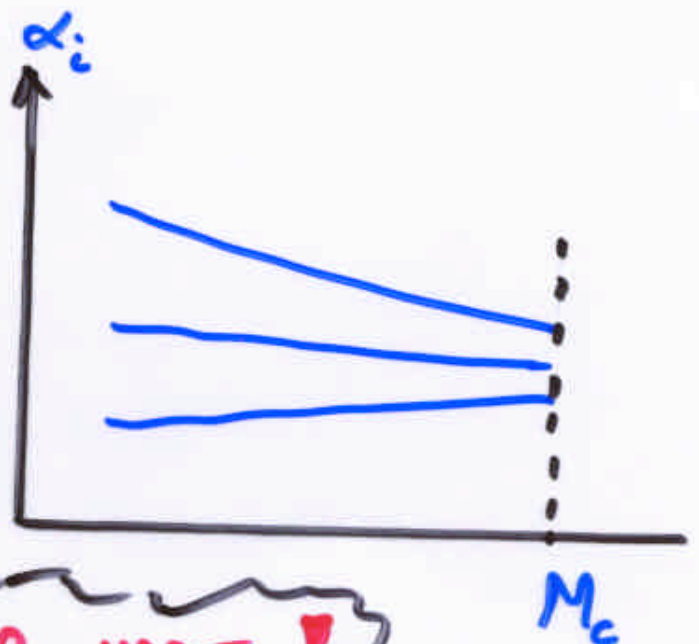


SUBSPACES WITH RESTRICTED GAUGE SYMMETRY:

$$\frac{1}{g_i^2} F_i^2$$

So dy:

$$\frac{1}{g_i^2} = \frac{R^d}{g_{4+d}^2} + \frac{R^{d'}}{\bar{g}_i^2}$$

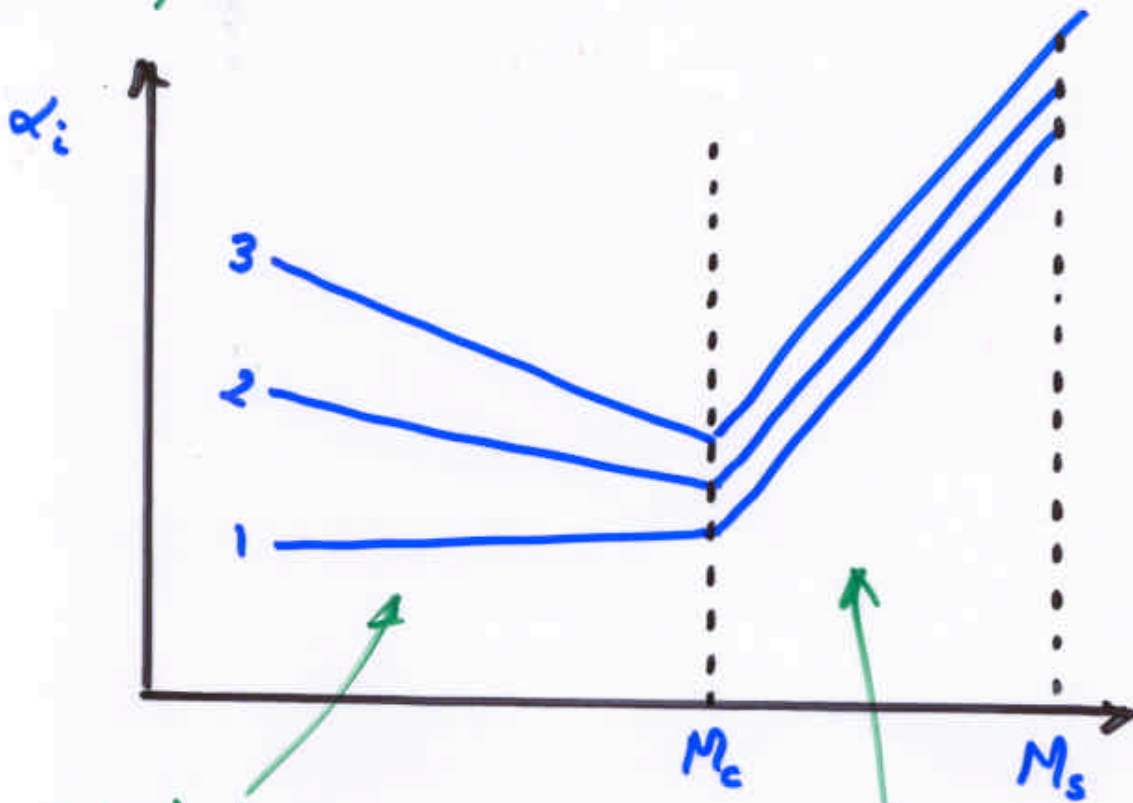


NO GAUGE COUP. UNIF!



# RADIATIVE CORR.

4D viewpoint



usual logs.

power law  
non-universal  
sensitive to  $M_s$

4+d GUT  
with  
RESTRICTED  
GAUGE  
SYMMETRY



EXPLICIT  
LOCAL  
BREAKING  
OF  
G



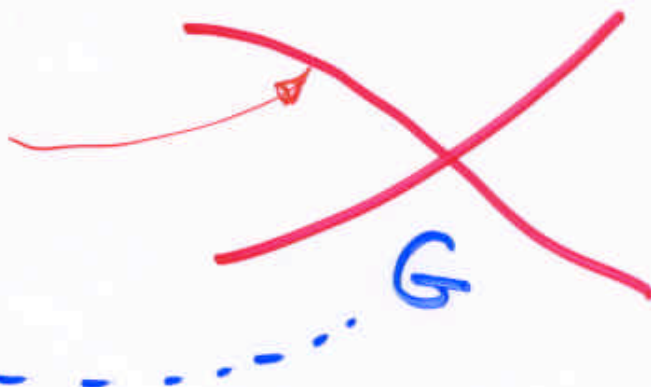
UNIF.  
 $\alpha_i$   
LOST

# NEW PHYSICS OF UNIFICATION

II.1

Hall, Nomura (01)

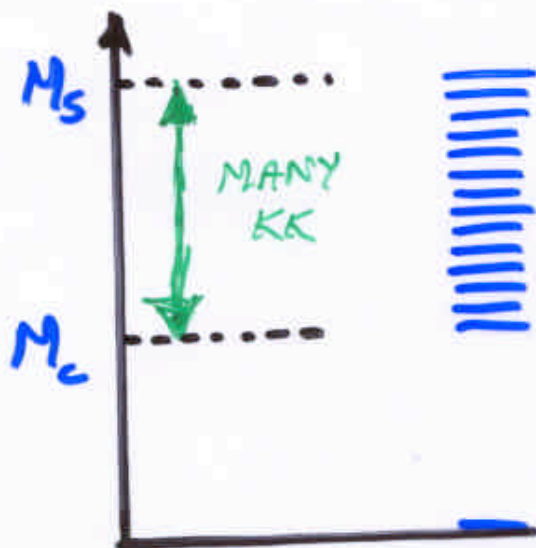
$$\frac{1}{g_i^2} = \frac{R^d}{g_{4+d}^2} + \frac{R^{d'}}{\tilde{g}_i^2}$$



## ① Large vol. of bulk

$$RM_S = \frac{M_S}{M_c} \Rightarrow 1$$

$M_S$  AT/CLOSE TO STRING COUPLING



## ② Local Breaking To 3-2-1 at POINTS

⇒ RELATIVE RUNNING

ABOVE  $M_c$  IS LOG



⇒ SENSITIVITY TO PHYSICS AT  $M_S$  IS  $\frac{1}{100}$

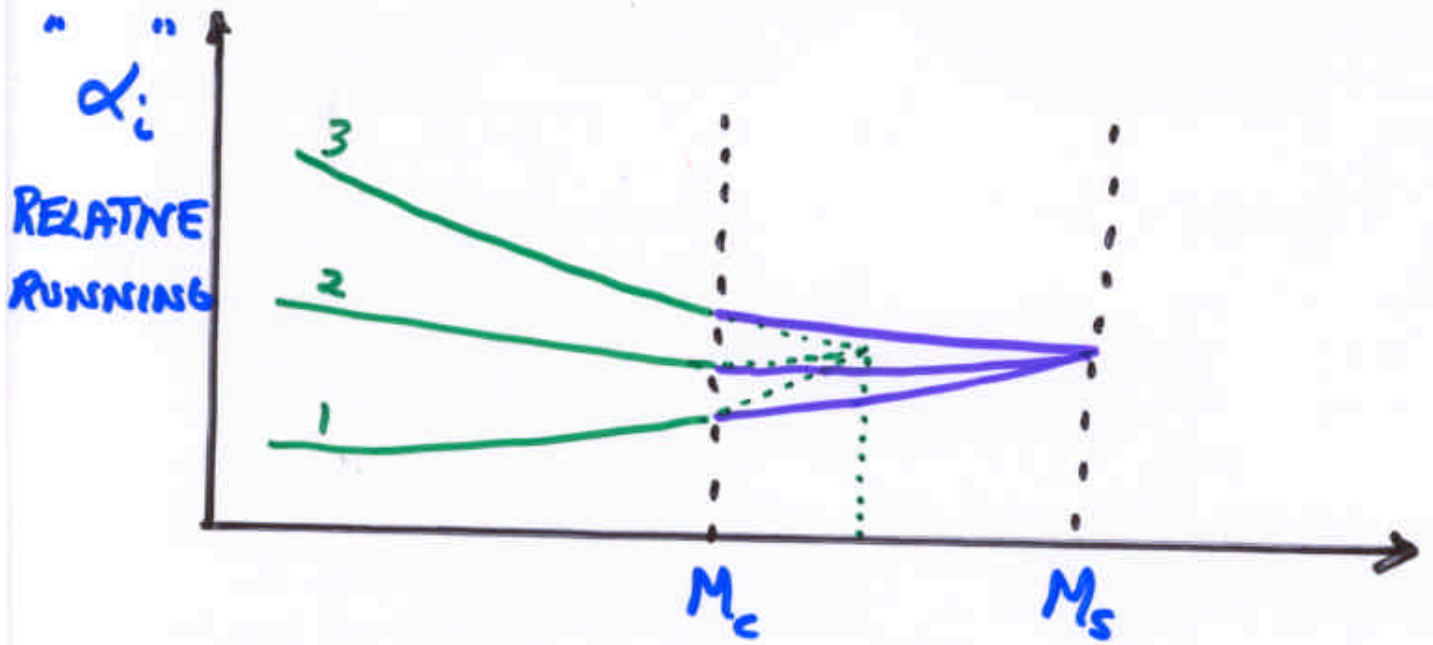
3-2-1

G

4-2-2



# HOW CAN 2 STAGE UNIF. BE PREDICTIVE



$$\frac{1}{\alpha_i(M_Z)} = b_i^{\text{asy}} \ln \frac{M_c}{M_Z} + b_i^{\text{KK}} \ln \frac{M_s}{M_c} + \frac{1}{\alpha_G}$$

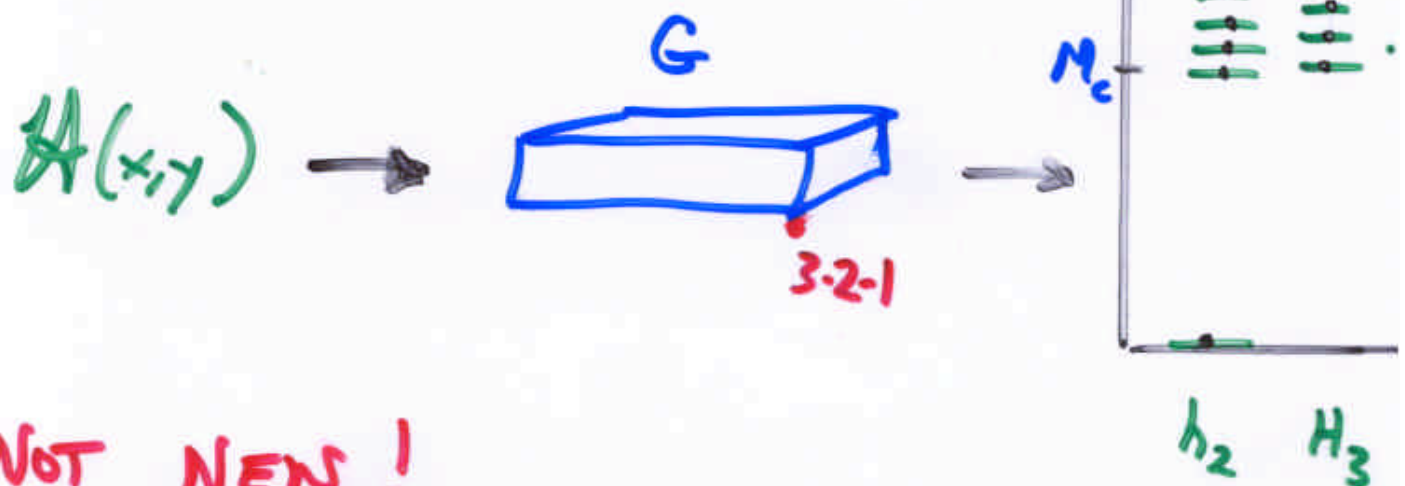
$\frac{M_s}{M_c}$  determined by strong coupling:

$$C \frac{g^2(M_c)}{16\pi^2} \left( \frac{M_s}{M_c} \right)^d \approx 1$$

$$\alpha_s = \alpha_s(d, \text{B.C.}, \Phi(\gamma))$$

# SPLIT MULTIPLETS: THE LIGHT $h_2$

The "machine" automatically creates split multiplets.



NOT NEWS!

Candelas, Horowitz, Strominger, Witten (85)

Dixon, Harvey, Vafa, Witten (85)

Ibanez, Kim, Nilles, Quevedo (87)

⋮

Kawamura (00)

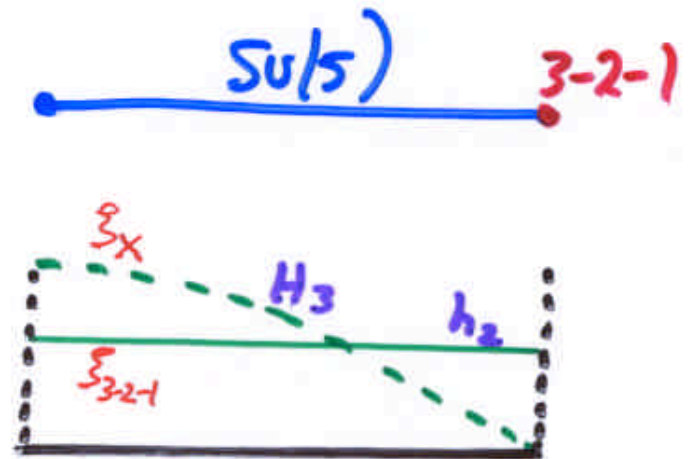
$SU(5)$

$3-2-1$

# d=5 PROTON DECAY

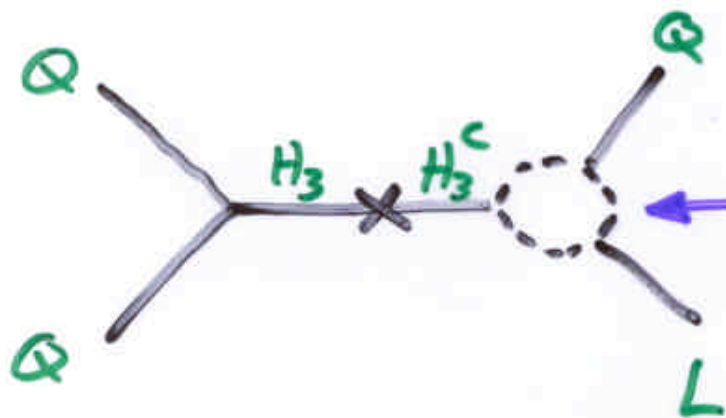
II.

Study masses for  $H_3$ :



$$\mathbb{Z} = H_3 \supset H_3^c \rightarrow H_3 \supset \frac{n}{R} H_3^c, \quad n \neq 1$$

N=2 partners



No coupling by R symmetry

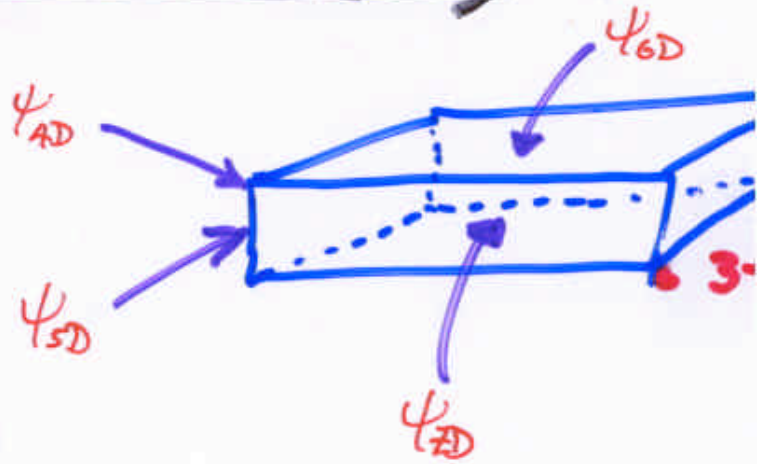
d=4,5  
ABSENT BY  
SYM.

d=6  
 $M_x = M_c = ?$

# QUARK-LEPTON MASS RELATIONS

II. 11

- DEPENDS ON LOCATION:



- Yukawas in 4d  $\delta^d(\bar{y}-\bar{y}_0) \psi_1 \psi_2 H$

•  $m_{12} \propto \frac{1}{\sqrt{V_1}} \frac{1}{\sqrt{V_2}}$  if both touch  $y_0$

\* HEAVIEST FERMIONS LIVE IN 4D & HAVE UNIFIED MASS REL (IF AWAY FROM ...)

\* LIGHTER FERMIONS LIVE IN BULK & DON'T HAVE UNIFIED MASS REL. (IF TOUCH ...)

### III

## THE MINIMAL MODEL

Conceptual framework



Calculable, predictive theories



$\alpha_s$  SUGGESTS  $SU(5)$  IN 5D III

Seek:  $\alpha_s(\text{BULK, B.C., G, H}) \rightarrow 0.117 \pm .002$   
(Exp. at  $M_Z$ )

Recall  $\alpha_s^{\text{SGUT}} \approx 0.130$

Find  $d < 3$

H in bulk (not part of V)

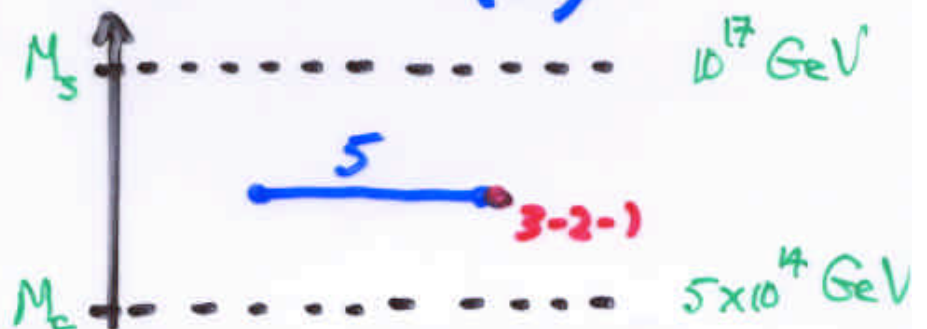
$$\Delta\alpha_s = -\frac{1}{2\pi} \alpha_s^2 \frac{12}{7m} \ln \frac{M_s}{M_c} \left\{ \begin{array}{l} T^2/2m \\ \text{or} \\ m=2 \\ \text{for } d= \end{array} \right.$$

Best if

$d=1 \Rightarrow \left\{ \begin{array}{l} M_s/M_c \approx 200 \\ G = SU(5) \end{array} \right.$

Simplest Model

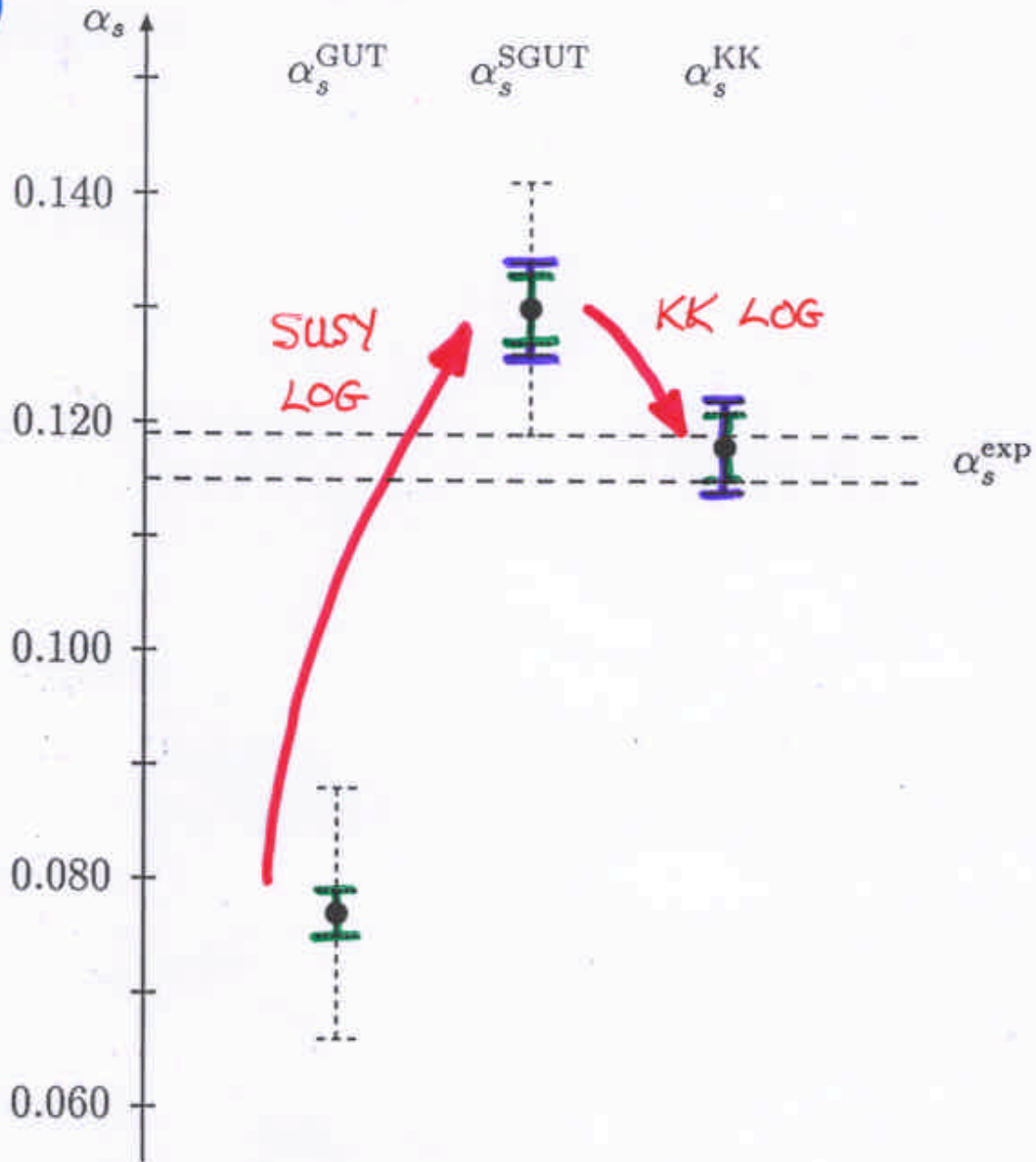
is selected over a large energy interval.





# THE $\alpha_s$ PREDICTIONS

$\alpha_s(M_Z)$



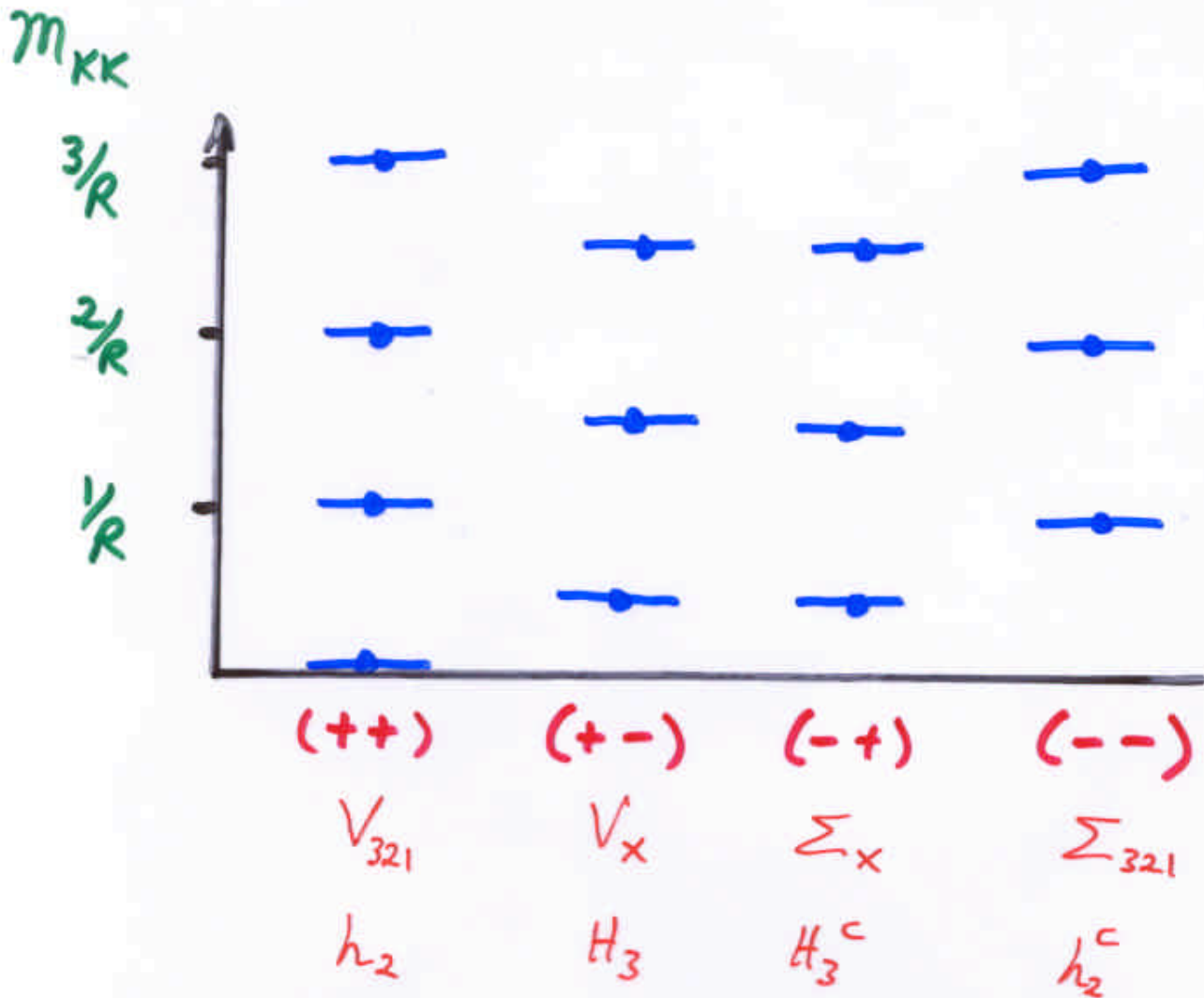
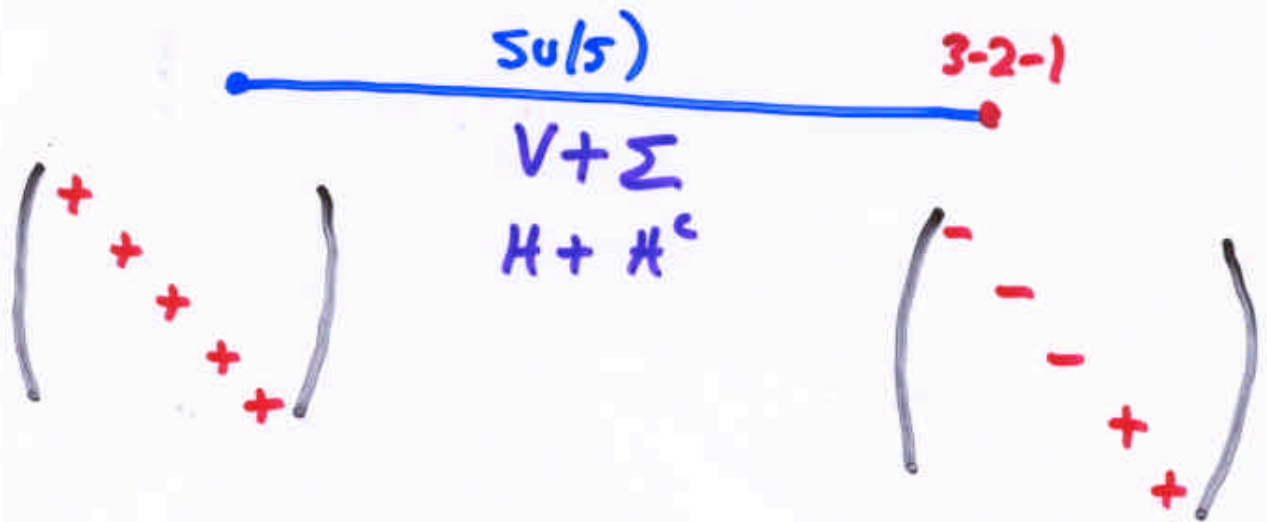
$\bar{\text{I}}$  uncertainty from non-log corr. at high scale

$\text{I}$  uncertainties from susy scale

- 5D Theory with large  $M_s/M_c$  is best fit to data
- Other theories not excluded

# THE KK MODES

III.



cf Kawamura (2000)

# MATTER LOCATION

"Boundary Matter"

"UNIFIED"

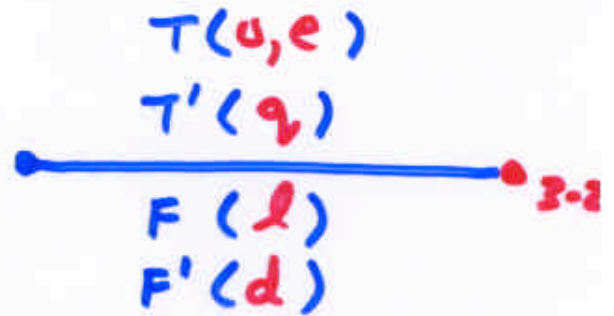
$$S(y) \left\{ T\bar{T}H + T\bar{F}H \right\}$$



"Bulk" Matter

"NON-UNIFIED"

$$S(y) \left\{ \begin{array}{l} T\bar{T}H + T\bar{T}'H + \dots \\ T\bar{F}H + T\bar{F}'H + \dots \end{array} \right\}$$



Boundary	$1/R^{1/2}$	SUL(S) REL.
Bulk	$1/R^{3/2}$	NO REL.

Hall, Nomura ph/0103  
Hebecker, March-Russell ph/0106

$T_3$  must be on boundary

$F_3$  on boundary for  $b/c$

# BOUNDARY INTERACTIONS

III.

4D, $N=1$ allows:	$FH, \bar{H}H$	NO
	$TTH, TF\bar{H}$	YES
	$TFF$	NO
	$TTTF$	NO

$N=2$  possesses continuous  $R$  sym.

$$\text{eg } [H^{(0)} \partial_y H^{(2)}]_F$$

Extend to boundary interactions:

$$T^{(1)} \quad F^{(1)} \quad H^{(0)} \quad \bar{H}^{(0)} \quad H^{(2)} \quad \bar{H}^{(2)}$$

Complete sol<sup>n</sup> to

- $\Delta_B, \Delta_L \neq 0$  at  $d=4,5$
- $h_2$  massless

## IV

# EXPERIMENTAL SIGNALS

- $d=6$   $p$  decay
- Theories of quark & lepton masses.

Hall, March-Russell, Okui, Smith [ph/0108161](#)

Hebecker, March-Russell [ph/0205143](#)

Talk by John March-Russell

If  $m$  supersym. breaking has high messenger scale

- Superpartner spectrum
- Lepton flavor violation

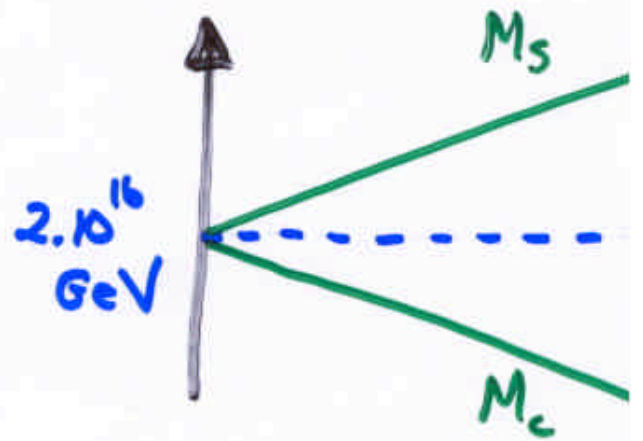


# $d = 6$ PROTON DECAY

IV.

①

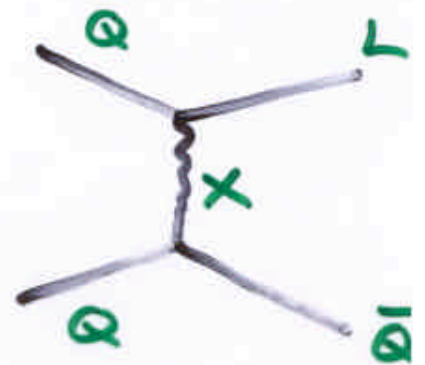
$$M_x = M_c$$



②

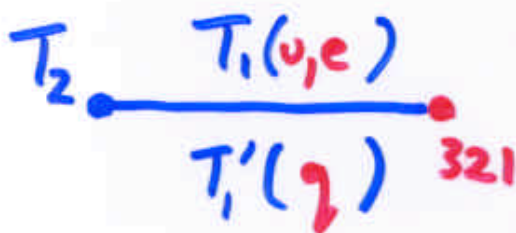


can give  
 $p \rightarrow e^+ \pi^0$



③

For large  $\frac{M_s}{M_c}$ ,  $T_1$  must be in bulk



$$p \rightarrow K^+ \nu$$

via CKM mixing.

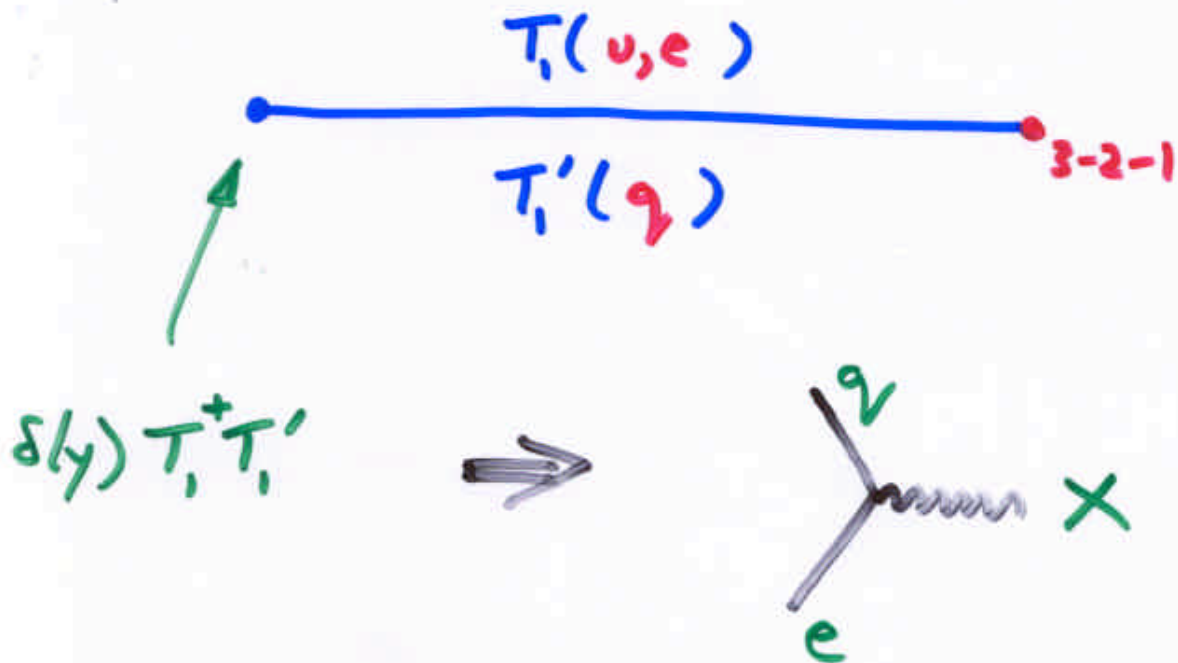
Nomura [ph/0108170](#)

Hebecker, March-Russell [ph/0204037](#)



# ④ Boundary Gauge Interactions

N.2



Minimal model gives

$$\tau_p \approx 10^{34} \text{ yrs for } l^+ \pi^0, l^+ K^0, \bar{\nu} \pi^+, \bar{\nu} K^0$$

$\swarrow \quad \searrow$   
 $e \text{ or } \mu$

~~SUSY~~ :  $M_{\text{mess}} > M_c$



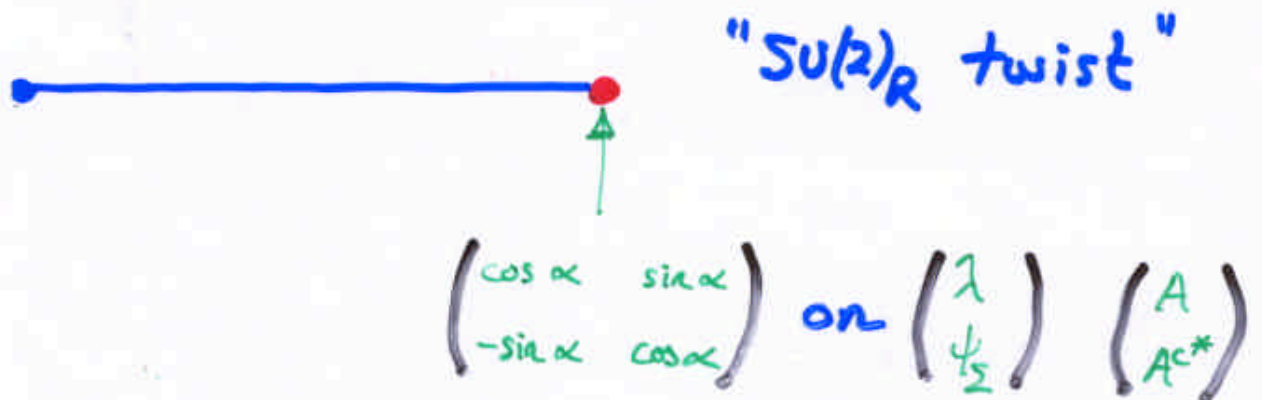
Various possibilities for  $T_2, F_1, F_2$

$U(3)_T \times U(3)_F \xrightarrow{\text{matter location}} \dots$

Expect:

- Non-universal squark/slepton masses
- FCNC from superpartner exchange

Ex: Susy breaking B.C.



$\alpha \sim 10^{-13}$  !

DYNAMICS OF F  
COMPONENT OF RADION

Marti Pomarol th/01/06/02

SOFT OPS. AT  $M_c$

All bulk superpartners:

$$\tilde{m} = \frac{\alpha}{R}$$

All boundary superpartners:

$$0$$

Tilinear A parameters:

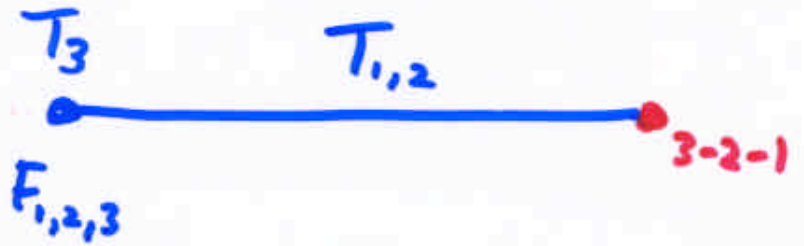
$$1, 2, 3$$

(counts # scalars in bulk,

# CONSEQUENCES

IV.5

Unique  
matter  
locations :



Param.

$\tilde{m}, \mu, B$

ENSB



$\tilde{m}, \tan \beta$



superpartner  
spectrum

# PREDICTIONS FOR

IV.6



3-2-1

$$\begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix}$$

Probe  
matter  
location:

Boundary  
squarks  
light

Boundary  
sleptons  
light.

$\tilde{m}$	300	400
$\tan \beta$	5	10
$\tilde{g}$	699	911
$\tilde{\chi}_1^\pm$	251	334
$\tilde{\chi}_2^\pm$	427	531
$\tilde{\chi}_1^0$	130	175
$\tilde{\chi}_2^0$	251	334
$\tilde{\chi}_3^0$	417	518
$\tilde{\chi}_4^0$	422	528
$\tilde{q}$	701	915
$\tilde{u}$	675	880
$\tilde{d}$	602	780
$\tilde{l}$	209	277
$\tilde{e}$	317	422
$\tilde{t}_1$	425	547
$\tilde{t}_2$	619	780
$\tilde{b}_1$	563	727
$\tilde{b}_2$	601	774
$\tilde{\tau}_1$	106	126
$\tilde{\tau}_2$	214	280
$h$	118	128
$A$	552	690
$H^0$	553	690
$H^\pm$	558	695
$\alpha_s(M_Z) \{\pm 0.003\}$	0.119	0.118
$m_b(M_Z) \{\pm 0.10\}$	3.37	3.26
$\text{Br}(\mu \rightarrow e\gamma)$	$6 \times 10^{-12}$	$8 \times 10^{-12}$
$\text{Br}(\mu \rightarrow 3e)$	$4 \times 10^{-14}$	$5 \times 10^{-14}$
$\text{Cr}(\mu \rightarrow e; {}^{48}\text{Ti})$	$4 \times 10^{-14}$	$5 \times 10^{-14}$
$\text{Br}(\tau \rightarrow \mu\gamma)$	$1 \times 10^{-8}$	$1 \times 10^{-8}$

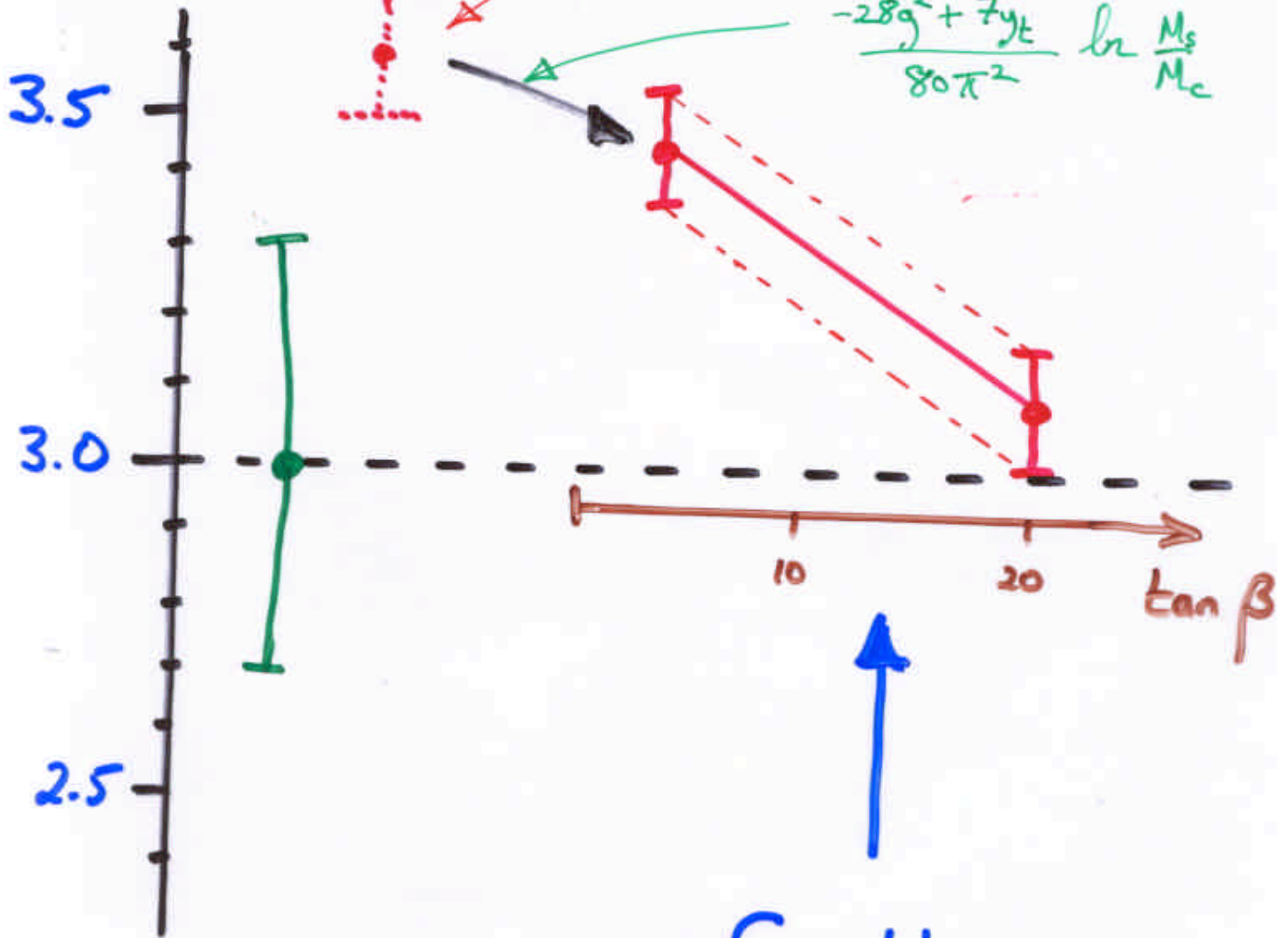
MASS  
EIGENVALUE  
IN GEV

includes  
susy  
threshold  
corrections



# $b/\tau$ MASS RELATION

$m_b (M_Z)$



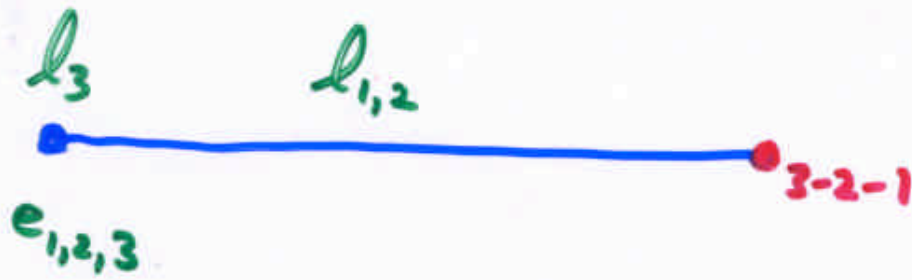
4D susyGUTs

$$\frac{-28g^2 + 7y_t^2}{80\pi^2} \ln \frac{M_s}{M_c}$$

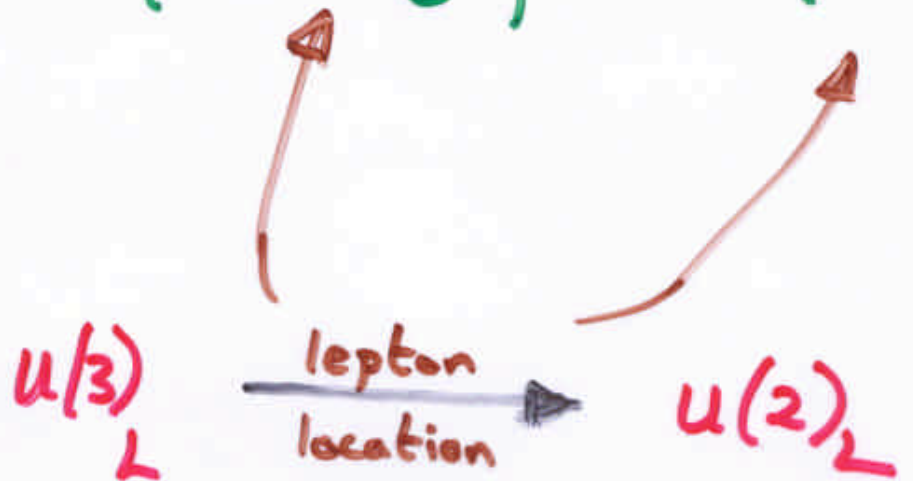
Complete success for moderate  $\tan \beta$ .



# LEPTON FLAVOR VIOLATION IV



$$m_E^2 = 0, \quad m_L^2 = \begin{pmatrix} \tilde{m}^2 & & \\ & \tilde{m}^2 & \\ & & 0 \end{pmatrix}, \quad A_E = \begin{pmatrix} 2 & & \\ & 2 & \\ & & 1 \end{pmatrix}$$

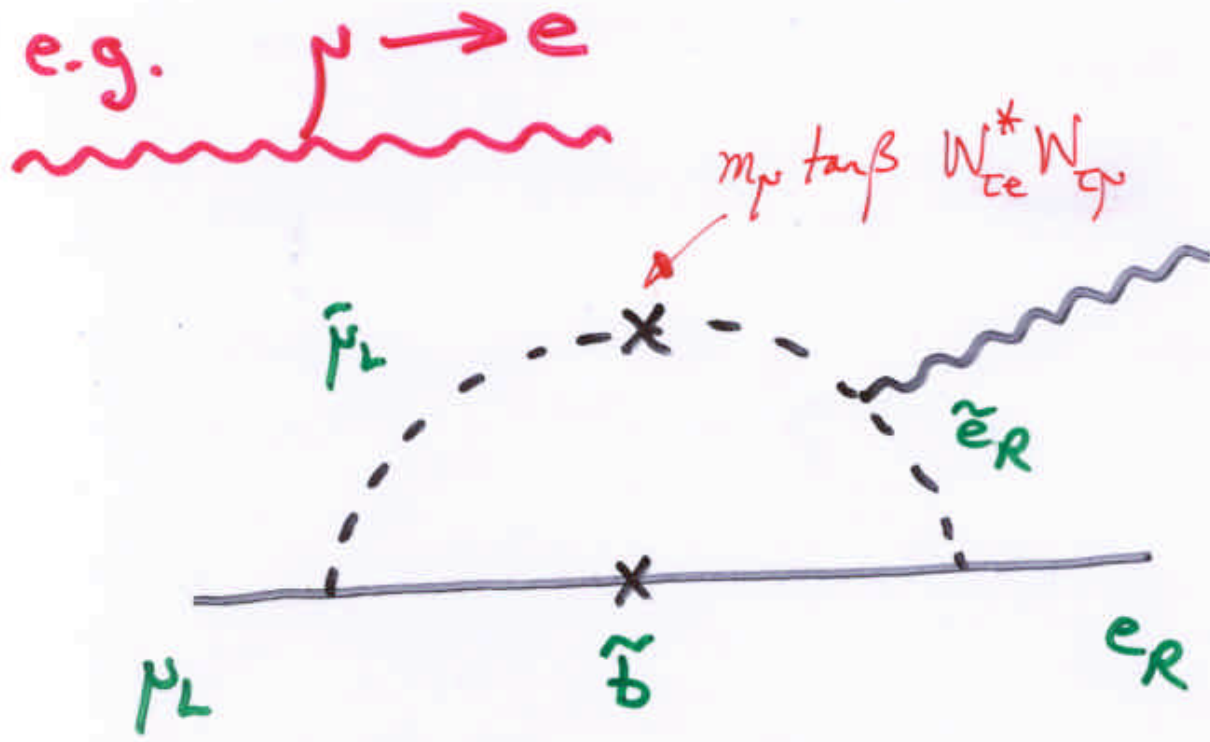


$O(1)$  tree-level effect !

Much larger than  $O\left(\frac{y_t^2}{16\pi^2}\right)$

4D susyGUT effect

Hall, Kostelecky, Raby (81)  
Barbieri, Hall (94)



$1.2 \times 10^{-11}$	$\xrightarrow{\text{PSI}}$	$10^{-14}$
$4.3 \times 10^{-12}$	$\xrightarrow{\text{MECO}}$	$10^{-16}$
$1.1 \times 10^{-6}$	$\xrightarrow{\text{B factories}}$	$10^{-7}$

VERY POWERFUL PROBE

$$\text{Br}(\mu \rightarrow e\gamma) \simeq \underline{3 \times 10^{-11}} \left(\frac{200 \text{ GeV}}{\tilde{m}}\right)^4 \left(\frac{|W_{\tau\mu}^e|}{0.04}\right)^2 \left(\frac{|W_{\tau e}^e|}{0.01}\right)^2 \left(\frac{\tan\beta}{5.0}\right)^2$$


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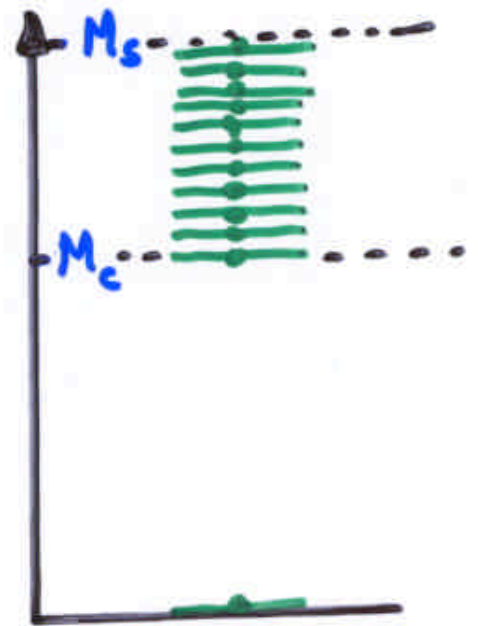
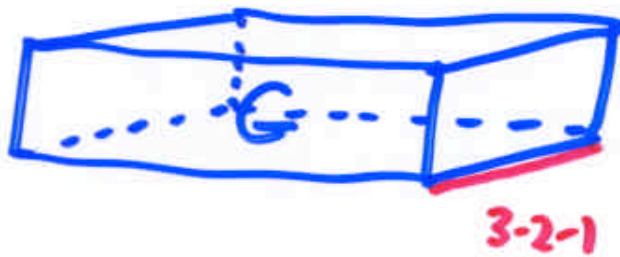

$$\text{Cr}(\mu \rightarrow e; {}^{48}\text{Ti}) \simeq \underline{2 \times 10^{-13}} \left(\frac{200 \text{ GeV}}{\tilde{m}}\right)^4 \left(\frac{|W_{\tau\mu}^e|}{0.04}\right)^2 \left(\frac{|W_{\tau e}^e|}{0.01}\right)^2 \left(\frac{\tan\beta}{5.0}\right)^2$$


---


$$\text{Br}(\tau \rightarrow \mu\gamma) \simeq \underline{5 \times 10^{-8}} \left(\frac{200 \text{ GeV}}{\tilde{m}}\right)^4 \left(\frac{|W_{\tau\mu}^e|}{0.04}\right)^2 \left(\frac{|W_{\tau\tau}^e|}{1.0}\right)^2 \left(\frac{\tan\beta}{5.0}\right)^2$$

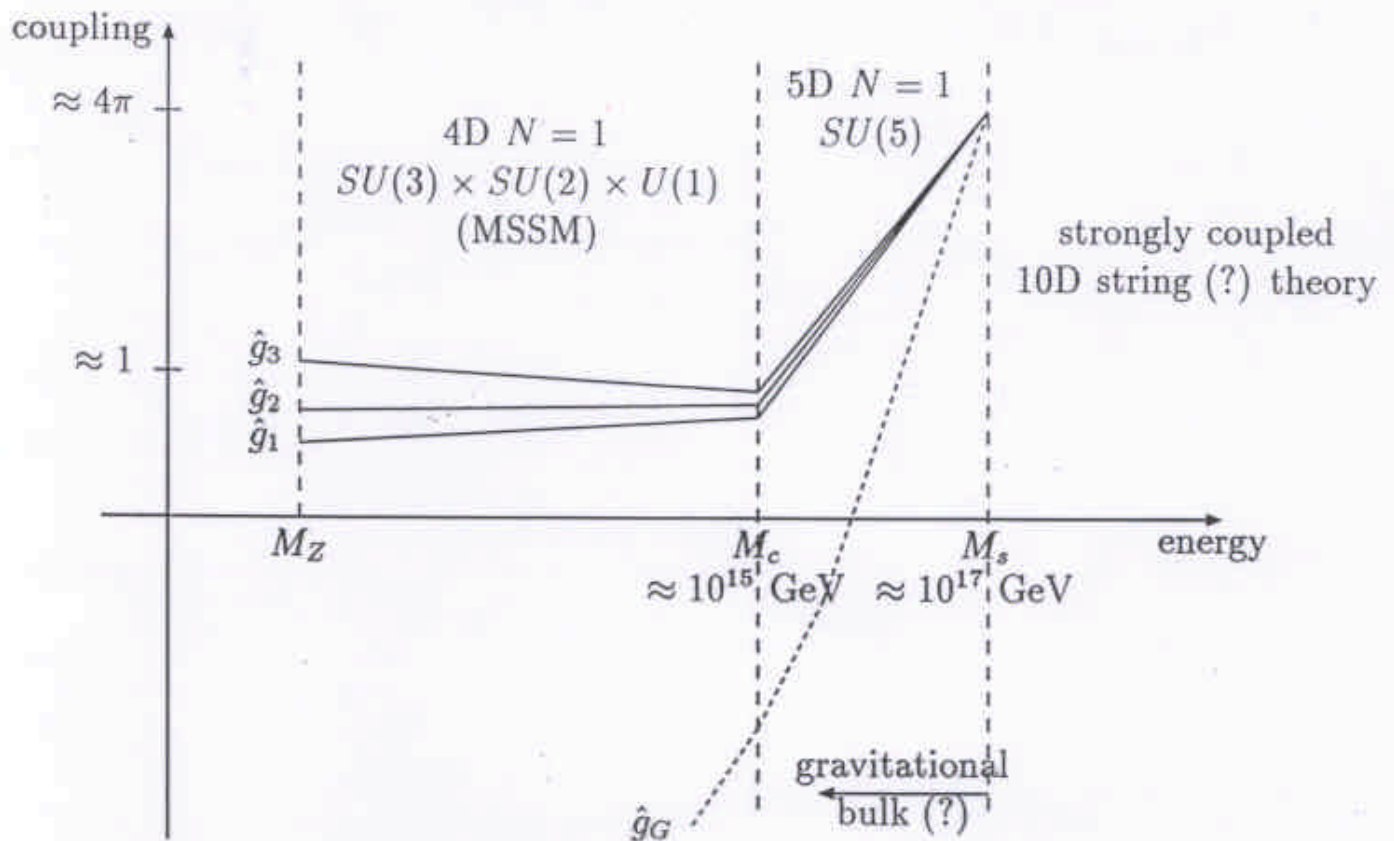
# CONCLUSIONS

- Alternative physics for unification at  $10^{16}$  GeV



- Solves dichotomy of 4D susy GUTs
  - $d=5$  p decay:  $u(1)_R$
  - $q-l$  mass rel. only for heavy  $g$
  - $h_2 - H_3$  splitting

# MINIMAL MODEL WORKS BEST



- $\alpha_s(M_Z) = .118 \pm .003$

Exp

$.117 \pm .003$

- $m_b(M_Z) = 3.3 - .02(\tan\beta - 10) \text{ GeV}$

$3.0 \pm 0$

- Some of flavor from  $M_s/M_c \approx 10^2$

- $d=6$   $p$  decay  $p \rightarrow \ell^+ \pi^0, \dots$

- \* Predictive superpartner spectra

- \*  $p \rightarrow e$  transitions

} large  $M_{\text{pl}}$