

①

DECONSTRUCTION,

G_2 HOLONOMY,

AND DOUBLET - TRIPLET

SPLITTING

②

I'LL WORK IN THE CONTEXT
OF SU(5) SUSY GUTS AND THEIR
CLOSE COUSINS ...

HIGGS BOSONS $2 \oplus \bar{2}$ OF SU(2) x U(1)

$$2 \oplus 3 = 5$$

$$\bar{2} \oplus \bar{3} = \bar{5}$$

TO GET THE ELECTROWEAK HIERARCHY
PLUS A LONG-LIVED PROTON,
WE NEED TO GIVE GUT MASSES
TO THE $3 \oplus \bar{3}$ - WHILE KEEPING
THE $2 \oplus \bar{2}$ LIGHT

③

THAT IS ALSO WHAT WE NEED
TO GET THE RIGHT SUSY-GUT
FORMULA FOR GAUGE COUPLINGS

BUT WHY ARE THE $2, \bar{2}$ LIGHT
IF THE $3, \bar{3}$ HAVE GUT MASSES?

AN OBVIOUS POSSIBILITY IS TO
FIND A GLOBAL SYMMETRY OF
THE LOW ENERGY WORLD THAT
ALLOWS THE $3-\bar{3}$ MASS

(i.e. A SUPERPOTENTIAL $\int d^2\theta 3-\bar{3}$)

AND FORBIDS THE $2-\bar{2}$ MASS

(4)

HOWEVER, THIS IS IMPOSSIBLE IN
FOUR-DIMENSIONAL $SU(5)$ THEORIES WITH
ONLY FINITELY MANY FIELDS
(GOODMAN & EW 1986)

MOST GENERAL GLOBAL SYMMETRY
ACTS ON THE 5 AS

$$e^{i\phi} \begin{pmatrix} e^{2i\alpha} & & & & \\ & e^{2i\alpha} & & & \\ & & e^{2i\alpha} & & \\ & & & e^{-3i\alpha} & \\ & & & & e^{-3i\alpha} \end{pmatrix}$$

GLOBAL
SYMMETRY
THAT COMMUTES
WITH $SU(5)$

$SU(5)$ GAUGE
TRANSFORMATION
THAT COMMUTES WITH
STANDARD MODEL

⑤

IT LIKEWISE ACTS ON THE $\bar{5}$

AS

$$e^{-i\phi} \tilde{\phi} \begin{pmatrix} e^{-2i\alpha} & & & & & \\ & e^{-2i\alpha} & & & & \\ & & e^{-2i\alpha} & & & \\ & & & e^{3i\alpha} & & \\ & & & & e^{3i\alpha} & \\ & & & & & e^{3i\alpha} \end{pmatrix}$$

$\tilde{\phi}$ CAN BE ANYTHING

SAME α AS FOR THE 5

THE $3-\bar{3}$ AND ~~2~~ $2-\bar{2}$ MASSES

BOTH TRANSFORM AS $e^{i(\phi-\tilde{\phi})}$

SO ONE IS ALLOWED IF AND ONLY IF THE OTHER IS

CAN WE AVOID THIS RESULT ⑥

BY ADDING MORE FIELDS?

LET'S TRY:

WE START WITH A $5 \oplus \bar{5}$

WITH $\phi, \tilde{\phi}$ CHOSEN SO THE $2 \oplus \bar{2}$

ARE LIGHT. THE TROUBLE IS THAT

THE $3 \oplus \bar{3}$ ARE ALSO LIGHT.

TO ELIMINATE THEM WE ADD A

NEW $5' \oplus \bar{5}'$ WITH QUANTUM

NUMBERS UNDER THE GLOBAL

SYMMETRY SUCH THAT

$3\bar{3}'$ AND $3'\bar{3}$ MASSES ARE ALLOWED

(7)

THIS ELIMINATES THE TRIPLETS

BUT NOW THE $2'$, $\bar{2}'$ MAKE

TROUBLE AS $2\bar{2}'$ AND $2'\bar{2}$

MASSSES ARE POSSIBLE, SO THAT

WE NO LONGER HAVE NATURALLY LIGHT

HIGGS BOSONS. TO AVOID THIS, WE

NEED TO ELIMINATE THE $2' \oplus \bar{2}'$ BY

ADDING A NEW $5'' \oplus \bar{5}''$ WITH

$2'\bar{2}''$ AND $2''\bar{2}'$ MASSES... BUT

NOW WE HAVE AN UNWANTED $3'' \oplus \bar{3}''$

SO WE NEED TO ADD A $5''' \oplus \bar{5}'''$

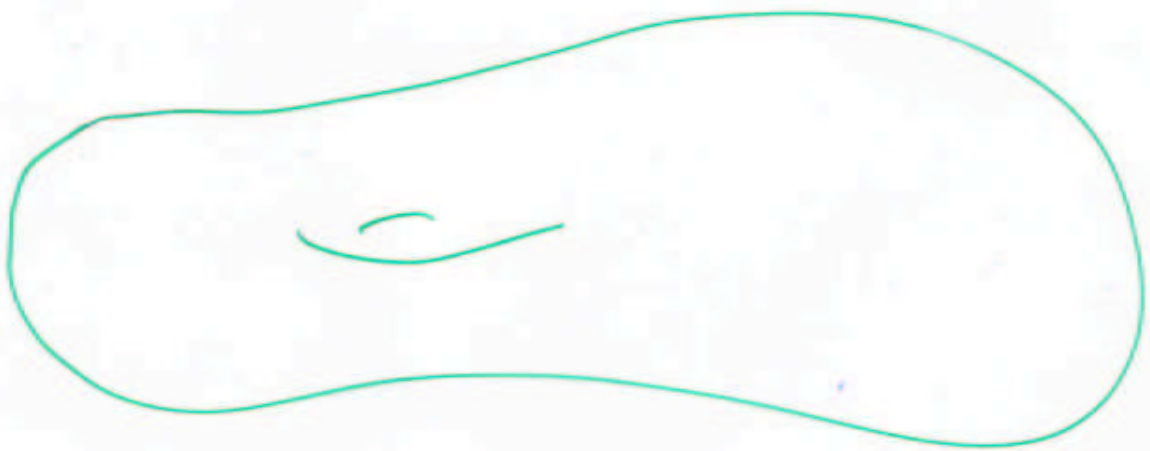
TO ELIMINATE IT.....

⑧

AS LONG AS WE HAVE ONLY
FINITELY MANY FOUR-DIMENSIONAL
 $SU(5)$ MULTIPLETS, WE'LL ALWAYS
RUN INTO TROUBLE AT THE END
OF THE CHAIN... BUT ALL
IS WELL IF WE ARE WILLING TO
USE INFINITELY MANY $5 \oplus \bar{5}$ PAIRS,
OF INCREASING MASSES

9

THIS INFINITE COLLECTION OF
FIELDS CAN HAVE A NATURAL
INTERPRETATION AS KALUZA-KLEIN
HARMONICS IN A HIGHER-DIMENSIONAL
THEORY



GUT SYMMETRY BREAKING BY
WILSON LINES + TOPOLOGICAL
TRICKS SO THE $5, \bar{5}$ TRANSFORM
DIFFERENTLY UNDER GLOBAL SYMMETRY ...

(9½)

(AS IN THE ORIGINAL

(ALABI-YAU CONSTRUCTIONS

Candelas, Horowitz,

Strominger & EW 1984

THOUGH IN THAT CASE

GLOBAL SYMMETRIES

WERE NOT USED)

(10)

FOR CALABI-YAU COMPACTIFICATION,
THIS APPROACH WAS WORKED OUT
IN THE MID-80'S ... A CHARACTERISTIC
PROPERTY IS THAT THERE IS NO ENERGY
SCALE AT WHICH THE MODEL IS A
FOUR-DIMENSIONAL GUT...

THE USUAL GUT ~~PREDICTIONS~~ PREDICTIONS
FOR

FERMION QUANTUM NUMBERS

GAUGE COUPLING UNIFICATION

HOLD JUST AS IN FOUR-DIMENSIONAL
GUT'S, BUT THERE ARE SOME DEPARTURES

9

FOR ONE THING, COUPLING

UNIFICATION IS MORE ROBUST

THAN IN 4-DIM'L GUTS, WHERE

IT CAN BE SPOILED BY TERMS

OF ORDER M_{GUT}/M_{string}

$$\text{Tr } \phi^n F_{\mu\nu} F^{\mu\nu} \quad n=1,2$$

WITH SYMMETRY BREAKING BY

WILSON LOOPS ϕ IS REPLACED BY

$$A_\alpha, \alpha > 4$$



AND ANYTHING LIKE

$$\text{Tr } \overbrace{F_{\alpha\beta}}^{\alpha\beta > 4} \overbrace{F_{\mu\nu} F^{\mu\nu}}^{\mu, \nu = 1 \dots 4}$$

$$\text{AS } F_{\alpha\beta} = 0$$

THIS ARGUMENT WORKS TO
 ALL ORDERS IN MGUT / MSTRING
 BUT ACTUALLY THE RESULT IS
 EXACT, AS SHOWN BY A CONFORMAL
 FIELD THEORY ARGUMENT.

(this does not depend on supersymmetry)

(Witten 1986)

ANOTHER DIFFERENCE : MAGNETIC MONOPOLES

4-DIM'L
 FIELD THEORY



$D^3 \times K$

HIGHER DIMENSIONS
 START ON $S^2 \times K$



FIND QUANTUM OF
 MAGNETIC CHARGE IS

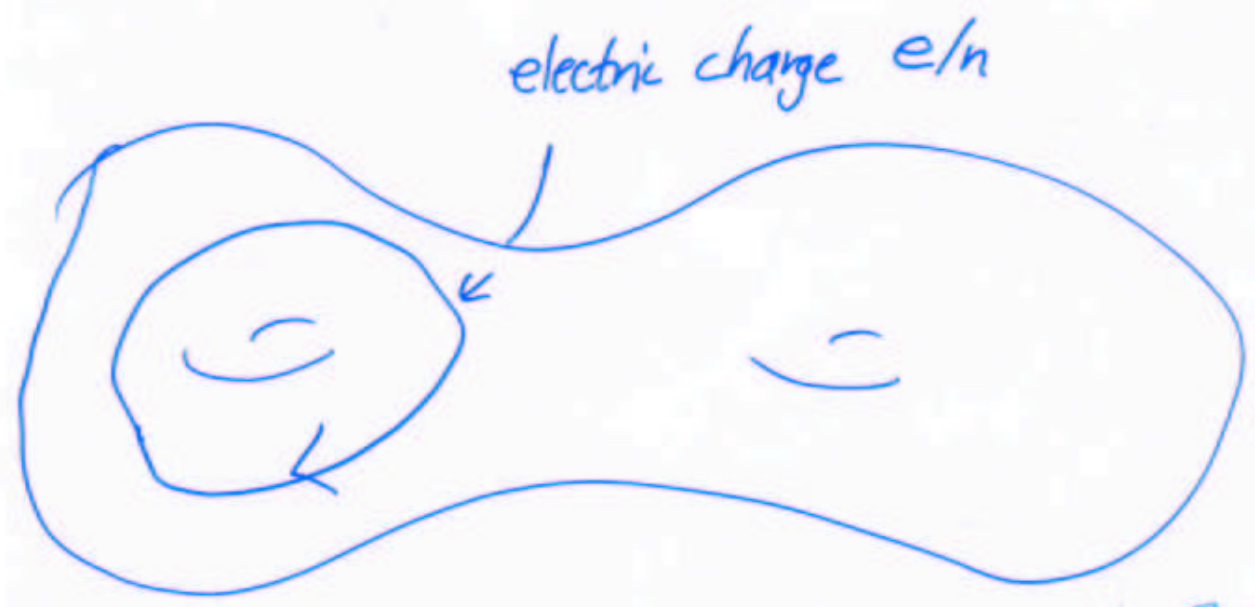
$$\frac{2\pi}{e} \cdot n$$

11/1/21

THIS RESULT WOULD BAFFLE A
4+k - DIMENSIONAL FIELD THEORIST -

WHERE ARE THE PARTICLES OF ELECTRIC
CHARGE ~~???~~ $\frac{e}{n}$?

IN HETEROTIC STRING THEORY THEY
EXIST: WRAPPED STRINGS



$$\pi_1(K) = \mathbb{Z} \times \mathbb{Z}$$

(11½)

NOTE:

THE USUAL CHARGE
QUANTIZATION (AND GUT-LIKE
STRUCTURE) DOES, OF
COURSE, HOLD FOR
PARTICLES OF ORDINARY
MASS

(12)

APART FROM THE HETEROTIC
STRING ON A CALABI-YAU, ANOTHER
WAY TO GET $N=1$ SUSY IN FOUR
DIMENSIONS AND A MODEL SOMEWHAT LIKE
THE STANDARD MODEL IS TO
COMPACTIFY M-THEORY FROM ELEVEN
TO FOUR DIMENSIONS ON A MANIFOLD OF
 G_2 HOLONOMY

$$\mathbb{R}^4 \times X$$

$X =$ A SEVEN-MANIFOLD OF G_2 HOLONOMY

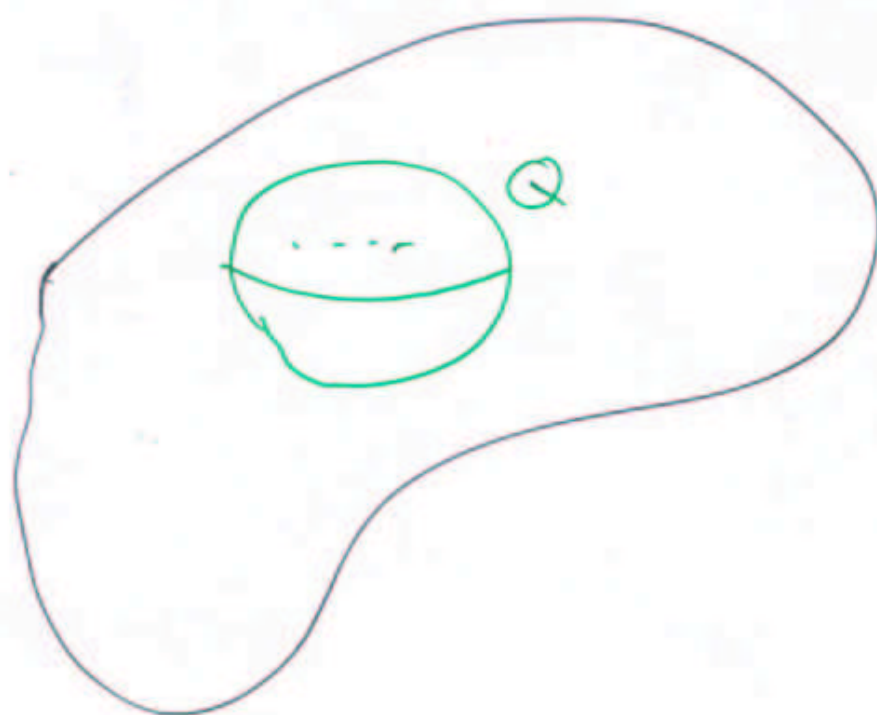
$G_2 =$ THE SMALLEST "EXCEPTIONAL" LIE
GROUP AND THE ONLY ONE THAT
CAN BE A HOLONOMY GROUP

(13)

THIS AVENUE WASN'T PURSUED 20 YEARS
AGO BECAUSE IN SUPERGRAVITY IT
DOESN'T WORK ... THAT IS, IF X
IS SMOOTH - NEEDED FOR SUPERGRAVITY
TO BE VALID - ONE GETS AN
ABELIAN GAUGE GROUP ONLY AND NO
CHIRAL FERMIONS.

BUT WITH THE MODERN UNDERSTANDING
OF HOW TO GENERATE GAUGE GROUPS
AND CHIRAL MATTER FROM SINGULARITIES
ONE CAN DO MUCH BETTER

(14)



X

X IS A SEVEN-MANIFOLD

OF G_2 HOLONOMY, AND Q IS A

THREE-MANIFOLD (THE "NORMAL"

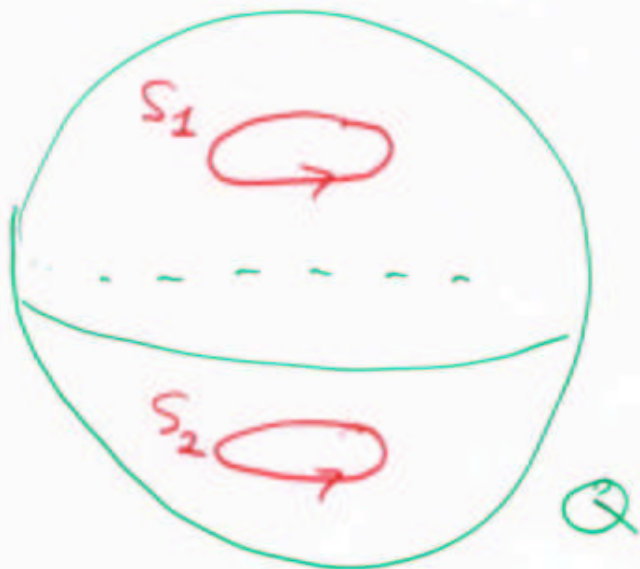
SPACE TO Q IS SINGULAR) ON

WHICH PROPAGATE $SU(5)$

(OR $SO(10)$, OR E_6) GAUGE FIELDS.

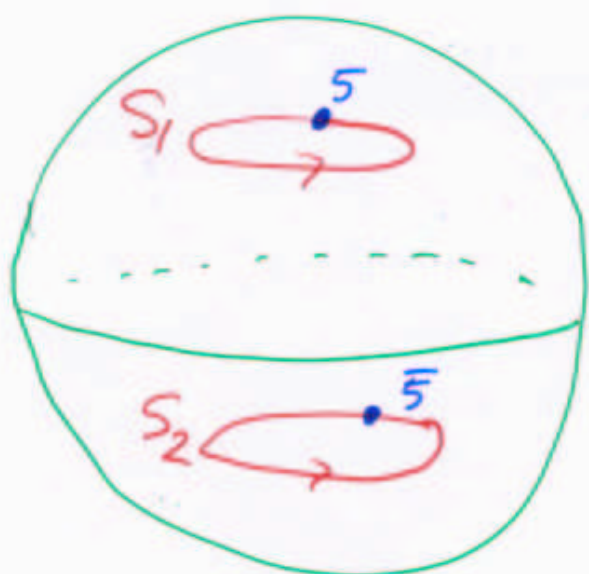
(15)

LET US JUST LOOK AT Q :



S_1 AND S_2 ARE TWO CIRCLES
(CONTAINING FURTHER SINGULARITIES)
ON WHICH CHIRAL SUPERFIELDS
ARE SUPPORTED.

(16)



IT TURNS OUT THAT IF WE
PLACE ONE HIGGS - THE 5 -
ON S_1 AND THE OTHER - THE $\bar{5}$ -
ON S_2 , THEN A CERTAIN DISCRETE
SYMMETRY (THAT ROTATES S_1 OR S_2)
TRANSFORMS THE HIGGSES
IN A WAY THAT CANNOT BE
ACHIEVED IN A FOUR-DIM'L SU(5) THEORY

(17)

NAMELY

$$5: \begin{pmatrix} 3 \\ 2 \end{pmatrix} \rightarrow e^{i\phi} \begin{pmatrix} e^{2i\alpha} & & & & \\ & e^{2i\alpha} & & & \\ & & e^{2i\alpha} & & \\ & & & e^{-3i\alpha} & \\ & & & & e^{-3i\alpha} \end{pmatrix}$$

AS EXPECTED

BUT (FOR EXAMPLE)

$$\bar{5}: \begin{pmatrix} \bar{3} \\ \bar{2} \end{pmatrix} \rightarrow e^{i\bar{\phi}} \begin{pmatrix} \bar{3} \\ \bar{2} \end{pmatrix}$$

i.e. JUST A GLOBAL SYMMETRY ON $\bar{5}$

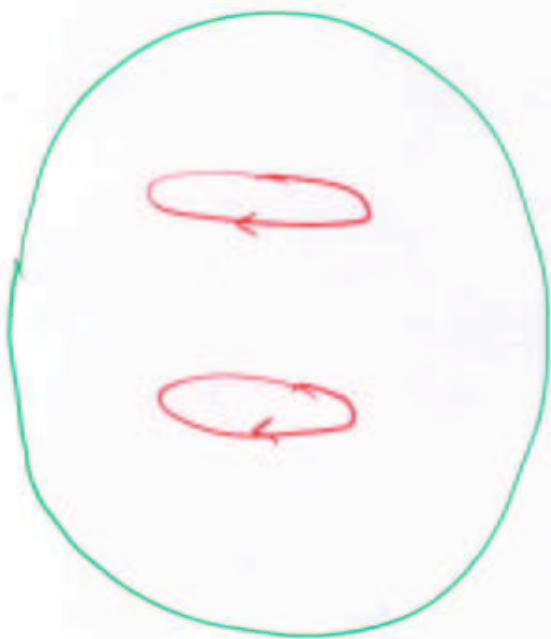
BUT A GLOBAL SYMMETRY PLUS A

GAUGE TRANSFORMATION ON THE 5

CLEARLY, WE CAN PICK $\phi, \bar{\phi}, \alpha$ TO
 ALLOW A $3\text{-}\bar{3}$ MASS TERM AND FORBID ONE
 FOR $2\text{-}\bar{2}$

I WON'T EXPLAIN THE DETAILS OF (18)
THE CONSTRUCTION, SINCE IT TURNS OUT
THAT THIS ASPECT CAN BE IMITATED
(via "deconstruction" — ARKANI-HAMED, COHEN, GEORGE
& HILL, POKORSKI, WANG)

BY A FOUR-DIMENSIONAL THEORY
IN WHICH THE GAUGE GROUP
IS NOT $SU(5)$ BUT $SU(5) \times SU(5)$



IN AN EXTREME VERSION
OF LATTICE GAUGE
THEORY, EACH CIRCLE
IS REPLACED BY ONE
POINT AND THE
REST IS IGNORED

NOW WE BREAK $SU(5) \times SU(5)$

TO $SU(3) \times SU(2) \times U(1) \times \Gamma$
underbrace
embedded diagonally in $SU(5) \times SU(5)$

WHERE Γ IS A FINITE SUBGROUP
OF THE HYPERCHARGE GROUP OF THE
SECOND $SU(5)$ TIMES A GLOBAL SYMMETRY

IF HIGGS BOSONS ARE NOT $\bar{5} \oplus 5$ OF $SU(5)$,

BUT

$$(\bar{5}, 1) \oplus (1, 5) \text{ OF } SU(5) \times SU(5)$$

THEN Γ IS THE SORT OF DISCRETE
SYMMETRY WE WANT - THAT CAN ALLOW
 $3 \cdot \bar{3}$ AND FORBID $2 \cdot \bar{2}$

(20)

IN THIS WAY WE MAKE A FOUR-DIMENSIONAL MODEL THAT REPRODUCES SOME RESULTS OF THE M-THEORY MODEL, BUT NOT OTHERS.

THE MOST STRIKING DIFFERENCE IS PERHAPS IN ELECTRIC CHARGE CONSERVATION - THE DECONSTRUCTED, FOUR-DIMENSIONAL MODEL HAS CONVENTIONAL QUANTIZATION OF ELECTRIC CHARGE, AND THE M-THEORY MODEL DOES NOT.

TO ENSURE THAT ALL QUARKS AND
LEPTONS CAN GET MASSES, IT IS MOST
STRAIGHTFORWARD TO TAKE THEM TO BE
THREE COPIES OF

$$(\bar{5}, 1) \oplus (10, 1)$$

SO BELOW THE GUT SCALE, THE
SPECTRUM IS SO FAR

$$\underbrace{(\bar{5}, 1) \oplus (1, \bar{5})}_{\text{HIGGS}} \oplus \underbrace{3 (\bar{5}, 1) \oplus 3 (10, 1)}_{\text{QUARKS AND LEPTONS}}$$

IF WE TOOK QUARKS AND LEPTONS TO
BE, SAY, $(\bar{5}, 1) \oplus (1, 10)$, WE'D GET
PECULIAR QUANTUM NUMBERS UNDER Γ_5
AND WE'D HAVE TROUBLE GIVING MASS TO ALL
QUARKS & LEPTONS

(22)

THERE IS A LOT OF FREEDOM

IN PRECISELY SPECIFYING THE GLOBAL SYMMETRY GROUP - COMING FROM THE CHOICE OF THE HYPERCHARGE SUBGROUP Γ AND HOW IT MIXES WITH GLOBAL SYMMETRIES OF THE DIFFERENT MULTIPLETS.

$\overline{55} H H$

WE CAN PICK THE DISCRETE SYMMETRY TO

- ALLOW $3 \cdot \overline{3}$, FORBID $2 \cdot \overline{2}$
- ALLOW QUARK, LEPTON, AND NEUTRINO MASSES
- FORBID PROTON DECAY BY OPERATORS OF DIMENSION FOUR OR FIVE

BUT OF COURSE WE DON'T HAVE A PURELY THEORETICAL REASON FOR THESE CHOICES.

NOW LET'S RETURN TO THE SPECTRUM, ⁽²³⁾
WHICH SO FAR WAS (BELOW GUT SCALE)

$$\underbrace{(\bar{5}, 1) \oplus (1, 5)}_{\text{HIGGS}} \oplus 3 \underbrace{(\bar{5}, 1) \oplus (10, 1)}_{\text{QUARKS, LEPTONS}}$$

THIS SPECTRUM IS ANOMALOUS; TO
CANCEL ANOMALIES, WE MIGHT ADD

$$(5, 1) \oplus (1, \bar{5}) \quad \text{OR} \quad (10, 1) \oplus (1, \bar{10})$$

HOW DO WE INTERPRET THESE FIELDS?

IT IS TEMPTING TO CONSIDER THEM AS
"MESSENGER FIELDS" OF GAUGE-MEDIATED
SUPERSYMMETRY BREAKING.

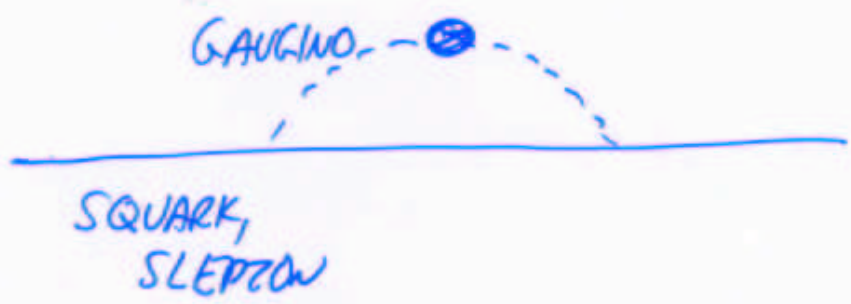
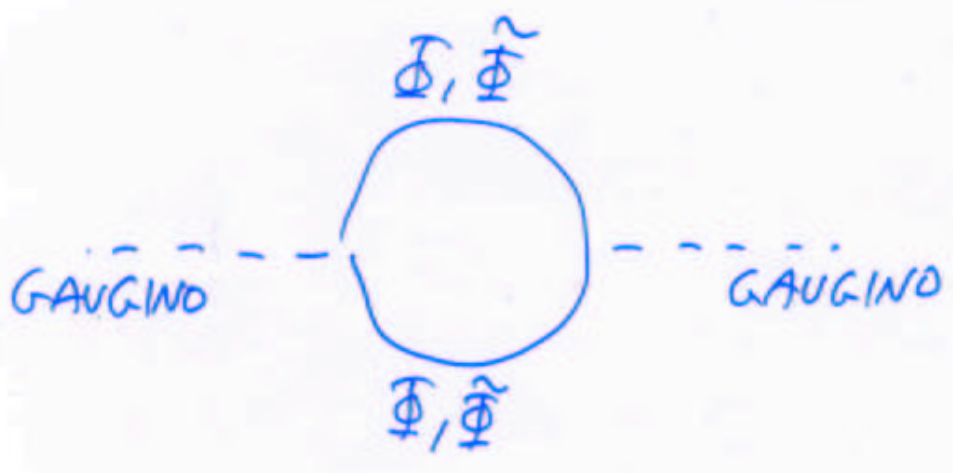
LET US CALL THESE MESSENGERS 24

$$\Phi, \tilde{\Phi}$$
$$(5, 1) \quad (1, \bar{5})$$

IN GAUGE-MEDIATED SUSY BREAKING
~~TRANSFORMS~~ ONE USUALLY COUPLES
THE MESSENGERS TO A CHIRAL
SUPERFIELD S OF THE
SUPERSYMMETRY BREAKING SECTOR

$$\int d^2\theta S \Phi \tilde{\Phi}$$

THEN CERTAIN ONE AND TWO-LOOP
DIAGRAMS GIVE MASSES TO
GAUGINOS AND SQUARKS, SLEPTONS



VIRTUE : NATURAL SQUARK, SLEPTON
DEGENERACY \Rightarrow AVOID FLAVOR
CHANGING PROCESSES

(26)

AT LOW ENERGIES, UNDER
THE STANDARD MODEL,

$$\Phi = 2 \oplus 3, \quad \tilde{\Phi} = \bar{2} \oplus \bar{3}$$

AND WE MIGHT INTRODUCE
SEPARATE S-FIELDS FOR THE
 $2 \cdot \bar{2}$ AND $3 \cdot \bar{3}$:

$$\int d^2\theta (S 2 \cdot \bar{2} + S' 3 \cdot \bar{3})$$

IN AN $SU(5)$ MODEL, ONE COULD

ARRIVE AT THIS STRUCTURE

BY COUPLING $\Phi, \tilde{\Phi}$ TO BOTH

A SINGLET AND ADJUNT OF $SU(5)$

→ COMPONENTS BEING S, S'

IN THE PRESENT CONTEXT,
 WE HAVE A MORE PRECISE
 RATIONALE FOR THIS. JUST AS
 FOR THE HIGGS BOSONS, SO FOR THE
 MESSENGERS $2 \cdot \bar{2}$ TRANSFORMS
 DIFFERENTLY FROM $3 \cdot \bar{3}$ UNDER THE
 GLOBAL SYMMETRY. SO S, S'
 MUST BE DIFFERENT.

MINIMAL
 IN ~~MINIMAL~~ GAUGE MEDIATION WITH
 ONLY ONE S FIELD, THE
 RATIO $\frac{\langle F_S \rangle}{\langle S \rangle}$ DETERMINES MASSES
 IN THE OBSERVED SECTOR
 AND ONE GETS UNIQUE PREDICTIONS FOR
 THE SPECTRUM.

(28)

HERE, WE HAVE THE TWO

PARAMETERS

$$\Lambda_2 = \frac{\langle F_S \rangle}{\langle S \rangle} \quad \text{AND} \quad \Lambda_3 = \frac{\langle F_{S'} \rangle}{\langle S' \rangle}$$

SO THE SPECTRUM ISN'T UNIQUE:

IT DEPENDS ON THEIR DIMENSIONLESS

RATIO $r = \frac{\Lambda_2}{\Lambda_3}$

Λ_2 CONTRIBUTES TO MASSES OF

SU(2) NONSINGLETS AND Λ_3 TO

THOSE OF SU(3) NONSINGLETS

SO, FOR EXAMPLE, BY INCREASING (29)
 r ONE ~~INCREASES~~ INCREASES THE MASS
OF LEFT-HANDED SLEPTONS AND
CHARGINOS RELATIVE TO RIGHT-HANDED
SLEPTONS.

FOR EXAMPLE, IN RUN 1 AT FERMI
THERE WAS A MUCH-DISCUSSED

$e^+e^- \rightarrow \gamma\gamma + \text{missing } E_T$ EVENT THAT
MIGHT BE A SIGNAL OF GAUGE
MEDIATED SUSY BREAKING, BUT NOT
IF ONE ASSUMES THE SPECTRUM OF
THE MINIMAL MODEL.

HERE WE HAVE A RATIONALE FOR
DEPARTING FROM THE SPECTRUM OF
THE MINIMAL MODEL ... IT IS

POSSIBLE TO GET A BETTER "FIT"

TO THE $e^+e^- \rightarrow \gamma\gamma + \cancel{E}_T$ EVENT

(AND ABSENCE OF CERTAIN OTHER
EVENTS) THAN THE MINIMAL

MODEL (S. THOMAS AND E.W.)

IN THE PROCESS, ONE CAN ALSO MAKE
SQUARKS LIGHTER AND THE REQUIRED
 μ PARAMETER LESS ... ALLEVIATING THE
SUSY FINE-TUNING PROBLEM....

30½

ONE CAN ALSO, FOR EXAMPLE,
MAKE GLUONS MUCH LIGHTER
THAN OTHER PARTICLES,
SOMEWHAT AS IN YESTERDAY'S
"PARTIAL GAUGE MEDIATION"

(31)

DEPENDING ON WHETHER

ONE GETS e^+e^- OR \cancel{E}_T FROM

LEFT- OR RIGHT-HANDED SLEPTON

PRODUCTION, ONE WILL OR WILL

NOT ALSO PREDICT

$$e\nu \text{ OR } \cancel{E}_T \simeq e \text{ OR } \cancel{E}_T$$

WHEN ONE SLEPTON IS A

NEUTRINO

IF ONE INCLUDES R-PARITY VIOLATION,

ONE CAN ALSO (ALLANACH et al.) TRY

TO EXPLAIN THE $\mu \gamma \cancel{E}_T$ EXCESS

SEEN AT RUN 1 AT FERMILAB.

(ACCORDING TO ALLANACH'S TALK, ONE CAN'T QUITE
MAKE A GOOD FIT TO BOTH)

(32)

HOPEFULLY THE LUMINOSITY

AT FERMILAB WILL SOON INCREASE

AND WE'LL LEARN WHICH

HINTS FROM RUN 1 ARE REAL.