

Lepton Flavor Violating Process in Bi-large Texture of Neutrino Mixing

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- A.Kageyama, S.K., N.Shimoyama, M.Tanimoto
Phys.Lett **B527**:206-214,2002
 - A.Kageyama, S.K., N.Shimoyama, M.Tanimoto
Phys.Rev. **D65**:096010,2002

I. Motivations

• Neutrino Oscillations [Super-K, SNO]

⇒ Bi-large neutrino mixing

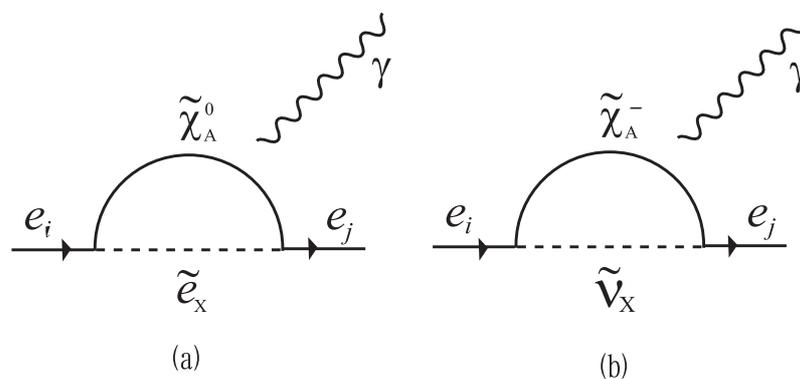
- $\sin^2 2\theta_{\text{sol}} = 0.6 - 0.85$ (best fit : $\theta_{\text{sol}} = 30^\circ$)
- $\sin^2 2\theta_{\text{atm}} \geq 0.92$ (best fit : $\theta_{\text{atm}} = 45^\circ$)
- $\Delta m_{\text{sol}}^2 = 5 \times 10^{-5} \text{eV}^2$ (best fit)
- $\Delta m_{\text{atm}}^2 = 2.5 \times 10^{-3} \text{eV}^2$ (best fit)

⇒ Small neutrino mass

- **See-saw mechanism** (introducing ν_{R})
- Small Lepton Flavor Violation in SM + ν_{R}
 $\text{BR}(\mu \rightarrow e\gamma) \sim 10^{-50}$

• Supersymmetry

- Solving Gauge Hierarchy Problem
- Gauge Coupling Unification
- **Large Lepton Flavor Violation (→ Sec.II)**
[F.Borzumati and A.Masiero]



MSSM + ν_R (see-saw)

In this framework, **LFV is related to neutrino mass spectrum and neutrino mixing matrix elements.**

Spectrums:

1. Hierarchical

- J.Hisano, T.Moroi, K.Tobe and M.Yamaguchi, Phys.Rev.D53(1996)2442
- J.Sato, K.Tobe and T.Yanagida, Phys.Lett.B498(2001)498

...

2. Degenerate

3. Inverse-hierarchical

→ **Sec.III**

Neutrino mixing matrix:

$$U_{\text{MNS}} \sim \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & U_{e3} < 0.2 \\ -\frac{1}{2} & \frac{1}{2} & \frac{1}{\sqrt{2}} \\ \frac{1}{2} & -\frac{1}{2} & \frac{1}{\sqrt{2}} \end{pmatrix} \quad (1)$$

We can compare between the theoretical predictions and experimental bound for $\mu \rightarrow e\gamma$ and $\tau \rightarrow \mu\gamma$:

$$\text{BR}(\mu \rightarrow e\gamma) < 1.2 \times 10^{-11} \quad (\text{MEGA;1999})$$

$$\text{BR}(\tau \rightarrow \mu\gamma) < 1.1 \times 10^{-6} \quad (\text{CLEO;2000})$$

for three types of neutrino mass spectrum.

II. LFV in MSSM + ν_R

[F.Borzumati and A.Masiero(1986)]

[J.Hisano,T.Moroi,K.Tobe and M.Yamaguchi(1996)]

Superpotential

$$W_{\text{lepton}} = \mathbf{Y}_e L H_d e_R^c + \mathbf{Y}_\nu L H_u \nu_R^c + \frac{1}{2} \nu_R^{cT} \mathbf{M}_R \nu_R^c$$

(\mathbf{Y}_e : diagonal, \mathbf{M}_R : diagonal , \mathbf{Y}_ν : non-diagonal) (2)

Soft SUSY breaking terms

$$\begin{aligned} -\mathcal{L}_{\text{soft}} = & (\mathbf{m}_{\tilde{L}}^2)_{ij} \tilde{L}_i^\dagger \tilde{L}_j + (\mathbf{m}_{\tilde{e}}^2)_{ij} \tilde{e}_{Ri}^* \tilde{e}_{Rj} + (\mathbf{m}_{\tilde{\nu}}^2)_{ij} \tilde{\nu}_{Ri}^* \tilde{\nu}_{Rj} \\ & + \mathbf{A}_{ij}^e H_d \tilde{e}_{Ri}^* \tilde{L}_j + \mathbf{A}_{ij}^\nu H_u \tilde{\nu}_{Ri}^* \tilde{L}_j + \dots \end{aligned} \quad (3)$$

We take universality condition:

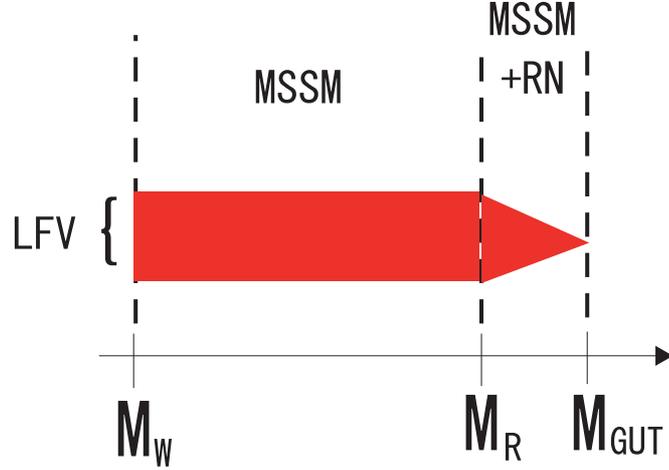
$$\begin{aligned} (\mathbf{m}_{\tilde{L}}^2)_{ij} &= (\mathbf{m}_{\tilde{e}}^2)_{ij} = (\mathbf{m}_{\tilde{\nu}}^2)_{ij} = \dots = \delta_{ij} m_0^2, \\ \mathbf{A}^\nu &= \mathbf{Y}_\nu a_0 m_0, \\ \mathbf{A}^e &= \mathbf{Y}_e a_0 m_0, \quad (a_0 = 0) \end{aligned} \quad (4)$$

at GUT scale ($\sim 10^{16}$ GeV).

There are **no flavor violations** at GUT scale.

This assumptions (4) are the **most conservative** one for lepton flavor violations.

However, lepton flavor violation are induced through **Renormalization Group Equations effect**. This is due to generation **off-diagonal elements** of $(\mathbf{m}_{\tilde{L}}^2)_{ij}$.



$$\mu \frac{d}{d\mu} (\mathbf{m}_{\tilde{L}}^2)_{ij} = \mu \frac{d}{d\mu} (\mathbf{m}_{\tilde{L}}^2)_{ij} \Big|_{\text{MSSM}} + \frac{1}{16\pi^2} [(\mathbf{m}_{\tilde{L}}^2 \mathbf{Y}_{\nu}^{\dagger} \mathbf{Y}_{\nu} + \mathbf{Y}_{\nu}^{\dagger} \mathbf{Y}_{\nu} \mathbf{m}_{\tilde{L}}^2)_{ij} + 2(\mathbf{Y}_{\nu}^{\dagger} \mathbf{m}_{\tilde{\nu}} \mathbf{Y}_{\nu} + \tilde{m}_{H_u}^2 \mathbf{Y}_{\nu}^{\dagger} \mathbf{Y}_{\nu} + \mathbf{A}_{\nu}^{\dagger} \mathbf{A}_{\nu})_{ij}] \quad (5)$$

Source of LFV comes from second and third lines including neutrino Yukawa coupling matrix \mathbf{Y}_{ν} .

$$(\Delta \mathbf{m}_{\tilde{L}}^2)_{ij} \simeq -\frac{(6 + 2a_0^2)m_0^2}{16\pi^2} (\mathbf{Y}_{\nu}^{\dagger} \mathbf{Y}_{\nu})_{ij} \ln \frac{M_X}{M_R}, \quad (6)$$

$$BR(\mu \rightarrow e\gamma) \propto |(\Delta \mathbf{m}_{\tilde{L}}^2)_{21}|^2 \tan^2 \beta, \quad (7)$$

$$BR(\tau \rightarrow \mu\gamma) \propto |(\Delta \mathbf{m}_{\tilde{L}}^2)_{32}|^2 \tan^2 \beta, \quad (8)$$

\mathbf{Y}_{ν} depends on $m_1, m_2, m_3, M_R, U_{MNS}$ (next section). **Therefore, discussing neutrino masses, mixings and M_R is crucial for estimating LFV.**

III. LFV in three types of Neutrino Spectrums

Neutrino Mass Spectrum

1. $m_1 \ll m_2 \ll m_3$ **Hierarchical**

$$m_1 \sim 0, \quad m_2 = \sqrt{\Delta m_{\text{sol}}^2}, \quad m_3 = \sqrt{\Delta m_{\text{atm}}^2}$$

2. $m_1 \sim m_2 \sim m_3$ **Degenerate**

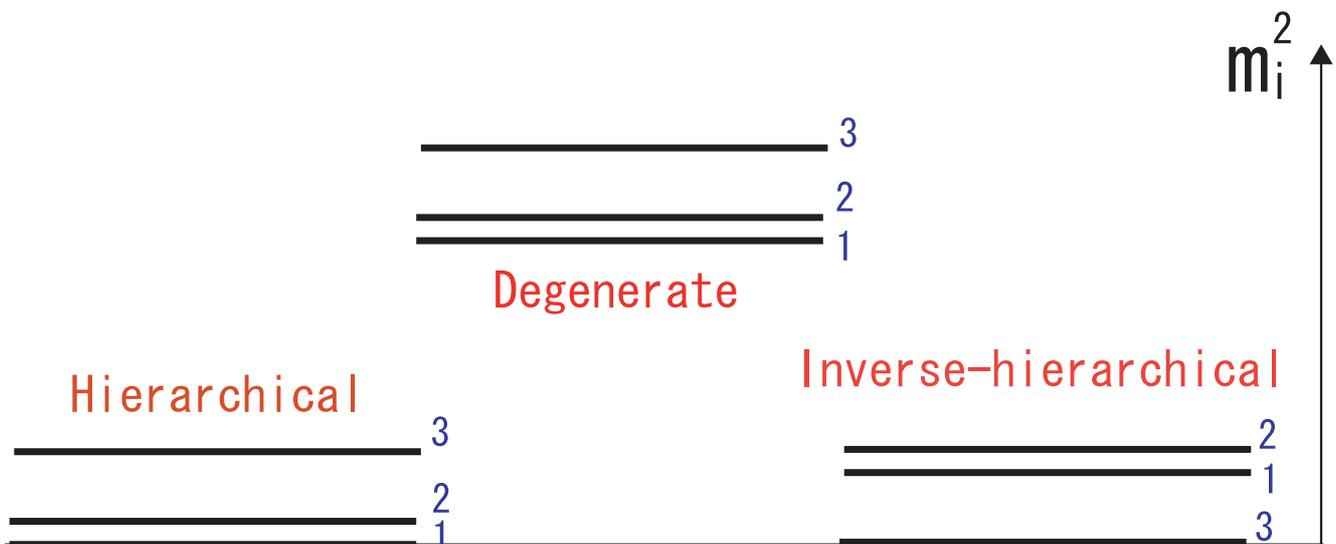
$$m_1 \equiv m_\nu, \quad m_2 = m_\nu + \frac{1}{2m_\nu} \Delta m_{\text{sol}}^2, \quad m_3 = m_\nu + \frac{1}{2m_\nu} \Delta m_{\text{atm}}^2$$

3. $m_1 \sim m_2 \gg m_3$ **Inverse-hierarchical**

$$m_1 = m_2 - \frac{1}{2m_2} \Delta m_{\text{sol}}^2, \quad m_2 \equiv \sqrt{\Delta m_{\text{atm}}^2}, \quad m_3 \simeq 0$$

Taking **Large Mixing Angle MSW** solution :

$$\Delta m_{\odot}^2 = 7 \times 10^{-5} \text{eV}^2, \quad \Delta m_{\text{atm}}^2 = 3 \times 10^{-3} \text{eV}^2$$



Neutrino Yukawa coupling

$$\mathbf{Y}_\nu = \frac{1}{v_u} \sqrt{\mathbf{M}_R^{\text{diag}}} \mathbf{R} \begin{pmatrix} \sqrt{m_1} & 0 & 0 \\ 0 & \sqrt{m_2} & 0 \\ 0 & 0 & \sqrt{m_3} \end{pmatrix} \mathbf{U}_{MNS}^T \quad (9)$$

\mathbf{R} is an orthogonal matrix ($\mathbf{R}\mathbf{R}^T = \mathbf{1}$), which is model dependent.

- [1] $M_{R1} = M_{R2} = M_{R3} = M_R \rightarrow$ **Simple formula**
- [2] $M_{R1} \neq M_{R2} \neq M_{R3} \rightarrow$ **Sec. V** (Discussion)

- ♣ **No CP violating phase**
- ♣ $U_{e3} \leq 0.2$ (**CHOOZ**)

In the case of [1],

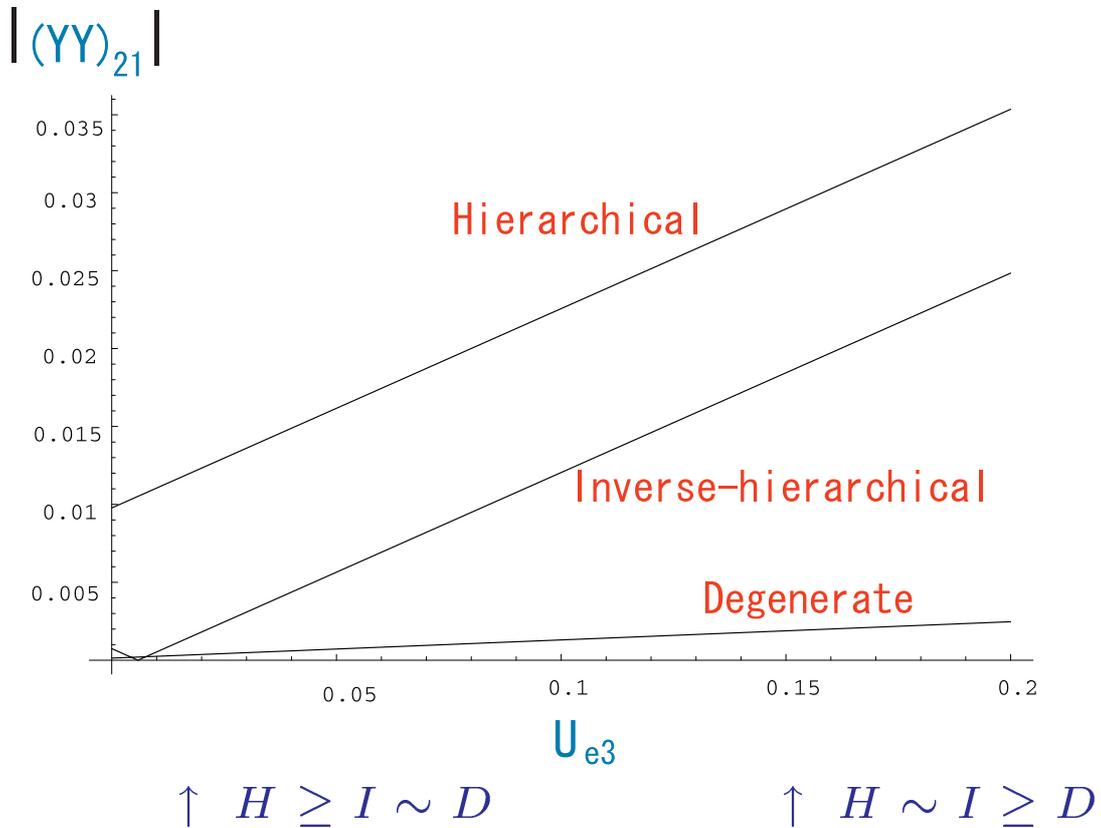
$$(\mathbf{Y}_\nu^\dagger \mathbf{Y}_\nu)_{21} = \frac{M_R}{v_u^2} [U_{\mu 2} U_{e 2}^* (m_2 - m_1) + U_{\mu 3} U_{e 3}^* (m_3 - m_1)]$$

$$(\mathbf{Y}_\nu^\dagger \mathbf{Y}_\nu)_{32} = \frac{M_R}{v_u^2} [U_{\tau 2} U_{\mu 2}^* (m_2 - m_1) + U_{\tau 3} U_{\mu 3}^* (m_3 - m_1)]$$

↑ **These are R independent !**

$\mathbf{Y}_\nu^\dagger \mathbf{Y}_\nu$ is expressed in terms of neutrino oscillation parameters for three types of neutrino spectrums:

$$\begin{aligned}
 (\mathbf{Y}_\nu^\dagger \mathbf{Y}_\nu)_{21} &\simeq \frac{M_R}{\sqrt{2}v_u^2} \sqrt{\Delta m_{\text{atm}}^2} \left[\frac{1}{\sqrt{2}} U_{e2}^* \sqrt{\frac{\Delta m_\odot^2}{\Delta m_{\text{atm}}^2}} + U_{e3}^* \right] \quad \text{Hier.} \\
 &\simeq \frac{M_R}{\sqrt{2}v_u^2} \frac{\Delta m_{\text{atm}}^2}{2m_\nu} \left[\frac{1}{\sqrt{2}} U_{e2}^* \frac{\Delta m_\odot^2}{\Delta m_{\text{atm}}^2} + U_{e3}^* \right] \quad \text{Deg.} \\
 &\quad (m_\nu = 0.3 \text{ eV}) \\
 &\simeq \frac{M_R}{\sqrt{2}v_u^2} \sqrt{\Delta m_{\text{atm}}^2} \left[\frac{1}{2\sqrt{2}} U_{e2}^* \frac{\Delta m_\odot^2}{\Delta m_{\text{atm}}^2} - U_{e3}^* \right] \quad \text{Inv.}
 \end{aligned}$$



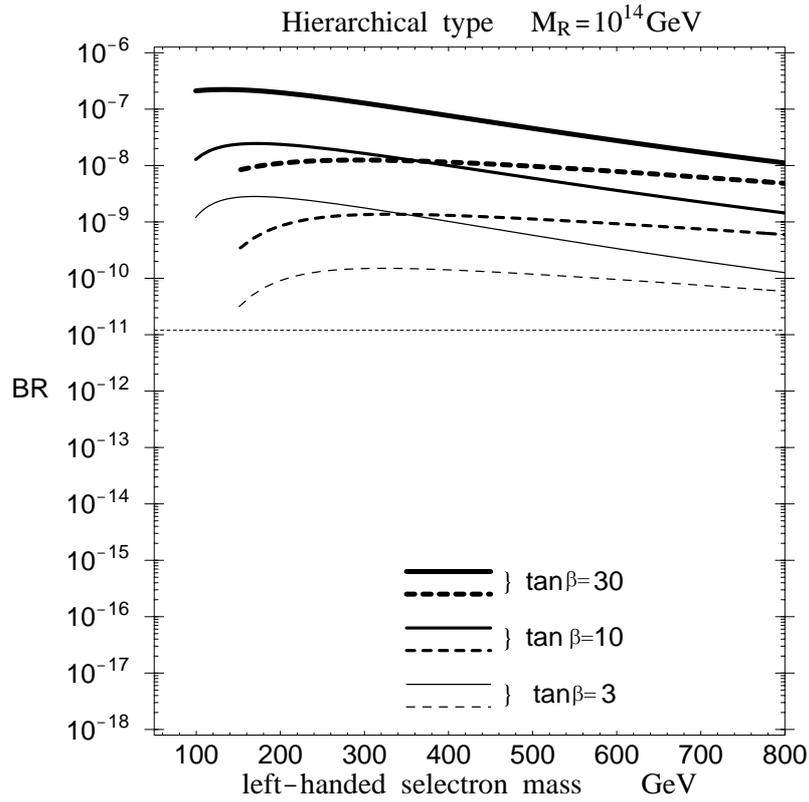


Figure 1: **Hierarchical** $\text{BR}(\mu \rightarrow e\gamma)$; $M_R = 10^{14} \text{ GeV}$, $U_{e3} = 0.2$

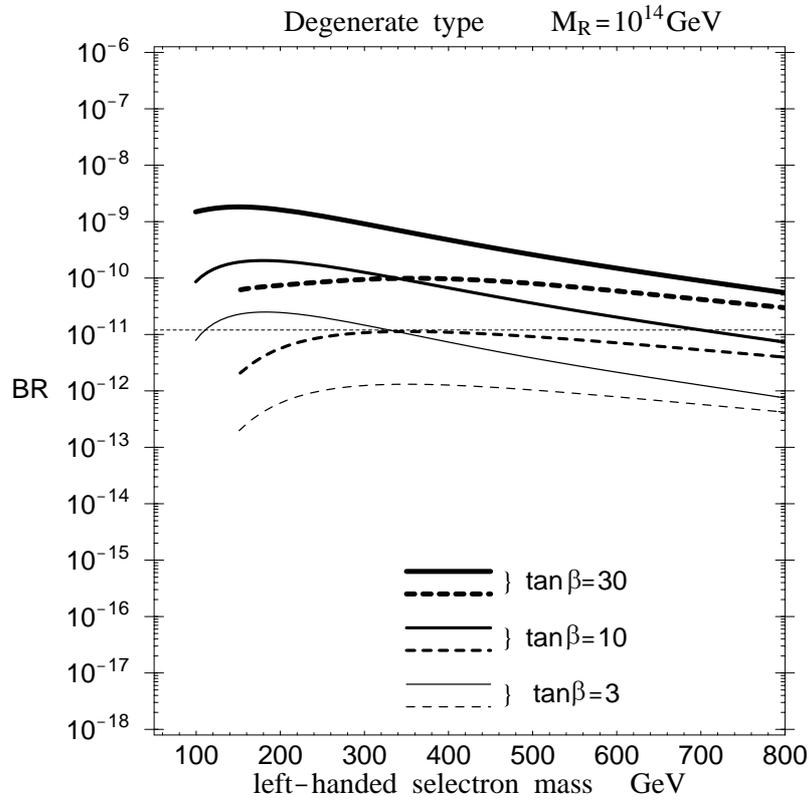


Figure 2: **Degenerate** $\text{BR}(\mu \rightarrow e\gamma)$; $M_R = 10^{14} \text{ GeV}$, $U_{e3} = 0.2$

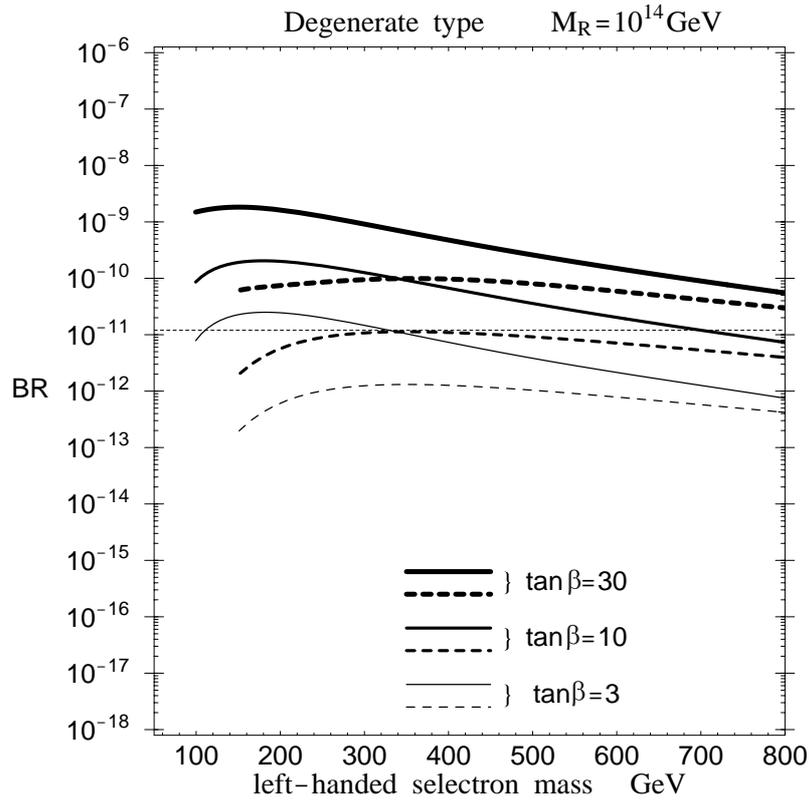


Figure 2: Degenerate $\text{BR}(\mu \rightarrow e\gamma)$; $M_R = 10^{14} \text{ GeV}, U_{e3} = 0.2$

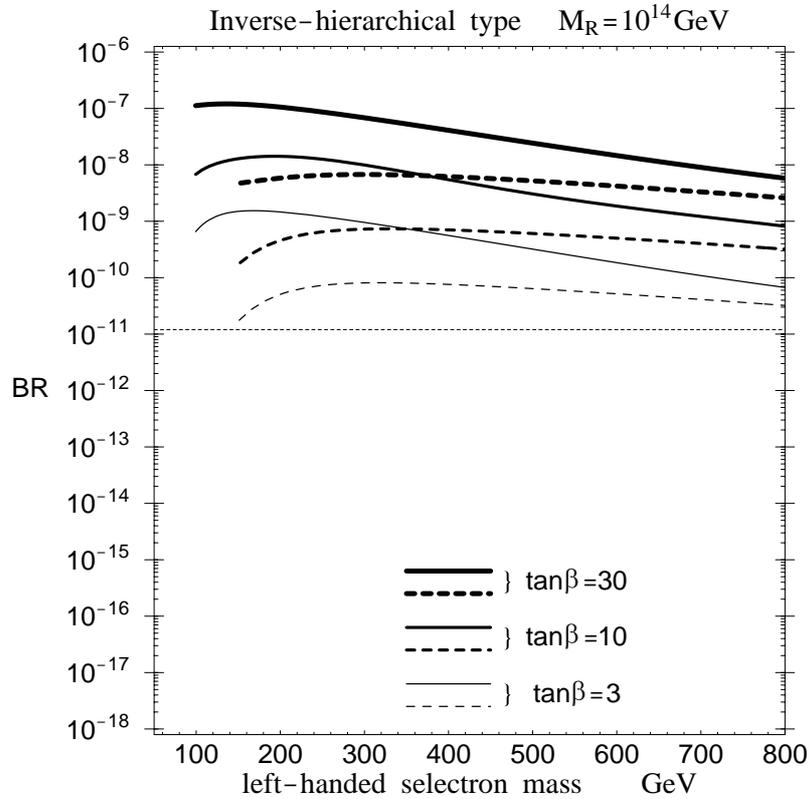


Figure 3: Inverse-hier. $\text{BR}(\mu \rightarrow e\gamma)$; $M_R = 10^{14} \text{ GeV}, U_{e3} = 0.2$

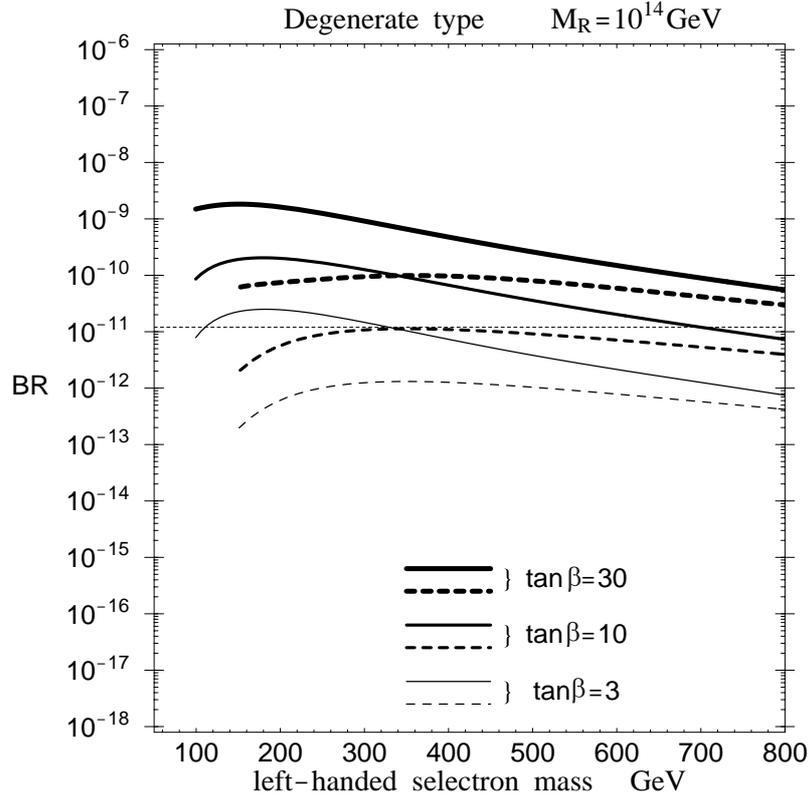


Figure 2: **Degenerate** $\text{BR}(\mu \rightarrow e\gamma)$; $M_R = 10^{14} \text{ GeV}$, $U_{e3} = 0.2$

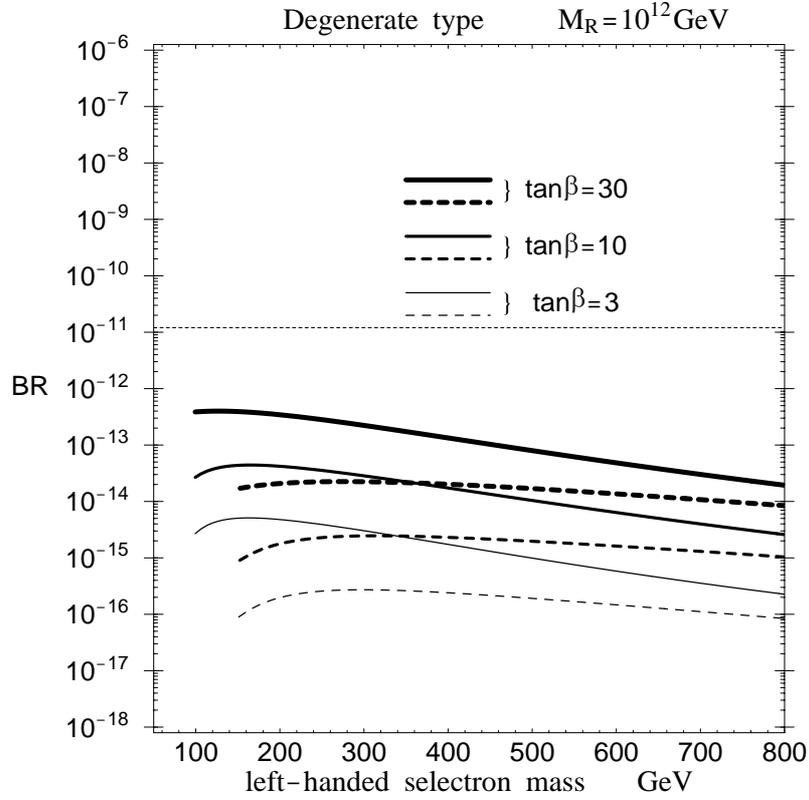


Figure 4: **Degenerate** $\text{BR}(\mu \rightarrow e\gamma)$; $M_R = 10^{12} \text{ GeV}$, $U_{e3} = 0.2$

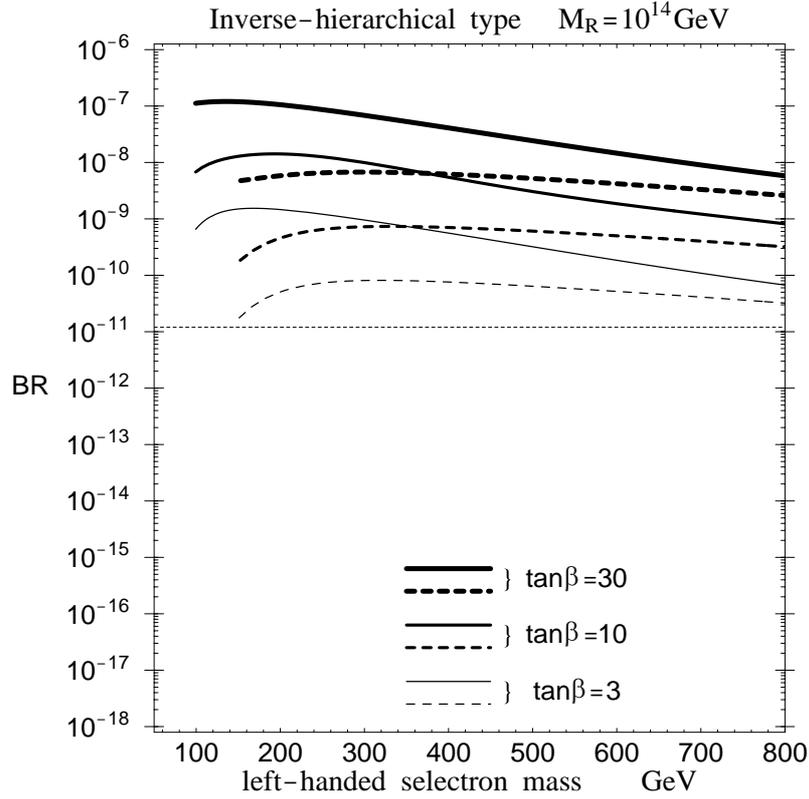


Figure 3: **Inverse-hier.** $\text{BR}(\mu \rightarrow e\gamma)$; $M_R = 10^{14} \text{ GeV}$, $U_{e3} = 0.2$

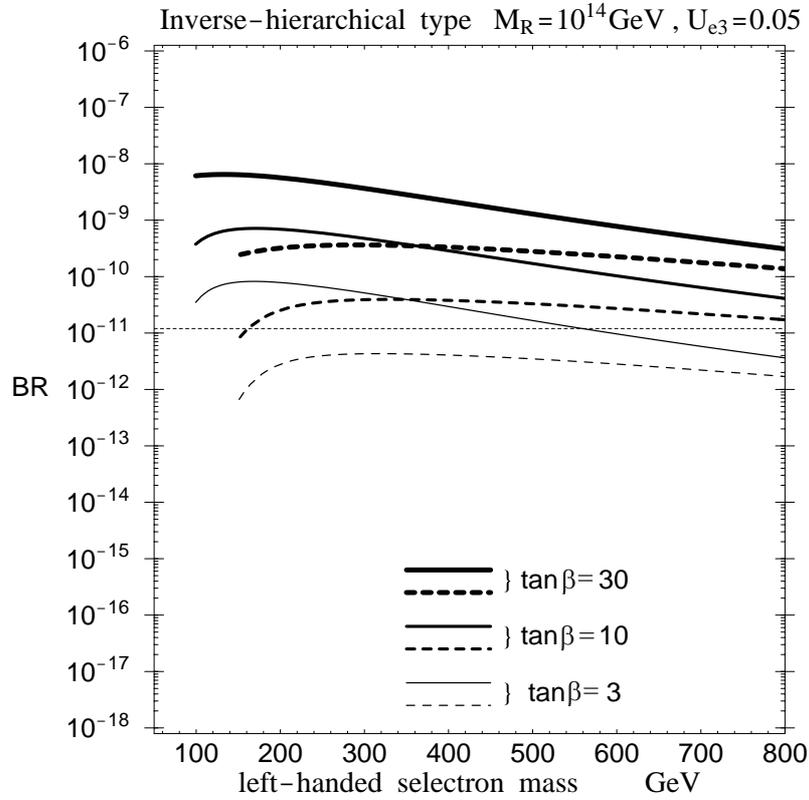


Figure 5: **Inverse-hier.** $\text{BR}(\mu \rightarrow e\gamma)$; $M_R = 10^{14} \text{ GeV}$, $U_{e3} = 0.05$

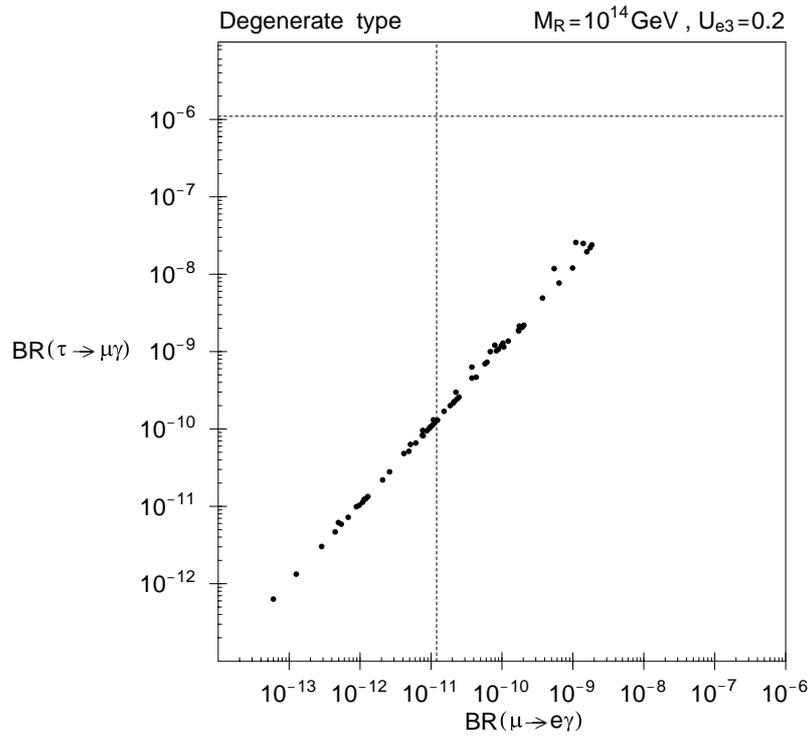


Figure 6: **Degenerate** $BR(\tau \rightarrow \mu\gamma)$ vs $BR(\mu \rightarrow e\gamma)$
 $M_R = 10^{14} \text{ GeV}, U_{e3} = 0.2$

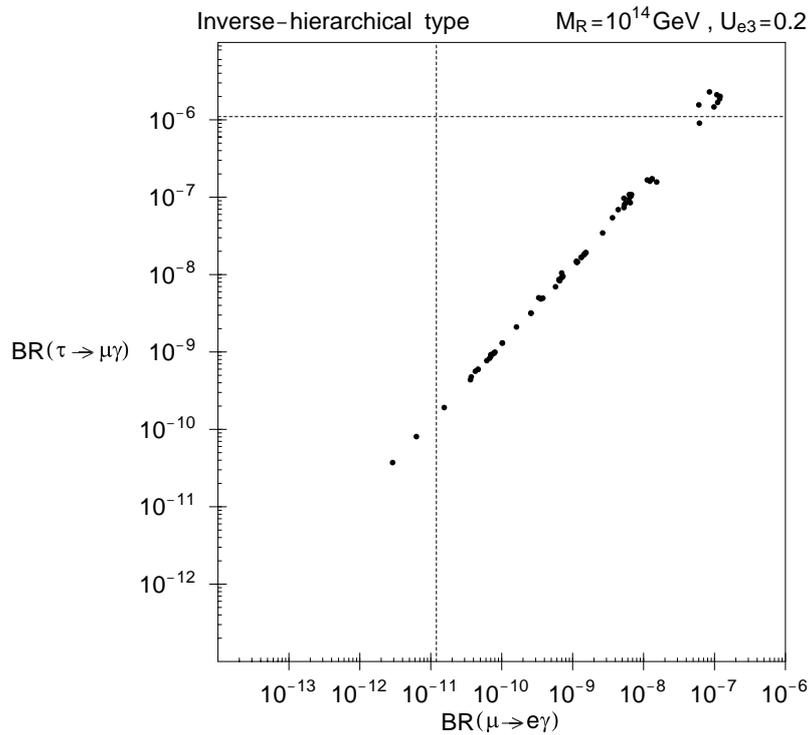


Figure 7: **Inverse-hier.** $BR(\tau \rightarrow \mu\gamma)$ vs $BR(\mu \rightarrow e\gamma)$
 $M_R = 10^{14} \text{ GeV}, U_{e3} = 0.2$

V. Summary and Discussion

We investigate lepton flavor violations ($\mu \rightarrow e\gamma$, $\tau \rightarrow \mu\gamma$) in **MSSM**+ ν_R in **bi-large** texture of neutrino mixing.

$$[1] M_{R1} = M_{R2} = M_{R3} = M_R$$

$$\text{BR}(\text{deg.}) < \text{BR}(\text{Inverse-hier.}) \sim \text{BR}(\text{Hier.}) \quad @ U_{e3} > 0.05$$

$\times 100$

$$[2] M_{R1} \neq M_{R2} \neq M_{R3}$$

$$\text{Degenerate} : \frac{\text{BR}_{M_{R1} \neq M_{R2} \neq M_{R3}}}{\text{BR}_{M_{R1} = M_{R2} = M_{R3} = M_R}} \leq 5.5$$

$$\text{Inverse - hierarchical} : \frac{\text{BR}_{M_{R1} \neq M_{R2} \neq M_{R3}}}{\text{BR}_{M_{R1} = M_{R2} = M_{R3} = M_R}} \leq 1.04$$

M_R , U_{e3} and m_ν are crucial quantities for estimating lepton flavor violations.

Future experiment

♣ $\text{BR}(\mu \rightarrow e\gamma) \rightarrow 10^{-14}$ @ PSI

♣ $\text{BR}(\tau \rightarrow \mu\gamma) \rightarrow 10^{-8}$ @KEK, SLAC

♣ Double beta decay : m_ν

♣ LBL experiment ($\nu_\mu \rightarrow \nu_e$) : U_{e3}

More precise test of LFV are possible in future.