# Search for Excited Leptons at HERA

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Searches for excited electrons and neutrinos have been performed using the complete HERA I and II data samples collected by the H1 detector at  $\sqrt{s} = 320$  GeV corresponding to an integrated luminosity of up to 435 pb<sup>-1</sup>. In absence of a signal, the limits on the ratio of the coupling to the compositeness scale derived extend the excluded region to higher masses than has been possible in previous searches.

#### 1 Introduction

Compositeness models [2] attempt to explain the hierarchy of masses in the Standard Model (SM) by the existence of a substructure within the fermions. Several of these models predict excited states of the known leptons. Excited leptons  $(F^*)$  are assumed to have the same electroweak SU(2) and U(1) gauge couplings, g and g', to the vector bosons, but are expected to be grouped into both left- and right-handed weak isodoublets with vector couplings. The existence of the right-handed doublets is required to protect the ordinary light leptons from radiatively acquiring a large anomalous magnetic moment via the  $F^*FV$  interaction (where V is a  $\gamma$ , Z, or W). Considering only the electroweak interaction, the phenomenological model describes this interaction by the Lagrangian density:

$$L_{F^*F} = \frac{1}{2\Lambda} \bar{F_R^*} \sigma^{\mu\nu} [gf\frac{\vec{\tau}}{2} \partial_{\mu} \vec{W_{\nu}} + g'f'\frac{Y}{2} \partial_{\mu} B_{\nu}]F_L + h.c.$$

where the new weights f and f' multiply the standard coupling constants g and g' corresponding to the weak SU(2) and the electromagnetic U(1) sectors respectively. The matrix  $\sigma^{\mu\nu}$  is the covariant bilinear tensor,  $\tau$  are the Pauli matrices,  $W_{\mu\nu}$  and  $B_{\mu\nu}$  represent the fully gauge invariant field tensors, and Y is the weak hypercharge. The parameter  $\Lambda$  has units of energy and can be regarded as the compositeness scale. The relative values of f and f' affect the size of the single-production cross section, their detection efficiencies and also the branching ratios of excited leptons.

Excited electrons and neutrinos may be produced in electron(positron)-proton collisions at HERA via *t*-channel  $\gamma(Z)$  or  $W^{\pm}$  gauge boson exchange. In the case of excited neutrinos, the cross section is much larger in  $e^-p$  collsions than in  $e^+p$  collsions due to the favourable valence u-quark and the helicity enhancement, specific to CC-like processes. Therefore the search for excited neutrinos uses only  $e^-p$  sample data with an integrated luminosity of 184 pb<sup>-1</sup>. In the case of excited electrons, both  $e^-p$  and  $e^+p$  collision modes are used, corresponding to total integrated luminosity of 435 pb<sup>-1</sup>.

### 2 Data analysis and results

Excited leptons  $(l = e, \nu)$  are searched for in the following decay channels:  $l^* \to l\gamma$ ,  $l^* \to lZ$ ,  $l^* \to lW$ . The final states resulting from the Z or W hadronic decays are taken into account

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Selection	Data	$\mathbf{SM}$	Efficiency $\times$ BR
$e^* \rightarrow \nu W_{\hookrightarrow qq}$	172	$175 \pm 39$	$\sim 40 \%$
$e^* \rightarrow eZ_{\rightarrow qq}$	351	$318 \pm 64$	$\sim 45 \%$
$e^* \rightarrow e\gamma$	112	$125 \pm 19$	60 - 70 %
$\nu^* \rightarrow \nu \gamma$	9	$15 \pm 4$	50 %
$\nu^* \rightarrow eW_{\rightarrow qq}$	198	$189 \pm 33$	30 - 40 %
$\nu^* \rightarrow \nu Z_{\hookrightarrow qq}$	111	$102 \pm 24$	40 %
$\nu^* \rightarrow eW_{\hookrightarrow \nu\mu}$	0	$0.54 \pm 0.04$	3-4.5~%
$\nu^* \rightarrow e W_{\hookrightarrow \nu e}$	0	$0.6 \pm 0.3$	4-6 %
$\nu^* \rightarrow \nu Z_{\hookrightarrow ee}$	0	$0.12\pm0.04$	2 %

Search for e<sup>\*</sup>,  $\nu^*$  HERA I+II ( $\sqrt{s} = 320$  GeV, preliminary)

Table 1: Observed and predicted event yields for the event classes of  $e^*$  and  $\nu^*$  searches. The selection efficiency for the signal multiplied by the branching ratio (BR) in each decay channel is also presented.

for both excited electrons and neutrinos and the Z or W leptonic decays are taken into account only for excited neutrinos. In the following, the selection criteria are described for the decay channels.

## The $\nu^* \rightarrow \nu \gamma$ channel

Candidate events are selected by requiring missing transverse momentum  $P_T^{miss} > 15$  GeV, where the photon is identified as an isolated electromagnetic (e.m.) cluster in the LAr calorimeter within a polar angle of 5° to 120°. The photon candidates measured within the acceptance of the central tracker ( $\theta^{\gamma} > 20^{\circ}$ ) are required to have no associated tracks. The neutral current (NC) and charged current (CC) backgrounds are reduced by imposing the longitudinal momentum balance  $E - P_Z > 45$  GeV for events with photon candidates at lower transverse momentum  $P_T^{\gamma} < 40$  GeV and by requiring the virtuality ( $Q_{\gamma}^2$ ) to satisfy  $\log(Q_{\gamma}^2) > 3.5$  GeV<sup>2</sup>. The background is further suppressed by rejecting events with a transverse momentum of the final hadronic in the calorimeter  $P_T^h < 5$  GeV.

### The $e^* \to e\gamma$ channel

Candidate events are selected with two isolated e.m. clusters in the LAr calorimeter of transverse energy greater than 20 GeV and 15 GeV, respectively, and with a polar angle between  $5^{\circ}$  and  $130^{\circ}$ . The sum of the energies of the two clusters has to be greater than 100 GeV. The background from NC is further suppressed by rejecting events with a total transverse energy of the two isolated e.m. clusters lower than 75 GeV.

## The $e^* \to eZ, \nu W$ and $\nu^* \to eW, \nu Z$ channels with $Z, W \to qq'$

These channels use subsample of events with at least two jets with high transverse momentum  $P_T^{j1(j2)} > 20(15)$  GeV reconstructed within  $5^\circ < \theta^{j1(j2)} < 130^\circ$ . The dijet invariant mass must be compatible with the relevant boson mass and should be closest to them.

Events with two high  $P_T$  jets and one electron: The channels  $e^* \to eZ_{\to qq}$  and  $\nu^* \to eW_{\to qq}$  are characterised by at least two high  $P_T$  jets and an energetic isolated e.m. cluster  $P_T^e > 10$  GeV ( $P_T^e > 20$  GeV for  $e^* \to eZ_{\to qq}$ ) in the polar angle  $5^\circ < \theta^e < 130^\circ$ . The

polar angle of the highest  $P_T$  jet resulting from W boson should be lower than 80°. The dijet invariant mass has to be greater than 60 GeV. If  $P_T^e < 65$  GeV, the dijet invariant mass must be greater than 75 GeV. In the case of the  $\nu^* \rightarrow eW_{\rightarrow qq}$  channel, to reduce the NC background, the polar angle of e.m. cluster must be lower than 90°. Furthermore, the background is reduced by requiring the virtuality computed from the e.m. cluster kinematics  $Q^2 > 2500 \text{ GeV}^2$  if  $P_T^e < 25 \text{ GeV}$  and by requiring a third jet with  $P_T > 5 \text{ GeV}$  to be present in the event if  $P_T^e > 50 \text{ GeV}$ .



Figure 1: The 95% C.L. limits obtained for coupling constants  $(f/\Lambda)$  as a function of the excited neutrino mass within two assumptions: f = -f' (left) and f = +f' (right). The observed limits from this annalysis using all H1  $e^-p$  data is presented by the yellow area. Values of the couplings above the curves are excluded. The orange-dark area corresponds to the exclusion domain published by the H1 experiment using 98/99 data and the dashed line to the exclusion limit from the L3 experiment at LEP [4].

Events with two high  $P_T$  jets and  $P_T^{miss}$ : The channels  $e^* \to \nu W_{\to qq}$  and  $\nu^* \to \nu Z_{\to qq}$ are characterised by at least two high jets and  $P_T^{miss} > 12$  GeV. In the case of  $e^* \to \nu W_{\to qq}$ channel, the ratio  $V_{ap}/V_p$  of transverse energy flow anti-parallel and parallel to the hadronic final state [3] is required to be lower than 0.3 to reject the photoproduction  $(\gamma p)$  background. The dijet invariant mass is required to be greater than 50 GeV. Furthermore, if  $P_T^{miss} < 65$  GeV the dijet is required to have an invariant mass above 65 GeV. In the case of the  $\nu^* \to \nu Z_{\to qq}$  channel, a dijet invariant mass greater than 60 GeV is required. In order to reduce CC background, the total hadronic system is required to have a polar angle above  $20^\circ$  and a third jet with  $P_T > 5$  GeV has to be present in the event if  $P_T^{miss} < 65$  GeV. The longitudinal balance  $E - P_Z > 25$  GeV is imposed if  $P_T^{miss} < 50$  GeV. In addition, events with  $P_T^{miss} < 30$  GeV are only accepted if the topological variable  $V_{ap}/V_p > 0.1$ .

The  $\nu^* \to eW, \nu Z$  channels with  $Z \to ee, W \to \nu e(\mu)$ 

Events with two electron and  $P_T^{miss}$ : These channels use a subsample of events with two high  $P_T$  isolated e.m. clusters  $P_T^{e1(e2)} > 20(15)$  GeV and a polar angle  $5^\circ < \theta^{e1(e1)} < 100^\circ(120^\circ)$ and  $P_T^{miss} > 12$  GeV. The clusters are required to have associated tracks if they are measured within the acceptance of the central tracker. Events with one muon and an electron: Candidate events containing an isolated muon plus an isolated electron, both having a high transverse momentum  $P_T^{e(\mu)} > 20(10)$  GeV and a polar angle  $5^{\circ} < \theta^{e(\mu)} < 100^{\circ}(160^{\circ})$  are selected. A cut  $P_T^{miss} > 12$  GeV is applied. The backgrounds are reduced by requiring the virtuality  $(Q_e^2)$  to satisfy  $\log(Q_e^2) > 3$  GeV<sup>2</sup>.

The observed number of events are compared to the expected SM background in table 1 for each search channel. A good overall agreement is found for all channels. No significant deviation is observed in the data. The selection efficiency for each decay channel for the both  $e^*$  and  $\nu^*$  search is presented also in the table 1.

#### **3** Interpretation and Conclusions



Figure 2: The 95% C.L. limits obtained for coupling constants  $(f/\Lambda)$  as a function of the excited electron mass within assumption: f = +f'. The observed limit from this analysis using 435  $pb^{-1}$  of H1 data is presented by the yellow area. The orangedark area corresponds to the exclusion domain published by the H1 experiment using HERA I data. The combined exclusion limit from LEP experiments is presented by the violet line. The result of the CDF [5] experiment at the Tevatron is also shown.

In absence of a signal for both excited electron and neutrino production, upper limits on the coupling  $f/\Lambda$  have been derived at 95% Confidence Level (C.L.) as a function of excited electron and neutrino masses. In case of excited neutrinos production, the obtained limits are displayed for the two assumptions f = -f' and f = +f' (figure 1). For f = -f'(maximal  $\gamma\nu*\nu$  coupling) and assuming  $f/\Lambda = 1/M_{\nu*}$ , excited neutrinos with masses below 211 GeV are excluded at 95% C.L. The limits on the ratio  $f/\Lambda$  also are given for the excited electron for the hypothesis f = +f' (figure 2). We do not consider the case f = -f', because the  $\gamma e^*e$  coupling constant would be equal to zero and the production cross section of the excited electron is very small. For this hypothesis and assuming  $f/\Lambda = 1/M_{e^*}$ , excited electrons with masses below 273 GeV are excluded at 95% C.L.

#### References

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