Search for ZH Production at D0 in $p\bar{p} \to \ell^+\ell^-b\bar{b}$ Events at $\sqrt{s}=1.96$ TeV

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We present a search for a low-mass standard model Higgs boson produced in association with a Z boson decaying to charged leptons at a center-of-mass energy of \sqrt{s} =1.96 TeV with the D0 detector at the Fermilab Tevatron collider. The search is performed in a large data set of events containing two opposite-sign leptons (electron, muon, tau) and one or two b-tagged jets. Recent improvements to the sensitivity, from increased lepton acceptance to optimized signal-to-background discrimination, will be discussed.

1 Introduction

The mass of a standard model Higgs boson is constrained by direct searches performed at LEP and measurements of the top quark and W boson masses [1]. Combining these results, the Higgs mass m_H must be less than 186 GeV at 95% confidence level.

If $m_H < 135$ GeV, then Higgs bosons are expected to decay primarily to $b\bar{b}$. At hadron colliders, the inclusive $b\bar{b}$ cross section is six orders of magnitude larger than the cross section for Higgs production, so it is not feasible to find evidence for low-mass Higgs bosons produced alone. Instead, we search for associated production of vector (W, Z) and Higgs bosons. Requiring leptonic decay of the W or Z dramatically reduces the multijet background and increases our sensitivity to the Higgs signal. The analysis discussed here is concerned exclusively with ZH production in the $\ell^+\ell^-b\bar{b}$ final state; the D0 collaboration has also completed similar analyses using $\ell\nu b\bar{b}$ [2] and $\nu\bar{\nu}b\bar{b}$ [3] final states.

The irreducible backgrounds in this search are Z production with heavy-flavor jets, top quark pair production, and diboson final states. Instrumental backgrounds include jets faking charged leptons and light jets faking heavy-flavor jets.

A more detailed description of this analysis may be found in [4].

2 The Tevatron and the D0 detector

The Tevatron is a proton-antiproton collider located at Fermilab near Chicago, IL. Collisions have a center-of-mass energy of $\sqrt{s}=1.96$ TeV. Fermilab's Accelerator Division continues to optimize the integrated luminosity produced by the Tevatron, and currently the accelerator is performing better than ever before. The total integrated luminosity delivered from RunII is over 7 fb⁻¹; results discussed here use up to 4.1 fb⁻¹ of data.

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D0 is a multi-purpose particle detector, one of two detectors located at collision points around the Tevatron. We have taken data with 90% average efficiency since the start of RunII. In this search, we employ every major component of D0 in order to identify muons, electrons, and heavy-flavor jets [5].

3 Event selection

3.1 Leptons

We strive for maximum Higgs signal acceptance, so our event selection is very loose. For muons, we require central track matches, $p_T > 10 \text{ GeV}$, and $|\eta| < 2$. Electron requirements are $p_T > 15 \text{ GeV}$ and $|\eta| < 2.5$. All identified leptons are also required to pass various tracking and calorimeter isolation criteria, and the invariant mass of the dilepton pair must match the Z boson resonance: $70 < m_{\ell\ell} < 130 \text{ GeV}$.

To be sure we accept as many Higgs events as possible, we select some electrons and muons that are not initially identified as such. In the inter-cryostat region (ICR) where there is little calorimeter coverage, we look for electrons that have been reconstructed as taus. In the various gaps in muon coverage, we look for isolated tracks. These additions improve our signal acceptance by 15%.

3.2 Jets

We select events with at least 2 jets, leading jet $p_T > 20$ GeV and second jet $p_T > 15$ GeV. Before tagging, S/B = 0.0003.

Using D0's neural net b-tagging algorithm [6], we require either two loose tags (inclusive) or one tight tag (exclusive), which improves S/B by factors of 20 and 10, respectively. The optimization of our final multivariate discriminant depends significantly on b-tag criteria, as do the expected signal and background yields, so we can improve our Higgs sensitivity by analyzing these orthogonal b-tag samples separately.

With the two highest- p_T tagged jets or the one tagged jet and highest- p_T untagged jet, we compute the invariant mass of the dijet system, which is the kinematical variable most sensitive to low-mass Higgs production.

3.3 Kinematic fit

With an ideal detector, we would have very little missing E_T in $ZH \to \ell^+\ell^-b\bar{b}$ events. Thus, events with a large p_T imbalance must result from either background processes or mismeasurement. Given our knowledge of the jet and lepton energy resolutions, we can make our measurements more precise and discriminate against backgrounds such as multijet production and $t\bar{t} \to \ell^+\ell^-\nu\bar{\nu}b\bar{b}$. To do this, we perform a constrained multi-dimensional fit on the p_T , η , and azimuthal angle ϕ of the two leptons and two candidate jets. We constrain the dilepton invariant mass to 91.2 ± 2.5 GeV and the vector sum of p_T to 0.0 ± 7.0 GeV.

Subsequently, we remove events with high kinematic fit χ^2 values to reduce instrumental background. As a result of the fit, the dijet mass resonance in Higgs events is more prominent, which translates directly to a 6-11% increase in Higgs sensitivity.

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4 Results

We use boosted decision trees (BDT) [7] to combine the discrimination power of several kinematical variables: the dijet mass and p_T , the dilepton p_T and colinearity, and many others. No evidence for ZH production is seen, so we compute upper cross section limits based on the shape of the BDT output, using a modified frequentist approach [8, 9]. The leading sources of systematic uncertainty are the Z+heavy-flavor cross section (20%), the jet energy scale (10%), and b-tagging efficiencies (10%). Assuming a Higgs mass of 115 GeV, we exclude ZH production above 9.1 times the standard model expectation at 95% confidence level. Limits assuming other Higgs masses are shown in Fig. 1.

Upon comparison to previous results [10], our limits have improved roughly 12% beyond

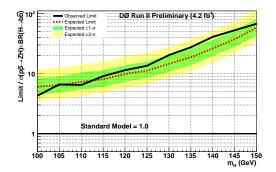


Figure 1: Expected and observed limits on ZH production, expressed as ratios to the standard model expectations.

what is expected with more data. This is due to our use of more optimal selection criteria and more sophisticated analysis techniques. We are currently investigating further improvements to the analysis, including the use of matrix-element discriminants, improved b-tagging, and further optimization of our multivariate discriminant.

References

- [1] LEP, Tevatron and SLD Electroweak Working Groups, arXiv:0911.2604 [hep-ex] (2009).
- [2] The D0 Collaboration, D0Note 5972-CONF (2009).
- [3] The D0 Collaboration, arXiv:0912.5285 [hep-ex] (2010).
- [4] The D0 Collaboration, D0Note 5876-CONF (2009).
- [5] V. M. Abazov et al. (D0 Collaboration), Nucl. Instrum. Methods in Phys. Res. A 565, 463 (2006); M. Abolins et al., Nucl. Instrum. and Methods A 584/1, 75 (2007); R. Angstadt et al., arXiv:0911.2522 [physics.ins-det].
- [6] T. Scanlon, FERMILAB-THESIS-2006-43.
- [7] L. Breiman et al., "Classification and Regression Trees" (Wadsworth, Belmont, CA, 1984); Y. Freund and R. E. Schapire, Machine Learning: Proceedings of the Thirteenth International Conference, Ed. L. Saitta (Morgan Kaufmann, San Francisco, 1996), p148.
- [8] T. Junk, Nucl. Instrum. Methods in Phys. Res. A 434, 435 (1999); A. Read, J. Phys. G 28, 2693 (2002).
- [9] W. Fisher, FERMILAB-TM-2386-E (2006).
- [10] The D0 Collaboration, D0Note 5570-CONF (2008).

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