

Forward photon spectrum in 7 TeV pp collisions measured by the LHCf experiment

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LHCf is an experiment designed to measure the energy and transverse momentum spectra of very forward neutral particles produced at the Interaction Point 1 of LHC. Its goal is to calibrate the hadron interaction models in air-shower simulations used for ultra-high-energy cosmic-ray (UHECR) experiments, and its results would also be useful to get a glimpse of the physics in inelastic hadronic collisions. LHCf has taken data in 2009-2010 pp runs at 0.9 and 7 TeV. Then the detector was removed from the experimental area. In this paper the forward photon spectrum obtained at 7 TeV will be presented. Results will be discussed compared to the MC models widely used in UHECR physics. Future prospects for the analysis and the data taking at 14 TeV will be also discussed.

1 Introduction

Cosmic rays with energy above 10^{15} eV have been detected on the Earth since several decades ago, but there still remain many uncertainties in understanding their properties. In particular, the properties of cosmic rays beyond 10^{18} eV, called ultra-high-energy cosmic ray (UHECR), have not been clarified yet. We do not know much about its chemical composition, the mechanism of its propagation, and thus its astrophysical origin. One reason is that it can be detected indirectly only with a particle cascade called air-shower, and that we must always care for systematic errors caused by the air-shower development.

If we try to understand the cosmic rays in a step-by-step way from the Earth side, we must understand first what the cosmic rays are, i.e., the chemical composition. For its determination, the maximum depth of the air-shower development is generally used. Recently, a measurement by Pierre Auger Observatory suggests that the chemical composition of UHECR has a gradual transition from a light chemical composition (proton) to a heavy one (iron) [1]. However, this result is different from a result by the HiRes experiment [2]. Furthermore, the depth of the shower maximum is severely dependent on the hadronic interaction models used in Monte Carlo simulations for the air-shower development. The models are based on the perturbative QCD, but they treat the non-linear effects only in phenomenology. Therefore, in order to reduce the uncertainty in the discussion of the chemical composition, we must reduce first the model dependence with experimental data obtained in colliders like Large Hadron Collider (LHC).

There are several key quantities for the air-shower development. It starts with the first interaction of the cosmic ray and a nucleus of the atmosphere, where the inelastic cross section of the first interaction is important. In LHC the TOTEM experiment has already measured the inelastic cross section [3]. Then, the energy of the cosmic ray is transferred inelastically

to secondary particles like baryons and mesons. These secondary particles in the forward direction make following particle cascades. Thus, the amount of transferred energy (inelasticity) and the energy spectrum of the forward particles are important here. The latter is the main goal of the LHCf experiment [5]. It measures the spectra of very forward neutral particles (photons, neutral pions and neutrons) emitted in high energy collisions in LHC. As for the spectrum of forward pions, available data with the highest energy is that for neutral pions by the UA7 experiment [4] in SPS, whose collision energy corresponds to 10^{14} eV in cosmic-ray energy. On the other hand, the LHCf experiment is designed to obtain data with a collision energy of 14 TeV, corresponding to 10^{17} eV in cosmic-ray energy. Thus, it will be expected to reduce the uncertainty of the model dependence close to the UHECR energy.

2 The LHCf detector

LHCf is installed at the CERN Large Hadron Collider in Geneva, Switzerland. The LHCf detector is located at 140m away from Interaction Point 1 (IP1; the ATLAS site) and at zero degree collision angle. The detector is installed in the instrumentation slots of the neutral particle absorbers (TAN). Inside the TAN the beam vacuum chamber makes a Y-shaped transition from a single beam tube facing the IP to two separate beam tubes joining to the arcs of LHC. The TAN instrumentation slot is in the crotch of the Y. Charged secondary particles from IP are swept aside by the inner beam separation dipole before reaching TAN, so only neutral particles are incident on the LHCf detectors. This location covers the pseudo-rapidity range from 8.4 to infinity. This instrumentation slot exists in each of both sides of IP1, thus we have two detectors located on opposite sides of IP1, named Arm1 and Arm2.

The overall concepts for the two Arms are the same (Fig. 1). Each of the two LHCf detectors consists of two small sampling calorimeters and four position sensitive layers inserted into the calorimeter layers. The calorimeter is made of 16 plastic scintillators sandwiched by tungsten absorbers, and it has a total length equivalent to 44 radiation lengths, and 1.55 interaction length. We call the calorimeters small & large towers. The four position sensitive layers are distributed among the layers of the calorimeters for determining the transverse shower positions. Arm1 utilizes scintillating fibers as the position-sensitive layers, while Arm2 uses silicon (Si) microstrip sensors. In addition the geometrical configurations of the two towers for each Arm are different for the purposes of redundancy and consistency check. In front of each detector, a Front Counter made of plastic scintillators is inserted. It provides useful trigger information by covering a larger aperture than the calorimeters. Many additional details of both Arms can be found in [5].

The detector can identify photons with the two calorimeters. It measures their energy spectrum beyond 100 GeV with less than 5% energy resolution, and their incident position with 0.2mm position resolution. If a photon is detected in each of the two calorimeters at the same time, it is expected to be reconstructed as an event from a neutral pion. This is the main target of the LHCf detector. Measured spectrum of neutral pions can be used to discriminate among the hadronic interaction models. Also, hadronic showers of high-energy neutrons can be measured with energy resolution of about 30%.

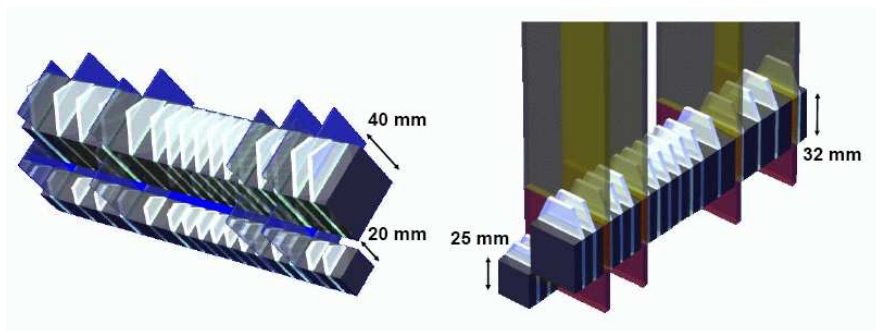


Figure 1: Schematic view of the LHCf detectors (Arm1 on the left, Arm2 on the right). Plastic scintillators (light blue color) are interleaved with tungsten blocks (dark gray color). Four position sensitive layers (scintillating fibers in Arm1, dark blue color; silicon micro-strip detectors in Arm2, purple color) are distributed in each calorimeter.

3 Data acquisition

The current LHCf detector was not designed to be a radiation hard detector. That is because the model discrimination requires just a short period during the early phase of the LHC commissioning before the high luminosity operations. For the LHC operation in 2009 and 2010, we have obtained data for pp collision with energies of $\sqrt{s} = 0.9$ and 7 TeV. The accumulated numbers of events are about 100k showers at 0.9 TeV, and about 400M showers at 7 TeV, corresponding to about 1M neutral pion events, for each of the two Arms. After the operation for several months, the detectors were removed from the slot in 2010.

In Fig. 2, we show an example of a candidate of a neutral pion event, i.e., one photon for each tower, obtained by Arm2. We can see sharp peaks from single photon events detected by the first two layers of the Si microstrip sensors. With the information we can reconstruct the invariant mass of the photon pairs. It is expected to be around the mass of the neutral pion (135 MeV) and will be used for an energy calibration. We have also obtained the energy spectra of the neutral pions for both of the Arms. A preliminary result can be seen in [6].

4 Photon analysis

The first physics result from the LHCf experiment has been published in 2011 [7]. It is dedicated to measurements of the single photon spectrum. Events by neutral hadrons are removed by a simple criteria based on the longitudinal development of the showers, obtained by the scintillators. The energy of the remained photons is determined from the same information on the light produced in the scintillators. We applied corrections to it: for non-uniformity of light collection efficiency, and for particles leaking out of the edges of the calorimeter towers. In this corrections we used the lateral positions of showers determined with the information by the position sensitive layers. The information by the position sensitive layers is also used to exclude ‘multi-hit’ events that have more than one showers inside the same tower. The two Arms have different geometrical configurations, thus in this analysis we have selected a common region for each tower in order to combine the two spectra without any acceptance correction.

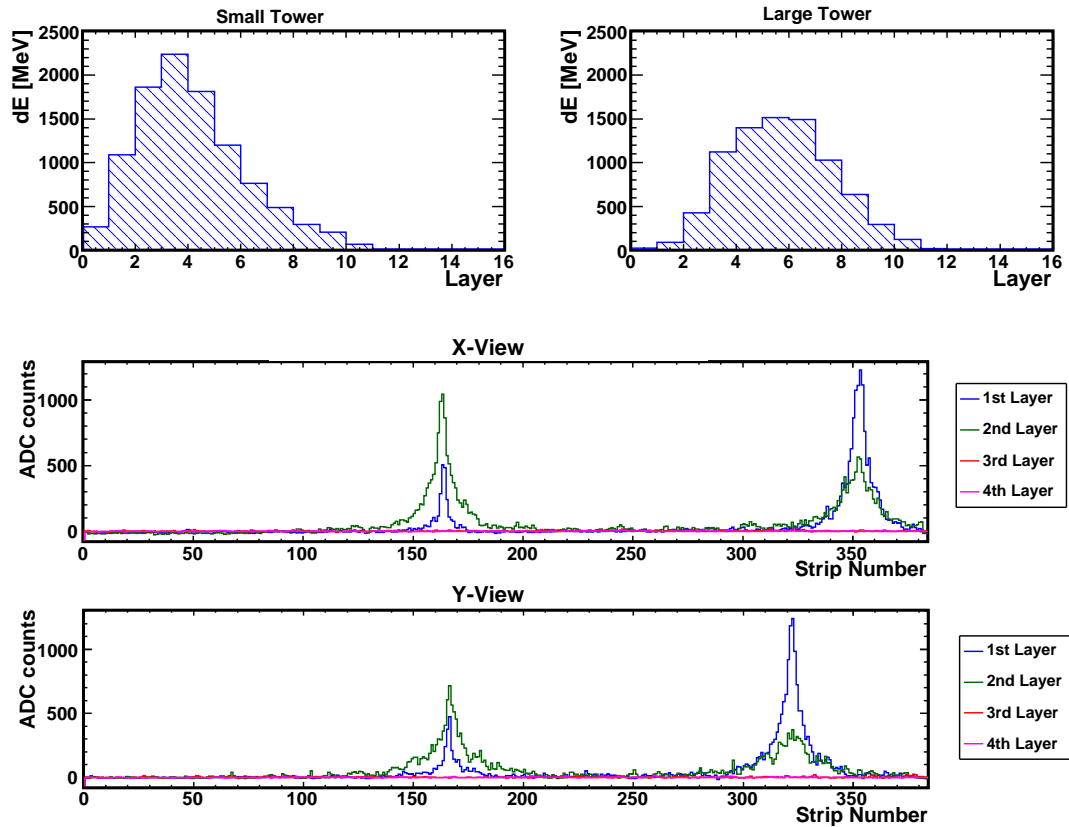


Figure 2: An example of data taken during 2010 operation [7] (a candidate for detection of two photons from the decay of a neutral pion). The upper two panels show the transition curves taken by the two calorimeters: the left is by the small tower, while the right is by the large tower. The lower two panels are the data from the Si microstrip sensors; The upper is the four x-view data, while the lower is the four y-view data. The identifiable peaks are from the 1st and 2nd layers, while the data from the 3rd and 4th layers are almost zero.

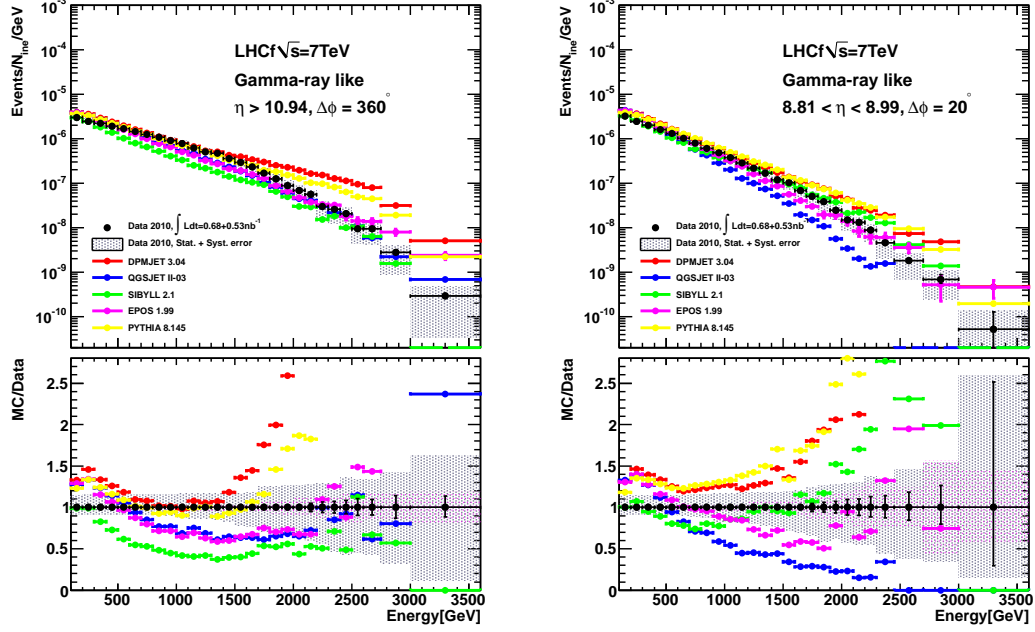


Figure 3: Comparison between the measured single-photon energy spectra (black dots) and the predictions of the following MC codes: DPMJET 3.04 (red), QGSJET II-03 (blue), SIBYLL 2.1 (green), EPOS 1.99 (magenta) and PYTHIA 8.145 (yellow), taken from [7]. Top panels show the spectra and bottom panels show the ratio of MC results to experimental data. Left and right panels refer to different pseudo-rapidity ranges. Error bars show the statistical error and gray shaded areas the systematic error for experimental data. Magenta shaded areas indicate the statistical error associated to MC simulations.

The region is pseudo-rapidity $\eta > 10.94$ and azimuthal range $\Delta\phi = 360^\circ$ for the small towers, while $8.81 < \eta < 8.99$ and $\Delta\phi = 20^\circ$ for the large towers. Further details about the analysis can be found in [7].

Figure 3 shows the single-photon energy spectra obtained from a data set taken in pp collisions at $\sqrt{s} = 7$ TeV in 2010. During the period for the data set, the integrated luminosity is estimated to be $\int Ldt = 0.68 \text{ nb}^{-1}$ and 0.53 nb^{-1} for Arm1 and Arm2, respectively. Multiplying it with an assumed inelastic cross section $\sigma_{\text{ine}} = 71.5 \text{ mb}$, we derived the number of inelastic collisions, N_{ine} , on the vertical axis. The black points are the energy spectra obtained by the combination of the two Arms, and they are compared with results predicted by MC simulations using different models: DPMJET 3.04[8], QGSJET II-03 [9], SIBYLL 2.1 [10], EPOS 1.9[11] and PYTHIA 8.145[12, 13]. Left and right panels refer to the selected regions as mentioned before.

As a result, we see that none of the model predictions nicely describe the LHCf data in the whole energy ranging from 100 GeV to 3.5 TeV. In particular there is a big discrepancy in

high energy region. Discussions about this result with the model developers have been already started. Improvements of the hadron interaction models are expected in near future.

5 Future prospects

Now we are working for the following analyses for topics such as neutral pions, 900 GeV photons, and Pt distribution of photons and hadron spectra at 7 TeV collisions. At the same time, we are preparing for foreseen data acquisitions with the proton-ion collision in 2012, and the pp collisions at $\sqrt{s} = 14$ TeV in 2014. A higher luminosity is expected for the 14 TeV runs than that in the 7 TeV runs. Thus we are planning a hardware upgrade of the plastic scintillators and the scintillating fibers to ones made of Gd_2SiO_5 (GSO), which are more radiation-hard than the plastic scintillators. The radiation hardness of the GSO scintillator was measured in a test beam, and we confirmed that it has properties good enough for our use [14].

Another upgrade is planned for the Arm2 detector. The current configuration of the Si microstrip sensors distributed in the plastic scintillators are not optimized for an energy reconstruction using the Si sensors. If we change the configuration, an improvement of the energy resolution only with the Si is expected, by a MC study, to be less than 10%. However, the Si sensors were originally not considered as a calorimeter, thus we require additional work for them. The first is the gain calibration using the test beam data and a MC simulation for the same configuration. We have obtained a gain factor from ADC counts to the energy deposit on the Si sensors. Then the incident particle energy can be reconstructed from the energy deposit, using a function obtained by a simulation for LHC configuration. A preliminary result with a MC data set shows that the energy resolution for photons is $\sim 15\%$ up to 1.5 TeV [15]. When the Si sensor is used as a calorimeter, it would help not only the calorimetry of the scintillators, but also a separation of the ‘multi-hit’ events, which could not be resolved only by the scintillators.

6 Conclusions

LHCf is a collider experiment dedicated for the cosmic-ray physics. It measures the energy spectrum of very forward particles generated in pp collisions in LHC. Its aim is to minimize systematic errors in air-shower simulations used in ultra-high-energy cosmic-ray experiments. In this paper we briefly surveyed our detector and the current status of the experiment. We have already done the data acquisition for pp collisions at $\sqrt{s} = 7$ TeV, and removed the detectors. The obtained photon spectra are not well described by the MC expectations, especially in high energies. Now we are working for following analyses and foreseen detector upgrade. The prospects about the hardware upgrade were also shown in this paper.

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