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On the Photoproduction of Multipion Systems

by

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Summary: Combining vector meson dominance and quark model with isospin independence of production reactions above the resonance region, we predict the cross-sections for multipion photoproduction from the corresponding $\pi-N$ data.

For many aspects in the description of the electromagnetic interaction of hadrons the vector meson dominance model (VDM) l) has met with considerable success $^{2)}$. Using the electromagnetic current $\mathbf{j}_{\mu}(\mathbf{x})$ as interpolating vector meson field

$$j_{\mu}(x) = \sum_{V} \frac{m_{V}^{2}}{2\gamma_{V}} \Psi_{\mu}^{V}(x) = \frac{m_{\rho}^{2}}{2\gamma_{\rho}} \rho_{\mu}(x) + \frac{m_{\omega}^{2}}{2\gamma_{\omega}} \omega_{\mu}(x) + \frac{m_{\phi}^{2}}{2\gamma_{\phi}} \phi_{\mu}(x)$$
(1)

(with constants γ_V) and assuming a slow variation of the corresponding matrix elements as we extrapolate from photon mass $k^2=0$ to vector meson mass $k^2=m^2_V$, the VDM gives a connection between γ -induced and V-induced hadronic reactions. Neglecting interferences between ρ , ω , and ϕ we thus have

$$\sigma(\gamma p \to x) = \sum_{V} \frac{\alpha \pi}{\gamma_{V}^{2}} \sigma(V p \to x)$$
 (2)

with $\alpha=1/137$; X denotes any compatible hadronic final state and $\sigma(Vp \to X)$ describes the cross-section of the V-p reaction with transversally polarized vector mesons.

In the present note we wish to apply (2) to obtain a prediction of multipion photoproduction processes from available π -N data. Since experiments on ρ^0 -photoproduction indicate a ratio $\vec{\gamma}_\rho^2$: $\vec{\gamma}_\omega^2$: $\vec{\gamma}_\phi^2$ of about 16:3:1/2, we shall here include only ρ and ω in (2). As coupling constant values we take $\gamma_\rho^2/4\pi=1/2$, $\vec{\gamma}_\omega^2=(3/16)\vec{\gamma}_\rho^2$, in accord with the mentioned ρ^0 -photoproduction data as well as with most other VDM applications 2).

To connect V-p with π -p processes, we appeal to the sum rule

$$\sigma_{\text{tot}}(\rho^{\circ}p) = \sigma_{\text{tot}}(\omega p) = \frac{1}{2} \left[\sigma_{\text{tot}}(\pi^{\dagger}p) + \sigma_{\text{tot}}(\pi^{\circ}p)\right] = \sigma_{\text{tot}}(\pi^{\circ}p)$$
 (3)

derived from the quark model $^{4,5)}$ with additivity of quark scattering amplitudes. This and the agreement of diffraction peak widths $^{3)}$ yield equal ρ^{0} -p and π^{0} -p elastic cross-sections; therefore we shall assume (3) to hold term by term

$$\sigma(\rho^{O}p \to x) = \sigma(\omega p \to x) = \sigma(\pi^{O}p \to x) \tag{4}$$

for incident energies well above threshold; (an "additive" quark model for inelastic reactions would presumably also lead to (4), since ρ^0 and π^0 have the same quark content). Thus we obtain

$$\sigma(\gamma p \to x) = 4.33 \times 10^{-3} \sigma(\pi^{0} p \to x)$$
 (5)

where X denotes any hadronic final state with more than two pions.

Beyond VDM and the extended sum rule (4) we lastly assume the isospin independence of π -N amplitudes for energies above the main resonance region (> 2 - 3 GeV/c). Neglecting interferences, this reduces the relation between π -N cross-sections of identical pion and nucleon number final and initial states but of different charge configurations to a relation of the corresponding Clebsch-Gordan coefficients. Thus we have e.g.

$$\sigma(\pi^{\circ}p \to p2\pi^{+}2\pi^{-}) / \sigma(\pi^{-}p \to p\pi^{+}2\pi^{-}\pi^{\circ}) = \frac{.201}{.358}, \text{ etc.}$$
 (6)

Isospin independence, suggested by statistical considerations, is rather well supported both by the near-equality of π^+ and π^- total cross-sections and by analyses of various multipion experiments $^{7,8)}$.

Our predictions obtained via (5) and isospin independence from the available π -N data⁹⁾ are shown in Figures I and II. The errors given are those from the π -N experiments, assuming here an exact value of the γ -V coupling constant.

In Fig. I we compare our result with the measured values $^{3,10)}$ for the reaction $\gamma p \rightarrow p \ 2\pi^{+}2\pi^{-}$, so far the only multipion photoproduction cross-section for which the energy dependence has been established. The agreement is seen to be very good.

Fig. II gives the prediction for the seven-prong $\gamma p \to p\pi 3\pi^{+}3\pi^{-}$ as well as the predicted behaviour of the no-fit channels $\gamma p \to p \ 2\pi^{+}2\pi^{-}\pi^{-}$, $p \ 2\pi^{+}2\pi^{-}2\pi^{-}$. It is seen that the smallness of the seven-prong cross-section is essentially due to isospin effects, not to a small $p \ 6\pi$ cross-section relative to $p \ 5\pi$.

In conclusion we find that the VDM successfully relates electromagnetic and hadronic multipion production. On one hand this gives additional experimental checks of the γ -V coupling constant, on the other one can generally apply

the method used here to evaluate contributions from no-fit events in multipion photoproduction. Finally we consider our results as an indication that perhaps quark concepts may eventually also lead to useful "sum-rules" for genuinely inelastic reactions.

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Note added in proof: After completion of this note, data on the seven-prong reaction have come in 10) and are found to agree with our predictions (see Fig. II).

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Figure Captions

- Fig. I: The predictions for $\sigma(\gamma p \to p \ 2\pi^+ 2\pi^-)$ as calculated from $\pi^- p$ data (•: $\pi^- p \to n \ 2\pi^+ 2\pi^-$: Δ : $\pi^- p \to p \ \pi^+ 2\pi^- \pi^0$: \square : $\pi^+ p \to p \ 2\pi^+ \pi^- \pi^0$:

 •: $\pi^+ p \to n \ 3\pi^+ \pi^-$): the measured $\sigma(\gamma p \to p \ 2\pi^+ 2\pi^-)$ is denoted by $\pi^+ \to 0$.
- Fig. II: The prediction for $\sigma(\gamma p \to p \ 3\pi^{+}3\pi^{-})$ (\Box), $\sigma(\gamma p \to p \ 2\pi^{+}2\pi^{-}2\pi^{0})$ (Δ), and $\sigma(\gamma p \to p \ 2\pi^{+}2\pi^{-}\pi^{0})$ (o) as calculated from $\pi^{+}p \to p \ 3\pi^{+}2\pi^{-}$, $p \ 3\pi^{+}2\pi^{-}\pi^{0}$ data. The measured $\sigma(\gamma p \to p \ 3\pi^{+}3\pi^{-})$ is denoted by \Box .





