

# DEUTSCHES ELEKTRONEN-SYNCHROTRON **DESY**

DESY 79/81  
December 1979



MEASUREMENTS OF THE PROPERTIES OF THE  $\gamma$  FAMILY FROM  $e^+e^-$  ANNIHILATION

by

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MEASUREMENTS OF THE PROPERTIES OF THE Y FAMILY FROM  $e^+e^-$  ANNIHILATION

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Invited talk given at the 1979 International Symposium on Lepton  
and Photon Interactions at High Energies, FNAL, August 23-29, 1979

Summary

Information on Y particles from  $e^+e^-$  annihilation experiments at DORIS is summarized. The mass differences in the Y family as well as the coupling to lepton pairs is found to be very similar to the  $\psi$  family. The charge of the constituent b quark is established to be  $1/3$ . The event structure reveals non trivial correlations between the decay particles. The three-gluon decay hypothesis assuming spin 1 gluons provides a quantitative explanation of the data. Gluon jets are found to be rather similar to quark jets for jet energies up to  $\sim 5$  GeV.

Introduction

Information on the properties of the Y family from  $e^+e^-$  annihilation comes from experiments at the  $e^+e^-$  storage ring DORIS. The data were taken about a year ago by three collaborations: DASPII (DESY-Dortmund-Heidelberg-Lund), DHHM (DESY-Hamburg-Heidelberg-München) and PLUTO (Aachen-DESY-Hamburg-Siegen-Wuppertal). The experiments were made possible by a great effort of the DORIS storage ring group to operate the  $e^+e^-$  ring at energies much higher than originally foreseen<sup>1</sup>.

In the first running period in April - May 1978 at total energies more than 9.0 GeV DASPII<sup>2</sup> and PLUTO<sup>3</sup> obtained resonance signals corresponding to the Y(9.46) as shown in fig. 1. PLUTO then left the south interaction region of DORIS to be ready for first measurements at PETRA and was replaced by the DESY-Heidelberg detector. The Y(9.46) was confirmed in June 1978 and finally in August 1978<sup>4,5</sup> also the Y'(10.02) was detected, fig. 2.

1. Parameters of the Y Particles

From the resonance signals several interesting parameters have been determined, the masses, the partial width  $\Gamma_{ee}$ , the hadronic width  $\Gamma_{had}$  and the charge of the constituent quark. The resonance signal at mass M with total width  $\Gamma_{tot}$  and hadronic width  $\Gamma_{had}$  is given by

$$\sigma_{res} = \frac{12 \pi M^2}{S} \frac{\Gamma_{had} \Gamma_{ee}}{(S - M^2)^2 + M^2 \Gamma_{tot}}$$

with  $S=4 E_0^2$  is the total center of mass energy squared. A single parameter is determined,  $\Gamma_{ee}$ , the partial width to  $e^+e^-$  pairs, if

$$\Gamma_{had} \gg \Gamma_{ee}.$$

The observed width of the Y(9.46),  $\sigma = 7.3$  MeV, and of the Y'(10.02),  $\sigma = 11$  MeV, is due to the finite energy spread of the  $e^+e^-$  beams in DORIS and compatible with values calculated from the  $e^+e^-$  beam dynamics.

Radiation of photons in the initial stage strongly reduces the resonance signal, in case of Y(9.46) at DORIS by a factor of about  $1/0.63$ . The values for  $\Gamma$  obtained are shown in table 1. The mass determinations are dominated by the uncertainty in absolute energy calibration of the DORIS ring, about 10 MeV at the Y(9.46) and 20 MeV at the Y'(10.02). The much larger uncertainty at 10 GeV as compared to 9.5 GeV results from strong saturation effects in the DORIS magnets near the end of the accessible energy range.

Using the masses of the Y(9.46) and Y'(10.02) as input it was possible from the high resolution data of

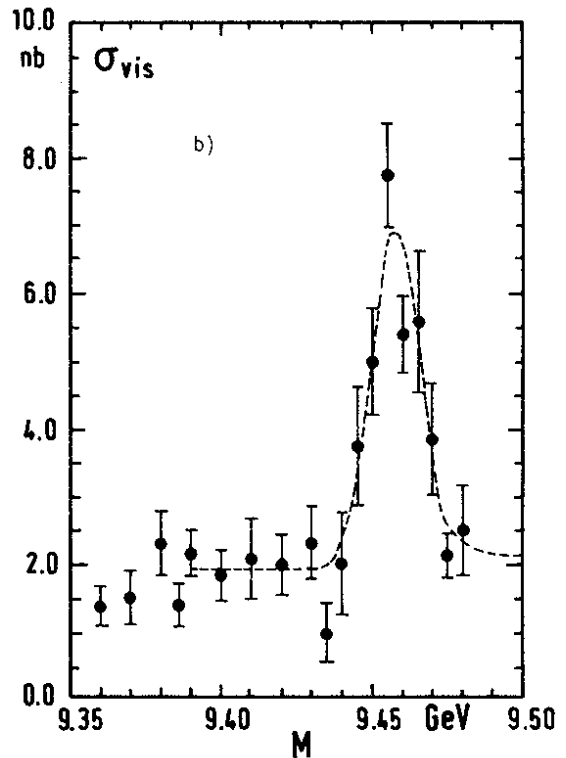
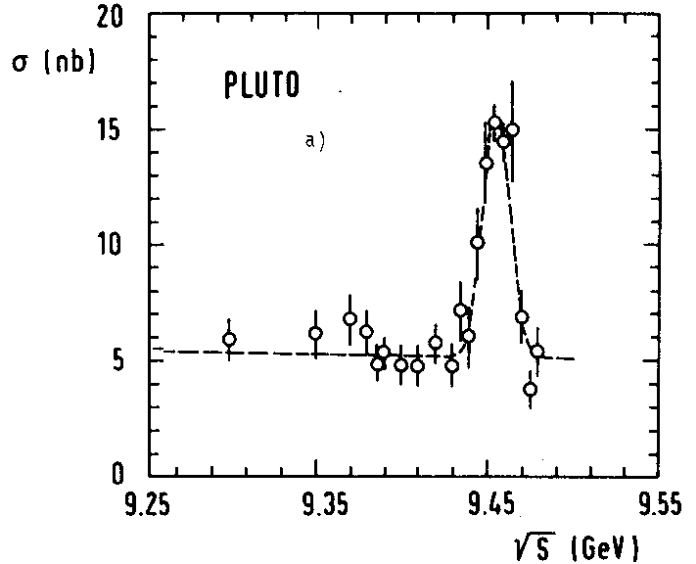


Fig. 1 The Y(9.46) signals from  
a) PLUTO and  
b) DASPII.

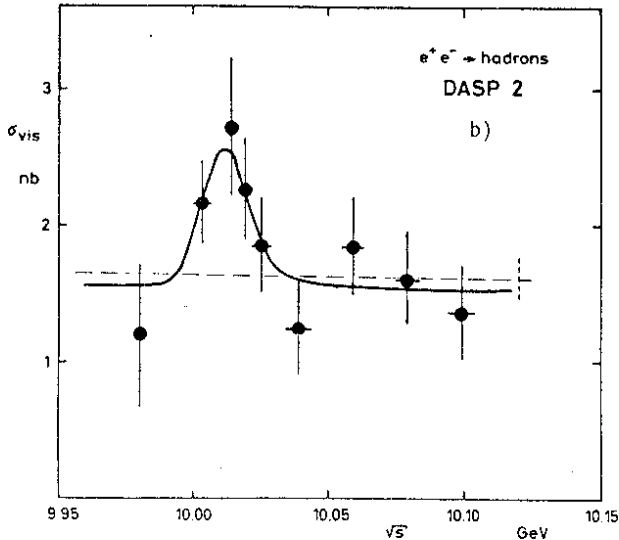
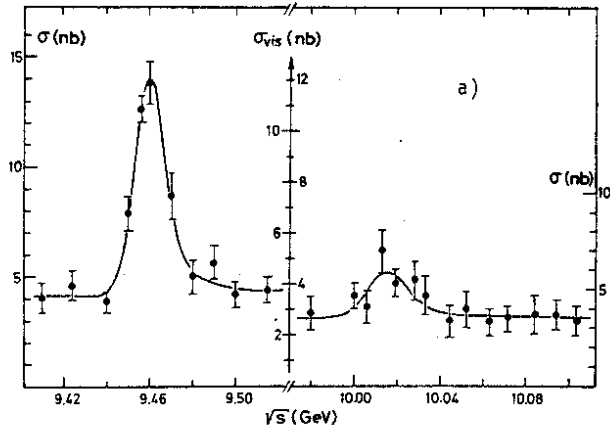


Fig. 2: The  $Y(9.46)$  and  $Y'(10.02)$  signal as  
a) seen by DHM and  
b) the  $Y'(10.02)$  from DASPII.

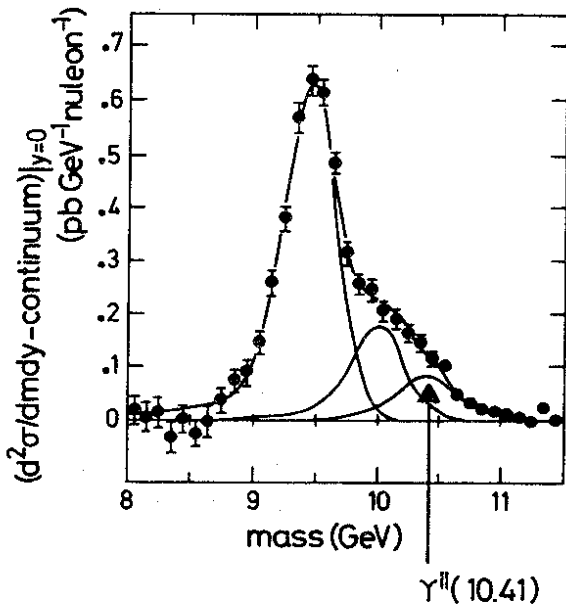
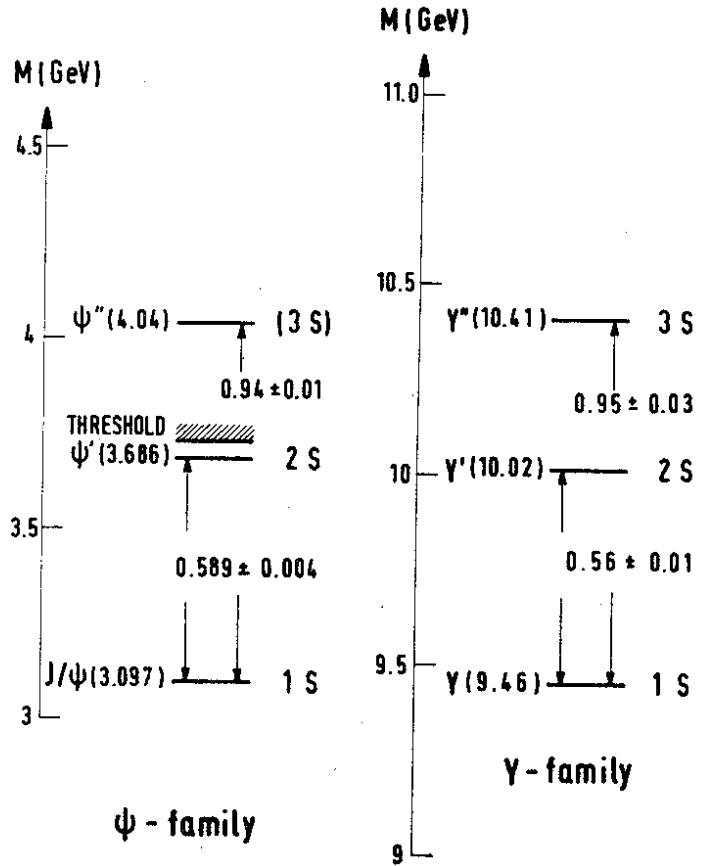


Fig. 3: High resolution data in  $M_{\mu^+\mu^-}$  from  
ref. 6 for the reaction  
 $p + N \rightarrow \mu^+\mu^- + \text{anything}$ .

TABLE 1

| Experiment  | Ref. | Mass(Gev)         | $\Gamma_{ee}$ (KeV) | $B_{\mu\mu}$ (%)       | $\sigma$ (MeV) |
|-------------|------|-------------------|---------------------|------------------------|----------------|
| PLUTO       | 15   | $9.456 \pm 0.01$  | $1.33 \pm 0.14$     | $2.2 \pm 2.0$          | $7.3 \pm 0.1$  |
| DHMM        | 4,17 | $9.46 \pm 0.01$   | $1.04 \pm 0.28$     | $1.0 \pm 3.4$<br>$1.0$ | $7.1 \pm 0.8$  |
| DASP 2      | 2,16 | $9.457 \pm 0.01$  | $1.50 \pm 0.40$     | $2.5 \pm 2.1$          | 7.6            |
| Mean values |      | $9.46 \pm 0.01$   | $1.28 \pm 0.27$     | $2.1 \pm 1.4$          | $7.3 \pm 0.2$  |
| DHMM        | 4    | $10.02 \pm 0.02$  | $0.32 \pm 0.13$     |                        | $12 \pm 4$     |
| DASP 2      | 5    | $10.01 \pm 0.02$  | $0.35 \pm 0.14$     |                        | 9              |
| Mean values |      | $10.015 \pm 0.02$ | $0.33 \pm 0.14$     |                        | $10 \pm 2$     |

the collaboration at Fermilab <sup>6</sup> to establish the existence of a  $Y''$  at a mass of  $10.41 \pm 0.05$  GeV, see fig. 3. It leads to the level scheme for the Y family sketched below and compared to the  $\psi$  family. The  $Y''(10.41)$  waits for observation in  $e^+e^-$  annihilation



experiments as well as the detection of the threshold for the new 'B particles', probably at energies just above the  $Y''(10.41)$  mass. A signal recently observed in the  $\psi K\pi$  final state at a mass of 5.30 GeV may be taken as evidence for the B particles <sup>7</sup>. It would place the threshold at an energy of 10.60 GeV.

A very detailed study of quark-antiquark binding potentials by Quigg, Rosner and Thacker <sup>8</sup> allows a de-

termination of the quark charge from the measured width of  $Y(9.46)$  and  $Y'(10.02)$  to  $e^+e^-$  pairs. Fig. 4 illustrates the result based on the average of the values for  $\Gamma_{ee}$  from table 1, clearly the charge assignment  $1/3$  is favoured.

Recent data from PETRA on the total annihilation cross section confirm this conclusion. The sum rule for  $\sigma_{tot}$  including first and second order QCD corrections reads <sup>9,10</sup>

$$R = \frac{3S}{4\pi\alpha^2} \sigma_{had} = 3 \sum_i (e_{q_i})^2 \times \left[ 1 + \frac{\alpha_s}{\pi} + (1.98 - 0.116 N_f) \left(\frac{\alpha_s}{\pi}\right)^2 \right]$$

where  $e_{q_i}$  are the quark charges and  $N_f$  is the number of quark flavours. From u, d, s, c quarks we get an asymptotic ( $\alpha_s/\pi$  very small) value of  $10/3$ , a b ( $1/3$ ) adds  $1/3$  while a b ( $2/3$ ) will add  $4/3$ . Higher order corrections to R including calculations to second order in  $\alpha_s/\pi$  are small, less than 7% for a value  $\alpha_s = 0.2$ . Fig. 5 shows data from the PLUTO collaboration in the range  $3.6 \text{ GeV} \leq E_{CM} \leq 31.6 \text{ GeV}$ , with the contribution from the  $\tau$  lepton subtracted. The average value of R above 13 GeV is  $3.88 \pm 0.22$  ( $\pm 15\%$ ) in good agreement with the b ( $1/3$ ) charge assignment. Table 2 summarizes the values for R obtained by the four groups working at PETRA <sup>11-14</sup>. The possible systematic scale shifts reported are  $\sim 10 - 15\%$ , however, since all four experiments determine the luminosity separately and are rather different in event definition and acceptance corrections the true error on R for the four experiments taken together is probably only  $\sim 5\%$ . The b ( $2/3$ ) charge assignment then is excluded at the 6 standard deviation levels.

TABLE 2

| Experiment | Ref. | Energy (GeV) | R    | $\Delta R$ |
|------------|------|--------------|------|------------|
| PLUTO      | 11   | 13 - 31.6    | 3.88 | 0.22       |
| JADE       | 13   | 22 - 31.6    | 4.14 | 0.26       |
| MARK J     | 12   | 13 - 31.6    | 4.11 | 0.17       |
| TASSO      | 14   | 13 - 31.6    | 3.98 | 0.20       |
| Mean value |      |              | 4.03 | 0.12       |
| Theory     | 10   |              | 3.90 |            |

A comparison of the width  $\Gamma_{ee}$  divided by the quark charge squared of the Y family with the values known for the lower lying vector mesons is given in fig. 6. It suggests a rather constant behaviour with vector meson mass and a ratio of about 2.5 for the values from the ground state ( $^1S$ ) to the first radial excitation ( $^3S$ )

The determination of  $\Gamma_{ee}$  uses the assumption  $\Gamma_{had} \gg \Gamma_{ee}$  where  $\Gamma_{had}$  is the decay width of the  $Y(9.46)$  into hadrons.  $\Gamma_{had}$  can be determined from a measurement of a resonance signal in  $e^+e^- \rightarrow \mu^+\mu^-$ . Assuming lepton universality the DORIS groups have attempted to measure  $B_{\mu\mu}$  <sup>15-17</sup>. The results are summarized in table 1. The uncertainty in  $B_{\mu\mu}$  is still so large that the upper limit on  $\Gamma_{had}$  is only the observed width of the  $Y(9.46)$

$$25 \text{ keV} < \Gamma_{had} < 17000 \text{ keV.}$$

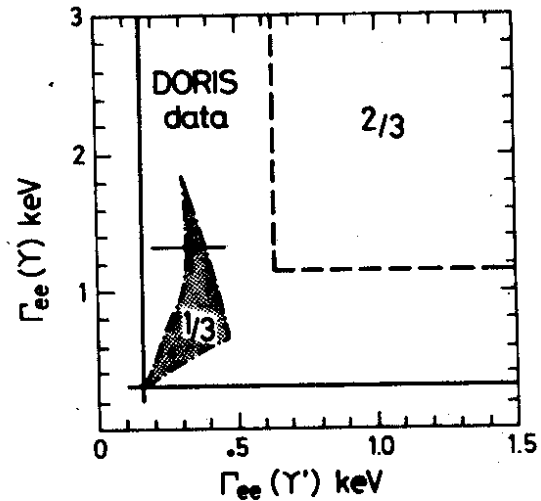


Fig. 4: Lower bounds for  $\Gamma_{ee}(Y)$  and  $\Gamma_{ee}(Y')$  for b-quark charges of  $1/3$  and  $2/3$ . The shaded area indicates the favoured region for charge  $1/3$ . (Ref. 8). The data point is from table 1.

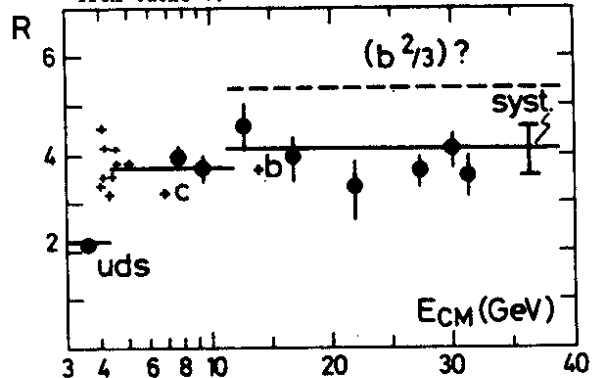


Fig. 5: Data from PLUTO on  $R(e^+e^- \rightarrow had)$ .

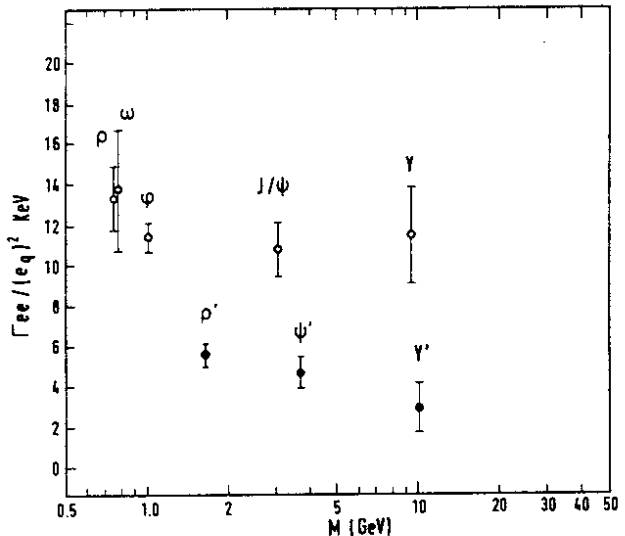


Fig. 6: Leptonic width divided by the quark charge squared for the known vector mesons.

A lower limit on  $\Gamma_{had}$  can also be derived using the event structure on and off resonance (see part 2 below). For the data from the PLUTO collaboration an increase of  $B_{\mu\mu}$  above 5 % leads to negative values for the lowest sphericity (thrust,  $\sin^2\delta$ ) bin in fig. 15 and therefore to a lower limit  $27 \text{ keV} < \Gamma_{had}$  in good agreement with the value obtained from  $B_{\mu\mu}$ .

## 2. Hadronic decays of the $Y$ and $Y'$

From the results on  $\Gamma_{had}$ ,  $\Gamma_{ee}$  and  $B_{\mu\mu}$  of the  $Y(9.46)$  we know that the electromagnetic contribution to the hadronic  $Y(9.46)$  decays, given by

$$\sigma_Y/\sigma_{res} = R \cdot B_{\mu\mu}$$

is less than  $\sim 20\%$ . It is the remaining part of the hadronic  $Y(9.46)$  decays that we are interested in and which will be discussed below. The experimental information is extracted as indicated in fig. 7. The signal at resonance contains three contributions

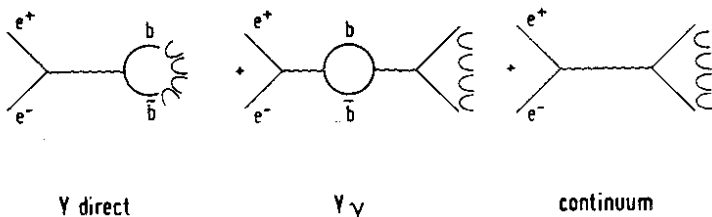
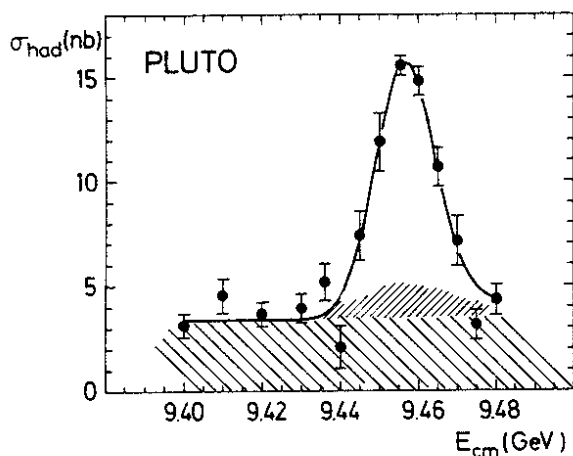


Fig. 7:  $Y(9.46)$  resonance signal from ref. 15. The three contributions to it ( $Y_{direct}$ ,  $Y_\gamma$  and continuum) are shown below and indicated by the shaded area.

the direct resonance decays ( $Y_{direct}$ ), the vacuum polarization term ( $Y_\gamma$ ) and nonresonant contributions (continuum). The second and third parts are identical in final state and are described by data just below resonance, since the energy variation in the final states going from 9.40 GeV to 9.46 GeV can safely be neglected. The main uncertainty is in the value of  $B_{\mu\mu}$ , however, the conclusions about properties of the hadronic  $Y(9.46)$

decays to be discussed below are only weakly dependent on the exact value of  $B_{\mu\mu}$ .

## 2.1. Models

The data on  $Y_{dir}$  will be compared with three models, different by the degree of correlation among the final state particles. One extreme is simple phase space, which has no other correlation between the particles than energy and momentum conservation. It provides a one (essentially) parameter comparison since the particle multiplicity has been fixed to correspond to the observed one. Also a ratio  $K/\pi = 0.30$  was put in.

Another extreme with very strong correlations is a two-jet model based on the description by Field and Feynman<sup>18</sup>. It has been extended beyond the original version by including heavy quarks, the charm quark ( $c$ ) and above  $\sqrt{s} = 11 \text{ GeV}$  the  $b$  quark. This model is known to give a fairly good fit to experimental data on hadronic jet properties in weak and electromagnetic current induced processes and in particular for the two-jet events in  $e^+e^-$  annihilation<sup>19</sup>.

However, there is also a very interesting model which supposedly provides a realistic framework for describing the  $Y(9.46)$  decay properties. In QCD a heavy quark-antiquark resonance system decays into 3 gluons in lowest order of the running coupling constant

$$\alpha_s = \frac{12\pi}{(33 - 2N_f) \ln S/\Lambda^2}$$

At  $Y(9.46)$  we expect a value for  $\alpha$  of  $\sim 0.2$ . Therefore, it seems not to be unreasonable to pursue a model for  $Y(9.46)$  decays based on the three-gluon decay matrix element (positronium analogue)<sup>20</sup>

$$\frac{1}{\Gamma} \frac{d\Gamma}{dx_1 dx_2} = \frac{6}{\pi^2 - 9} \times$$

$$\left| \frac{x_1^2(1-x_1)^2 + x_2^2(1-x_2)^2 + x_3^2(1-x_3)^2}{x_1^2 x_2^2 x_3^2} \right|$$

where  $x_i$  are the scaled (massless) gluon momenta,  $x_i = E_i(\text{gluon})/E(\text{beam})$  and  $x_1 + x_2 + x_3 = 2$ . The matrix element is rather flat over the  $x_1 x_2 x_3$  Dalitz plot (see fig. 24a) with variations of about 20%. The most frequent final state is two fast gluons  $x_1, x_2 \approx 1$  and  $x_3 \ll 1$  while the more spectacular 'mercedes star' structure  $x_1 = x_2 = x_3 = 2/3$  is the most unlikely one<sup>20</sup>. The structure of the average three-gluon event for the  $Y(9.46)$  drawn to scale is shown below.

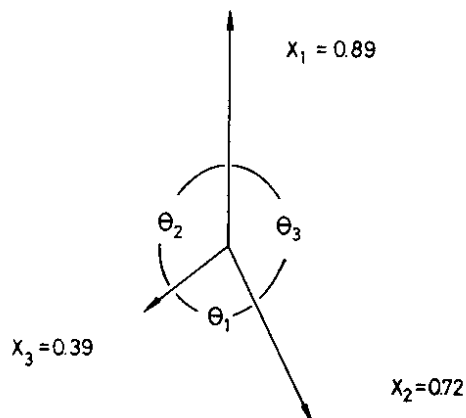


Fig. 8: The structure in momentum space of the average three-gluon event according to the positronium formula (ref. 20).

The fragmentation properties of the gluons have been assumed to be the same as for quarks. This assumption seems plausible since due to insufficient energy of the gluons no long tree with hard quark and gluon vertices can develop before the hadronization stage has been reached. The hadron distributions for each type of jet are certainly dominated by effects at the confinement radius as sketched below.

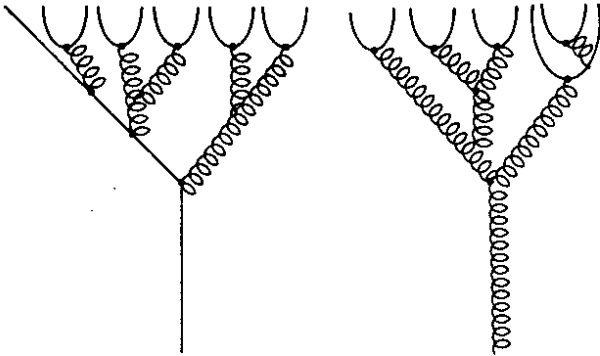


Fig. 9: Fragmentation trees of quarks and gluons at relatively low energy. After at most one hard radiation the confinement level is reached.

## 2.2. Two-Jet Analysis

The data in the vicinity of the resonances and at higher energies are clearly very well described by a dominant two-jet structure in good agreement with phenomenological parametrizations of quark jets à la Field-Feynman<sup>21</sup>. The  $Y(9.46)$  data are not described by a two-jet quark model. This is clearly seen for data from PLUTO using charged and neutral particles to get various measures of 'jettiness' sphericity (SPH)<sup>22</sup>, fig. 10, thrust ( $T$ )<sup>23</sup>, fig. 11, and the average energy weighted opening angle ( $\sin^2\theta$ )<sup>24</sup>, fig. 12. The dotted line shows the result of the Field-Feynman model, and also shown are the values for the phase space and the three-gluon model<sup>25</sup>. The value for phase space is much higher than the  $Y(9.46)$  data, while the three-gluon model is very close.

Fig. 13 shows data from DASPII with no subtraction procedure applied<sup>26</sup>. Since only angles and no momenta are measured 'pseudo' quantities have been defined to describe the jet properties. Clearly going across the resonance a change away from two-jet structure is observed.

A more detailed comparison is possible for distributions in sphericity or thrust. Fig. 14 shows the result from DHHM for both  $Y(9.46)$  and  $Y'(10.02)$  data and for the continuum<sup>27</sup>. Fig. 15 gives distributions from PLUTO.

The results can be summarized as follows:

(a) The continuum data are well described by the quark-jet parametrization from Field-Feynman.

(b) A simple phase space does not give a fit to the data from the  $Y(9.46)$ .

(c) The three-gluon model certainly is a rather good parametrization of the  $Y(9.46)$  (and  $Y'(10.02)$ ) data.

Two more arguments can be given to support these conclusions. Fig. 16 shows the sphericity (calculated from charged particles only) for different charged particle multiplicity classes<sup>25</sup>. It is seen that the  $Y(9.46)$  data is in between phase space and the two-jet model by about the same amount for each multiplicity class separately. It shows that essentially all events and not just a subclass from the  $Y(9.46)$  are different from the two extremes phase space and two-jet model

while there is again good agreement with the three-gluon model. Furthermore, from Fig. 17, we see that the

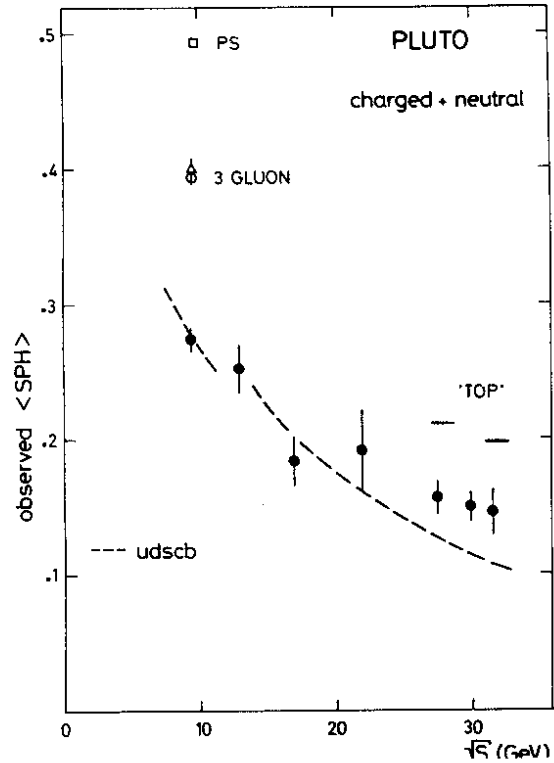


Fig. 10: The average sphericity vs. energy calculated from charged and neutral particles. The dotted line is the result of the Field-Feynman model. PS indicates phase space and 3 gluon the result of a three-gluon model prediction.

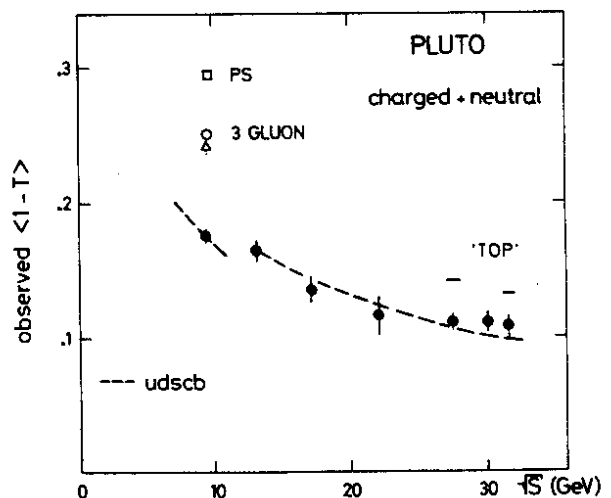


Fig. 11: The average thrust vs. energy, otherwise the same as fig. 10.

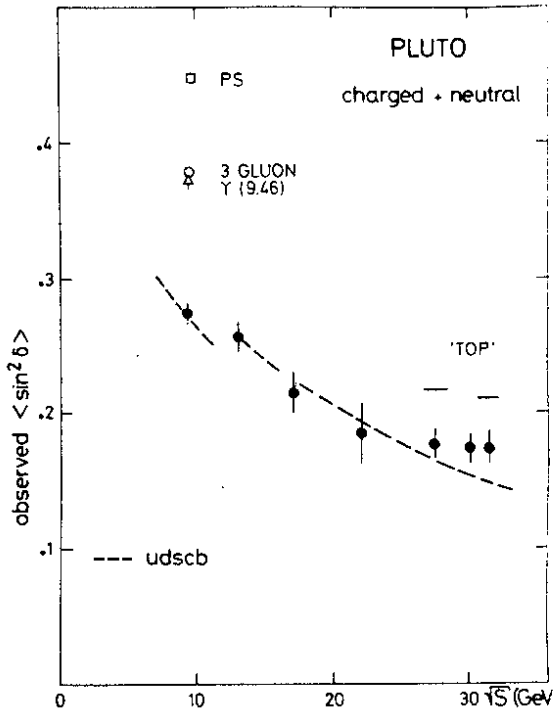


Fig. 12: The average energy weighted opening angle (squared) vs energy, otherwise the same as fig. 10.

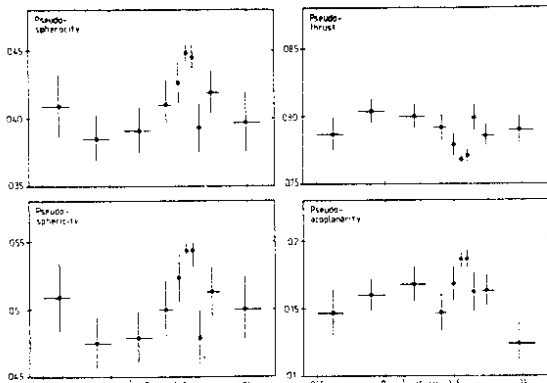


Fig. 13: Various jet quantities, calculated from the observed tracks in DASPII, ref. 26.

Fig. 14: Sphericity distributions for the  $\Upsilon(9.46)$ ,  $\Upsilon'(10.02)$  and the nearby continuum from ref. 27. For comparison the results from two-jet, three-gluon jet and phase-space calculations.

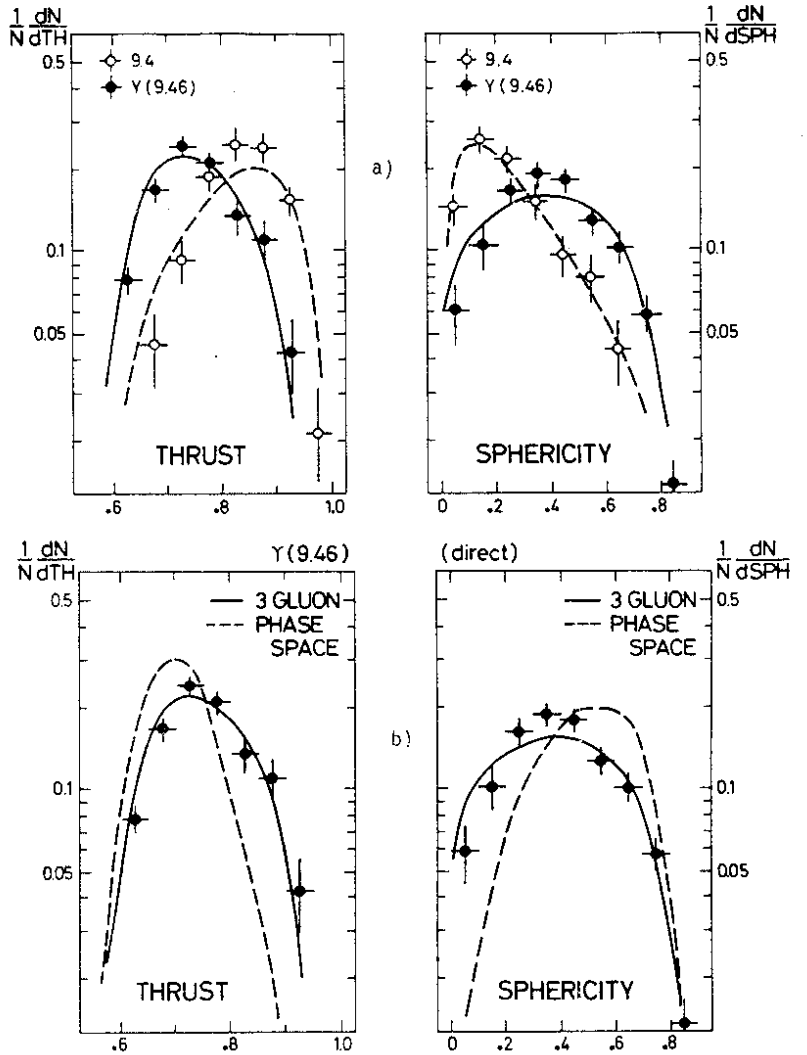
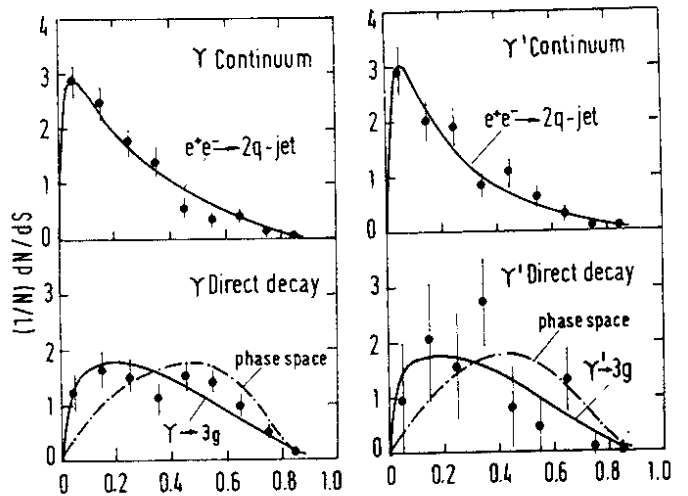


Fig. 15: Sphericity and thrust distribution for a)  $\Upsilon(9.46)$   $\blacklozenge$  and continuum (9.4)  $\circ$ , the dashed line is the two-jet model, and for b)  $\Upsilon(9.46)$   $\blacklozenge$  and phase space (dashed line). The full line is the result of the three-gluon model. The data is from PLUTO.





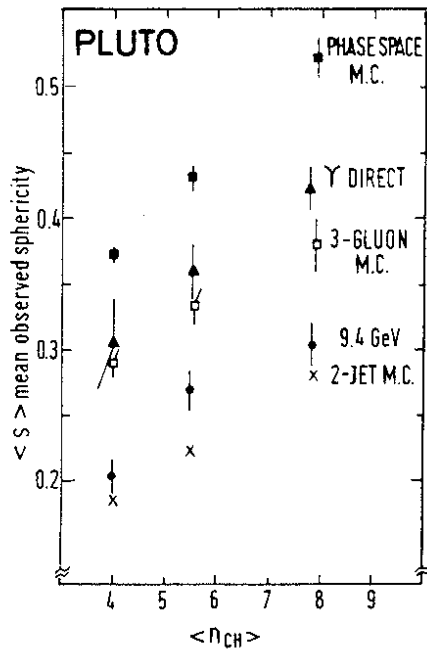


Fig. 16: Average sphericity (charged particles) for various multiplicity classes from ref. 25.

sphericity values for the  $Y(9.46)$  and the  $J/\psi(3.1)$  are about the same and in addition the values from phase space and at 3.6 GeV are very close to the  $J/\psi(3.1)$  value<sup>25</sup>. Going to higher energies the sphericity of phase space has to rise irrespective of the details of the phase space model just from the increase in particle multiplicity. It is found, however, that the sphericity of  $Y(9.46)$  is significantly lower than phase space for the same multiplicity. This shows in a model independent way that one finds experimentally nontrivial momentum correlations at  $Y(9.46)$ . In the three-gluon model this is due to the structure of the three-gluon matrix element which predicts three-gluon jets<sup>20</sup>. The gluon fragmentation properties are taken similar to quark fragmentation resulting in a very good agreement for  $Y(9.46)$  decay data and the three-gluon model.

Finally the fact that the sphericity on  $Y(9.46)$  is higher than the two-quark jet data just off resonance excludes 'black box' models in general, where the  $b\bar{b}$  system turns into two light quarks as sketched below.

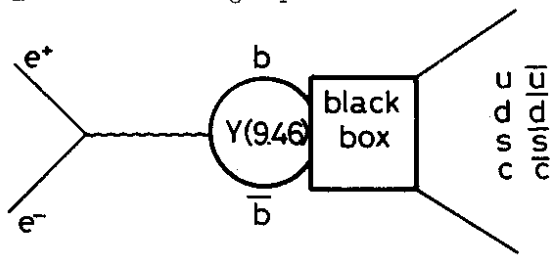


Fig. 18: Possible transitions of the  $b\bar{b}$  resonance into low mass quark pairs.

The success of the three-gluon model using two-jet quantities calls for a three-jet analysis. The PLUTO collaboration has used a three-jet quantity (Triplcity) defined analogous to collinear thrust<sup>28</sup>. The final state hadrons with the momenta  $\vec{p}_1 \dots \vec{p}_N$  are

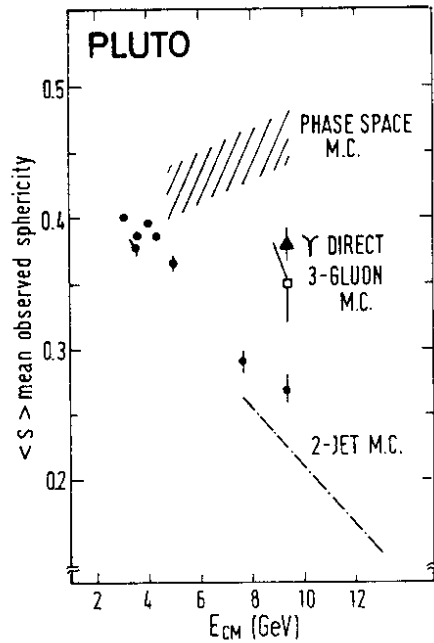


Fig. 17: Average sphericity calculated from charged particles only  $\sqrt{s}$  energy, ref.25.

grouped into 3 non-empty classes  $C_1, C_2, C_3$  with the total momenta  $\vec{P}(C_i) = \sum \vec{p}_i, i = 1, 2, 3, i \in C_i$ . Triplcity  $T_3$  is then defined by

$$T_3 = \max \left\{ \left| \sum_{C_1} p_i \right| + \left| \sum_{C_2} p_i \right| + \left| \sum_{C_3} p_i \right| \right\} / \sum_{i=1}^3 |p_i|.$$

The limiting values are

$$3 \cdot \sqrt{3}/8 < T_3 < 1.$$

Momentum conservation puts the three-jet vectors into a plane. One gets two equivalent Dalitz plots having either  $x_1 < x_2 < x_3$  the three-jet energies or the  $\theta_1 < \theta_2 < \theta_3$ , the angles between the three jets (see fig. 8). The ordering of  $x_i, \theta_i$  reduces the large triangle to the small one, fig. 19, and the characteristic event structures are indicated for the three corners of the triangles<sup>29</sup>.

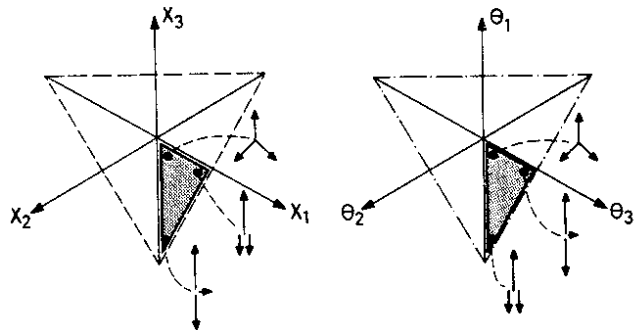


Fig. 19: Event structures in  $x_i$  and  $\theta_i$  as defined in fig. 8

Projections of the event density in the triangles on  $x_1, x_3$ , Fig.20 and on  $\theta_1, \theta_3$ , fig. 21, are compared to the three-gluon model. The agreement again is very good. For comparison phase space is shown (dotted line) again to demonstrate that phase space is still not very much separated from the three-gluon model at this low energy although quantitatively it is not a description for the  $Y(9.46)$  data.

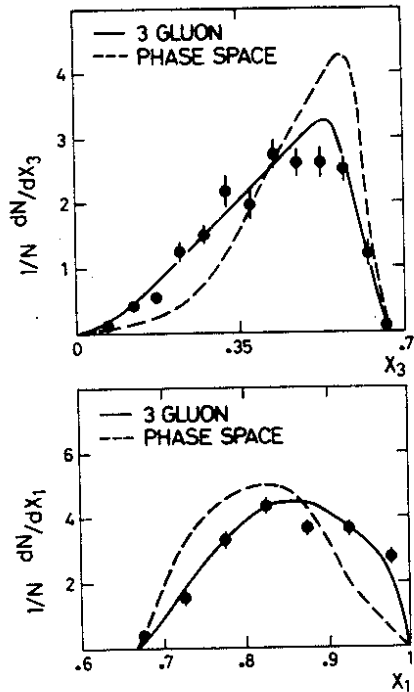


Fig. 20: Results of the triplicity analysis for  $x_1, x_3$  on  $Y(9.46)$  and compared to the three-gluon model (full line) and phase space (dashed line).

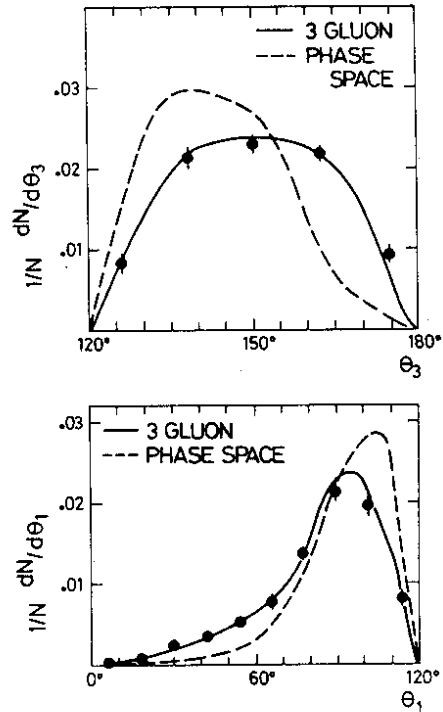


Fig. 21: Results of the triplicity analysis for  $\theta_1, \theta_3$  on  $Y(9.46)$ , otherwise as in fig. 20.

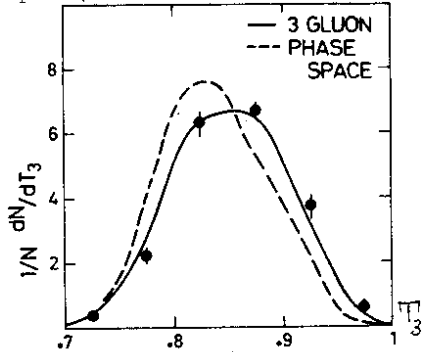


Fig. 22: The triplicity distribution on  $Y(9.46)$ .

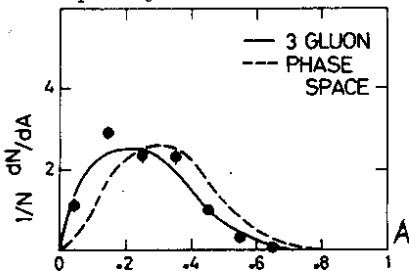


Fig. 23: The aplanarity distribution on  $Y(9.46)$ .

The triplicity ( $T_3$ ) distribution alone does not give a very strong discrimination against phase space, see fig. 22. This holds true also for the aplanarity<sup>28</sup>, fig. 23.

Table 3 shows average values for three jet quantities for the PLUTO analysis of the  $Y(9.46)$  in compari-

son to the phase space and the three-gluon model. The off resonance data has also been analysed resulting in sufficiently different distributions in  $\theta_1, \theta_3$  and  $x_1, x_3$  (not shown)<sup>28</sup>.

Model Comparison

Table 3

|                            | Y Direct data               | 3-Gluon MC                  | Phase Space MC              |
|----------------------------|-----------------------------|-----------------------------|-----------------------------|
| $\langle T \rangle$        | $0.715 \pm 0.004$           | $0.712 \pm 0.003$           | $0.671 \pm 0.003$           |
| $\langle T_3 \rangle$      | $0.858 \pm 0.002$           | $0.850 \pm 0.002$           | $0.838 \pm 0.002$           |
| $\langle x_1 \rangle$      | $0.855 \pm 0.004$           | $0.853 \pm 0.003$           | $0.819 \pm 0.002$           |
| $\langle x_2 \rangle$      | $0.722 \pm 0.004$           | $0.724 \pm 0.003$           | $0.700 \pm 0.002$           |
| $\langle x_3 \rangle$      | $0.423 \pm 0.006$           | $0.422 \pm 0.005$           | $0.481 \pm 0.004$           |
| $\langle \theta_1 \rangle$ | $84.1^\circ \pm 1.0^\circ$  | $85.5^\circ \pm 0.8^\circ$  | $93.2^\circ \pm 0.6^\circ$  |
| $\langle \theta_2 \rangle$ | $125.6^\circ \pm 0.7^\circ$ | $124.3^\circ \pm 0.5^\circ$ | $122.9^\circ \pm 0.4^\circ$ |
| $\langle \theta_3 \rangle$ | $150.3^\circ \pm 0.6^\circ$ | $150.2^\circ \pm 0.5^\circ$ | $144.0^\circ \pm 0.4^\circ$ |

The three-jet analysis does not prove, however, the existence of three-jet structure unambiguously since possibly other more arbitrary final-state configurations can be constructed which numerically agree with the result of this analysis. However, the agreement demonstrated with the numbers summarized in table 3 is most likely not accidental and can be taken as strong support of the three-gluon decay hypothesis for the  $Y(9.46)$ .

### 3. Gluon Spin

The analysis of the  $Y(9.46)$  decay structure reveals quantitative agreement with the three-gluon decay matrix element. It is calculated for massless spin 1 gluons. The data therefore support the spin 1 assignment to the gluon. A natural question then is, can the data be described by spin 0 gluons as well? Koller and Krasemann have provided an expression for a three-gluon

matrix element with spin 0 gluons<sup>30</sup>.

$$\frac{1}{\Gamma} \frac{d\Gamma}{dx_1 dx_2} = \frac{1}{\ln M/A}$$

$$\frac{(x_1 - x_2)(x_3 - x_2)(x_1 - x_2) + \text{even perm.}(1,2,3)}{x_1^2 x_2^2 x_3^2}$$

### SPIN 1 GLUON

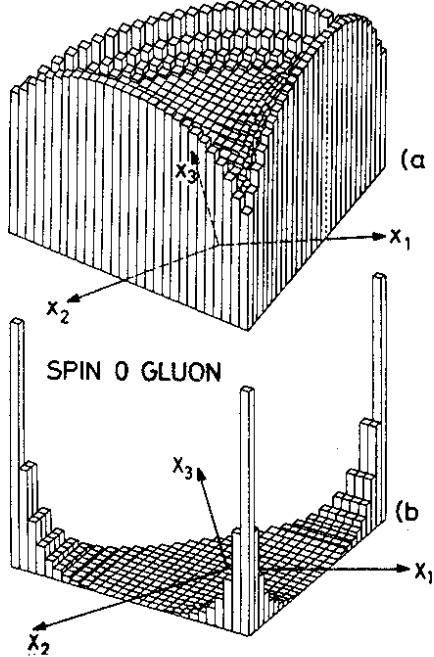


Fig. 24: Graphical representation of the three-gluon matrix element for a) spin 1 and b) spin 0 from ref. 30.

Fig. 24 shows the Dalitz plot density distribution for spin 1 and spin 0. In the spin 0 case there is a peaking in the corners of the triangle. The corners correspond to  $x_1 = x_2$ ,  $x_2 = x_3$  and  $x_2 = x_1$  and the final state is dominantly two gluons of about equal energy and the third one with very little energy. It follows that one should expect a dominant two-jet structure. This is not supported by the data and provides a falsification of the spin 0 case. Spin 1 does, however, not necessarily follow uniquely. Although the true spin 1 gluon matrix element is in good agreement with the data any phase space like distribution of three low mass clusters will give a very similar distribution<sup>31</sup>. The analogue of this statement was already observed by Ore and Powell in their paper<sup>32</sup> on the three-photon decay of positronium.

For a positive spin 1 determination a measurement of the angular distribution of the fastest gluon ( $x_1$  in fig. 8) with respect to the  $e^+$  beam axis is most appropriate<sup>33</sup>. The general form is

$$d\sigma/d \cos\theta = A + B(x_1)\cos^2\theta.$$

$B(x_1)$  is shown in fig. 25. The statistics of the data samples in the DORIS experiments is not sufficient for such a detailed comparison. Integration over the probability distribution for  $x_1$  predicts for the coefficient of  $\cos^2\theta$  a value of 0.39. The spin 0 case, however, gives -0.995<sup>30</sup>. The data is shown in fig. 26<sup>25</sup>. While spin 0 again is excluded by the data only consistency with a value 0.39 can be seen, not an unambiguous determination. It should be remarked here that this analysis is sensitive to the exact value of  $B_{\mu\mu}$ , since the

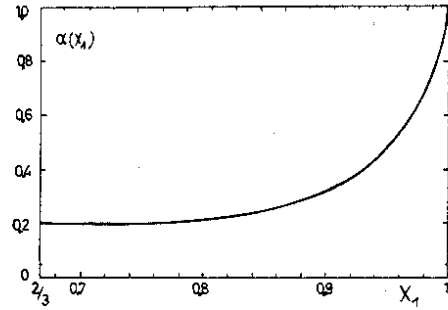


Fig. 25: The coefficient  $B(x_1)$  vs  $x_1$  for a spin 1 three-gluon decay of the  $Y(9.46)$  (see text).

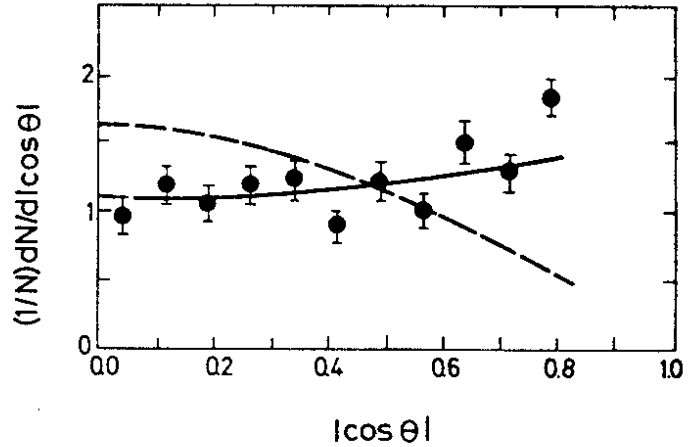


Fig. 26: Data on the angular distribution of the sphericity axis (for charge sphericity) with the beam and compared to the predictions for spin 1 (full line) and spin 0 (dashed line), refs.25 and 30.

continuum and the vacuum polarisation term have a value - 1 for the coefficient of  $\cos^2\theta$ .

While it seems rather unsatisfactory to try to construct a  $Y(9.46)$  decay model based on spin 0 gluons in the first place, it is gratifying to see the one model actually worked out in detail to be falsified<sup>30</sup>. Spin 1 on the other hand remains to be proven.

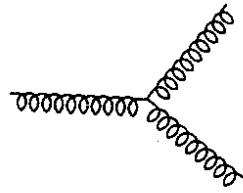
#### 4. Fragmentation Properties of the Gluon

There are several specific features of the gluon that should reflect into the properties of the gluon jets and make them different from quark jets<sup>31</sup>

(a) The gluon should have a softer fragmentation function which means a higher multiplicity than a quark jet, asymptotically by a factor 9/4.

(b) In addition the average  $\langle p_{\perp} \rangle$  of hadrons with respect to the gluon-jet axis should be broader.

Both (a) and (b) are due to the existence in QCD of a gluon selfcoupling, the three-gluon vertex, the most



characteristic feature of QCD. It provides asymptotically for a gluon jet more hard branches in its development and leads to the predictions (a) and (b) above.

#### 4.1. Multiplicities and $p_{\perp}$

Given the existence of the three-gluon coupling qualitatively one should see a rather spectacular increase in  $\langle n \rangle$  going from the continuum (two-quark jets) to the top of a heavy quark-antiquark resonance, both because of the softer fragmentation of the gluons and because of the development of the third-jet. In case of the  $Y(9.46)$  a change in  $\langle n_{ch} \rangle$  has actually been observed by the three DORIS groups. The evidence is shown in fig. 27 from DHHM and fig. 28 from DASPII. Corrected numbers are shown in table 4 from PLUTO and DHHM. The increase is significant, but not very spectacular. In

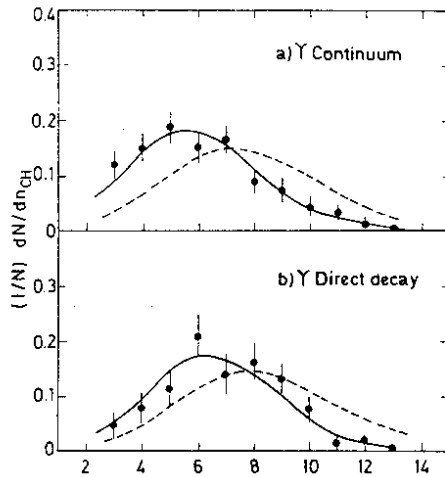


Fig. 27: Observed charged multiplicity  
a) near and  
b) at the  $Y(9.46)$ .  
For comparison the result of the two-jet and three-jet calculations from ref.27.

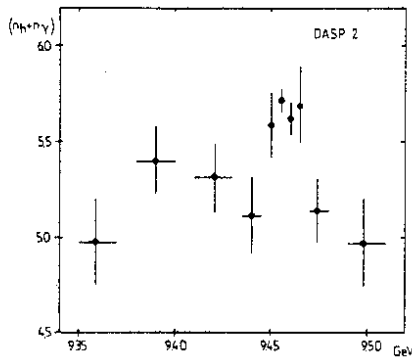


Fig. 28: Observed particle multiplicity in the  $Y(9.46)$  region from ref. 35.

the framework of the three-gluon model, assuming no difference in the gluon jet and the quark jet multiplicities from the average energy of the three different gluon jets (see fig. 8) one gets

$$\langle n_{ch} \rangle_{Y(9.46)} = \frac{1}{2} \sum_i n_{ch}(2 E_i)$$

where  $n_{ch}(2 E_i)$  is the quark multiplicity at two times the gluon energies  $E_i$ . Evaluated using the data compiled in ref. 34 one expects  $\langle n_{ch} \rangle_{Y(9.46)} = 7.6 \pm 0.5$  in agreement with the data in table 4. The observations on  $Y(9.46)$  are thus consistent with equal fragmentation properties of quark and gluon jets and the development of a third jet in support of the three-gluon jet decay model.

A fragmentation model for the gluon where the gluon decays to a quark pair which then fragment independently would lead to higher multiplicity ( $\sim 10$ ) for  $Y(9.46)$ <sup>36</sup>. This option is, therefore, not supported by the multiplicity data.

Larger  $p_{\perp}$  with respect to the gluon-jet axis will lead to a bigger  $\langle p \rangle_{out}$  measured against the plane of the three-gluon events compared to the prediction based on the  $p_{\perp}$  of quark jets. The numbers obtained are  $0.132 \pm 0.003$  for the  $Y(9.46)$  data and  $0.140 \pm 0.006$  for the gluon<sup>25</sup>. No increase over the quark-jet value is observed.

Table 4

| Exp.    | ref. | continuum     | $Y(9.46)$     | $Y'(10.02)$   | $Y(9.46)_{cont.}$ |
|---------|------|---------------|---------------|---------------|-------------------|
| DHHM    | 27   | $6.1 \pm 0.2$ | $6.9 \pm 0.2$ |               | 1.13              |
| DASPII  | 35   | 5.15          | 5.88          |               | 1.14              |
| PLUTO   | 25   | $6.3 \pm 0.4$ | $8.0 \pm 0.3$ |               | 1.27              |
| DHHM    | 27   | $5.9 \pm 0.2$ |               | $7.5 \pm 0.6$ | 1.27              |
| Expect. | 36   | 6.2           |               | 7.6           | 1.23              |

#### 4.2. Charged Particle Momentum Spectra

Since different processes for particle production on and off resonance - three-gluon jets versus two-quark jets - have been proposed a comparison of (charged) particle momentum spectra can give further details on the two processes.

The data are shown in fig. 29 and fig. 30. There are some differences seen (fig. 29) not unexpected since the multiplicities are different. The relative particle yields (fig. 30), however, do not show strong variations.

A momentum analysis of charged particle spectra from heavy quarkonium decays was proposed as a sensitive test for the presence of the three-gluon coupling<sup>37</sup>. Only two resonances, the  $J/\psi(3.1)$  and the  $Y(9.46)$  are available so far. The measured momentum spectra at  $J/\psi(3.1)$  are shown in fig. 31a and are compared to the nearby continuum data ( $E_{CM} = 3.6$  GeV), fig. 31b. They are not very different while the data on and off resonance at 9.46 are certainly different, see fig. 29.

The momentum spectra from  $J/\psi(3.1)$  and  $Y(9.46)$  corrected for detector effects are shown in Fig. 32. A moments analysis following ref. 37 has been performed, the results for ratios of moments for  $A = 0.5$  GeV, a value preferred by inclusive lepton scattering analysis<sup>38</sup>, the data require a  $A$  of  $\sim 0.7$  GeV, see fig.33.

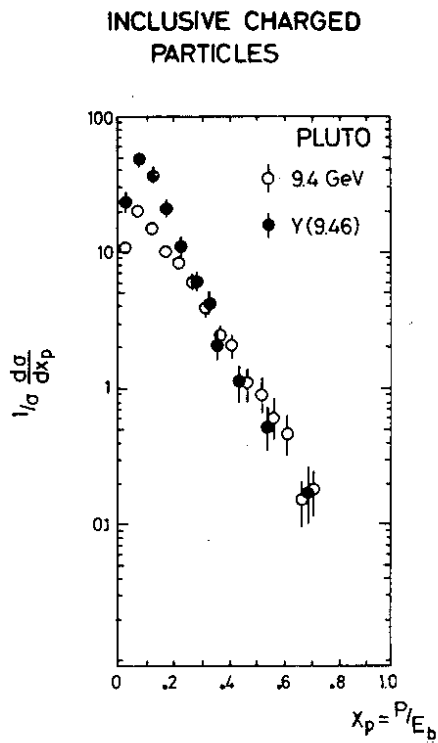


Fig. 29: Charged particle momentum spectra for the Y(9.46) and the nearby continuum.

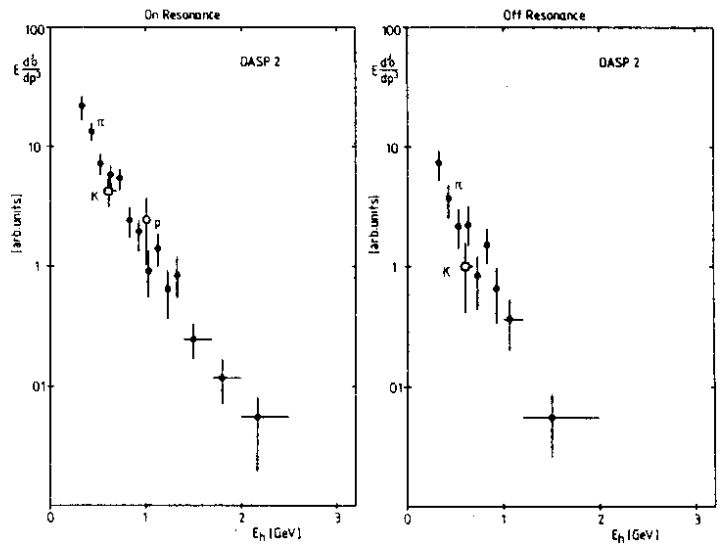


Fig. 30: The momentum spectrum of charged particles for the Y(9.46) and the continuum, for pions  $\pi$ , kaons  $K$  and  $p$  from ref. 35.

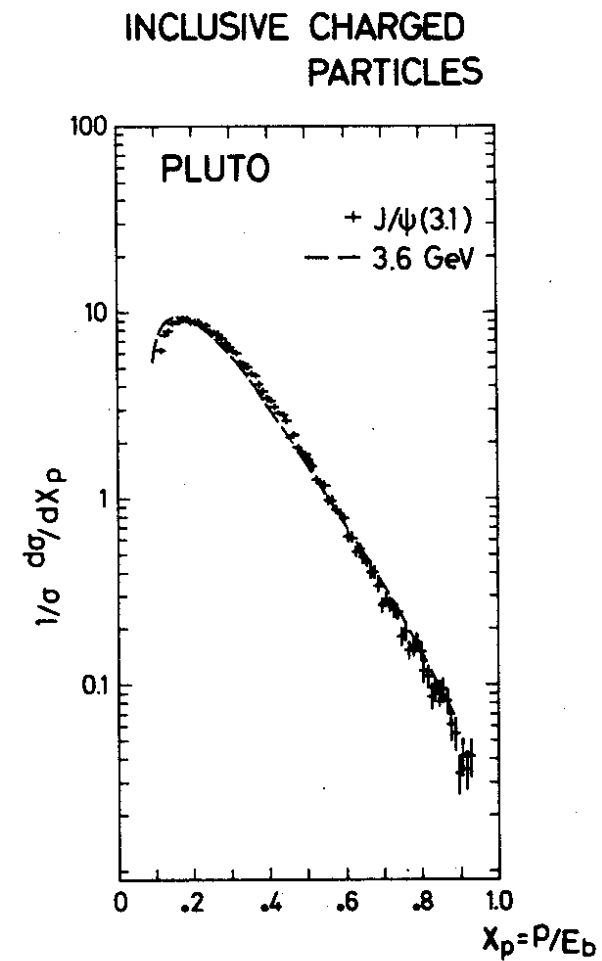
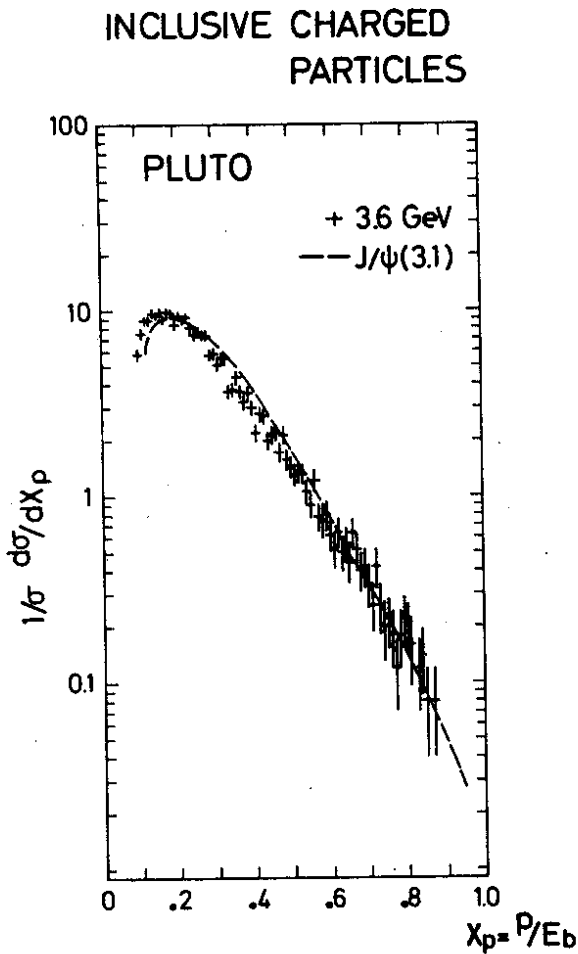


Fig. 31: Inclusive charged particle spectra  
 a) at the J/ψ(3.1) and  
 b) at 3.6 GeV.

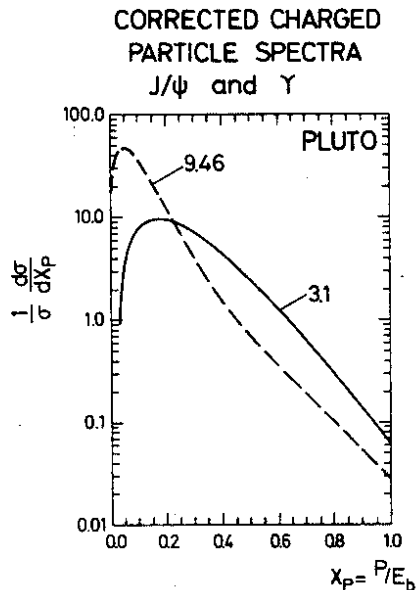


Fig. 32: Corrected data on inclusive charged particle momentum spectra on  $J/\psi(3.1)$ , full line, and  $Y(9.46)$ , dashed line.

Certainly the  $J/\psi(3.1)$  resonance is not a good reference point for such an analysis. The jet structure is not at all apparent at this low energy, see for example the sphericity values in fig. 17 where both the data from the  $J/\psi(3.1)$  and the off resonance data (3.6) are still close to phase space, this means energy and momentum conservation dominate the final state once the multiplicity is fixed. At  $Y(9.46)$  differences show up which makes it most likely a good reference point once a higher resonance (toponium) were found. Still qualitatively an inspection of fig. 33 shows the 'QCD' trend for the data. This sheds some light on the difficulties of performing a QCD analysis based on the energy dependence of particle momentum spectra since the range of validity for such an analysis is not well defined. In  $e^+e^-$  annihilation the comparison on and off resonances at least provides some basis for a judgement on this question.

#### 4.3. Higher resonances and Three-Jet Structure

Finally a remark about a resonance at higher energies, e.g. a 'toponium' at 30 GeV. The triplicity distribution expected using the successful three-gluon model for the  $Y(9.46)$  discussed is shown in fig. 34a together with phase space at 30 GeV in comparison to the situation at  $Y(9.46)$ , fig. 34b. The difference is impressive and indicates that very significant gluon-jets studies can be done on such a resonance.

Concerning the existence proof of three-jet events with the help of energy-flow diagrams<sup>39</sup> (e.g.  $q\bar{q}g$  events at 30 GeV<sup>40</sup>) I want to show an energy-flow diagram of phase-space events at 30 GeV where triplicity has been used to give an orientation for each event, fig. 35. A 'three-jet' structure is apparent. One must conclude that just the existence of a three-arm structure does not prove the three-jet structure even at this high energy, since as we see in the most trivial particle correlation the analysis can impose the three-jet structure one is looking for.

#### MOMENTS - ANALYSIS $J/\psi(3.1)$ and $Y(9.46)$

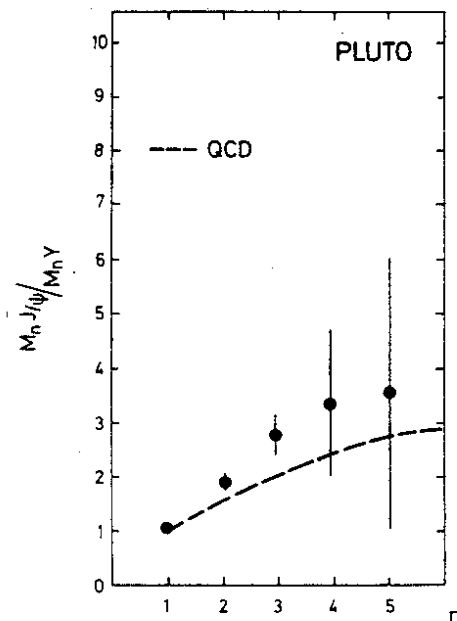


Fig. 33 The ratio of the moments of the charged particle spectra calculated from the fits to the PLUTO data shown in fig. 32. The dashed line shows the prediction of ref. 37 for  $\Lambda = 0.5$  GeV.

#### 30 GeV MODEL

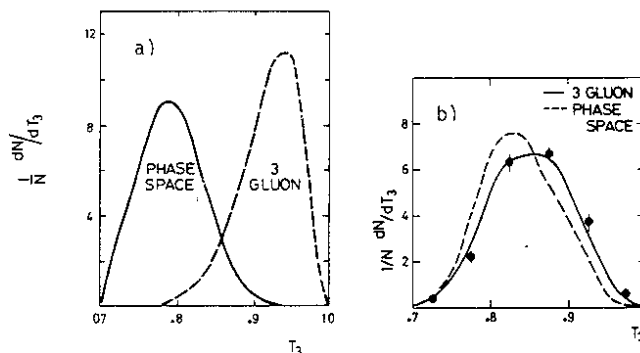


Fig. 34: A calculation of the triplicity distribution for  
a) a hypothetical resonance at 30 GeV (dashed line) as compared to phase space (full line) and  
b) compared to  $Y(9.46)$

#### ENERGY - FLOW 30 GeV

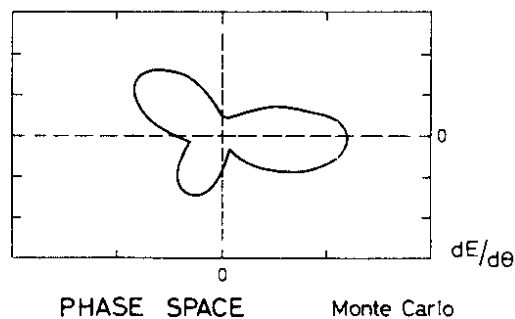


Fig. 35:

### Conclusions

(1) The constituent quark of the  $Y(9.46)$  is most likely the b-quark with charge  $1/3$ . A final proof has to come from the observation of open b-flavour and its decay modes.

(2) The hadronic decay of the  $Y(9.46)$  is consistent with a three-jet structure and the lowest order three-gluon decay matrix element. This is based on the two- and three-jet analysis of the  $Y(9.46)$  data.

(3) A 3 spin 0 gluon decay of the  $Y(9.46)$  is ruled out.

(4) Gluon jets are very similar to quark jets at jet energies  $\leq 5$  GeV.

(5) A resonance at high energies (toponium family) would be a great place for definite QCD tests.

### Acknowledgement

I am very grateful to Fermilab for the warm hospitality extended to me. Many thanks to my scientific secretary, R. Wagner, for his very efficient help. Also I want to thank H. Harari for critical and useful comments. I want to thank G. Flügge and E. Lohrmann for reading the manuscript.

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