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by

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The Formation of Excited  $\infty_2^+$  Ions by Vacuum Ultraviolet Radiation

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Excited  $\varpi_2^+$  ions are formed by irradiating  $\varpi_2$  molecules with the synchrotron radiation. The excitation function of  $\varpi_2^+$  ( $\tilde{B}^2 \Sigma_u^+$ ) shows a pronounced structure which can be explained by an interaction of the Rydberg states  $np\pi_u$  converging to  $\tilde{C}^2 \Sigma_g^+$  with the continuum of  $\tilde{B}^2 \Sigma_u^+$ .

The formation of excited  $\operatorname{CO}_2^+$  ions by photon molecule reaction is of special interest in the physics of planetary atmospheres<sup>1)</sup>. In order to study the excitation mechanism we have used the synchrotron radiation as a primary light source. The light intensity of the synchrotron at the exit slit of the first monochromator amounted to  $10^9$  photons/Å sec at about 600 Å. (This value is related to a beam current of the accelerator of 30 mA). Towards 400 Å the intensity drops down to approximately  $10^8$  photons/Å sec<sup>2)</sup>. The wavelength resolution was about 3.5 Å. The pressure within the collision chamber was varied between  $10^{-3}$  torr and  $5 \cdot 10^{-2}$  torr.

The fluorescence radiation was observed with a second monochromator which had a wavelength resolution of about 20  $\mathring{\Lambda}$  in the region from 2000  $\mathring{\Lambda}$  to 4000  $\mathring{\Lambda}$ . The signal was detected with a cooled multiplier (type EL1 6256 S).

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Fig. 1 shows the emission spectrum of  $CO_2$  which was obtained at a primary photon energy of 19 eV. The radiation components can be classified according to ref. 3. The vibrational structure with  $\Delta v$  constant is not resolved. For example the component designated by 2,0,0 - 0,0,0 contains the transitions 3,0,0 - 1,0,0 , 4,0,0 - 2,0,0 and 5,0,0 - 3,0,0. A similar picture was obtained by measuring the spectrum of the Martian atmosphere<sup>1)</sup>. Fig. 2 gives the excitation function of the transition  $\tilde{B}^2 \Sigma_u^+ - \tilde{X}^2 \Pi_g$ at about 2900 Å. We have found a sharp structure at 679 Å and 662.5 Å. The electron configuration of the ground state of  $CO_2$  can be written as

KKK 
$$(\underline{(\sigma_g)^2}, (\sigma_u)^2, (\sigma_g)^2, (\sigma_u)^2, (\sigma_u)^4, (\pi_g)^4, \mathbf{\Sigma}_g^+$$

The removal of a nonbonding  $\sigma_u$  electron of type b results in the state  $\tilde{B}^2 \Sigma_u^+$  of the ion. On the other hand one gets a Rydberg series converging to the state  $\tilde{C}^2 \Sigma_g^+$  if a  $\sigma_g$  electron of the same type is raised to a np $\pi_u$  orbital<sup>5)</sup>.

This apparent emission series is well known from absorption measurements<sup>6</sup>. Since the window structure at 679 Å and 662.5 Å in the excitation function of  $\tilde{B}^2 \Sigma_u^+$  coincides, with respect to energy, with the states  $4p\pi_u$  and  $5p\pi_u$  we conclude that it is due to an interaction of the Rydberg states with the continuum corresponding to  $\tilde{B}^2 \Sigma_u^+$ . This interpretation is strengthened by the fact that we have found a similar structure in N<sub>2</sub>0 (to be published). A window structure in N<sub>2</sub>0 was found by ref. 7 who measured the total fluorescence with four filters. We would like to thank our colleagues of the Synchrotron Radiation Group at DESY for valuable help in our experimental work. Especially we are indebted to the Munich Group for making available to us their monochromator.

#### References

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## Figure Captions

- Fig. 1 Emission spectrum of  $OO_2^+$  obtained by irradiating  $OO_2$  with 19 eV photons. The dotted part of the spectrum is taken from another measurement.

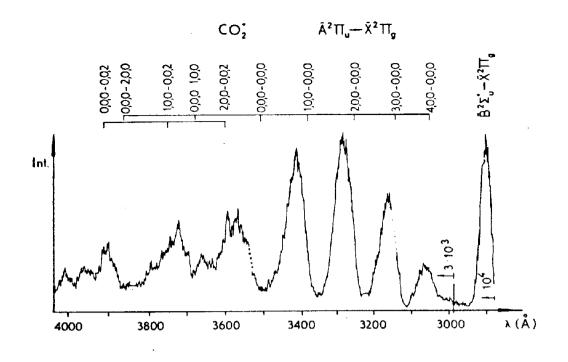


Fig. 1

