

An Yb-fiber frequency comb phase-locked to microwave standard and optical reference*

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We present a fully stabilized Yb-fiber frequency comb locked to a microwave standard and an optical reference separately. The carrier-envelope offset frequency is generated by a standard $f-2f$ interferometer with 40 dB signal-to-noise ratio. The offset frequency and the repetition rate are stabilized simultaneously to the radio frequency reference for more than 30 hours, and the fractional Allan deviation of the comb is the same as the microwave standard of 10^{-12} at 1 s. Alternatively, the comb is locked to an ultra-stable optical reference at 972 nm using an intracavity electro-optic modulator, exhibiting a residual integrated phase noise of 458 mrad (1 Hz–10 MHz) and an in-loop tracking stability of 1.77×10^{-18} at 1 s, which is significantly raised by six orders comparing to the case locked to the microwave frequency standard.

Keywords: Yb-fiber, frequency comb, low-noise, optical reference

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1. Introduction

Low-noise optical frequency combs are desired for many scientific applications, such as optical atomic clocks, high-resolution spectroscopy, fundamental physical constant measurement, and attosecond science.^[1–6] Over the past two decades, several strategies were used to build optical frequency combs, such as Ti:sapphire laser combs, all-solid-state laser combs, fiber-laser combs, electro-optic combs, and microresonator combs.^[7–11] In particular, fiber-laser frequency combs have become one of the most widely used optical frequency combs due to their robustness, low-cost, and long-term stability. Previous works have shown the great potential of fiber-laser optical frequency combs.^[12–14] In terms of average power, a fiber-laser comb with submillihertz linewidth and higher than 10 W average power was achieved based on an Yb-fiber system in 2008^[15] and then the average power was increased to 80 W employing the linear chirped-pulse amplification.^[16] Several effective noise suppression measures have been implemented, promoting the fiber-laser comb to the same noise level as that of the Ti:sapphire laser comb.^[17,18] Moreover, with all-fiber structure and compact design, the fiber frequency comb can operate outside the well-controlled optical laboratory or even in the space environment for a long time.^[19,20]

Femtosecond mode-locked fiber lasers are the key element of the fiber optical frequency comb and directly affect the stability and accuracy of the total system. Early experiments

indicated that fiber lasers have large phase noise introduced by the high gain high loss cavity dynamics and high nonlinearity during the propagation in the fiber.^[21] Despite daunting technical challenges, the past decade has witnessed remarkable progress in different fiber mode-locking schemes.^[12,22] Among these laser sources, the fiber laser using nonlinear polarization evolution (NPE) technique is a good choice due to its broader spectra output and low phase noise generation. With the help of intracavity grating pairs, zero net cavity group-delay dispersion (GDD) could be easily realized for getting the lowest noise and the narrowest linewidth for the carrier-envelope offset (CEO) frequency (f_{ceo}).^[23]

Frequency combs have millions of optical modes, whose frequencies can be determined by the equation in the frequency domain: $f_n = n f_{\text{rep}} + f_{\text{ceo}}$, where n is the integer of mode number. The repetition rate, f_{rep} , represents the spacing between the comb teeth. The carrier-envelope offset frequency, f_{ceo} , determines the absolute position of the comb structure. According to this formula, a frequency comb has only two degrees of freedom. Generally, f_{rep} and f_{ceo} are stabilized to a radio frequency (RF) reference and the frequency stability is limited to the accuracy of the RF standard and phase-locked loop (PLL). For further improving the frequency comb stability, it is necessary to replace the RF reference with an optical reference. In this way, f_n is well stabilized by locking the heterodyne beat frequency between the comb tooth and optical reference with an intracavity electro-optic modulator

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(EOM). The EOM is used as a fast actuator for cavity length fine modulation, which can provide MHz feedback bandwidth with a fixed point near direct current (DC). The combination of a high-bandwidth EOM and a low-bandwidth, long-range piezoelectric transducer (PZT) actuated mirror allows for tight stabilization of the frequency comb to an optical reference over a large dynamic range.^[24]

In this paper, we demonstrate a home-built 250-MHz Yb-fiber frequency comb locked to the radio frequency standard and the 972 nm ultrastable continuous wave (CW) laser separately. The NPE mode-locked Yb-fiber laser output is launched into three branches to detect the repetition rate, the CEO frequency, and the heterodyne beat frequency between the comb and the CW laser, respectively. Both f_{rep} and f_{ceo} are phase locked to the radio frequency standard over 30 hours without any temperature control. Furthermore, we show the scheme of frequency comb phase locked to an ultrastable CW laser with Hz level linewidth. In this case, the integrated phase noise of 458 mrad is obtained from 1 Hz to 10 MHz integrating range. The in-loop frequency instability is measured to be 1.77×10^{-18} at 1 s gate time, which represents six orders of

magnitude improvement compared with phase locking to the radio frequency standard.

2. Phase locking to the RF standard

Figure 1 shows the schematic of the stabilized 250-MHz Yb-fiber frequency comb. The seed source is a home-built femtosecond Yb-fiber ring oscillator based on NPE mode-locking technique, which consists of two major sections. The fiber section provides gain and nonlinearity, while the free space section compensates for the dispersion and rotates the polarization to achieve mode-locking. A bulk EOM is placed in the cavity for locking the frequency comb to an optical reference. The oscillator directly delivers near zero dispersion pulses at 250-MHz repetition rate with 150-mW average output power. The measured output spectrum is centered at 1045 nm with 51 nm bandwidth (full width at half maximum, FWHM), corresponding to the transform-limited pulse duration of 32 fs, as shown in Fig. 2(a). For suppression of atmosphere fluctuation, acoustic noise, and mechanical vibration, the oscillator is set in an aluminum sheet metal box with foam on the side for insulation and sound attenuation.

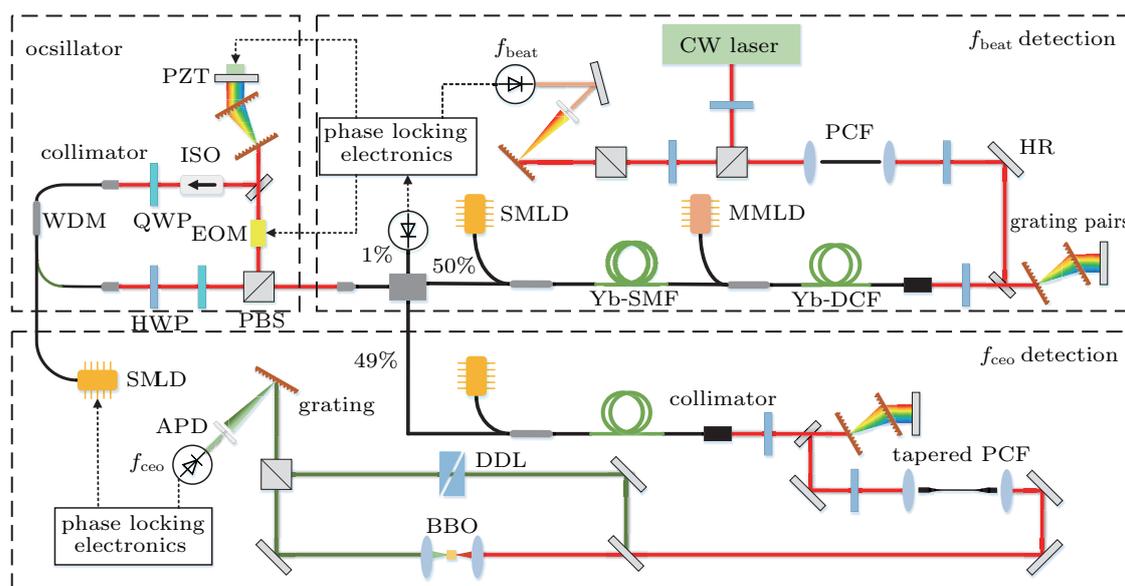


Fig. 1. Schematic of the experimental setup. SMLD, single-mode laser diode; WDM, wavelength division multiplexer; HWP, half waveplate; QWP, quarter waveplate; ISO, isolator; EOM, electro-optic modulators; PZT, piezoelectric transducer; SMF, single mode fiber; DCF: double cladding fiber; PCF, photonic crystal fiber; PBS, polarization beam splitter; DDL, dispersion delay line; APD, avalanche photodetector.

The pulses are split by a fiber power splitter to three replicas with 1% for detecting the repetition rate, 49% for spectral expansion and f_{ceo} signal detection, and 50% for obtaining the heterodyne beat frequency to phase lock to the optical reference. The f_{ceo} signal is usually obtained by self-referencing technique. To overcome the relatively low power, we use a one-stage fiber amplifier that amplifies the 30 mW laser output from the 49% port to 300 mW and then compress the pulses to ~ 80 fs pulse width using a pair of transmission grating pairs, as shown in Fig. 2(b). The compressed pulse is then launched

into a tapered photonic crystal fiber (PCF) for supercontinuum generation. The tapered PCF is based on a commercial PCF, which has a 3.6- μm mode field diameter with zero dispersion around 1030 nm. By tapering this fiber, the pitch of the fiber is decreased to 2 μm and its zero-dispersion wavelength is blue-shifted. The total length of the tapered fiber is 14 cm with 6 cm taper transition. Octave-spanning spectrum broadening is obtained when coupling the 200 mW pulses into the tapered PCF. Figure 3(a) shows the octave-spanning spectrum spanning from 500 nm to 1200 nm. The intensity of short wave-

length and long wavelength is high enough to make it easy to generate f_{ceo} signal. A free-running f_{ceo} signal with 40 dB signal to noise ratio (SNR) in a 100 kHz resolution bandwidth (RBW) is obtained by a standard $f-2f$ interferometer, as shown in Fig. 3(b).

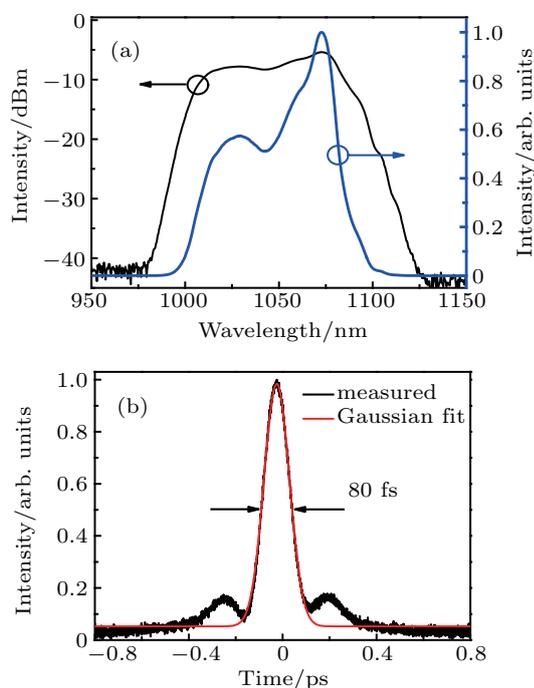


Fig. 2. (a) Measured oscillator output spectrum in linear (blue) and log (black) scales. (b) Measured autocorrelation traces of the compressed pulses after one-stage amplifier.

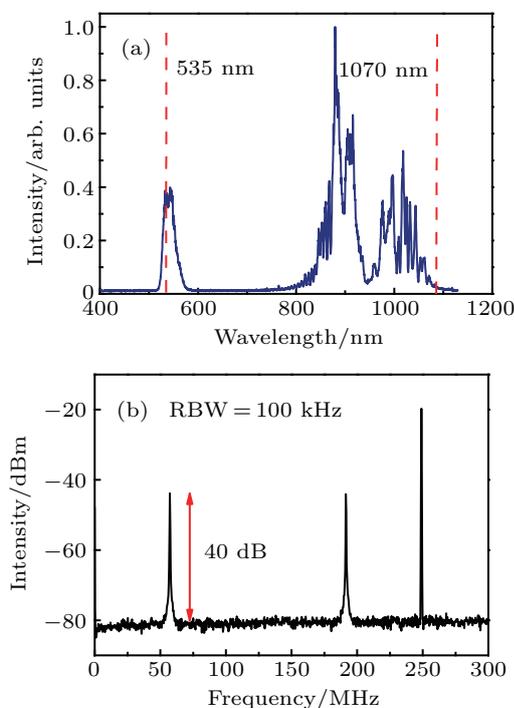


Fig. 3. (a) Octave-spanning spectrum after the tapered photonic crystal fiber. The two dotted lines indicate the spectral components used for generating the f_{ceo} signal. (b) Detection of free-running f_{ceo} signal in 100 kHz resolution bandwidth.

Phase locking to the RF standard is achieved by two

phase-locked loops electronics, one loop acting on the PZT to stabilize f_{rep} and the other loop controlling the pump LD current to lock f_{ceo} . Both references are provided by an RF synthesizer synchronized to an Rb clock.

For f_{ceo} stabilization, the beat signal detected by an avalanche photodetector passes through a bandpass filter and mixes with the RF reference at 20 MHz to get an error signal. This error signal is then processed by a proportion-integration-differentiation (PID) and used to stabilize the f_{ceo} signal by adjusting the pumping current of the Yb-fiber oscillator. Figure 4(a) shows the long-term shift and stability of the f_{ceo} frequency, which operated fully-stabilized without phases slip for 30 hours until the measurement was stopped. The standard deviation of the locked f_{ceo} signal is 6 mHz. The relative fractional Allan deviation of the signal normalized to 20 MHz and to the optical frequency is calculated to be 2.8×10^{-10} and 1.99×10^{-17} at 1 s averaging time, as shown in Fig. 4(c), which is limited by the accuracy of the synthesizer used in the PLL.

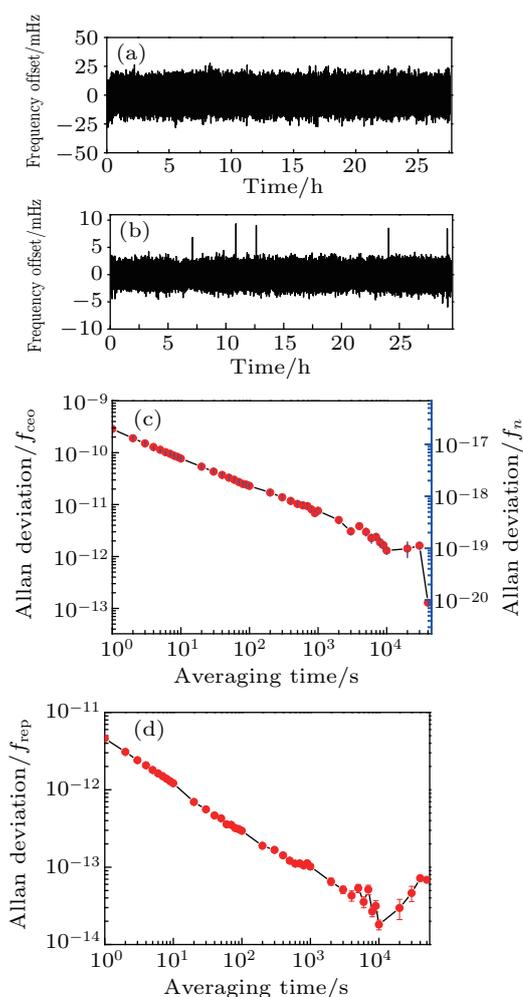


Fig. 4. Long-term frequency shift and Allan deviation of the comb locked to an RF reference. (a) Residual fluctuations of the f_{ceo} signal with a 1 s gate. The standard deviation for f_{ceo} is 6 mHz. (b) Residual fluctuations of the f_{rep} signal with a 1 s gate. The standard deviation for f_{rep} is 1 mHz. (c) Allan deviation normalized to the 20 MHz standard and optical frequency. (d) Allan deviation of the repetition rate.

Synchronously, the repetition rate is phase locked by controlling the PZT in the oscillator, which has a resonant frequency of 110 kHz and a full dynamic range of 9.1 μm . Here we select the fourth harmonic of the repetition rate as the input signal to improve the locking accuracy. In the experiment, the f_{rep} signal was also stabilized for more than 30 hours without any temperature control. The long-term frequency shift curve of the locked f_{rep} signal is shown in Fig. 4(b). The standard deviation of the locked f_{rep} signal is 1 mHz with an in-loop tracking stability of $4.46 \times 10^{-12} \text{ s}^{-1}$ (normalized to the repetition rate of 250 MHz). We have verified that the true feedback bandwidth of the PZT in the experiment is about 100 Hz, which is limited by the resonant frequency of the PZT and the phase-locked loop bandwidth. For the high-frequency noise suppression, it is necessary to introduce fast actuators into the cavity, such as EOM.

3. Phase locking to the optical reference

To further improve the stability of the Yb-fiber optical frequency comb, we stabilize the comb to an optical standard by using an intracavity EOM. A 972 nm CW laser, which is stabilized to an F-P cavity with cavity finesse measured to be 208600, is used as an ultra-stable optical reference. The CW laser power for beat signal detection is about 100 mW. The stability between the comb and the CW laser is accomplished by phase locking the heterodyne beat frequency between one of the comb modes and the CW laser to a microwave reference and feedback to the intracavity EOM.

While EOM is designed to modulate the phase, polarization, or amplitude of light propagating through it, we are only interested in using it to modulate the phase. According to the basis of Pockel's effect (linear electro-optic effect), the index of refraction of the crystal along the optical axis is a function of the applied electric field. The index change of EOM in the cavity means the change of the total cavity length. Compared with the PZT, the intracavity EOM has higher feedback bandwidth but smaller travel range. Therefore long-term phase locking to the optical standard needs the combination of high-bandwidth intracavity EOM and low-bandwidth, long-range PZT. In the experiment, a homebuilt EOM made by a 7-mm MgO-doped lithium niobate crystal was used as a fast actuator. We chose the short LiNO_3 crystal in order to reduce the effect of high-order dispersion while increasing the corresponding voltages. In the experiment, the bandwidth of the electronics for f_{rep} stabilization in optical frequency is about 1 MHz and the voltage range applied to the EOM is $\pm 15 \text{ V}$.

Since the output spectrum of the oscillator does not contain the 972 nm wavelength component, we need to expand the spectral coverage. Pulses generated from the 50% port

of the oscillator output are first stretched to 7.2 ps and then amplified in a two-stage fiber amplifier. The first stage amplifier is constructed with 30 cm single-clad gain fiber (Yb-125 from Coractive) and pumped with a 700 mW single-mode butterfly pump diode. The second stage is built using 1.2 m double-clad Yb-doped fiber (Yb10-128P from Coractive) and pumped by a multimode pump diode. With about 5 W pump power, we obtain 2 W average output power and then a compressor consisting of two identical fused-silica transmission diffraction-gratings (1000 lines/mm) de-chirps the amplified pulses to 135 fs. The average output power after the transmission grating pair is about 1.2 W, corresponding to 4.8 nJ single compressed pulse energy. These compressed pulses are then launched into a 23 cm length photonic crystal fiber with a zero-dispersion wavelength around 1000 nm to produce a supercontinuum spectrum as shown in Fig. 5(a). The maximum output power after the PCF is 620 mW with a high intensity at 972 nm wavelength, which benefits to obtaining high SNR heterodyne beat note signal (f_{beat}). Figure 5(b) shows the heterodyned beat note signal detected by an avalanche photo-detector with a SNR of 42 dB at 100 kHz RBW, which supports long-term electronic locking.

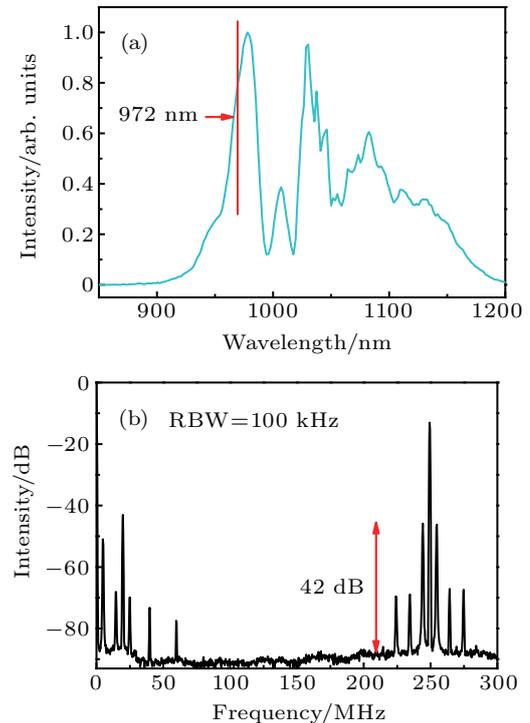


Fig. 5. (a) Spectrum expansion for 972 nm after the photonic crystal fiber. (b) Detection of free-running f_{beat} signal in 100 kHz resolution bandwidth.

Phase-locking the comb to an optical frequency reference other than an RF standard can increase the frequency stability and narrow the linewidth. In the experiment, the error signal was feed backed to the combination of intracavity EOM and PZT for long-term stability. The phase-locking results of f_{beat}

are shown in Fig. 6. Once the lock is achieved, the linewidth of f_{beat} becomes extremely narrow (RBW = 1 kHz) and two large peaks are obtained at around 320 kHz, which shows the actual bandwidth of the EOM in the servo loop. The relative narrowed linewidth of the controlled f_{beat} signal is measured to be 1.24 mHz, as shown in the inset of Fig. 6(a). The measured power spectral density (PSD) and integrate noise

of phase (IPN) of the locked f_{beat} are presented in Fig. 6(b). The integrated phase noise from 1 Hz to 10 MHz is about 458 mrad. Furthermore, we investigate the long-term stability of the locked f_{beat} signal. The standard deviation of the frequency offset is 0.54 mHz, corresponding to the relative frequency stability of 1.77×10^{-18} at an averaging time of 1 s and improved to 4.4×10^{-20} with an averaging time of 1000 s.

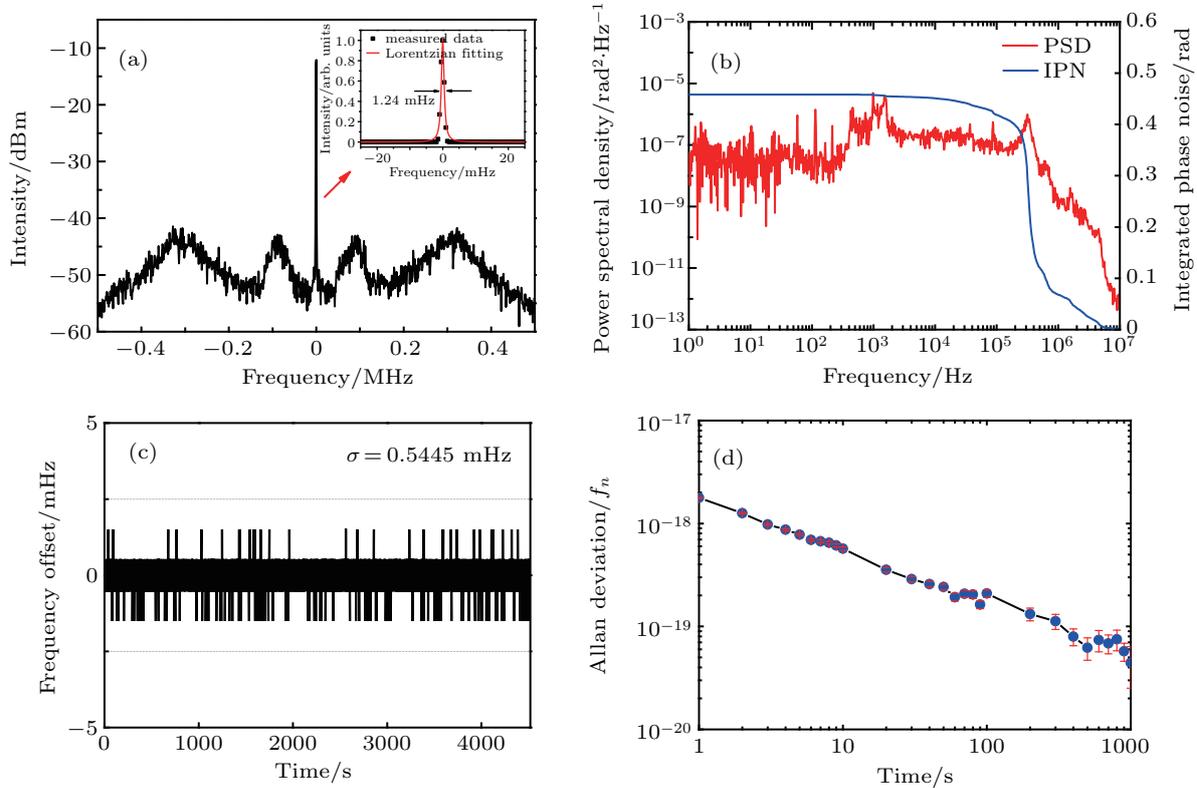


Fig. 6. (a) Frequency spectrum of stabilized heterodyne beat frequency under 1 kHz RBW. The inset: linewidth of the heterodyne beat frequency. (b) Phase noise PSD of f_{beat} (red line) and the IPN (blue line) ranging from 1 Hz to 10 MHz. (c) Residual fluctuations of the f_{beat} signal with a 1 s gate. The standard deviation for f_{beat} is 0.54 mHz. (d) Allan deviation of the f_{beat} signal.

4. Conclusions

We have demonstrated a fully phase locked fiber frequency comb based on a homebuilt 250-MHz NPE-mode-locked Yb-fiber laser oscillator, in which the RF standard and the ultrastable 972 nm laser are selected as the reference. The oscillator delivered 150-mW ultrashort pulses centered at 1045 nm with 51 nm bandwidth. A free running f_{ceo} with a SNR of ~ 40 dB in a 100 kHz resolution bandwidth was obtained by the $f-2f$ interferometer. Both f_{rep} and f_{ceo} were phase locked to the RF standard over 30 hours without any temperature control. Furthermore, phase locking to an optical reference was achieved by employing the intracavity EOM as a fast actuator. The integrated phase noise is approximately 458 mrad from 1 Hz to 10 MHz. The experimental result also shows the phase stability with Allan deviation of 1.77×10^{-18} at 1 s averaging time. Further studies will be devoted to improve the reliability and compactness of this optical frequency

comb and make it a powerful tool for dual-comb research.

References

- [1] Udem T, Holzwarth R and Hänsch T W 2002 *Nature* **416** 233
- [2] Schliesser A, Picqué N and Hänsch T W 2012 *Nat. Photon.* **6** 440
- [3] Picqué N and Hänsch T W 2019 *Nat. Photon.* **13** 146
- [4] Kolachevsky N, Fischer M, Karshenboim S G and Hänsch T W 2004 *Phys. Rev. Lett.* **92** 033003
- [5] Ludlow A D, Zelevinsky T, Campbell G K, Blatt S, Boyd M M, Miranda M H G, Martin M J, Thomsen J W, Foreman S M, Ye J, Fortier T M, Stalnaker J E, Diddams S A, LeCoq Y, Barber Z W, Poli N, Lemke N D, Beck K M and Oates C W 2008 *Science* **319** 1805
- [6] Porat G, Heyl C M, Schoun S B, Benko C, Dörre N, Corwin K L and Ye J 2018 *Nat. Photon.* **12** 387
- [7] Matos L, Kleppner D, Cuzuko O, Schibli T R, Kim J, Ippen E P and Kaertner F X 2004 *Opt. Lett.* **29** 1683
- [8] Pekarek S, Südmeyer T, Lecomte S, Kundermann S, Dudley J M and Keller U 2011 *Opt. Express* **19** 16491
- [9] Hundertmark H, Wandt D, Fallnich C, Haverkamp N and Telle H R 2004 *Opt. Express* **12** 770
- [10] DelHaye P, Schliesser A, Arcizet O, Wilken T, Holzwarth R and Kippenberg T J 2007 *Nature* **450** 1214
- [11] Carlson D R, Hickstein D D, Zhang W, Metcalf A J, Quinlan F, Diddams S A and Papp S B 2018 *Science* **361** 1358

- [12] Kim J and Song Y 2016 *Adv. Opt. Photon.* **8** 465
- [13] Li Y H, Kuse N, Rolland A, Stepanenko Y, Ranzewicz C and Fermann M E 2017 *Opt. Express* **25** 18017
- [14] Luo D P, Liu Y, Gu C L, Wang C, Zhu Z W, Zhang W C, Deng Z J, Zhou L, Li W X and Zeng H P 2018 *Appl. Phys. Lett.* **112** 061106
- [15] Schibli T R, Hartl I, Yost D C, Martin M J, Marcinkevicius A, Fermann M E and Ye J 2008 *Nat. Photon.* **2** 355
- [16] Ruehl A, Marcinkevicius A, Fermann M E and Hartl I 2010 *Opt. Lett.* **35** 3015
- [17] Newbury N R and Swann W C 2007 *Josa B* **24** 1756
- [18] Pang L H, Han H N, Zhao Z B, Liu W J and Wei Z Y 2016 *Opt. Express* **24** 28993
- [19] Sinclair L C, Coddington I, Swann W C, Rieker G B, Hati A, Iwakuni K and Newbury N R 2014 *Opt. Express* **22** 6996
- [20] Lezius M, Wilken T, Deutsch C, Giunta M, Mandel O, Thaller A, Schkolnik V, Schiemangk M, Dinkelaker A, Kohfeldt A, Wicht A, Krutzik M, Peters A, Hellmig O, Duncker H, Sengstock K, Windpassinger P, Lampmann K, Hulsing T, Hänsch T W and Holzwarth R 2016 *Optica* **3** 1381
- [21] Cingöz A, Yost D C, Allison T K, Ruehl A, Fermann M E and Ye J 2011 *Opt. Lett.* **36** 743
- [22] Chen H W, Chang G, Xu S, Yang Z and Kärtner F X 2012 *Opt. Lett.* **37** 3522
- [23] Nugent-Glandorf L, Johnson T A, Kobayashi Y and Diddams S A 2011 *Opt. Lett.* **36** 1578
- [24] Hudson D D, Holman K W, Jones R J, Cundiff S T and Ye J 2005 *Opt. Lett.* **30** 2948