

# Evaluation of inertia effects in planar squeeze flow inside soft, porous layers

G Lupu<sup>1</sup>, P Turtoi<sup>2</sup>, T Cicone<sup>2</sup>

<sup>1</sup> Military Equipment and Technologies Research Agency, Aeroportului 16, Clinceni, Ilfov, 077025, Romania

<sup>2</sup> Department of Machine Elements & Tribology, University POLITEHNICA of Bucharest, Spl. Independenței 313, Bucharest, 060042, Romania

**Abstract.** The resistance to flow of a fluid squeezed out from a soft and porous structure under compression, generates high load capacity. The process is dependent on permeability, which is variable during compression of the porous layer. This mechanism (named Ex-Poro-Hydrodynamic Lubrication – XPHD) plays an important role in a wide range of tribological applications (biolubrication, shock absorbers, etc.). A literature search reveals that Darcy's law is intensively used for modeling flow through porous media. However, for high porosity materials and relatively high velocity flows that may occur during impact, the inertia effects could be significant. The full squeeze solution of high velocity flow inside deformable porous materials including inertia effects and permeability variation with the level of compression could not be found in the literature. A numerical study for planar squeeze flow of a Newtonian liquid inside a soft porous layer, based on the Darcy flow model including Forchheimer correction, is presented in this paper. The porous material is compressed at high speed between two perfectly rigid, parallel and impermeable discs. Finite differences technique is used to solve the system of equations. Pressure distribution, load and Reynolds pore based number variation during compression are obtained and compared with similar values predicted by simple analytical models based on Darcy law. The results define the limits of Darcy approach in XPHD lubrication.

## 1. Introduction

Lubrication with fluid film is the main topic of numerous studies with various applications. A particular application based on fluids imbibed inside a highly deformable porous media is emerging and there are studies that prove higher load capacity. Load is generated by the resistance to flow inside the porous matrix with its variable (decreasing) permeability during compression. This mechanism (named Ex-Poro-Hydrodynamic Lubrication – XPHD [1]) shows great potential for squeeze dampers. For squeeze flow inside porous materials, Darcy law is widely used. However, for high velocity flow inertia effects could be important. It was already proven that Darcy law is less accurate when pore based Reynolds number  $Re_p$  becomes greater than unity [2]÷[4]. There is no well established limit or transition zone between laminar to turbulent flow regime. According to Zeng & Grigg [5] the transitional flow occurs for  $Re_p=1÷100$ , while Boomsma & Poulikakos [6] proposed a  $Re_p=5$ .

First proposal for modification of Darcy law by introducing a quadratic term was made by Dupuit [7], and Forchheimer [8] by analogy with Navier-Stokes equations. Presently the equation proposed by Ward [9] and Beavers et al. [10], based on a series of experimental results, is widely used:



$$\frac{dp}{dr} = -\frac{\eta}{\phi}u - \frac{C_f \cdot \rho}{\sqrt{\phi}}u^2 \quad (1)$$

where  $C_f$  is a nondimensional Forchheimer coefficient, experimentally determined,  $\phi$  is material permeability,  $\eta$  dynamic viscosity,  $\rho$  density. Some studies tried to find a universal value for this coefficient. Ward [9] stated that  $C_f=0.55$  can be used for all permeable materials. Later experimental results reported that  $C_f$  is different for various classes of materials. Beavers and Sparrow [11] found  $C_f=0.078 \div 1.32$  for porous media in the form of latticework of metallic fibers and water. For particles beds, Schwartz and Probst [12] used  $C_f=0.26$ . Other studies tried to find a connection between  $C_f$

and porosity. Ergun and Orning [13] proposed for gases that  $\frac{C_f}{\sqrt{\phi}} = b'' \frac{(1-\varepsilon)}{\varepsilon^3}$  where  $b''$  is a factor of proportionality. According to Irmay [14], for any fluid, the dimensionless coefficient is defined as  $C_f = \alpha \cdot (\beta \cdot \varepsilon^3)^{-1/2}$ , and later, Bear [3] suggested:  $\alpha=0.6$  and  $\beta=180$ . Thus  $C_f=0.045 \cdot \varepsilon^{-3/2}$ .

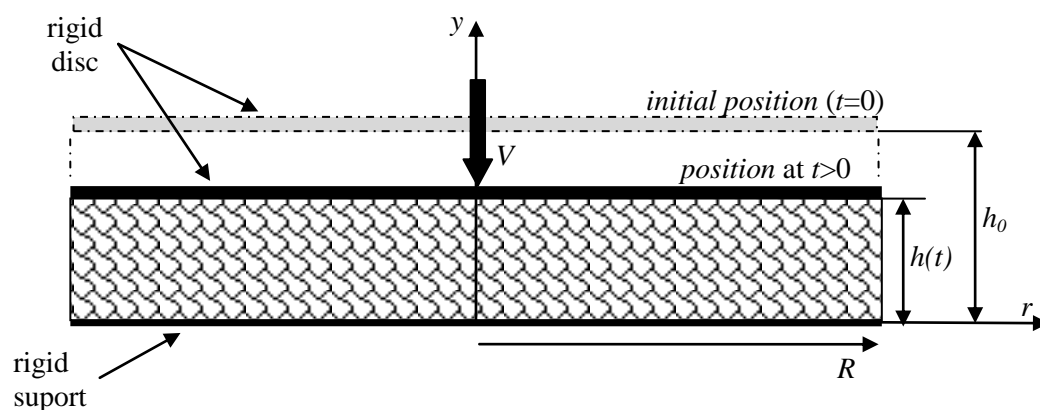
Despite a relatively great number of papers dedicated to inertia effects in porous media, no attempt could be found for the case of a highly deformable media with time variable porosity and permeability.

This study is focused on a particular application of XPHD lubrication for damping devices subjected to medium impact speed, when inertia effect can be significant. It is of major interest to evaluate the limits of Darcy flow model when squeezing speed is high. In the first stage, the analysis is performed for constant speed squeeze in a disc-on-plane configuration.

## 2. Problem formulation and main assumptions

A simple case of squeeze of an imbibed porous material between two circular parallel discs is considered for this study. The upper rigid disc approaches with constant speed  $V=dh/dt$  and compresses the material with initial thickness  $h_0$ . The porous material has the same radius  $R$  as the rigid discs. The model is axisymmetric with radial flow and atmospheric pressure on the outer boundaries. The geometry of the model is presented in figure 1.

The porous material is highly deformable (compression forces generated by the solid structure are neglected), homogeneous and isotropic. A reticulated structure with open pores, saturated with liquid, is considered. The fluid is Newtonian and the flow is laminar, isothermal and isoviscous.



**Figure 1.** The geometry of the axisymmetric planar squeeze flow

The equation of mass conservation in cylindrical coordinates  $(r,u)$  is:

$$\frac{1}{r} \frac{d}{dr}(ru) + \frac{dv}{dy} = 0 \quad (2)$$

Using Kozeny-Carman equation [15] permeability-porosity correlation is:

$$\phi = \frac{D \cdot \varepsilon^3}{(1 - \varepsilon)^2} \quad (3)$$

where  $D$  is a complex parameter, function of the characteristic size of the porous structure.

The material deformation is produced only in the normal direction of the porous layer and the conservation of the solid fraction is used  $h \cdot (1 - \varepsilon) = \text{ct}$ . Thus introducing dimensionless layer thickness  $\bar{h} = h / h_0$  equation (3) becomes:

$$\phi = \frac{D \cdot (\bar{h} - 1 + \varepsilon_0)^3}{\bar{h} \cdot (1 - \varepsilon_0)^2} \quad (4)$$

The following dimensionless parameters are introduced:  $\bar{y} = y / h_0$ ,  $\bar{v} = v / V$ ,  $\bar{u} = u / V$ ,  $\bar{p} = p \cdot D / (\eta VR)$ . Permeability  $\phi$  is also replaced with dimensionless form (eq. (4)).

The dimensionless form Darcy-Forchheimer (D-F) equation (1) in dimensionless form is:

$$\frac{d\bar{p}}{d\bar{r}} = - \frac{\bar{h} (1 - \varepsilon_0)^2}{(\bar{h} - 1 + \varepsilon_0)^3} \bar{u} - (C_f \cdot K) \frac{1}{\sqrt{\frac{(\bar{h} - 1 + \varepsilon_0)^3}{\bar{h} (1 - \varepsilon_0)^2}}} \bar{u}^2 \quad (5)$$

where  $\bar{y} = y / h_0$ ,  $\bar{v} = v / V$ ,  $\bar{u} = u / V$ ,  $\bar{p} = p \cdot D / (\eta VR)$  and  $K = \rho V \sqrt{D} / \eta$ . If we integrate eq. (2) with respect to radius  $r$ , while considering the boundary conditions (at  $y=0$ ,  $v=0$  and at  $y=h$ ,  $v=V$ ) we obtain eq. (6) in a dimensionless form:

$$\frac{d}{d\bar{r}} (\bar{r} \bar{h} \bar{u}) = - \bar{r} \left( \frac{R}{h_0} \right) \quad (6)$$

### 3. Numerical solution

There is no analytical solution possible for these equations and consequently, a numerical approach based on finite differences method is proposed. Because there is no flow on  $\bar{y}$  direction, the domain is split only in radial direction in small elements. Pressure and radial flow velocity is assumed to be constant across the thickness of the porous layer. An explicit first-order FD scheme with forward derivatives is used for the discretization of the eq. (5) and (6). A one-dimensional grid (along  $\bar{r}$ ) based on  $i=1 \dots N+1$  equally-spaced nodes is used and leads to the following set of discrete equations:

$$\bar{p}_i = \bar{p}_{i+1} + \left( f \cdot \bar{u}_i + C_f K \frac{1}{\sqrt{1/f}} \bar{u}_i^2 \right) / \Delta \bar{r} \quad (7)$$

$$\bar{u}_i = \left( \bar{u}_{i-1} \cdot \bar{r}_{i-1} + \frac{\Delta \bar{r}}{h_k} \cdot \frac{\bar{r}_i + \bar{r}_{i-1}}{2} \cdot \frac{R}{h_0} \right) / \bar{r}_i \quad (8)$$

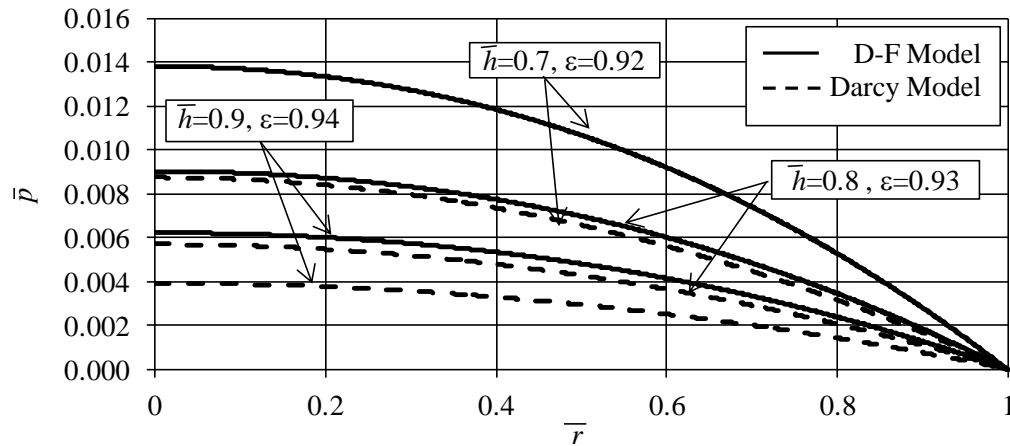
where  $f = \frac{h_k (1 - \varepsilon_0)^2}{(h_k - 1 + \varepsilon_0)^3}$ ,  $\bar{p}_i$  and  $\bar{u}_i$  refer to pressure and flow velocity, respectively, on each node and  $k$  is the subscript that counts for time step.

Gauss-Seidel iterative method without relaxation is used. At each time step pressure distribution is integrated and load calculated. The solution convergence is checked by: the relative variation in force  $\bar{F}$  between two iterations must be inferior to  $10^{-9}$ . Grid independence solution was satisfied for  $N > 1000$ .

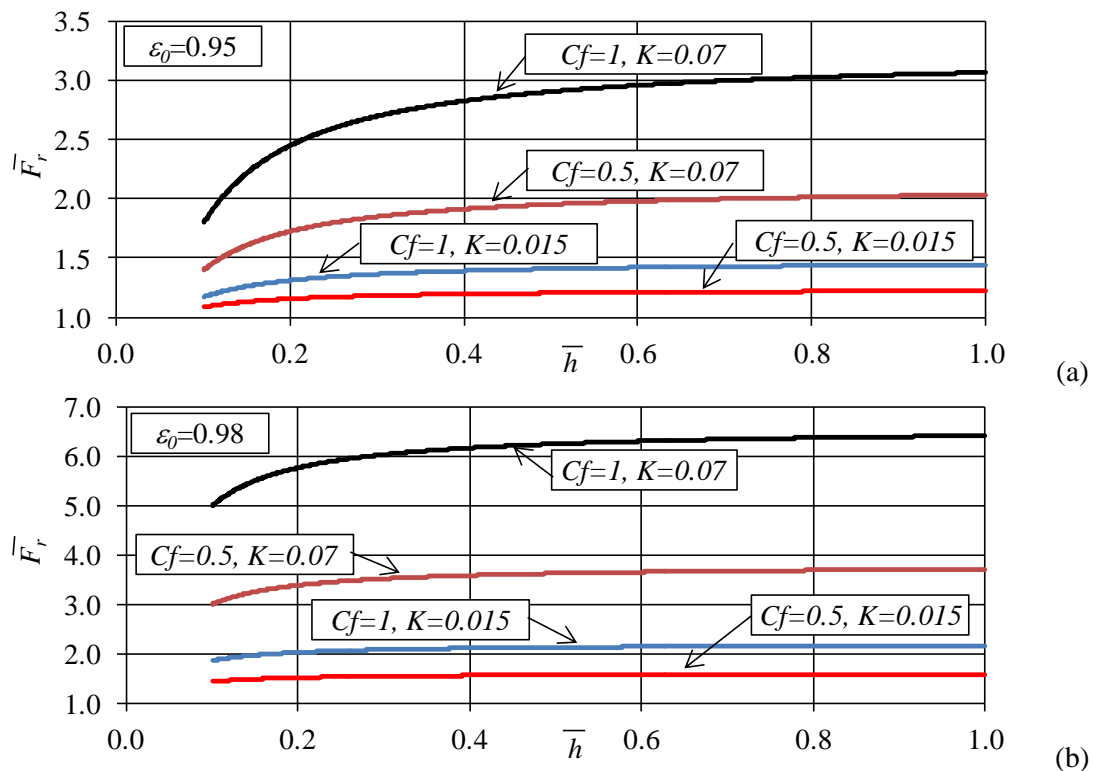
### 4. Results and discussion

The present analysis is focused on a particular application presented in [16]. A high porosity material (knitted fabric made of polyester yarns) and glycerin are considered for impact squeeze. First, an analysis of the values of  $K$  for various conditions reveals a large interval, spanning from  $10^{-3}$  for low fluid velocity and glycerin to  $10^2$  for high fluid velocity and water. For a squeeze speed of 1 m/s the coefficient  $K=0.015$  was found. A second less viscous fluid (e.g. synthetic oil), was considered which, for same speed, yields to a coefficient  $K=0.07$ . The algorithm allows determination of the pressure and

flow velocity distribution with radius when the porous imbibed material is compressed and the fluid is squeezed out including the inertia influence. Dimensionless pressure distribution on the radius of the porous disc during compression for different values of  $\bar{h}$  is presented in figure 2. According to continuity equation, variation of dimensionless flow velocity in radial direction is linear and reaches the maximum value when the pressure gradient is maximum at the outer boundary of the porous disk.



**Figure 2.** Variation of dimensionless pressure during compression for different values of dimensionless thickness  $\bar{h}$  ( $\epsilon_0=0.95$ ,  $K=0.015$ )

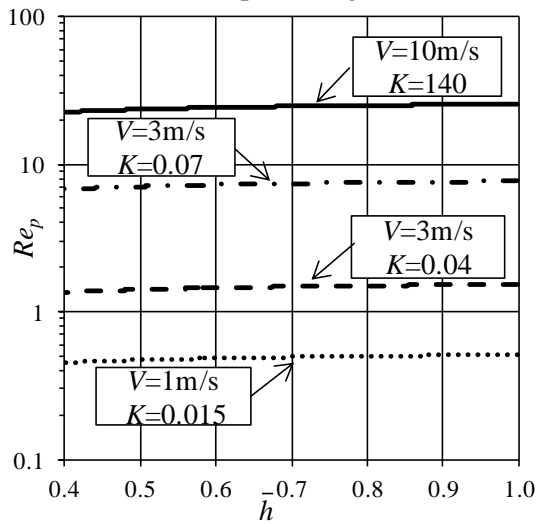


**Figure 3.** Variation of relative load  $\bar{F}_r$  for different values of  $C_f$ ,  $K$  and initial porosity a)  $\epsilon_0=0.95$ , b)  $\epsilon_0=0.98$

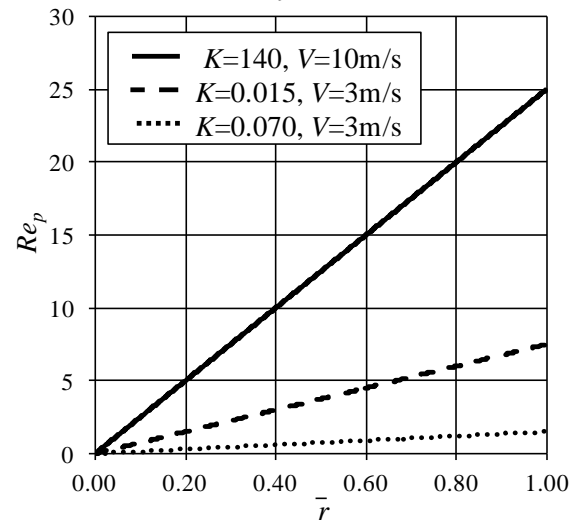
Further, by numerical integration the load  $\bar{F}$  can be determinate. To better understand the influence of inertia effects, relative load  $\bar{F}_r = \bar{F} / \bar{F}(C_f = 0)$  was calculated by dividing the load obtained using the proposed model for different values of  $C_f$  and  $K$  to the load obtained for  $C_f=0$  when no inertia effects are considered. Figure 3 shows the variation of relative load for two different initial

porosities: a)  $\varepsilon_0=0.95$ , b)  $\varepsilon_0=0.98$ . One can see that for low material thickness inertia effects decrease. The drop is accentuated for low thickness due to the variation of porosity in accordance with conservation of solid phase assumption. The load also increases with increasing  $C_f$ . If a less viscous fluid when  $K$  is increasing to 0.07, the inertia effects increase and also the estimated load. Further, to validate the numerical model, the variation of load  $\bar{F}$  was compared at limit, when Forchheimer coefficient  $C_f=0$ , with the analytical solution of the Darcy flow model proposed by Radu [17] for disc on plane squeeze. The difference is very small and can be neglected.

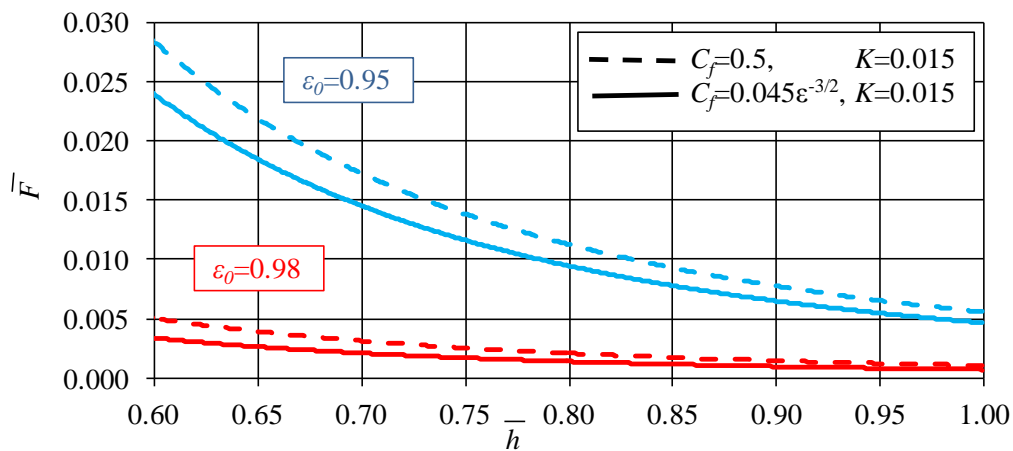
An analysis of pore based Reynolds number variation  $Re_p = \rho u \phi^{1/2} / \eta$  is presented in figure 4 for the maximum value of the flow velocity at the outer boundary of the domain. For a low squeeze speed of 1 m/s and  $K=0.015$  the  $Re_p$  exceeds the laminar flow limit. When considering the same fluid ( $K=0.04$ ), if the speed increases, the  $Re_p$  increases as well. Furthermore, when a less viscous fluid is used,  $Re_p$  exceeds the limit 10. This is in accordance with the assumption that for high squeeze flow inertia must be taken into account. The variation of  $Re_p$  on the radius is presented in figure 5. It can be seen that the variation is linear from centre, as no radial flow occurs, gradually increasing towards the exterior of the disc where the pressure gradient is maximum and also the flow velocity.



**Figure 4.** Variation of  $Re_p$  on compression for highest radial flow velocity  $u_{max}$



**Figure 5.** Variation of  $Re_p$  on the radius of the disc for  $\bar{h}=0.8$



**Figure 6.** Load variation during squeeze

The effect of Forchheimer coefficient  $C_f$  when is variable with porosity according to definition of Irmay [13] was studied in parallel with the case of constant  $C_f$  (figure 6). One can see that when  $C_f$  is variable the predicted force is lower and the difference becomes higher when the material is more compressed. Two cases for different initial porosities  $\varepsilon_0=0.95$  and  $\varepsilon_0=0.98$  are used.

## 5. Conclusions

An evaluation of the inertia effects on force generated during high constant velocity squeezing out of a liquid imbibed in a highly porous and soft layer has been performed. The analysis has been performed numerically using finite differences method.

Pressure distribution was determined and influence of inertia effects on load was evaluated for two values of Forchheimer coefficient. The dimensionless results show that for relatively low squeeze speed inertia effects could be important especially for high porous materials and low viscosity fluids. It is also evident from these results that at high rates of compression (near the end of squeeze process) the inertia contribution vanishes.

An improved numerical solution should consider the dependence of friction factor in Forchheimer term with Reynolds number.

Due to variable porosity specific to XPHD applications, a variable Forchheimer coefficient with porosity can be more appropriate to predict load. Further theoretical work appears necessary in order to evaluate inertia effects in the case of impact squeeze where fluid velocity decline with time due to damping effect.

## 6. References

- [1] Pascovici MD, Cicone T, 2003, Squeeze-film of unconformal, compliant and layered contacts *Tribol Int* **36** 791–799;
- [2] Gartling DK, Hickox CE, and Givler RC, 1996 Simulation of coupled viscous and porous flow problems *Int J Comput Fluid D.* **7** 1 23-48;
- [3] Bear J 1972 Dynamics of Fluids in Porous Materials *American Elsevier*;
- [4] Pearson JRA, Tardy PMJ, 2002 Models for flow of non Newtonian and complex fluids through porous media *J Non-Newton Fluid.* **102** 2 447-473;
- [5] Zeng Z, Grigg R, 2006 A criterion for non-Darcy flow in porous media *Transport Porous Med.* **63** 1 57–69;
- [6] Boomsma K, Poulikakos D, 2002 The effects of compression and pore size variations on the liquid flow characteristics in metal foams *J Fluids Eng* **124** 263–72;
- [7] Dupuit J, Etudes Theoretique et pratiques sur le mouvement des eaux, *Paris : Dunod* ;
- [8] Forchheimer P, 1901 Wassebewegung durch Boden *Z Ver Deutsch Ing* **45** 1782-1788;
- [9] Ward J, 1964 Turbulent flow in porous media *J Hydraul Div am Soc Civ Eng* **90**(HY5) 1-1;
- [10] Beavers GS, Sparrow EM, Rodenz DE, 1973 Influence of bed size on the flow characteristics and porosity of randomly packed of spheres *J Appl Mech* **40** 655-660;
- [11] Beavers GS, Sparrow EM, 1969 Non-Darcy flow through fibrous porous media *J Appl Mech* **36** 4 711-714
- [12] Schwartz J, Probstein RF, 1969 Experimental study of slurry separators for use in desalination *Desalinisation* **6** 239-266;
- [13] Ergun S, Orning AA 1949, Fluid flow through randomly packed columns and fluidized beds *Industrial and Engineering Chemistry* **41** 6 1179-1184;
- [14] Irmay S, 1958 On the theoretical derivation of Darcy and Forchheimer formulas *Eos Trans AGU* **39** 702-707;
- [15] Sheidegger AE, 1974 The physics of flow through porous media. *University of Toronto Press*;
- [16] Cicone T, Pascovici M, Melciu C, Turtoi P, 2019 Optimal porosity for impact squeeze of soft layers imbibed with liquids *Tribol. Int.* **138** 140–149;
- [17] Radu M, 2015 Modelling and simulation of the squeeze process during impact of highly compressible porous layers imbibed with fluids *PhD Thesis University POLITEHNICA of Bucharest (in Romanian)*;