



PAPER

Origin of exceptional magneto-resistance in Weyl semimetal TaSb₂

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E-mail: spatnaik@mail.jnu.ac.in**Keywords:** Weyl semimetal, Magnetoresistance, Kohler's scaling, Topological insulators**Abstract**

We study magneto-transport properties in single crystals of TaSb₂, which is a topological semimetal. In the presence of magnetic field, the electrical resistivity shows onset of insulating behaviour followed by a plateau at low temperature. Such resistivity saturation is generally assigned to topological surface states but we find that aspects of extremely large magneto resistance and resistivity plateau are well accounted by classical Kohler's scaling. In addition, magneto-resistance in TaSb₂ shows non-saturating field dependence. Evidence for anomalous Chiral transport is provided with observation of negative longitudinal magneto-resistance. Shubnikov-de Haas oscillation data reveal two dominating frequencies, 201 T and 455 T. At low temperature, the field dependence of Hall resistivity shows non-linear behaviour that indicates the presence of two types of charge carriers in consonance with reported electronic band structure. Analysis of Hall resistivity implies extremely high electron mobility.

The nascent classification of materials in accordance with the topological states of quantum matter has yielded a radically new technological paradigm [1, 2]. This relates to the discovery of several material systems belonging to topological insulators (TI) and topological semimetals (TSM) that exhibit extremely large magnetoresistance (XMR) [3–9]. Quite generally, the study of electrical resistance in the presence of external magnetic field renders deep insight into electronic transport mechanism [10] that collaterally leads to several technological applications such as in magnetic sensors, magnetic switches and magnetic storage devices. To bring in a prospective, the reported [6, 11–19] magnetoresistance in TSMs could be two orders of magnitude higher than the giant magnetoresistance (GMR) or colossal magnetoresistance (CMR) observed in metallic thin films [20], perovskite manganites [21] or Cr-based chalcogenide spinels [22]. At the core of such exceptional magneto-resistance is the peculiar band structure of three-dimensional TIs and TSMs that yields conducting surface states. Recent studies reveal extremely large magnetoresistance (XMR) in Dirac semimetals (with linear band crossing at the Fermi level e.g. Na₃Bi [23] and Cd₃As₂ [24]) that extends to Weyl semimetals like TaAs family [25, 26]. Such properties are also seen in compensated layered semimetals like WTe₂ [27] and MoTe₂ [28]. The XMR exhibited by topological materials make them very interesting from the prospective of technological applications. However, several recent works in a wide variety of these novel materials have raised a fundamental question; whether the XMR can be explained by classical magneto-resistance theories without considering the topological aspects [15, 29].

In terms of theories for topological insulators, the associated symmetry principles restrict the quantum states to be robust against disorder due to time reversal symmetry invariance [30]. Experimentally this is manifested as a low temperature plateau in electrical resistivity. The plateau at low temperature is understood to have origin in conducting surface states that negate the insulating bulk behaviour, and are protected against backscattering due to time-reversal symmetry (TRS) invariance. The observation of large magneto-resistance in rare-earth monpnictides like R(Sb, Bi) (R = La, Y, Pr etc) [18, 19, 31, 32] and transition-metal dipnictides TM₂ (T = Nb, Ta and M = Sb, As) [33, 34] have been analysed under such perspectives. In particular, TM₂ dipnictides are peculiar in the sense that at zero magnetic field they behave as weak topological insulators but with the application of external field, they can be classified as Type-II Weyl materials [35]. Several recent works have inferred that the plateau in these materials can also be explained by classical magneto-resistance theories without

invoking topological surface states [15, 18, 31, 33, 36, 37]. In this paper, we present a detailed study of magneto-transport behaviour in single crystalline TaSb₂. We observe XMR and metal-insulator like transition at low temperatures under external magnetic field. The resistivity plateau is observed at temperatures below 13 K. Our data on XMR and resistivity plateau is well accounted by Kohler's scaling. Further we observe unambiguous signatures of negative longitudinal magneto-resistance that can be assigned to Weyl phase in TaSb₂ in the presence of external magnetic field. Moreover, clear Shubnikov-de Haas (SdH) oscillations are observed revealing two major Fermi pockets.

Single crystals of TaSb₂ were synthesized by a two-step iodine vapour transport technique. In the first step, polycrystalline TaSb₂ was synthesized by solid-state reaction method. Stoichiometric amounts of Ta powder (3 N, Alfa Aesar) and Sb shots (6 N, Alfa Aesar) were reacted together by heating at 700 °C for 3 days. The polycrystalline sample was then vacuum sealed with Iodine (50 mg cm⁻¹ [3]) in quartz ampoule and put in a tubular furnace with sample at 1000 °C and temperature gradient of ~100 °C across the sealed tube for one week. Shiny needle like single crystals were obtained in the size range 2–6 mm. The crystallinity and structure of the sample were analyzed using single crystal x-ray diffractometer (Bruker D8 Quest single crystal x-ray diffractometer). The crystal structure of the sample was also studied using high resolution transmission electron microscope (HRTEM). Magnetotransport measurements in the temperature range 2–300 K and field range 0–6 T were performed in Cryogenic made cryogen free magnet (CFM) system. SdH oscillation measurements were performed using Cryogenic 14 T Physical Properties Measurement System (PPMS). Single crystal x-ray diffraction data were collected using a microfocus anode (Mo) and a CMOS detector (PHOTON 100).

The XRD data analysis suggests a monoclinic (C_{12/m1}) structure with lattice constants $a = 10.39(\pm 0.03)$ Å, $b = 3.66(\pm 0.01)$ Å, $c = 8.42(\pm 0.03)$ Å, $\beta = 121.38(\pm 13)$ and $V = 273(\pm 2)$ Å³ which are in agreement with the reported data for TaSb₂ [17]. Inset (i) in figure 1(a) shows the HRTEM image of the single crystal. The lattice fringes reveal excellent crystalline nature of the specimen. The inter-planar spacing is ~0.286 nm which corresponds to the orientation of (111) atomic planes of the monoclinic TaSb₂. Main panel of figure 1 shows the temperature dependent longitudinal resistivity $\rho(T)$ of TaSb₂ sample in the absence of external field. The contact geometry used for this measurement is included in the Inset (ii) of figure 1(a). The sample displays metallic behaviour and electrical resistivity decreases monotonically down to 2 K, with $\rho(2\text{ K}) = 0.54\ \mu\Omega\text{-cm}$. Such low residual resistivity value reflects negligible defect scattering in the sample. The residual resistivity ratio ($\text{RRR} = \rho(300\text{ K})/\rho(2\text{ K})$) of the sample is observed to be ~260. Three different crystal of about same size were studied and here we are reporting the data from highest RRR specimen. We note that the magnitude of resistivity is lower than that in good metals and this is clear reflection of role played by topological states where impurity and defect scattering become oblivious to current flow. At low temperatures, $\rho_{xx}(T)$ can be fitted with $\rho_{xx} = \rho_{xx}(0) + AT^n$ with $n \sim 2.56$, where $\rho_{xx}(0)$ is the residual resistivity ($=0.538\ \mu\Omega\text{cm}$) and $A = (4.29 \times 10^{-10}\ \Omega\text{cmK}^{-n})$. Such extremely low resistivity has also been observed in other topological semimetals such as WTe₂⁶ and LaSb⁴⁵. The exponent value n is reflective of dominant scattering mechanism in the system with limiting values $n = 2$ for strong electron-electron scattering and $n = 5$ for conventional electron-phonon scattering processes. The observed value for TaSb₂ would imply dominance of limiting electron-electron scattering as was observed in NbP [38].

An extensively studied aspect of magneto-transport in TSM is the reported negative magnetoresistance when magnetic field is applied parallel to current direction. Such observations are a common feature in several Weyl and Dirac semimetal such as TaAs, Cd₃As₂, and ZrSiS. Dirac semimetals are essentially gapless semiconductors with linear dispersion that become Weyl semimetals when Dirac point splits into two Weyl points due to either spatial inversion or time reversal symmetry breaking. The chiral transport current between the Weyl points is not conserved in this case leading to what is referred to as Adler-Bell-Jackiw (ABJ) anomaly [39]. The experimental manifestation of this is negative magnetoresistance (for $B \parallel I$). Figure 1(b) shows the transverse ($B \perp I$) magnetoresistance at several temperatures with current along a-axis and magnetic field along c-axis. At 2 K and 6 T, the sample shows extremely large magnetoresistance ($\text{MR} = [\rho(B) - \rho(0)] / \rho(0) \times 100$ with $\rho(B)$ and $\rho(0)$ being electrical resistivity at applied B field and 0 field, respectively), with magnitude $\sim 3.55 \times 10^4\%$, without any trace of saturation. With increase in temperature, MR decreases to 300% at 50 K, 6 T. Inset in figure 1(b) compares the field dependence of transverse and longitudinal magnetoresistance (LMR) at 2 K. In the longitudinal configuration, magnetic field and current are applied along the a-axis of the sample. This is figuratively elucidated in the inset of figure 1(b). We observe negative magnetoresistance above ~5.5 T below which a parabolic field dependence is seen. The possible reasons for the negative LMR could be (i) the magnetism in the sample that can be ruled out in TaSb₂, (ii) Improper contact geometry (non-uniform current) may also give rise to negative LMR [40]. Importantly, the specimen was needle-like thus eliminating such possibilities, and (iii) The emergence of Weyl points with application of magnetic field leading to Chiral anomaly in current transport. In agreement with band structure of TaSb₂, we find the negative LMR observation to be a clear signature of ABJ anomaly. A recent theoretical report [35] confirms the possibility of hidden Weyl points in

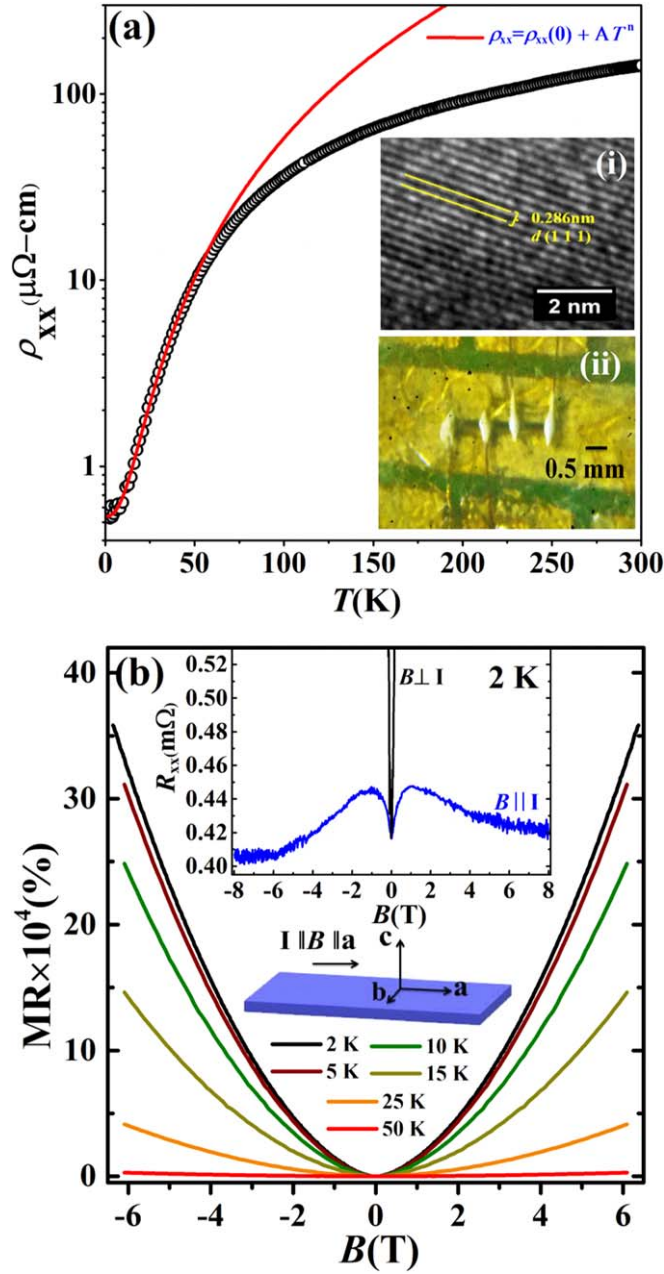


Figure 1. (a) Temperature dependence of resistivity in zero applied magnetic field. Inset (i) shows high resolution TEM image showing the planes (111). Solid line is a fit to the relation $\rho_{xx}(T, B) = \rho_{xx}(0 T) + AT^n$. Inset (ii) shows picture of contact configuration. (b) Field dependence of resistivity at various temperatures. Inset shows the field dependence of transverse ($B \perp I$) and longitudinal resistivity ($B \parallel I$) at 2 K.

TaSb₂ that appear under applied magnetic field resulting in anisotropic chiral anomaly and consequent Type-II Weyl semimetal characterization [41].

In figure 2(a) we show Kohler's scaling with regard to constant temperature field scans. Kohler's rule gives a classical description of electronic motion that can provide insight into the MR behaviour in the sample. According to Kohler's scaling: $MR = \alpha(B/\rho_0)^m$ where α and m are sample dependent constants, B is the applied field strength and ρ_0 is the zero field resistivity. Employing the Kohler's rule: $MR = \alpha(B/\rho_0)^m$, $MR(B)$ data are fitted for all temperatures against $\alpha(B/\rho_0)^{1.78}$ ($\alpha = 1.13 \times 10^4$). Inset (i) of figure 2(a) shows MR as a function of B/ρ_0 (at 2 K) from which α and m were determined. The collapse of MR data at all temperatures to a single line in the Kohler plot (main frame of figure 2(a)) implies that the sample shows similar power law for MR at all temperatures. This means similar scattering mechanism is followed by the carriers at all temperatures. Inset of figure 2(b) shows resistivity $\rho_{xx}(T)$ in applied perpendicular magnetic fields ranging from 0 to 6 T. We note that in the presence of magnetic field, above ~ 100 K, TaSb₂ shows metallic temperature dependence similar to the zero field $\rho_{xx}(T)$ behaviour. In the presence of magnetic field, the resistivity shows sharp upturn and drastic increase at temperatures below ~ 100 K. Moreover, below around 13 K, the resistivity starts to saturate, leading

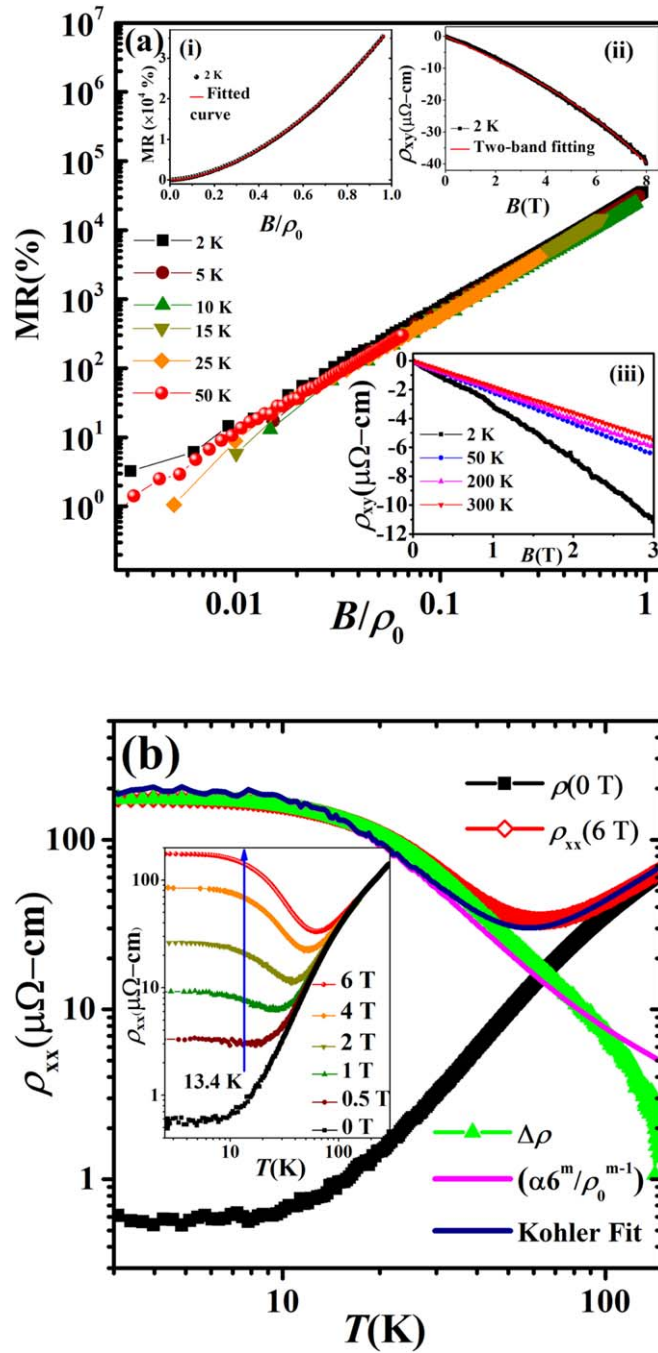


Figure 2. (a) Kohler scaling of MR behaviour, $MR = \alpha(B/\rho_0)^m$. Inset (i) shows the MR fitted with $MR = \alpha(B/\rho_0)^m$ with $m = 1.78$ at 2 K. Inset (ii) Two band model fitting of Hall resistivity at 2 K. Inset (iii) Field dependence of Hall resistivity at temperatures 2 K, 50 K, 200 K and 300 K. (b) Temperature dependence of $\rho_{xx}(0\text{ T})$, $\rho_{xx}(6\text{ T})$ and their difference $\Delta\rho$ are shown. Solid lines are fit to $\alpha B^m/\rho_0^{m-1}$ and Kohler scaling. Inset in (b) shows temperature dependence of ρ_{xx} at various magnetic fields.

to a plateau-like behaviour in the $\rho_{xx}(T)$ curve. The temperature where the sample starts showing saturation remains unchanged at all applied magnetic fields. Such sharp upturn in resistivity below a particular temperature has been seen in several topological semimetals. The associated XMR [7, 31, 39, 42–44] has been widely reported. The microscopic explanation of such phenomena has been attempted with theories that include, (a) magnetic field induced metal to insulator (MIT) transition, (b) electron-hole compensation, and (c) high mobility transport in metallic surface states of topological materials. However, as recently reported for LaSb [29], the upturn in resistivity and its eventual saturation can also be ascribed to classical magnetoresistance theories involving Kohler scaling without invoking topological surface states. Kohler scaling can also be written as: $\rho_{xx}(T, B) = \rho_0 + \alpha B^m/(\rho_0)^{m-1}$. Since ρ_0 is the only temperature dependent term in this equation, the temperature variation of $\rho_{xx}(T, B)$ is mainly governed by ρ_0 . Evidently, the second term ($\Delta\rho = \alpha B^m/(\rho_0)^{m-1}$) and ρ_0 have opposite dependence on temperature, leading to a minimum in total resistivity, $\rho_{xx}(T, B)$, at a

particular temperature. The main panel of figure 2(b) shows the temperature dependent behaviour of $\rho_{xx}(6\text{ T})$, $\rho_0 (= \rho_{xx}(0\text{ T}))$ and $\Delta\rho (= \rho_{xx}(6\text{ T}) - \rho_0)$. The Kohler fitting to $\rho_{xx}(6\text{ T})$ is shown by dark blue line and the fitting parameters α and m are taken as 1.13×10^4 and 1.78 respectively. Evidently, the $\rho_{xx}(6\text{ T})$ in this figure fits well in the entire temperature range down from 2 K to 150 K that includes the plateau region as well. The correct crossover temperature of $\sim 13\text{ K}$ is also verified from Kohler scaling. At low temperature, since ρ_0 is small and independent of temperature variation, it implies $\Delta\rho \gg \rho_0$ and $\rho_{xx}(T, B) \approx \Delta\rho \propto 1/\rho_0^{m-1}$. In summary, the low temperature emergence of plateau and XMR in TaSb₂ can be explained with the help of Kohler's scaling. In conjunction with the observation of ABJ anomaly, we conclude that Kohler's analysis is applicable to Magneto-resistance in Weyl semimetal and there is no apparent contradiction between Kohler scaling and existence of surface transport in TSMs. We emphasize that in the case of Dirac semimetal Cd₃As₂, the magnetic field can lead to a bulk gap due to the magnetic-field-induced rotational symmetry breaking. Recent work by Liu *et al* confirms anisotropic magneto-resistivity in TaSb₂ that imply breaking of two fold rotational symmetry in the presence of magnetic field [41]. There are two key results of our data analysis; 1. TaSb₂ is a Weyl semimetal that is proven by longitudinal negative magneto-resistance and, 2. Kohler's scaling is followed for transverse magneto-resistance data. The case is similar to WTe₂ which is a well known Weyl semimetal [6] that follows Kohler's scaling [45]. The point that we want to emphasize is that Kohler's scaling takes into account magneto-resistance of normal metals that does not require peculiar surface states of topological systems. Therefore, while the edge modes are present in single crystals of TaSb₂, one does not need to invoke this to explain metal-insulator like upturn on application of magnetic field and its eventual saturation at low temperature. Simply said, TaSb₂ is one more example (like WTe₂) that questions the general assignment of topological states to its exceptional magneto-resistance.

The Hall resistivity measurements were performed to investigate the carrier type, concentration and mobility. The field dependence of Hall resistivity ρ_{xy} is shown in the inset (iii) of figure 2(a). The negative sign of ρ_{xy} indicates possible dominance of electronic transport although a large difference in mobilities can also give rise to same result. The Hall resistivity at 2 K is fitted with the two-band model [46] (inset ii) to evaluate the charge carrier concentration and mobility:

$$\rho_{xy} = \frac{B}{|e|} \frac{(n_h \mu_h^2 - n_e \mu_e^2) + (n_h - n_e)(\mu_h \mu_e)^2 B^2}{(n_h \mu_h + n_e \mu_e)^2 + (n_h - n_e)^2 (\mu_h \mu_e)^2 B^2}$$

here $n_e(n_h)$ and $\mu_e(\mu_h)$ are density and mobility of electrons (holes), respectively. The constraint: $\rho_{xx}(B = 0) = 1/e(n_e \mu_e + n_h \mu_h)$ is used to fit the Hall resistivity. The obtained values of electron and hole carrier concentrations are estimated to be $\sim 1.02 \times 10^{18}\text{ cm}^{-3}$ and $1.01 \times 10^{18}\text{ cm}^{-3}$, respectively, revealing TaSb₂ to be a compensated semimetal. From the fitting parameters, the obtained electron and hole mobilities are estimated as $\sim 5.1 \times 10^4\text{ cm}^2\text{V}^{-1}\text{s}^{-1}$ and $1.36 \times 10^4\text{ cm}^2\text{V}^{-1}\text{s}^{-1}$, respectively. In essence, the resistivity upturn and low temperature saturation can be well accounted by strong temperature dependence of the high mobilities of the charge carriers in this compensated semimetal [45].

Further, the magnetoresistance measurements show clear Shubnikov-de Haas (SdH) oscillations at low temperatures and high magnetic fields (figure 3(a)). Inset (i) shows zoomed data between 11 T-13 T. The oscillation component (dR) of the MR is extracted by subtracting a higher order polynomial fit from the high field oscillation data. The oscillations can be identified from the plots of dR versus B^{-1} , as shown in the inset (ii) of figure 3(a). In figure 3(b) we show Fast Fourier transformation (FFT) of the data shown in inset (ii) of figure 3(a). The oscillation frequencies are identified as 201 T (β) and 455 T (γ). From these SdH oscillation frequencies, the extremal cross-sectional area, A_F , of the Fermi surface can be extracted using the Onsager relation: $F = (\Phi_0/2\pi^2)A_F$. The frequency of $F = 201\text{ T}$ corresponds to $A_F = 0.019\text{ \AA}^2$ and $F = 455\text{ T}$ corresponds to $A_F = 0.043\text{ \AA}^2$. Also, we obtain Fermi wave vector, $\kappa_F (= (A_F/\pi)^{1/2}) = 0.078\text{ \AA}^{-1}$ and 0.117 \AA^{-1} corresponding to frequencies 201 T and 455 T, respectively. The FFT amplitude decreases with increasing temperature. In the inset of figure 3(b) we show FFT amplitude with increasing temperature for β and γ peaks. From the temperature damping of SdH oscillation amplitude, the effective quasiparticle mass (m^*) can be extracted. The value of m^* is extracted from the fit of temperature dependence of FFT amplitude with Lifshitz-Kosevich (LK) equation: $\Delta\rho/\rho \propto AT/\sinh(AT)$, where $A = 2\pi^2\kappa_F m^*/\hbar e B$. The LK fit is shown in the inset of figure 3(b). The value of m^* obtained from the LK fit for frequency 201 T is $0.17m_e$ and that for frequency 455 T is $0.15m_e$, where m_e is the free electron mass. The Fermi velocity $v_F = \hbar\kappa_F/m^*$ is estimated to be $5.29 \times 10^5\text{ m s}^{-1}$ and $7.94 \times 10^5\text{ m s}^{-1}$, respectively corresponding to frequencies 201 T and 455 T. The SdH oscillatory component was further analysed using the expression: $dR \propto \cos[2\pi(F/B - \gamma^*)]$ [9], here F is the frequency of oscillation and γ^* is the Onsager phase. From corresponding Landau fan diagram (not shown), the obtained γ^* was estimated to be zero confirming non-trivial Berry phase in TaSb₂ [17].

To summarize, we present a magneto-transport study in the topological semimetal TaSb₂. At 2 K and 6 T a large transverse MR ($= 3.55 \times 10^4\%$) is observed without any sign of saturation. The magnetic field induced turn-on behaviour and plateau like feature at low temperature are well explained with Kohler scaling.

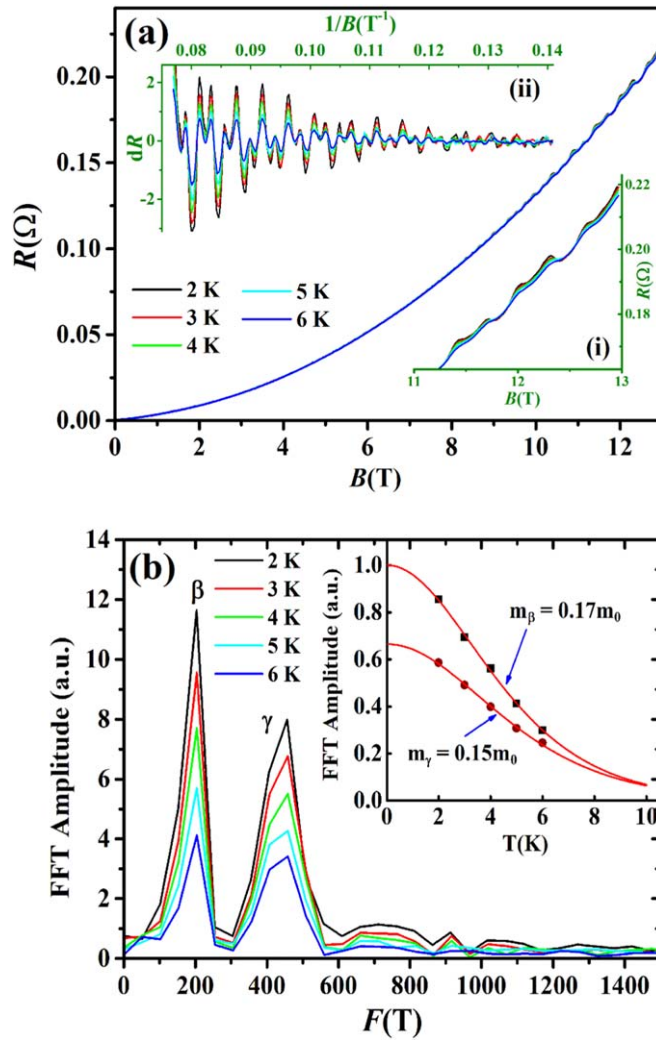


Figure 3. (a) Field dependence of resistivity at low temperatures (2–6 K) in the field range 0–13 T. Inset (i) shows the clear view of the SdH oscillations in the magnetic field range 11–13 T. Inset (ii) shows the oscillatory component (dR) of the resistivity at all temperatures (2–6 K) after subtracting the background. (b) FFT is plotted as a function of frequency at temperatures 2–6 K. Two dominant frequencies are obtained at 201 T and 455 T. Inset shows the FFT amplitude corresponding to frequencies 201 T and 455 T as a function of temperature. Solid lines show the fit of the data with Lifshitz-Kosevich formula.

Significantly we find evidence for negative longitudinal magnetoresistance that signifies the presence of topological Weyl points. Non-trivial Berry phase is also indicated from analysis of SdH oscillation measurements. The high field SdH oscillations show two dominant frequencies at 201 T and 455 T. The Hall measurements confirm compensated semi-metallic behaviour with exceptionally high mobilities. We find that applicability of classical Kohler's scaling to magneto-resistance and low temperature plateau in TaSb₂ is surprisingly accommodative with underlying topological Weyl states.

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