



Estimating the Temperature and Density of a Spicule from 100 GHz Data Obtained with ALMA

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Abstract

We succeeded in observing two large spicules simultaneously with the Atacama Large Millimeter/submillimeter Array (ALMA), the Interface Region Imaging Spectrograph (IRIS), and the Atmospheric Imaging Assembly (AIA) on board the *Solar Dynamics Observatory*. One is a spicule seen in the IRIS Mg II slit-jaw images and AIA 304 Å images (Mg II/304 Å spicule). The other one is a spicule seen in the 100 GHz images obtained with ALMA (100 GHz spicule). Although the 100 GHz spicule overlapped with the Mg II/304 Å spicule in the early phase, it did not show any corresponding structures in the IRIS Mg II and AIA 304 Å images after the early phase. It suggests that the spicules are individual events and do not have a physical relationship. To obtain the physical parameters of the 100 GHz spicule, we estimate the optical depths as a function of temperature and density using two different methods. One is using the observed brightness temperature by assuming a filling factor, and the other is using an emission model for the optical depth. As a result of comparing them, the kinetic temperature of the plasma and the number density of ionized hydrogen in the 100 GHz spicule are ~ 6800 K and $2.2 \times 10^{10} \text{ cm}^{-3}$. The estimated values can explain the absorbing structure in the 193 Å image, which appear as a counterpart of the 100 GHz spicule. These results suggest that the 100 GHz spicule presented in this Letter is classified to a macrospicule without a hot sheath in former terminology.

Unified Astronomy Thesaurus concepts: [The Sun \(1693\)](#); [Solar chromosphere \(1479\)](#); [Solar radio emission \(1522\)](#); [Solar spicules \(1525\)](#); [Solar ultraviolet emission \(1533\)](#)

Supporting material: animation

1. Introduction

A spicule is one of the building blocks of the solar atmosphere and a key phenomenon for understanding the heating of the corona and chromosphere. It is believed that spicules provide the energy and mass for forming the hot atmospheric layers and solar wind. Observations of spicules on the solar limb have been carried out since the 1870s (Secchi 1870), and visible and ultraviolet chromospheric lines (e.g., H α , Ca II K) are used for most observations. Due to the recent chromospheric observations with a high spatial resolution from solar observing satellites, our knowledge of a spicule is rapidly renewing (e.g., Parenti et al. 2002; de Pontieu et al. 2007; Okamoto & De Pontieu 2011; Zhang et al. 2012). On the other hand, it is hard to derive physical parameters of the spicules from these chromospheric lines because of deviating from the local thermodynamic equilibrium (LTE) for these lines. Some authors (e.g., Beckers 1972; Tsiropoula & Schmieder 1997; Alissandrakis et al. 2018) determined temperature and density of spicules and their distribution, but the results of them depend on the complicated forward modeling of the radiation from a non-LTE medium.

Theoretical studies of spicules have also been widely performed, such as the slow shock model (e.g., Osterbrock 1961; Hollweg 1982; Shibata & Suematsu 1982; Suematsu et al. 1982;

Sterling & Hollweg 1988; Sterling & Mariska 1990; Hansteen et al. 2006; Iijima & Yokoyama 2015) and Alfvén wave model (e.g., Hollweg 1982; Mariska & Hollweg 1985; de Pontieu & Haerendel 1998; Kudoh & Shibata 1999; Matsumoto & Suzuki 2012, 2014; Mathioudakis et al. 2013). Realistic radiative MHD simulation models have also been developed recently (e.g., Martínez-Sykora et al. 2017; Iijima & Yokoyama 2017), though more detailed physical analyses are necessary to understand the basic mechanism of spicule formation even for simulation results (see also Sterling 2000 and Tsiropoula et al. 2012 for a review of spicule theories). To confirm whether the physical processes in the simulations occur in the solar atmosphere, comparisons between the simulations and the observations are essential. For this purpose, it is required to derive physical quantities of spicules from the observations, but it is not so easy as mentioned above.

Since the millimeter waves emitted from the chromospheric plasma satisfy the LTE condition, it is relatively easier than with other chromospheric lines to derive the physical parameters. On the other hand, achieving a high spatial resolution with millimeter waves requires an interferometer with long baselines, and no one had investigated spicules with these wavelength ranges until recently. The Atacama Large Millimeter/submillimeter Array (ALMA; Wootten & Thompson 2009) is the largest

interferometer in the world for observing astronomical objects with millimeter and submillimeter waves, and started scientific solar observations in 2016. The spatial resolution of ALMA for observing the Sun is $\sim 2''$ with 100 GHz and is useful for investigating spicules. The ALMA data already show spicules and plasmoid eruptions (Shimojo et al. 2017b; Nindos et al. 2018; Yokoyama et al. 2018).

We obtained the observing time in ALMA Cycle 4 and succeeded in observing spicules with ALMA as well as the Interface Region Imaging Spectrograph (IRIS; De Pontieu et al. 2014), and the Atmospheric Imaging Assembly (AIA; Lemen et al. 2012) on board the *Solar Dynamics Observatory* (SDO; Pesnell et al. 2012) simultaneously. In this Letter, we present results of the coordinated observation, estimate the temperature and density of the spicules from the ALMA data, and discuss what phenomena in the past observations correspond to spicules observed at 100 GHz.

2. Observations

The ALMA observatory started offering solar observations with Band 3 and Band 6 at Cycle 4, which is the fifth open-use observing period from 2016 October to 2017 September. We succeeded in obtaining the observing time in Cycle 4, and our observation was carried out between 14:22 and 16:13 UT on 2017 April 26. We used the Band 3 receiver and the Time Division Mode (TDM) of the correlator. Hence, the frequency of the first local oscillator is 100 GHz (Andreani et al. 2016).

We synthesized one solar image from the visibility data obtained in each integration set whose duration is 2 s in the correlator, i.e., the time cadence of the images is 2 s. We used the full spectral range obtained with the TDM to enhance the signal-to-noise ratio in an image. Therefore, a synthesized image shows a distribution of intensity at 100 GHz with a bandwidth of about 8 GHz. To obtain synthesized images with high dynamic range, we carried out the self-calibration only for the phase after the standard calibration described in Shimojo et al. (2017a). The self-calibration has five steps with a different accumulating period to obtain the solution from 10 minutes for the first step to 2 s for the last step. The CLEAN deconvolution method is used for the image synthesis with the Clark algorithm and Briggs weighting scheme on CASA 5.4.0 (McMullin et al. 2007). During the five-step self-calibration, we used the CLEAN model created in the previous self-calibration as the initial guess of the synthesis. The field of view of the synthesized images is about $1'$. The sizes of the synthesized beam are $2''.67$ and $1''.88$ along the major and minor axes, respectively. Although the single-dish observations were carried out with the interferometric observations for obtaining the full-Sun images (White et al. 2017), we did not combine the synthesized images and full-Sun images because it is not established how to deal with the solar limb in the full-Sun image for the combining. Hence, in this study, we only deal with relative brightness temperatures. They are roughly the same as the differences from the averaged brightness temperature of the field of view. To remove jitter motion caused by the remaining phase error, we spatially shift the images based on the offsets estimated by the autocorrelation of the images.

The coordinated observation with IRIS was executed during the observing period. IRIS obtained slit-jaw images of Mg II (2796 Å) bands with a cadence and a spatial resolution of 40 s and $0''.40$, respectively. Although the spectra of the Mg II lines are obtained with IRIS during the period, we do not use the

spectral data because the slit positions do not cover the phenomena described in the next section. We also use 304 Å (He II) and 193 Å (Fe XII) band images taken with AIA. They provide context filtergrams with a cadence of 12 s. The spatial resolution is $1''.5$.

For the coalignment of the images, we compared between the AIA 304 Å images and the IRIS Mg II images, and aligned the Mg II images to the AIA 304 Å images. We aligned the ALMA images with the AIA ones only by using the coordinate information in the data header because the pointing accuracy of ALMA (within $0''.6$ rms) is sufficient for the current analysis (ALMA Partnership 2016).

3. The Spicules Observed with ALMA, IRIS, and AIA/SDO

The panels in Figure 1 show spicules in the 100 GHz, 304 Å, 193 Å, and Mg II images. The fields of view of them are the same. In the Mg II image, the solar limb is located around $X = 956''$, and elongated structures of most spicules cannot be recognized beyond $X = 963''$. In the 304 Å image, we cannot see the counterpart of the photosphere, and the limb at 304 Å is located around the apexes of the spicules seen in the Mg II images. The limb in the 100 GHz image is located around $X = 960''$ and is very similar to that in the 193 Å image, as reported by Yokoyama et al. (2018). We found two large spicules whose apexes are significantly higher than the others seen in the Mg II images. The region that these large spicules appear in is indicated by the red boxes in Figure 1. We concentrate on the phenomena seen in this region and describe them in Figure 2 and its animation.

A large spicule appears at around 14:37:50 UT above the limb of the Mg II and 304 Å images (Figure 2 and its animation). Because the shapes of the spicules in these images resemble each other, we call the large spicule hereafter “Mg II/304 Å spicule.” In the 304 Å images, the Mg II/304 Å spicule reached the maximum height at around 14:39 UT, and the height from the limb seen in the Mg II images is $\sim 20''$ ($\sim 15,000$ km) and the maximum width is about $4\text{--}5''$ (2900–3700 km). The rising velocity of the Mg II/304 Å spicule is about 70 km s^{-1} (the green line in Figure 3). We cannot find the counterpart of the spicule in 100 GHz and 193 Å images.

About one minute later from the start of the Mg II/304 Å spicule (14:38:40 UT), another large spicule appeared above the limb in the 100 GHz images (Figure 2 animation). We call the spicule the “100 GHz spicule” in this Letter. As pointed out in Shimojo et al. (2017a), the brightness just above the solar limb in the synthesized images is decreasing gradually with the distance from the limb. The brightness profile does not show the actual brightness profile of the limb that should be steeply decreasing at the limb. It is caused by the lack of baselines. In other words, the number of ALMA’s antennas is not enough for synthesizing images of the solar limb. To remove the influence of the artificial structure, we estimated the brightness temperature of the 100 GHz spicule from subtracting the brightness at the pre-event, from the brightness of the events, pixel by pixel. As a result, the peak enhancement of the brightness temperature caused by the spicule is about 240 K. This signal level is much larger than the noise, i.e., ~ 2.6 K (Shimojo et al. 2017a). An absorbing structure appeared in the 193 Å images, and its shape is similar to the spicule in the 100 GHz image. The rising velocity of the spicule seen in the 100 GHz and 193 Å images is about 40 km s^{-1} (the yellow

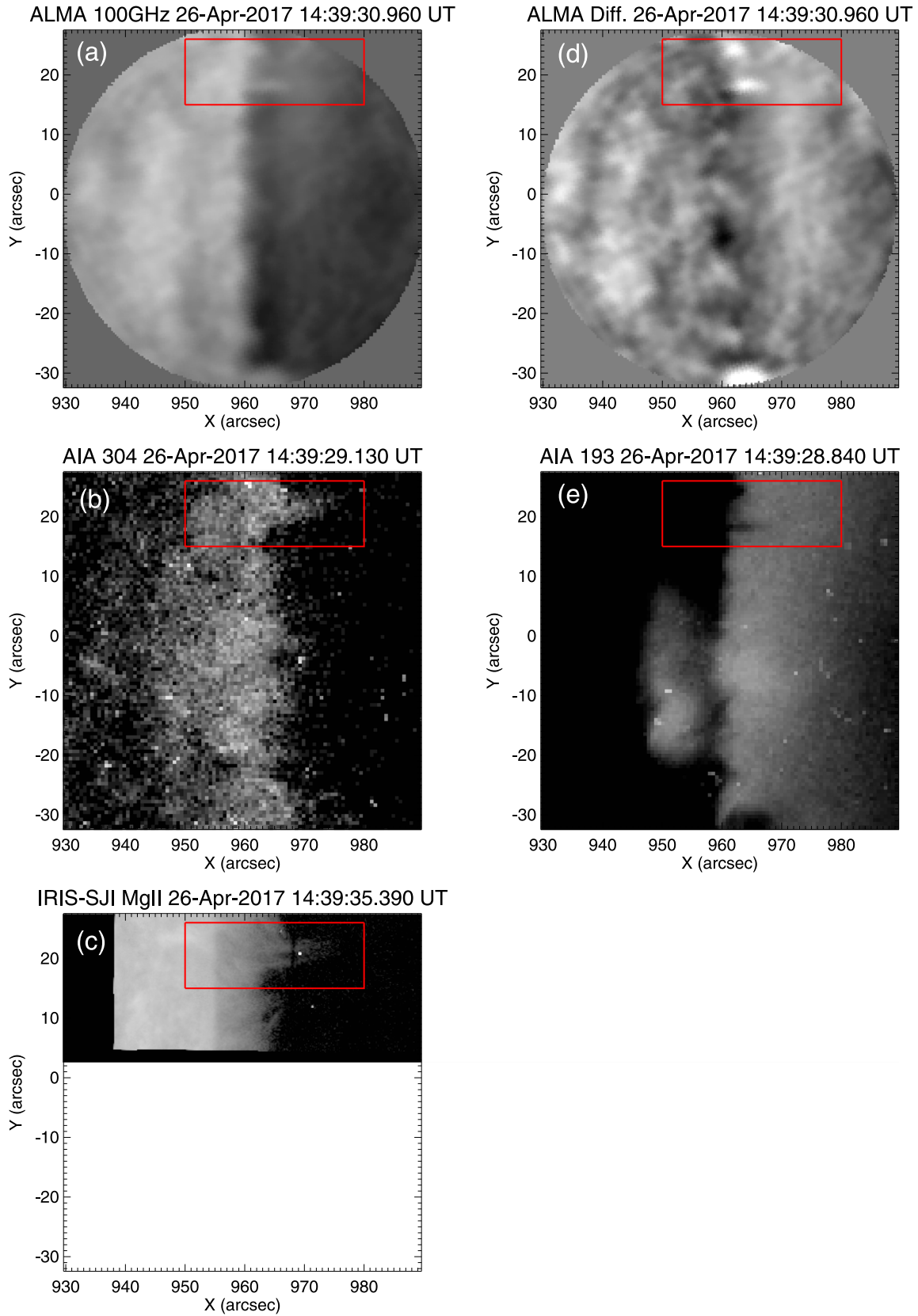


Figure 1. Joint observations with ALMA, IRIS, and AIA/SDO. Left column, from top: (a) ALMA 100 GHz, (b) AIA 304 Å band, and (c) IRIS slit-jaw image at Mg II 2796 Å band. Right column, from top: (d) the ALMA 100 GHz image subtracted from the averaged image created from the images obtained within 1.5 hr (ALMA difference image), and (e) AIA 193 Å band. The red boxes in the panels indicate the field of view of the panels in Figure 2. The X and Y in the panels present the heliocentric coordinates (SOL-X/Y).

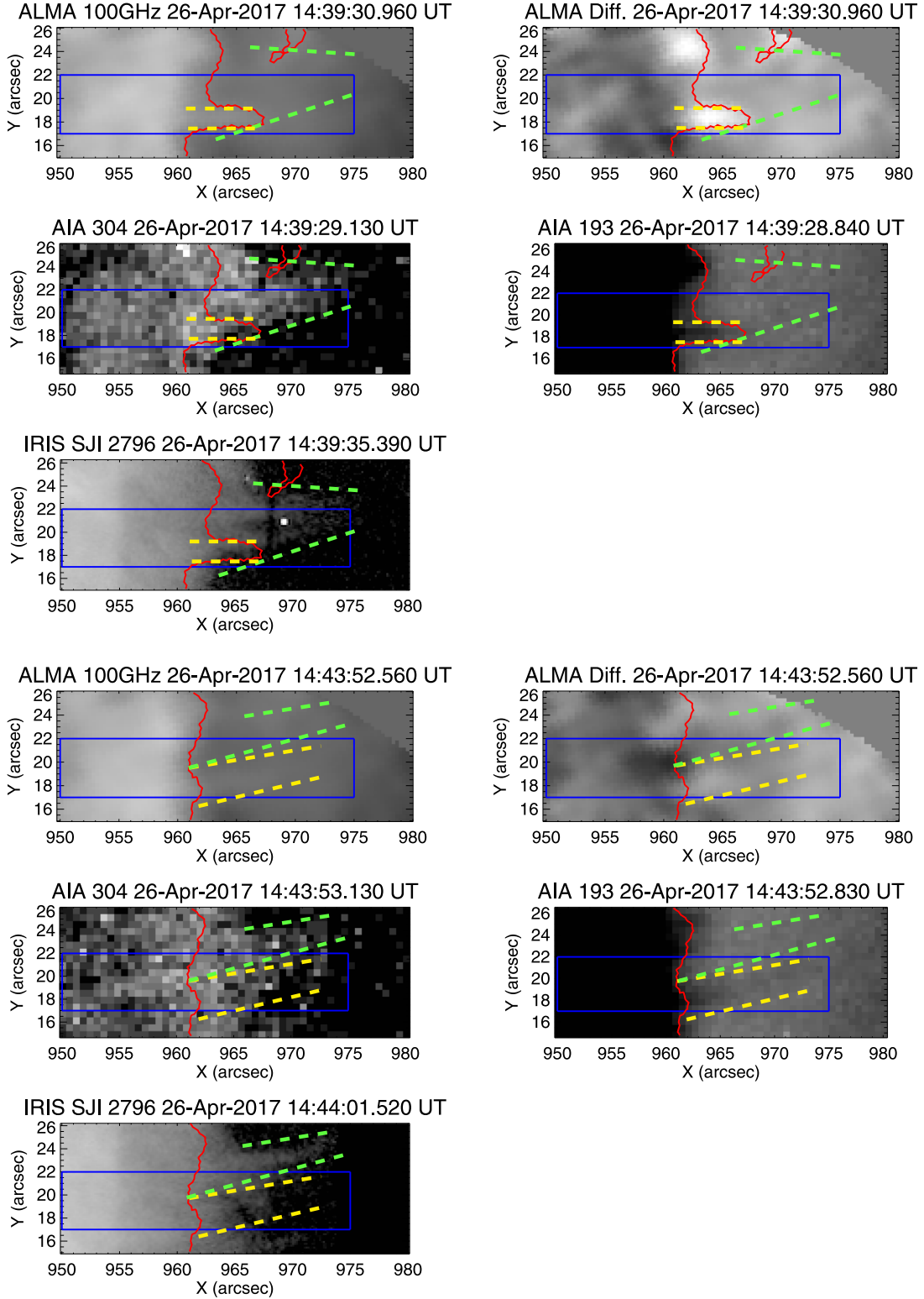


Figure 2. Close-up images around the spicules. The order of the panels is the same as that in Figure 1. Upper panels show the images at 14:39:30 UT, and lower panels show the images at 14:43:52 UT. The blue boxes show the area for creating the time-slice image (in Figure 3). The red contours indicate the 50 Jy beam^{-1} ($\sim 170 \text{ K}$) level in the 100 GHz images. The green dashed lines indicate the position of the $\text{Mg II}/304 \text{ \AA}$ spicule, and the yellow dashed lines indicate the position of the 100 GHz spicule. The animation of the figure covers from 14:35:35 to 14:45:35 UT, and the period is repeated twice. In the second one, the rising phase of the spicule in the Mg II images and 304 \AA images and the rising phase of the spicule in the 100 GHz and 193 \AA images are indicated as green and blue lines, respectively. The real-time duration is 40 s.

(An animation of this figure is available.)

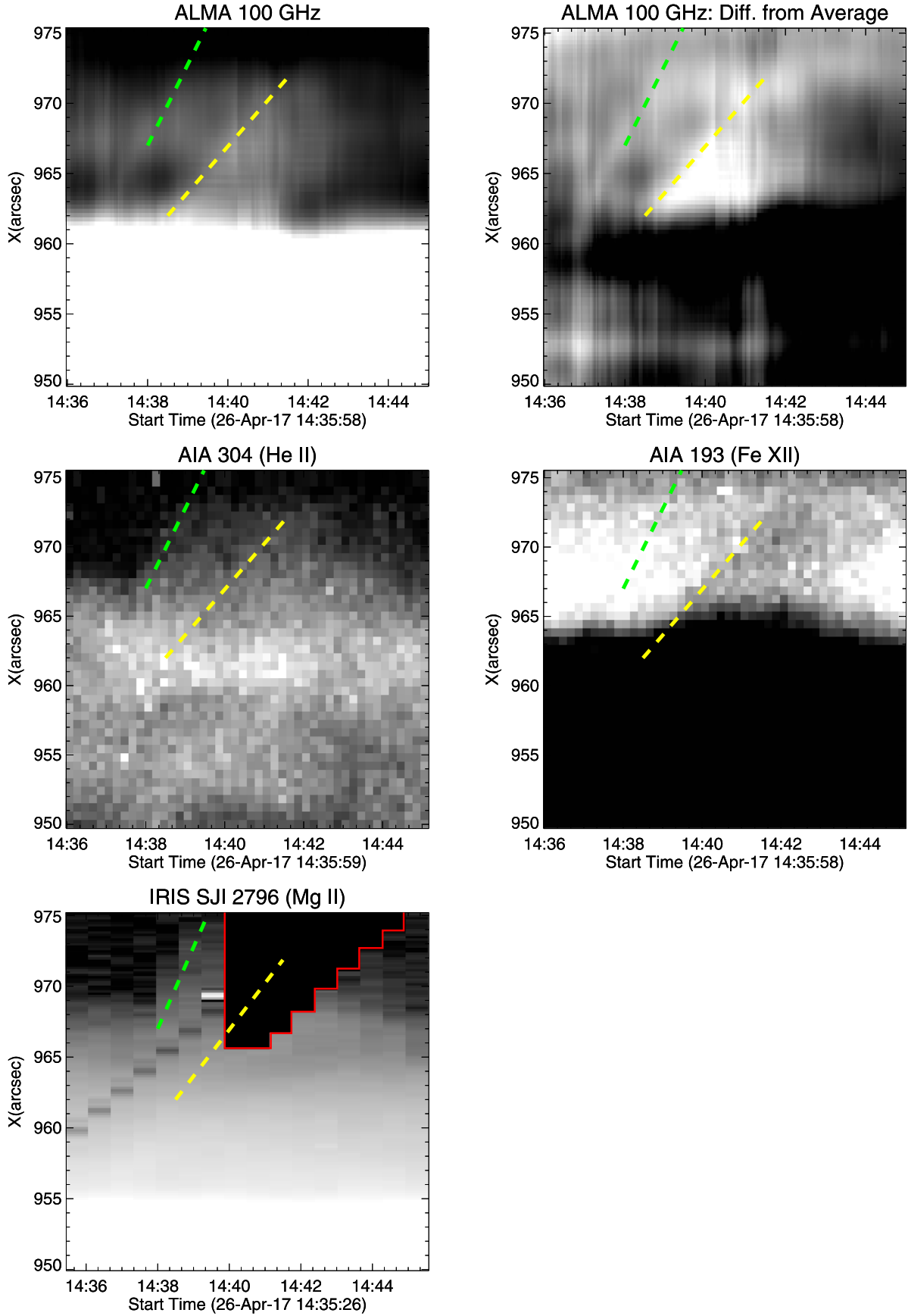


Figure 3. Time-slice images created from the data of the areas indicated by the blue boxes in Figure 2. The order of the panels is the same as that in Figure 1. The green and yellow dashed lines indicate the spicule in the AIA 304 Å and IRIS 2796 Å images and the spicule in the ALMA 100 GHz and AIA 193 Å images, respectively. The slope of the green lines corresponds with 70 km s^{-1} , and the slope of the yellow line reveals 40 km s^{-1} . The red lines in the IRIS image indicate the area that is not observed.

line in Figure 3). The 100 GHz spicule reached the maximum height of $\sim 15''$ ($\sim 11,000$ km) from the limb seen in the Mg II images at 14:41 UT. The obtained width in 100 GHz image is $2''$ (~ 1500 km), but the actual width can be narrower than the value because the width is comparable to the size of the synthesized beam.

The 100 GHz spicule is overlapped by the Mg II/304 Å spicule at 14:38:40 UT, as shown in the upper panels of Figure 2. Nevertheless, we cannot find any remarkable enhancements in Mg II and 304 Å images that correspond to the 100 GHz spicule. After 14:43 UT (the lower panels in Figure 2), there is no overlapped region between the 100 GHz spicule and Mg II/304 Å spicule. Moreover, the apparent velocity of the 100 GHz spicule (~ 40 km s $^{-1}$) is significantly slower than the Mg II/304 Å spicule (~ 70 km s $^{-1}$), as shown in Figure 3. The facts would suggest that the spicules are individual events and do not have a physical relationship.

Considering the typical formation temperatures of the Mg II band (5000–16,000 K; De Pontieu et al. 2014) and 304 Å band ($\sim 50,000$ K; Lemen et al. 2012), the temperature of the 100 GHz spicule should be lower than about 10,000 K because we cannot find its counterpart in Mg II and 304 Å images. Can the plasma with such temperature and density be explained by the observed structure at 100 GHz? To answer the question, we estimate the optical depths of the spicule at 100 GHz from observations and an emission model, and examine the temperature and density of the 100 GHz spicule. Before the estimations, we must consider a filling factor because the spatial resolution of the 100 GHz images is not enough for resolving the spicule. The spatial resolution of AIA is not enough either, so that we assume that the width of the 100 GHz spicule is $0''.5$. Since the apparent width of the 100 GHz spicule is $2''$, the filling factor of the spicule in the 100 GHz image is roughly 0.25. Thus, we use 960 K ($=240$ K/0.25) as the brightness temperature of the 100 GHz spicule. From the observed enhancement of the brightness temperature, the optical depth of the spicule as a function of electron temperature ($\tau = -\log_e(1 - T_b/T)$: optical depth, T : electron temperature, T_b : brightness temperature) is shown by the black solid line in Figure 4, based on the following assumptions: (1) the background emission is negligibly small because the phenomenon locates in the off-limb region, (2) the temperature range of the 100 GHz spicule is lower than 10,000 K, and (3) the emission from the spicule is the thermal free-free emission (electron-proton free-free) from the medium satisfied the LTE condition. On the other hand, by assuming the temperature and density of the spicule and the line-of-sight diameter of the spicule is $0''.5$ (same as the width), we derived the optical depth of the thermal free-free emission from the emission model (e.g., Lang 1980; Dulk 1985; Yokoyama et al. 2018). The model functions indicate the colored lines in Figure 4. The black line estimates from the observations and the colored line of the model cross when we assume the kinetic temperature of the plasma and the number density of ionized hydrogen to be ~ 6800 K and 2.2×10^{10} cm $^{-3}$, respectively, and the values indicate the physical parameters of the 100 GHz spicule.

Zhang et al. (1998) presented that the absorption at the limb in 195 Å images shows the chromosphere, where the number density of neutral hydrogen is 10^{10} cm $^{-3}$. Therefore, for explaining the absorbing counterpart of the 100 GHz spicule in the 193 Å images, the number density of neutral hydrogen should be similar or larger than the number density of the electrons in the temperature range that the thermal free-free

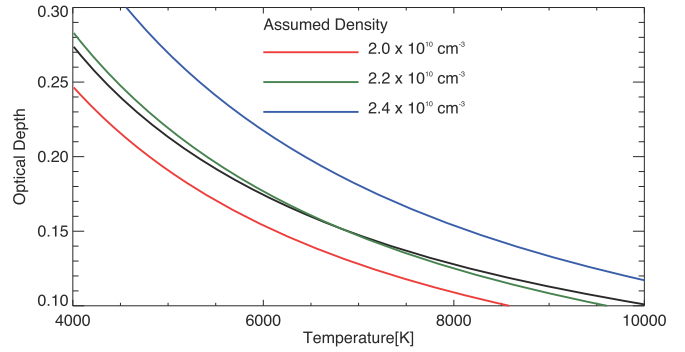


Figure 4. Estimated optical depths as a function of assumed temperatures. The solid black line indicates the optical depth at 100 GHz that is estimated from the brightness temperatures considering the filling factor, and assumed temperatures of the spicule. The colored lines reveal the theoretical optical depth at 100 GHz assuming the temperature, density, diameter of the spicule, and the emission via the thermal free-free interaction. The color indicates the difference of the assumed density.

emission is dominant. According to the lower right panel of Figure 1 in Rutten (2017), the temperature of the 100 GHz spicule should be 4000–7000 K. The temperature range is consistent with the temperature derived above, and thus justifies our estimation. Moreover, the total number density of hydrogen, which is the sum of the number densities of ionized and neutral hydrogen, is assumed to be 10^{11} cm $^{-3}$ in the lower right panel of Figure 1 in Rutten (2017). Hence, the mass density of the 100 GHz spicule might be about $\sim 10^{-13}$ g cm $^{-3}$, which is consistent with the previous results of the infrared observations (Sterling 2000).

4. Discussion

The spicules described in this Letter would be categorized into “macrospicules” formerly because their sizes and velocities are larger than the typical values of spicules. Therefore, it is appropriate to compare our results with the properties of macrospicules, rather than those of spicules. Macrospicules are most often visible in EUV transition-region lines, such as in He II 304 Å, and consist of a cool core and hot sheath (e.g., Bohlin et al. 1975; Withbroe et al. 1976; Parenti et al. 2002; Loboda & Bogachev 2019). Since the Mg II/304 Å spicule has a cool component revealed in the Mg II image and a hotter component shown in the He II 304 Å images, it has the typical properties of macrospicules described in the preceding papers.

On the other hand, for the 100 GHz spicule, we do not find its counterpart in Mg II/304 Å images. Habbal & Gonzalez (1991) presented that the brightness temperature of macrospicules at 4.8, 8.5, and 15 GHz can be explained by an empirical model that macrospicules consist of a cool core at ~ 8000 K surrounded by a hot sheath at $1\text{--}2 \times 10^5$ K. The temperature of the cool core is similar to the 100 GHz spicule, and there is no significant signal of hotter lines around the 100 GHz spicule. Hence, the 100 GHz spicule described in this Letter is a macrospicule without a hot sheath that is already reported by some authors (e.g., Bohlin et al. 1975; Madjarska et al. 2011).

The height of the limb in 100 GHz images is similar to the average heights of the typical spicules seen in Mg II images, and the limb seen in 304 Å images. It suggests that the 100 GHz limb is covered by multiple spicules in Mg II and

He II images. Hence, we cannot identify a hot sheath around the 100 GHz spicules, but we cannot argue that a spicule without a hot sheath is common in 100 GHz images. To reveal the properties of typical spicules seen in millimeter-wave images, we need a resolution high enough for resolving spicules, which will be realized with the higher observing frequency and/or long baselines of ALMA.

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