

Correlations in flavor-changing tqZ couplings constrained via experimental data

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Abstract – Possible non-standard tqZ couplings, where $q = c$ or u , originated from general flavor-changing neutral current interactions are studied model-independently using the effective Lagrangian consisting of several $SU(3) \times SU(2) \times U(1)$ invariant dimension-6 operators. After the electroweak symmetry breaking, these operators are recombined to form four kinds of independent terms whose coefficients are complex in general. The Lagrangian could therefore include up to eight independent coupling parameters. Through searches for experimentally allowed regions of these parameters, it is found that some correlations exist among the signs and sizes of those couplings.

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Searches for Flavor-Changing Neutral Current (FCNC) are a quite attractive mission for future collider experiments: The existence of physics beyond the standard model is strongly indicated if new phenomena originated from FCNC are observed, because the event probability of such phenomena is too tiny to detect within the standard-model framework [1]. In exploring such rare processes, the top quark is expected to play an especially important role, since it decays without being affected by non-perturbative effects thanks to its short lifetime [2,3] in contrast to the other heavy quarks.

We therefore studied top-quark FCNC processes model-independently in our latest paper [4] and derived constraints on tqZ ($q = u/c$) couplings, which induce FCNC interactions, using the effective Lagrangian¹. We however did not deal with any correlation among those coupling constants there, so we focus on this issue in this letter and study how these couplings are related with each other, aiming for deeper understanding of our results.

¹We have given a detailed list of preceding works by other authors in [4].

In our analysis, we use the following effective Lagrangian to describe the general tqZ interactions [5–8]:

$$\mathcal{L}_{tqZ} = -\frac{g}{2 \cos \theta_W} \left[\bar{\psi}_q(x) \gamma^\mu (f_1^L P_L + f_1^R P_R) \psi_t(x) Z_\mu(x) + \bar{\psi}_q(x) \frac{\sigma^{\mu\nu}}{M_Z} (f_2^L P_L + f_2^R P_R) \psi_t(x) \partial_\mu Z_\nu(x) \right], \quad (1)$$

where g and θ_W are the $SU(2)$ coupling constant and the weak mixing angle, $P_{L/R} \equiv (1 \mp \gamma_5)/2$, $f_{1/2}^{L/R}$ stand for the non-standard couplings parameterizing contributions from relevant $SU(3) \times SU(2) \times U(1)$ gauge invariant dimension-6 effective operators [4,7]. We treat these coupling parameters as complex numbers independent of each other in order to perform analyses as model-independently as possible. Thus, the resultant tqZ couplings are expressed by up to eight independent parameters.

Using the above Lagrangian, we can derive the theoretical partial decay width as an eight-variable function $\Gamma_{tqZ}^{\text{th}}(f_{1/2}^{L/R})$. On the other hand, the experimental partial decay width $\Gamma_{tqZ}^{\text{exp}}$ is obtained by the product of the branching ratio $\text{Br}(t \rightarrow qZ)$ and the top-quark total decay width Γ_t : $\Gamma_{tqZ}^{\text{exp}} = \text{Br}(t \rightarrow qZ) \times \Gamma_t$. Then, allowed

Table 1: Allowed minimum and maximum values of the tcZ coupling parameters for $\Gamma_{tcZ}^{\text{th}} < 3.0 \times 10^{-4}$ in the case that $\text{Re}(f_1^L)$ is fixed to its maximum value 3.4×10^{-2} .

	f_1^L		f_1^R	
	$\text{Re}(f_1^L)$	$\text{Im}(f_1^L)$	$\text{Re}(f_1^R)$	$\text{Im}(f_1^R)$
Min.	3.4×10^{-2}	-4.0×10^{-3}	-4.0×10^{-3}	-4.0×10^{-3}
Max.	(Fixed)	4.0×10^{-3}	4.0×10^{-3}	4.0×10^{-3}
	f_2^L		f_2^R	
	$\text{Re}(f_2^L)$	$\text{Im}(f_2^L)$	$\text{Re}(f_2^R)$	$\text{Im}(f_2^R)$
Min.	-3.0×10^{-3}	-3.0×10^{-3}	-2.5×10^{-2}	-3.0×10^{-3}
Max.	3.0×10^{-3}	3.0×10^{-3}	-2.2×10^{-2}	3.0×10^{-3}

Table 2: Allowed minimum and maximum values of the tcZ coupling parameters for $\Gamma_{tcZ}^{\text{th}} < 3.0 \times 10^{-4}$ in the case that $\text{Re}(f_2^R)$ is fixed to its minimum value -2.8×10^{-2} .

	f_1^L		f_1^R	
	$\text{Re}(f_1^L)$	$\text{Im}(f_1^L)$	$\text{Re}(f_1^R)$	$\text{Im}(f_1^R)$
Min.	2.6×10^{-2}	-5.0×10^{-3}	-5.0×10^{-3}	-5.0×10^{-3}
Max.	3.1×10^{-2}	5.0×10^{-3}	5.0×10^{-3}	5.0×10^{-3}
	f_2^L		f_2^R	
	$\text{Re}(f_2^L)$	$\text{Im}(f_2^L)$	$\text{Re}(f_2^R)$	$\text{Im}(f_2^R)$
Min.	-4.0×10^{-3}	-4.0×10^{-3}	-2.8×10^{-2}	-4.0×10^{-3}
Max.	4.0×10^{-3}	4.0×10^{-3}	(Fixed)	4.0×10^{-3}

regions for $f_{1/2}^{L/R}$ are obtained by varying their real and imaginary parts at the same time and searching for the parameter space that satisfies $\Gamma_{tqZ}^{\text{th}}(f_{1/2}^{L/R}) < \Gamma_{tqZ}^{\text{exp}}$.

Let us briefly show the result for the tcZ couplings as an example (see ref. [4] for more detailed results): their allowed regions at 95% confidence level for $\Gamma_t = 1.322 \text{ GeV}$ [9]² and $\text{Br}(t \rightarrow cZ) < 2.3 \times 10^{-4}$ [10] (thus, we are to solve $\Gamma_{tcZ}^{\text{th}} < 3.0 \times 10^{-4}$) are derived as

$$|\text{Re}/\text{Im}(f_1^{L/R})| \leq 3.4 \times 10^{-2}, \quad |\text{Re}/\text{Im}(f_2^{L/R})| \leq 2.8 \times 10^{-2}.$$

Here, we should comment on the meaning of the *allowed* regions. This means that if we give one parameter a value outside its allowed range, we can no longer reproduce the current experimental data no matter how we vary the other parameters.

²The direct measurement of the total decay width of the top quark is consistent with the prediction by the standard model, but the measured one still has a large uncertainty. Therefore, we use the standard-model value here instead of the experimental value (see ref. [4] for related discussions). Since we are focusing on the rare decays, this replacement does not bring any significant problem into our analysis.

Now, using these results and carrying out some further computations, we investigate if there is a certain relationship among the couplings. Let us start with the $t \rightarrow cZ$ process³. We first set one of the couplings to its *maximum* value and vary all the others. When $\text{Re}(f_1^L)$ is taken as such a fixed parameter, the allowed regions of the remaining couplings are derived as in table 1. From this table, we can see a relation between $\text{Re}(f_1^L)$ and $\text{Re}(f_2^R)$: the sign of $\text{Re}(f_2^R)$ is opposite to that of $\text{Re}(f_1^L)$ and the size of $\text{Re}(f_2^R)$ is of the same order as $\text{Re}(f_1^L)$. Next, in table 2, we show similar bounds to those in table 1 but in the case that $\text{Re}(f_2^R)$ is set to its *minimum* value. There we find the same kind of correlation between $\text{Re}(f_1^L)$ and $\text{Re}(f_2^R)$ again.

It must be meaningful here to point out the following fact: While we are now observing the relation between $\text{Re}(f_1^L)$ and $\text{Re}(f_2^R)$ assuming that they take the maximal/minimal values, the other parameters can also have

³The latest upper bound on $\text{Br}(t \rightarrow cZ)$ in the published paper [11] is not 2.3×10^{-4} but a bit weaker 2.4×10^{-4} . We however use the former value because the difference is tiny and also because it is crucial for our purpose to provide results that can be directly compared with what we gave in [4].

Table 3: Allowed minimum and maximum values of the tuZ coupling parameters for $\Gamma_{tuZ}^{\text{th}} < 2.2 \times 10^{-4}$ in the case that $\text{Re}(f_1^L)$ is fixed to its maximum value 2.9×10^{-2} .

	f_1^L		f_1^R	
	$\text{Re}(f_1^L)$	$\text{Im}(f_1^L)$	$\text{Re}(f_1^R)$	$\text{Im}(f_1^R)$
Min.	2.9×10^{-2}	-4.0×10^{-3}	-4.0×10^{-3}	-4.0×10^{-3}
Max.	(Fixed)	4.0×10^{-3}	4.0×10^{-3}	4.0×10^{-3}
	f_2^L		f_2^R	
	$\text{Re}(f_2^L)$	$\text{Im}(f_2^L)$	$\text{Re}(f_2^R)$	$\text{Im}(f_2^R)$
Min.	-4.0×10^{-3}	-4.0×10^{-3}	-2.2×10^{-2}	-4.0×10^{-3}
Max.	4.0×10^{-3}	4.0×10^{-3}	-1.8×10^{-2}	4.0×10^{-3}

 Table 4: Allowed minimum and maximum values of the tuZ coupling parameters for $\Gamma_{tuZ}^{\text{th}} < 2.2 \times 10^{-4}$ in the case that $\text{Re}(f_2^R)$ is fixed to its minimum value -2.4×10^{-2} .

	f_1^L		f_1^R	
	$\text{Re}(f_1^L)$	$\text{Im}(f_1^L)$	$\text{Re}(f_1^R)$	$\text{Im}(f_1^R)$
Min.	2.2×10^{-2}	-4.0×10^{-3}	-4.0×10^{-3}	-4.0×10^{-3}
Max.	2.7×10^{-2}	4.0×10^{-3}	4.0×10^{-3}	4.0×10^{-3}
	f_2^L		f_2^R	
	$\text{Re}(f_2^L)$	$\text{Im}(f_2^L)$	$\text{Re}(f_2^R)$	$\text{Im}(f_2^R)$
Min.	-4.0×10^{-3}	-4.0×10^{-3}	-2.4×10^{-2}	-4.0×10^{-3}
Max.	4.0×10^{-3}	4.0×10^{-3}	(Fixed)	4.0×10^{-3}

some allowed space though its size is one order of magnitude smaller. In fact, even if those other parameters had no allowed area, the allowed space for the correlated pair would not change drastically.

On the other hand, if we assume as an extreme case that only one non-standard coupling exists in the tqZ interactions, the allowed region of that coupling becomes rather small: In the case that only $\text{Re}(f_1^L)$ or $\text{Re}(f_2^R)$ exists in the tcZ couplings, we can get the corresponding allowed region as $|\text{Re}(f_1^L)| \leq 1.9 \times 10^{-2}$ or $|\text{Re}(f_2^R)| \leq 1.5 \times 10^{-2}$.

We performed these analyses for all the remaining parameters, too. The results are summarized as follows:

- Relational expressions $\text{Re}/\text{Im}(f_{1/2}^{L/R}) = -C \text{Re}/\text{Im}(f_{2/1}^{R/L})$ hold, where the maximum ranges of $f_1^{L/R}$ and $f_2^{L/R}$ are given with $0.93 \lesssim C \lesssim 1.1$ and $0.65 \lesssim C \lesssim 0.73$, respectively, by substituting the maximum or minimum values of $f_2^{R/L}$ and $f_1^{R/L}$ in the right-hand side of the expressions.
- Three or more non-standard couplings cannot take large values (within the allowed regions) at the same time.

- The allowed region could be roughly twice larger when two non-standard coupling constants exist than in the case that only one non-standard coupling constant exists.

In the same way, we then analyzed the $t \rightarrow uZ$ process. The main results coming from the current upper bound $\text{Br}(t \rightarrow uZ) < 1.7 \times 10^{-4}$ [10,11] are shown in table 3 and table 4, which are to be compared with table 1 and table 2. There we can clearly see parameter correlations quite similar to those in the tcZ couplings. That is, what we found here is common to both the tcZ couplings and the tuZ couplings.

Finally, we would like to present important comments on 1) to what extent our results depend on the upper limit of the branching ratio of $t \rightarrow qZ$ and 2) how model-independent and unique our analyses are:

- 1) We used here $\text{Br}(t \rightarrow qZ)$ presented in [10,11], but it is expected to be improved by half at High-Luminosity Large Hadron Collider [12]. Therefore we performed the same computations assuming that reduced upper bound. Of course, the allowed parameter ranges are thereby narrowed, which we already

showed in the previous work [4], but interestingly enough we found that the coupling-parameter correlations and related results are little affected and still hold even there.

- 2) We worked in the framework of the effective Lagrangian consisting of dimension-6 operators and treated all their coefficients as independent complex numbers, which led to eight-parameter analyses. That is, our results were derived by varying those eight parameters at the same time and comparing the resultant Γ_{tqZ}^{th} with the corresponding experimental bound. This is quite in contrast with many preceding analyses, where only some coefficients are treated as free parameters at once fixing the others. For example, the CMS Collaboration carried out detailed studies in [13] but they dropped some terms in the effective Lagrangian from the beginning. The ATLAS Collaboration also performed comprehensive analyses in [11], but what they gave are constraints on the absolute values of four coefficients assuming that only one parameter has a non-zero value in each analysis. In such limited analyses, we cannot take into account possible cancellations among contributions from different parameters. Indeed the existence of the correlations shown here tells us that it is true.

In conclusion, we have studied possible non-standard tqZ couplings and correlations among them in the framework of the effective Lagrangian. This Lagrangian can incorporate up to eight independent coupling parameters describing FCNC interactions. The allowed regions of these couplings satisfying the current experimental limits get larger when several numbers of couplings exist than when only one coupling exists. It was found that the allowed region becomes the largest when there are relations as $\text{Re}/\text{Im}(f_{1/2}^{L/R}) = -C \text{Re}/\text{Im}(f_{2/1}^{R/L})$ with $C \simeq 0.65 \sim 0.73$ ($C \simeq 0.93 \sim 1.1$) in the case that the maximum or minimum value of $f_1^{R/L}$ ($f_2^{R/L}$) is substituted in the right-hand side. Tables 1–4 also tell us that the non-standard couplings except for the two correlated ones can take non-zero values as well though their sizes are one order of magnitude smaller than those of the correlated ones. However,

even if they had no allowed regions, the current experimental limits could be realized by the existence of two correlated couplings alone.

Our results are from the current experimental data, but it is quite interesting to note that the parameter correlations and related results are little affected even if the experimental limits on $\text{Br}(t \rightarrow qZ)$ are improved, *e.g.*, by half at future facilities. Since this analysis was performed in a very general framework and does not depend on any special assumptions, the results pointed out here will be useful information for constructing a specific model inducing rather strong FCNC interactions.

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