



PAPER

A parameter sensitivity study for simulating DNA damage after proton irradiation using TOPAS-nBio

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23 April 2020Hongyu Zhu^{1,2,3}, Aimee L McNamara^{1,4}, Jose Ramos-Mendez⁵, Stephen J McMahon⁶ , Nicholas T Henthorn⁷, Bruce Faddegon⁵, Kathryn D Held^{1,4}, Joseph Perl⁸, Junli Li^{2,3}, Harald Paganetti^{1,4} and Jan Schuemann^{1,4} ¹ Department of Radiation Oncology, Massachusetts General Hospital, Boston, MA 02114, United States of America² Department of Engineering Physics, Tsinghua University, Beijing 100084, People's Republic of China³ Key Laboratory of Particle and Radiation Imaging (Tsinghua University), Ministry of Education, Beijing 100084, People's Republic of China⁴ Harvard Medical School, Boston, MA 02114, United States of America⁵ Department of Radiation Oncology, University of California San Francisco, CA 94143, United States of America⁶ Centre for Cancer Research and Cell Biology, Queens University Belfast, Belfast, United Kingdom⁷ Division of Molecular and Clinical Cancer Sciences, Faculty of Biology, Medicine and Health, University of Manchester, Manchester, United Kingdom⁸ SLAC National Accelerator Laboratory, Menlo Park, CA, United States of AmericaE-mail: jschuemann@mgh.harvard.edu**Keywords:** Monte Carlo, DNA damage, proton, sensitive study, TOPAS-nBio, Geant4-DNA, nucleus model**Abstract**

Monte Carlo (MC) track structure simulation tools are commonly used for predicting radiation induced DNA damage by modeling the physical and chemical reactions at the nanometer scale. However, the outcome of these MC simulations is particularly sensitive to the adopted parameters which vary significantly across studies. In this study, a previously developed full model of nuclear DNA was used to describe the DNA geometry. The TOPAS-nBio MC toolkit was used to investigate the impact of physics and chemistry models as well as three key parameters (the energy threshold for direct damage, the chemical stage time length, and the probability of damage between hydroxyl radical reactions with DNA) on the induction of DNA damage. Our results show that the difference in physics and chemistry models alone can cause differences up to 34% and 16% in the DNA double strand break (DSB) yield, respectively. Additionally, changing the direct damage threshold, chemical stage length, and hydroxyl damage probability can cause differences of up to 28%, 51%, and 71% in predicted DSB yields, respectively, for the configurations in this study.

1. Introduction

With the recognition that DNA is a significant target for causing biological effects after ionization irradiation (Hall and Giaccia 2018), modeling DNA damage mechanistically is a critical research topic. Different Monte Carlo (MC) track structure simulation tools as well as detailed DNA models have been developed to predict levels of DNA damage (Nikjoo *et al* 2001, 2006, 2016, Friedland *et al* 2003, 2017, Meylan *et al* 2017, Lampe *et al* 2018b, Tang *et al* 2019, Sakata *et al* 2019) by modeling the direct damages caused by initial physical interactions in the physical stage and indirect damages caused by subsequent chemical interactions of radiation induced reactive oxygen species in the physico-chemical and the chemical stages. The simulation of DNA damage allows researchers to further their understanding of the induction of cellular damage after ionizing radiation in order to interpret, analyze or design experiments (Campa *et al* 2005).

Differences in the simulation tool and geometry structure intrinsically affect the DNA damage result. Furthermore, different studies adopted different parameter values, which could also significantly alter simulation results. For instance, some studies (Nikjoo *et al* 2001, Meylan *et al* 2017, Mokari *et al* 2018, Lampe *et al* 2018b) assumed that a direct strand break (SB) is formed with a fixed threshold of 17.5 eV, i.e. a direct SB will be formed if the accumulated energy deposition within one sugar phosphate or hydration shell exceeds this threshold. In contrast, other studies (Friedland *et al* 2003, 2011, 2017, Sakata *et al* 2019) used a

Table 1. Simulation parameters adopted in different MC studies.

Reference	Code	Physics constructor	Chemical stage length	·OH damage probability ^c	Direct damage threshold
Nikjoo <i>et al</i> (2001)	PITS	—	1 ns	0.13 ^a	17.5 eV
Friedland <i>et al</i> (2017)	PARTRAC	—	100 ns	0.65 ^b	5–37.5 eV
Friedland <i>et al</i> (2011)	PARTRAC	—	2.5 ns	0.65 ^b	5–37.5 eV
Friedland <i>et al</i> (2003)	PARTRAC	—	10 ns	0.13 ^a	5–37.5 eV
Sakata <i>et al</i> (2019)	Geant4-DNA	Option4	2.5 ns	0.4 ^b	5–37.5 eV
Mokari <i>et al</i> (2018)	Geant4-DNA	Default	1 ns	0.65 ^b	17.5 eV
Lampe <i>et al</i> (2018a)	Geant4-DNA	Option4	1 ns	0.4 ^b	17.5 eV
Meylan <i>et al</i> (2017)	Geant4-DNA	Default	2.5 ns	0.4 ^b	17.5 eV

^aDefined between ·OH and whole DNA volume.

^bDefined between ·OH and deoxyribose ('backbone').

^cThe ·OH damage probability of $P_{\text{OH-DNA}} = 0.13$ is equivalent to $P_{\text{OH-backbone}} = 0.65$ (Friedland *et al* 2003).

linear acceptance threshold ranging from 5 to 37.5 eV, i.e. the probability to form a direct SB linearly increases from 0 when less than 5 eV energy deposition was accumulated to 1 when more than 37.5 eV energy deposition was accumulated in the target DNA volume. Many other parameters (see table 1), such as the length of the chemical stage and the probability of chemical interactions to cause DNA damages also vary over different studies.

Differences can also be found in handling of interactions among studies. Friedland *et al* assumed hydroxyl radicals were scavenged after an encounter with histones using a radius of 4.5 nm as the interaction distance, while other chemical species were assumed to not interact with the DNA or with histones by preventing these species from diffusing into the volumes occupied by these molecules (2003). Sakata *et al* and Lampe *et al* used a parameter of radical kill radius (r_{kill}) in the chemical stage, which defines a radius from the DNA volume beyond which chemical molecules will be terminated and would not be simulated to avoid modeling radical tracks that will not interact with the DNA; this subsequently speeds up the simulation runtime (Lampe *et al* 2018a, Sakata *et al* 2019).

Differences in the parameters as well as the handling of interactions can lead to significant uncertainties in the simulation of DNA damages. Lampe *et al* studied the impact of physics models, direct damage threshold, the hydroxyl damage probability and the radical kill distance on the DNA damage after irradiating a test geometry filled with DNA segments with electrons (2018a).

Lai *et al* studied the impact of cross-section data, energy cutoff in transportation, direct damage threshold, etc on electron induced DNA damage yields (Lai *et al* 2020). Margis *et al* examined the impact of the physics model, target size as well as the tracking cutoff energy on direct DNA damages induced by electrons with a microdosimetric method (Margis *et al* 2020).

In this study, TOPAS-nBio was used for a sensitivity study of physics and chemistry models as well as three key parameters that could affect direct and indirect DNA damage results. A previously developed full nucleus model (Zhu *et al* 2019) was used for the description of the DNA geometry. DNA damages were simulated with different physics constructors, chemistry models, direct damage threshold values, chemical stage time length, and the hydroxyl damage probability. The impacts of these parameters on the DNA damages were quantified and compared with the ratio of DNA damage yield. In addition to the differences in the above listed parameters, the nucleus geometry is another important aspect when studying DNA damage. Different research groups have developed their own nucleus geometries and implemented them in their own respective codes (Nikjoo and Girard 2012, Meylan *et al* 2017, Friedland *et al* 2017, Tang *et al* 2019, Sakata *et al* 2019). However, the uncertainties due to the chosen DNA geometry were not considered in this work.

2. Methods

2.1. Simulation tool and geometry setup

TOPAS-nBio (Schuemann *et al* 2019a) is a MC track structure simulation application layered on Geant4-DNA (Incerti *et al* 2010a, 2018), which extends the functionality of TOPAS (Perl *et al* 2012), an advanced and user-friendly MC tool based on Geant4 (Agostinelli *et al* 2003, Allison *et al* 2006, 2016) to support simulation studies at the patient scale. TOPAS-nBio offers an extensive library of advanced biological geometries ranging from the micrometer scale (e.g. cells and organelles) down to the nanometer

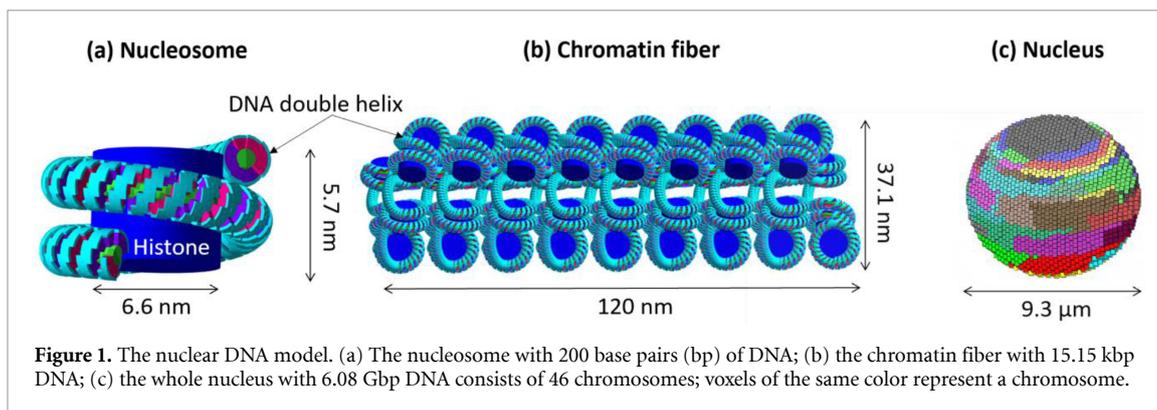


Table 2. Average number of primary protons to deposit a dose of 1 Gy within the nucleus.

Proton energy (MeV)	Average number of primary protons
0.5	6.3
0.6	7.5
0.8	9.9
1	12.1
1.5	16.9
2	21.1
5	43.0
10	76.0
20	139.4
50	312.0

scale (e.g. DNA molecules and proteins) to support studies at the cellular and sub-cellular level (McNamara *et al* 2018).

In this study, a previously developed full nucleus model (figure 1) (Zhu *et al* 2019), which was built using a fractal pattern based on a 3D Hilbert space-filling curve (Lieberman-Aiden *et al* 2009; McNamara *et al* 2018), was used to describe the DNA geometry. The spherical nucleus model has a diameter of 9.3 μm, and it was organized in a hierarchical pattern in the order of DNA double helix, nucleosomes, chromatin fibers, and chromosomal territories. The double helix structure was formed with a half-cylindrical base volume with a radius of 0.5 nm and a quarter-cylindrical sugar-phosphate backbone volume with a radius of 1.15 nm (Henthorn *et al*, 2017). Considering that ionizations in the hydration shell could also lead to DNA damage (Friedland *et al* 2003, 2017, Meylan *et al* 2017), a 0.16 nm thick quarter-cylindrical shell surrounding the backbone volume was modeled as the hydration shell. The DNA double helix wraps around the cylindrical histone volume to form the nucleosome (figure 1(a)), which is the unit of chromatin fibers (figure 1(b)). Each chromatin fiber has 51 nucleosomes and contains 15 150 base pairs (bp) of DNA. The nucleus was voxelized and consists of 14 328 voxels; each voxel was filled with chromatin fibers folded according to a continuous 3D Hilbert space-filling curve. The resulting nucleus (figure 1(c)) includes 46 chromosomes and was filled with 6.08 Gbp (Giga base pairs) of DNA; the DNA density of the nucleus was 14.4 Mbp μm⁻³ and agrees with the typical DNA density of a mammalian nucleus of 15 Mbp μm⁻³ (Ghosh and Jost, 2018).

DNA damage induction after irradiation by 0.5–50 MeV monoenergetic protons were simulated with the full nucleus model and scored in the SDD (Standard for DNA Damage) format for analysis (Schuemann *et al* 2019b). The nucleus was placed at the center of a cubic world volume with a side length of 14 μm, and the primary protons were sampled on the surface of the nucleus and directed inside the nucleus with a random direction to mimic the random relative orientation between the incident particle and chromatin fiber in experimental irradiations. If particles or radicals stepped outside the world boundary, they were considered to be terminated and were not reflected back into the world volume.

To achieve acceptable statistical accuracy, 100 runs were performed for each simulated particle energy. Each simulation run deposited a dose of ~1 Gy within the nucleus and the average number of histories to deposit a dose of 1 Gy for different energies was listed in table 2. The statistical uncertainties of DSB yields

were smaller than 2% (except for simulations with 10 ns chemical stage length, which were very time consuming, and the statistical uncertainty was limited to 5% for these simulations).

2.2. Physics and chemistry models in TOPAS-nBio

Most of the interaction models and cross sectional data used in MC track structure tools were derived based on water, and were considered as a near accurate description for biological tissue which is mostly water. The simulation of DNA damage includes three stages, starting from interactions in the physical stage, which lasts about 1 femtosecond. In this stage, primary and secondary particles interact with the DNA volume according to predefined interaction processes within the limits of the energy thresholds of each modeled physics process. The Geant4 (Geant4-DNA) release (version 10.5) used for this analysis provides nine physics constructors: G4EmDNAPhysics (default) and G4EmDNAPhysics_option1-8 for simulations at the DNA scale. These physics constructors are based on three main constructors (G4EmDNAPhysics_option2 (opt2), G4EmDNAPhysics_option4 (opt4), and G4EmDNAPhysics_option6 (opt6)) that are recommended for simulations in liquid water (Incerti *et al* 2018). Opt2 is an accelerated version of the default physics constructor; it uses a cumulated version of differential cross sections to describe the physical interactions while the default constructor uses non-cumulated differential cross sections. Opt4 includes an updated set of cross sections for electron-impact excitation and ionization in liquid water and provides an alternative elastic scattering model. Opt6 is an implementation of the CPA100 track structure code to Geant4-DNA. More details about different physics constructors in Geant4-DNA can be found elsewhere (Bordage *et al* 2016; Bernal *et al* 2015, Tran *et al* 2015, Incerti *et al* 2010b, 2018). All of these physics constructors are available in TOPAS-nBio.

Besides the physics constructors available in Geant4-DNA, TOPAS-nBio provides a flexible physics constructor, TsEmDNAPhysics, which uses the same physics processes and associated models as G4EmDNAPhysics. However, users can easily modify the interaction models in TsEmDNAPhysics constructor with a few command lines in the parameter file.

The pre-chemical stage follows the physical stage with a time scale of a few picoseconds and is followed by the chemical stage that can last up to microseconds. In the pre-chemical stage, the chemical species are generated by ionized or excited water molecules via dissociative decay or auto-ionization, and then the initial chemical species diffuse and react in the chemical stage. As mentioned before, potential reflections of chemical species at the geometry boundary have not been considered as this process is not currently supported in Geant4-DNA and TOPAS-nBio. Geant4-DNA provides a chemistry model (G4EmDNAChemistry) to describe the production, diffusion, and interaction of all chemical species (Karamitros *et al* 2014). TOPAS-nBio provides a chemistry model (TsEmDNAChemistry) to simulate the production of water radiolysis and the subsequent chemical interactions in the pre-chemical and chemical stage with reviewed and updated parameters, including the electron solvated thermalization distance, diffusion coefficients, and reaction rates in the G4EmDNAChemistry model. The updated values of parameters were adopted from RITRACKS (Plante and Devroye 2017), and this update was recently included in Geant4-DNA as G4EmDNAChemistry_option1 (Shin *et al* 2019). The TsEmDNAChemistry model was validated by calculating time-dependent G-values (number of molecular species per 100 eV of energy deposit) for electrons, protons and α particles covering a range of LET (0.05–230 keV μm^{-1}) (Ramos-Méndez *et al* 2018).

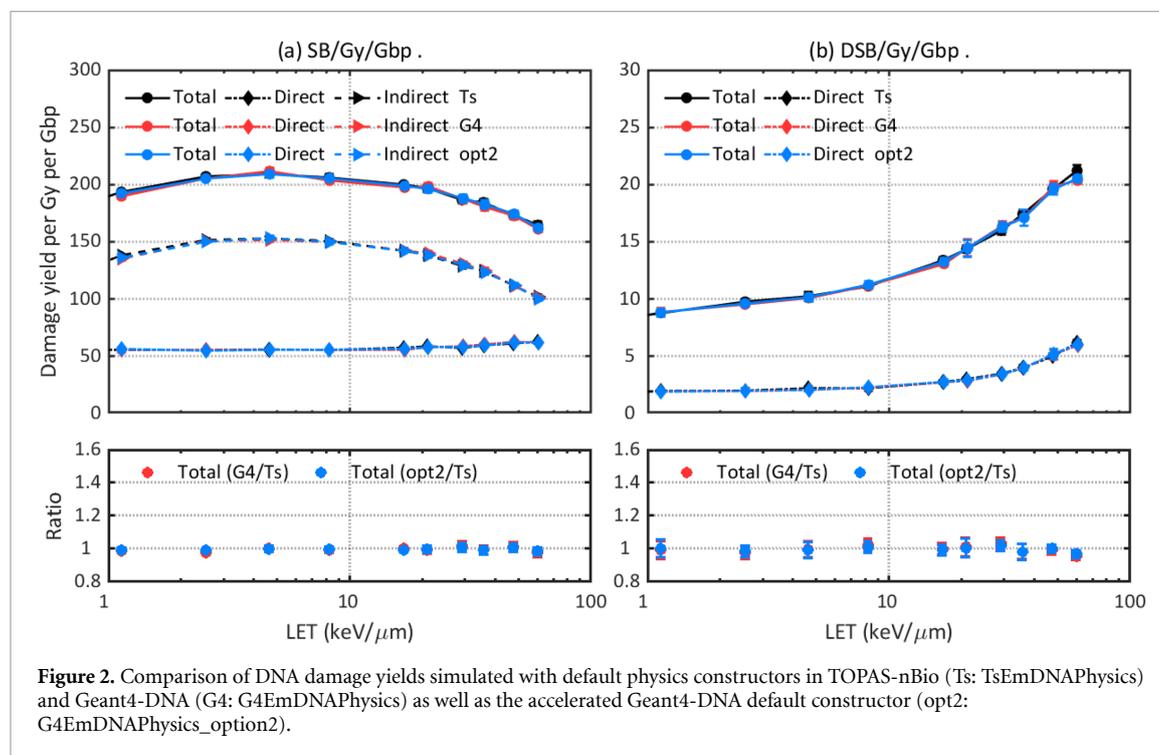
2.3. Induction of DNA damage

Initial DNA damages caused by the physical and chemical interactions within the nucleus were quantified as SBs in this work. SBs may result from the direct interactions of primaries and secondaries with the backbone volume and neighboring hydration shell or the indirect interactions of the chemical species with the backbone volume. Following our previous work (Zhu *et al* 2019), all radiolysis products and their interactions were simulated. The table of considered chemical species and their reactions can be found in the reference (Ramos-Méndez *et al* 2018). When chemical species diffuse into DNA volumes (base, backbone, hydration shell, and histone), they were considered to react and be scavenged by the DNA-related volumes. However, only interactions between the hydroxyl ($\cdot\text{OH}$) and the DNA backbone were regarded to be able to cause an indirect SB. No radical kill radius was applied in this work.

If two SBs were located on the opposite sides of the DNA double helix strand and separated by less than 10 base pairs, it was assumed to be a double strand break (DSB). Otherwise, the two SBs were classified as two single strand breaks (SSBs). DSBs were further classified as direct, indirect or hybrid DSBs according to the process that initiated SBs comprising the DSB. Following other studies, the DNA damages were counted after each history, i.e. inter-track damages were not considered assuming that proton tracks act

Table 3. Parameter set investigated in this study.

Parameter type	Value for sensitivity study	Default value
Physics constructors	G4EmDNAPhysics_option2/4/6	TsEmDNAPhysics
Chemistry model	G4EmDNAChemistry, TsEmDNAChemistry	TsEmDNAChemistry
Direct damaged threshold	17.5 eV, 5–37.5 eV	17.5 eV
Chemical stage length	1 ns, 2.5 ns, 10 ns	1 ns
·OH damage probability	$P_{\cdot\text{OH-backbone}} = 0.4$, $P_{\cdot\text{OH-backbone}} = 0.65$	$P_{\cdot\text{OH-backbone}} = 0.4$



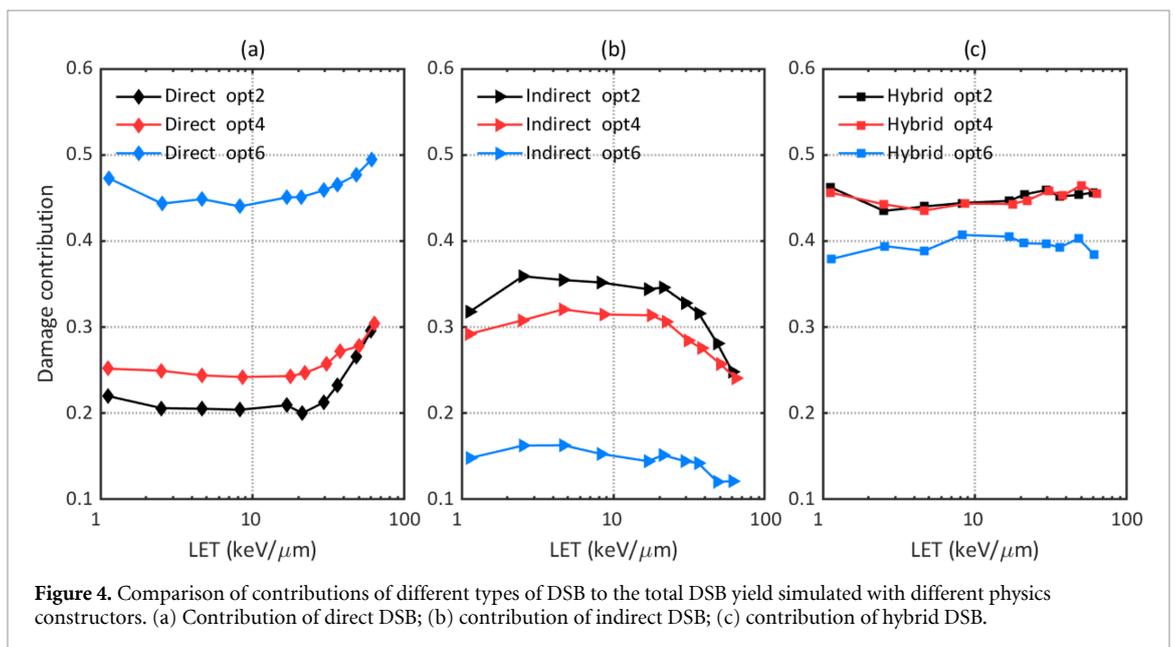
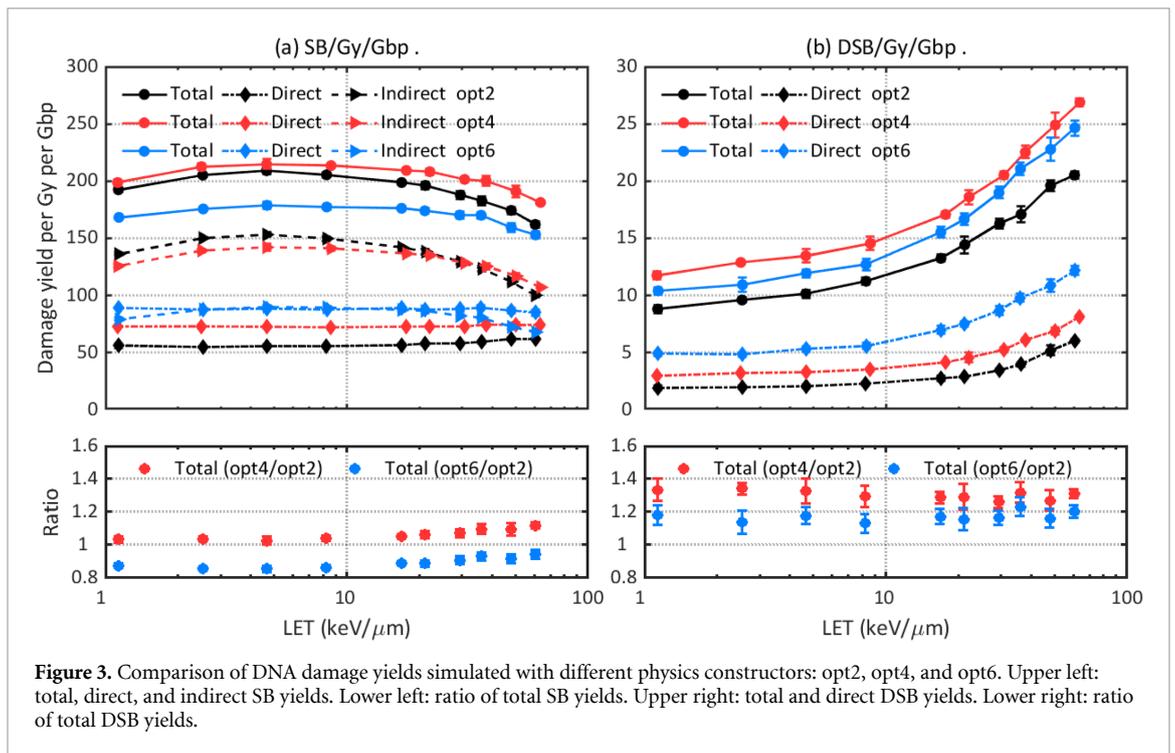
independently (Meylan *et al* 2017; Henthorn *et al*, 2017). The complexity of DSBs was not considered in this analysis.

2.4. Sensitivity study for parameters involved in DNA damage simulations

Differences in physics and chemistry models as well as the selection of other relevant parameter values across different studies can cause uncertainties in the DNA damage modeling. In this investigation, sensitivity studies were performed with TOPAS-nBio for physics and chemistry models, and three key parameters as listed in table 3. For each parameter, DNA damages were simulated with the values in the second column in table 3 while values for other parameters were fixed at the default values (third column in table 3). For instance, for the sensitivity study for physics constructors, DNA damages were simulated with opt2, opt4, and opt6, while for all of these simulations, the chemistry model was TsEmDNAChemistry, and values of direct damage threshold, chemical stage length and ·OH damage probability were fixed at thresholds of 17.5 eV, 1 ns and $P_{\cdot\text{OH-backbone}} = 0.4$, respectively.

3. Results and discussion

DNA damage results were simulated with three recommended physics constructors (opt2, opt4, and opt6) and default physics constructors in Geant4-DNA (G4EmDNAPhysics) as well as TOPAS-nBio (TsEmDNAPhysics). Figure 2 shows that TsEmDNAPhysics (Ts), G4EmDNAPhysics (G4), and opt2 give statistically equivalent damage yields. Changing to other physics constructors resulted in different total damage yields and different contributions from direct or indirect damage. Opt4 gave the highest total SB yield while opt6 gave the highest direct SB yield but the lowest indirect SB yield. Figure 3(b) shows that opt4 and opt6 caused up to 34% and 23% more DSBs, respectively, than opt2, and figure 4 shows that contributions of direct, indirect and hybrid DSBs to the total DSBs were also different when simulating with different physics constructors.



The main reason for the differences is the adopted process models and cross section data for electrons (an overview of these model differences can be found in (Incerti *et al* 2018; Bernal *et al* 2015)). Dose point kernel (DPK) simulations with different physics constructors were considered as benchmarks for the accuracy of electron elastic and inelastic scattering models (Incerti *et al* 2018). DPK is the average energy deposited per unit distance along the radius of a sphere with its center being the starting point of a particle source, and it describes spatial distribution of energy transferred from the particle source to the medium. The comparison of DPK obtained with different physics constructors showed opt2 is the most diffusive model (i.e. it has longer tail toward large radius values) while opt4 and opt6 deposit energies more concentrated in space (Kyriakou *et al* 2016, Bordes *et al* 2017, Shin *et al* 2018). A more diffuse model results in less instant recombination of initial pre-chemical and chemical species. Shin *et al* (2019) showed that opt2 simulates a more spread out distribution of electron interactions. This produces $\cdot\text{OH}$ radicals sufficiently separated to lead to higher G-values at the initiation of the chemical stage. Hence, more interactions between $\cdot\text{OH}$ and the DNA structure are expected, and our results show opt2 indeed gave higher indirect SB yield.

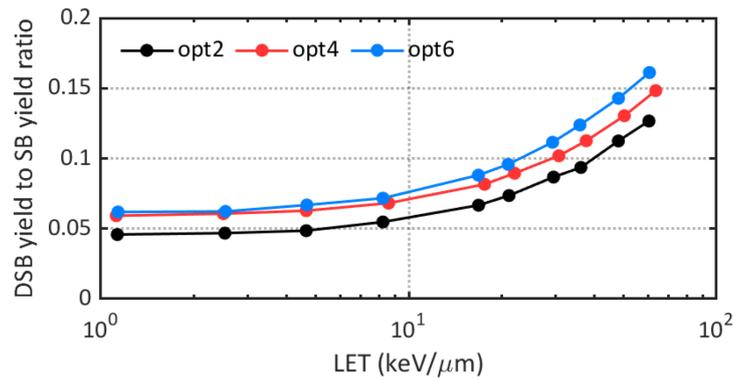


Figure 5. Comparison of DSB yield to SB yield ratios obtained with different physics constructors.

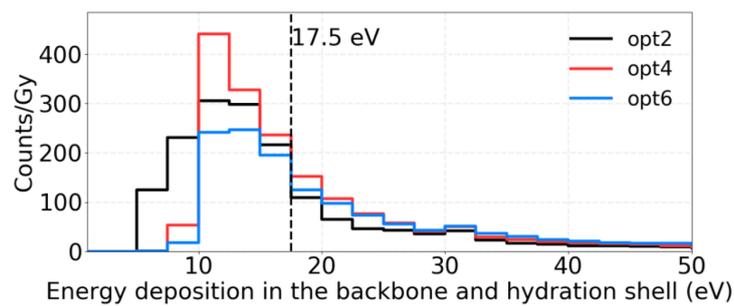


Figure 6. Energy deposition histograms of secondary electrons from 0.5 MeV primary protons in the backbone and hydration shell simulated with different physics constructors. The counts were normalized to the dose to the nucleus.

Table 4. Counts of energy depositions of secondary electrons from 0.5 MeV primary protons simulated with different physics constructors. The counts were normalized to the dose to the nucleus.

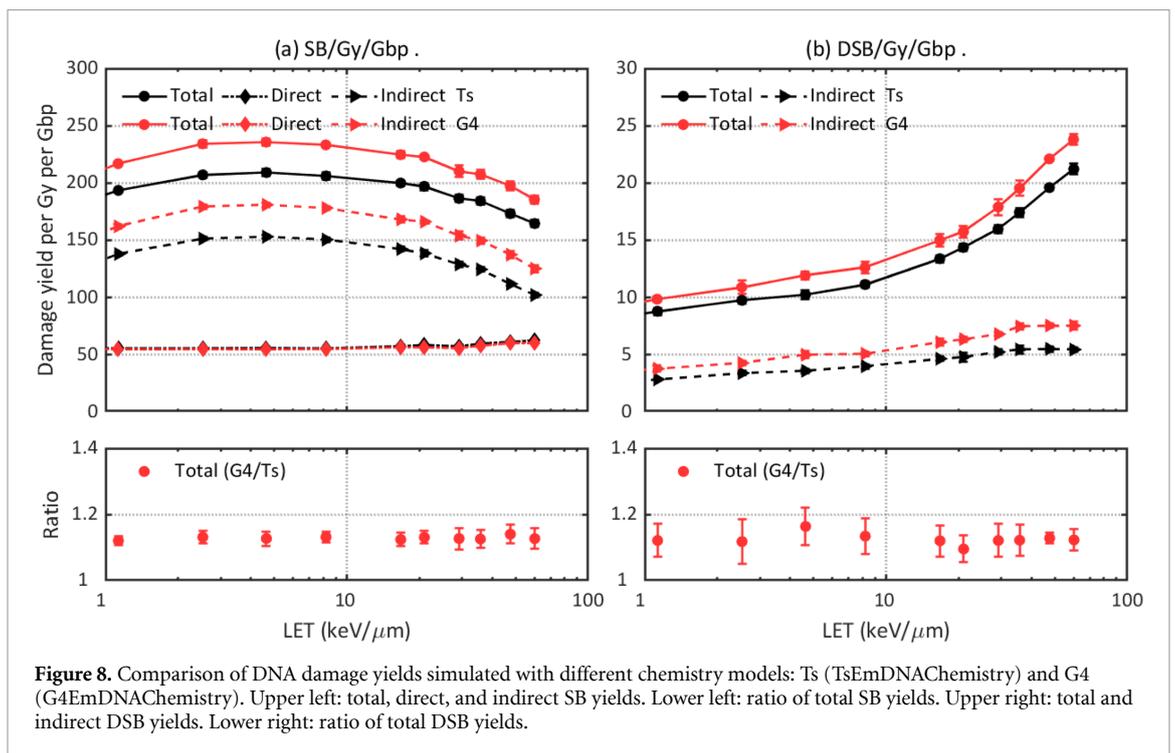
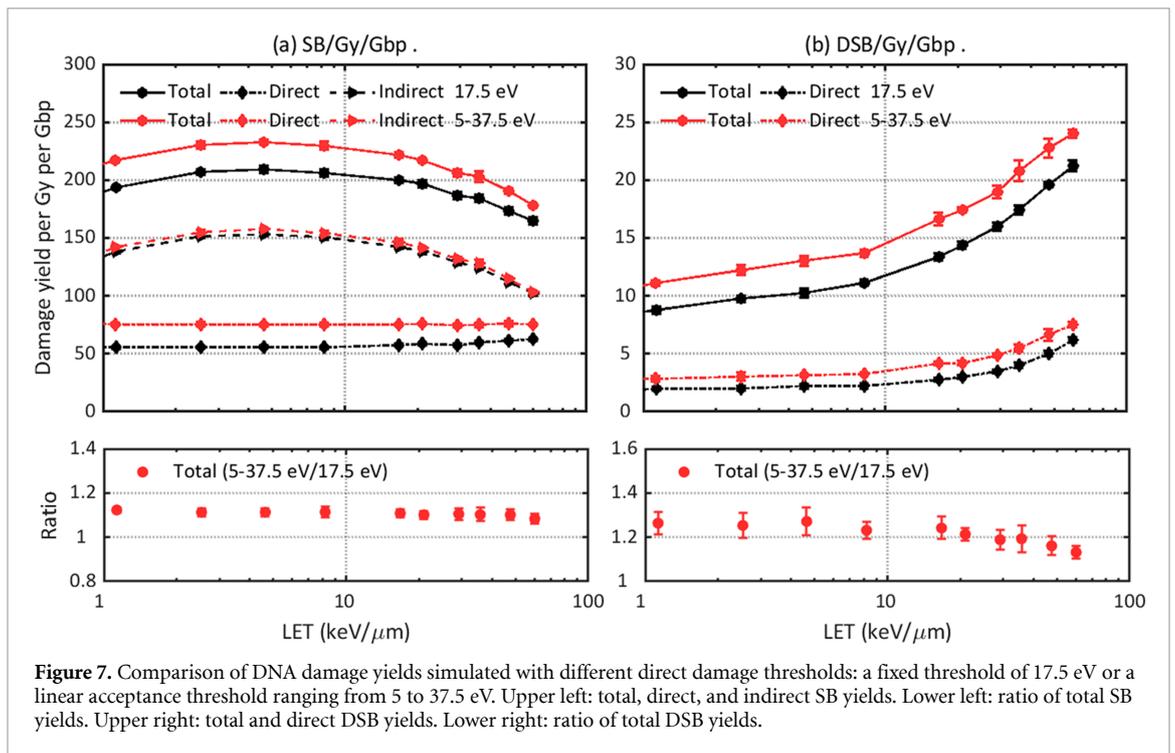
	Opt2	Opt4	Opt6
Counts/Gy (>0 eV) in the whole nucleus	1650.6	1729.8	1401.4
Counts/Gy (>17.5 eV) in the backbone and hydration shell	474.9	671.0	700.1

Figure 5 shows the total DSB yield to total SB yield ratios for different constructors, and opt2 gave the lowest DSB yield to SB yield ratios for the whole LET region, which is consistent with the diffusive characteristic of opt2 as SBs should be located close enough to form a DSB.

Our previous calculations (Zhu *et al* 2019) showed that more than 80% of the direct DNA damages after proton irradiation were caused by secondary electrons. Figure 6 and table 4 show counts of energy depositions of secondary electrons from 0.5 MeV primary protons in the backbone and hydration shell simulated with different physics constructors. In this case, opt6 resulted in the highest number of energy deposition events beyond the 17.5 eV damage threshold, which leads to the highest direct SB yield. On the other hand, opt2 and opt4 gave higher total counts of energy depositions, which then led to more excited water molecules (H_2O^*) and ionized water molecules (H_2O^+) in the pre-chemical stage, subsequently resulting in more radiolysis products that could lead to indirect damages.

Accordingly, another critical parameter that could affect the DNA damage results in the physical stage is the energy threshold to produce direct damage. Figure 7 compares the DNA damage results simulated with different energy thresholds: a fixed threshold of 17.5 eV and a linear acceptance threshold ranging from 5 to 37.5 eV. In the configuration of this work, up to 16% more SBs and 28% more DSBs were observed with the linear acceptance threshold. This trend is consistent with other works (Meylan *et al* 2017, Lampe *et al* 2018a, Sakata *et al* 2019).

Following the physical stage, the ionized and excited water molecules will produce initial chemical species through dissociative decay or auto-ionization leading to chemical interactions. The intrinsic difference between the Geant4-DNA and TOPAS-nBio chemistry models results in different G-values and hence



indirect damage yields. Figure 8 shows the difference of DNA damage simulated with G4EmDNACchemistry and TsEmDNACchemistry. The G4EmDNACchemistry model results in ~10% more SBs and DSBs than the TsEmDNACchemistry model in the investigated LET region. This is because only hydroxyl radicals were assumed to be able to cause indirect damages and the yield of hydroxyl radicals was affected by the diffusion and reaction rates in the two models. TsEmDNACchemistry produced ~10% fewer hydroxyl molecules than G4EmDNACchemistry at 1 ns (Ramos-Méndez *et al* 2018), thus resulting in fewer indirect damages.

The impact of the time cut of the chemical stage and the probability of hydroxyl radicals to induce DNA damage are shown in figures 9 and 10, respectively. One expects a higher damage yield with a longer chemical stage length because with a longer time chemical species will travel further before they were terminated,

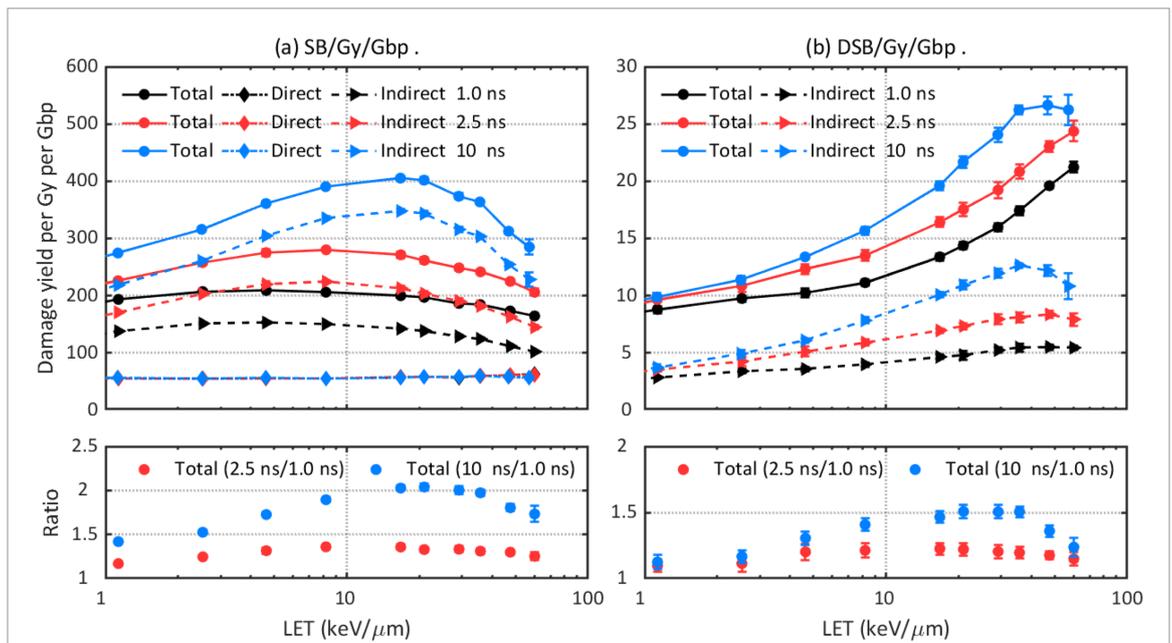


Figure 9. Comparison of DNA damage yields simulated with different chemical stage lengths: 1 ns, 2.5 ns, and 10 ns. Upper left: total, direct, and indirect SB yields. Lower left: ratio of total SB yields. Upper right: total and indirect DSB yields. Lower right: ratio of total DSB yields.

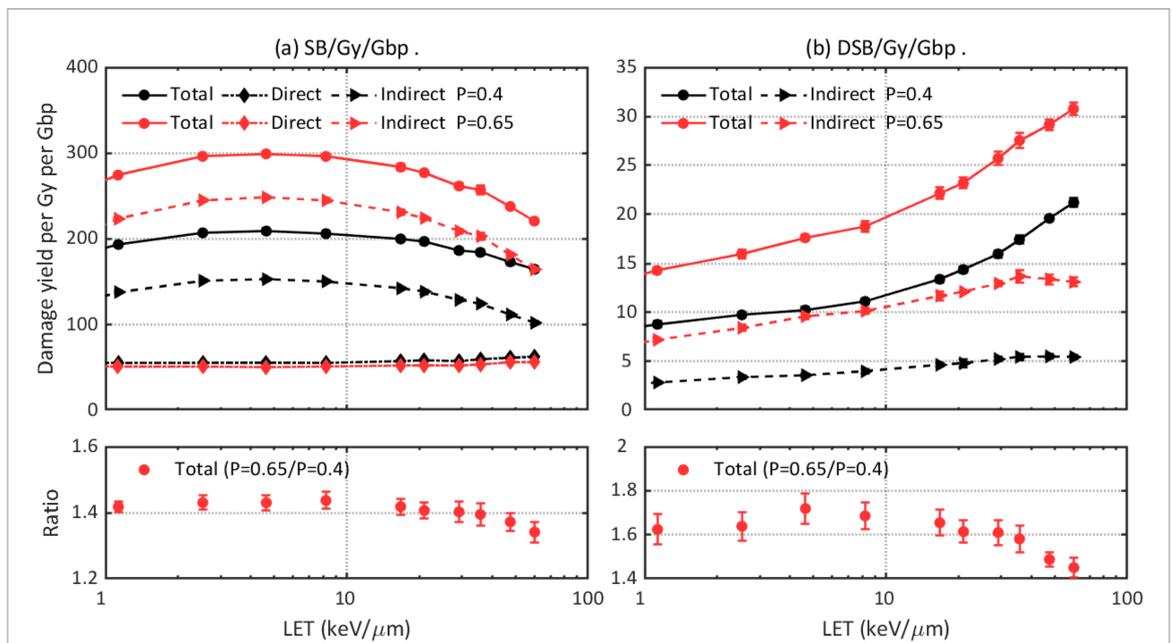


Figure 10. Comparison of DNA damage yields simulated with different hydroxyl damage probabilities: 0.4 and 0.65. Upper left: total, direct, and indirect SB yields. Lower left: ratio of total SB yields. Upper right: total and indirect DSB yields. Lower right: ratio of total DSB yields.

leading to a higher probability to interact with the DNA volume and a higher indirect damage yield. As shown in figure 9, the 2.5 ns long chemical stage results in up to 36% more SBs and 23% more DSBs than the 1 ns long chemical stage. The 10 ns long chemical stage produces up to 104% more SBs and 51% more DSBs than the 1 ns long chemical stage. As DNA volumes were the only scavengers in the simulation and the typical time scale for scavenging hydroxyl in cellular media is on the order of 1 ns (Lampe *et al* 2018a), a 1 ns long chemical stage length may be preferable for DNA damage simulations. As for the hydroxyl damage probability, it is shown that up to 44% more SBs and 71% more DSBs are predicted when a higher damage probability was adopted.

4. Summary

In this work, we studied the impact of physics and chemistry models as well as three parameters on the DNA damage yields after proton irradiation with a whole nucleus model implemented in TOPAS-nBio. This study aims at providing an estimation for the relative change in DNA damage yields, rather than absolute values, when different models and parameter values are adopted. However, the differences caused by varying DNA geometry configurations and simulation tool on damage results were not evaluated in this work.

The difference in physics constructors intrinsically affect the physical stage interactions, which will sequentially affect the production and interaction of radiolysis in the subsequent pre-chemical and chemical stages. Our simulations show that using G4EmDNAPhysics_option4 results in the highest DNA damage yields and caused up to 34% more DSBs than the G4EmDNAPhysics_option2 model. The G4EmDNAChemistry model predicted ~16% higher DSB yields than the TsEmDNAChemistry model due to different diffusion coefficients and reaction rates. Differences in direct damage threshold, chemical stage length, and hydroxy damage probability can cause differences of up to 28%, 51%, and 71% in predicted DSB yields, respectively, for the DNA geometry and irradiation modalities investigated here.

This study highlights the uncertainties in DNA damage simulations caused by differences in simulation parameters and demonstrates the impact when comparing DNA damage yields from different simulation studies. However, the uncertainty due to the geometry is not considered in this work. A recent study shows that different chromatin compaction forms lead to different DNA damage yields and a change in the contributions from direct and indirect reactions (Tang *et al* 2019). In addition, simulation performance critically depends on the adopted cross sections which are prone to uncertainties arising from physics models. For example, it was shown that including the exchange and correlation effects in inelastic electron scattering at low energies (below 10 keV) lead to better agreement with experimental cross section data (Emfietzoglou *et al* 2013, 2017). The underlying approximations made in cross section calculations may have a sizeable effect on DNA damage calculation. As a result, comparisons of relative effect sizes may be much more stable than absolute predictions as many of the above mentioned uncertainties cancel out.

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